

The interaction function and lattice duals

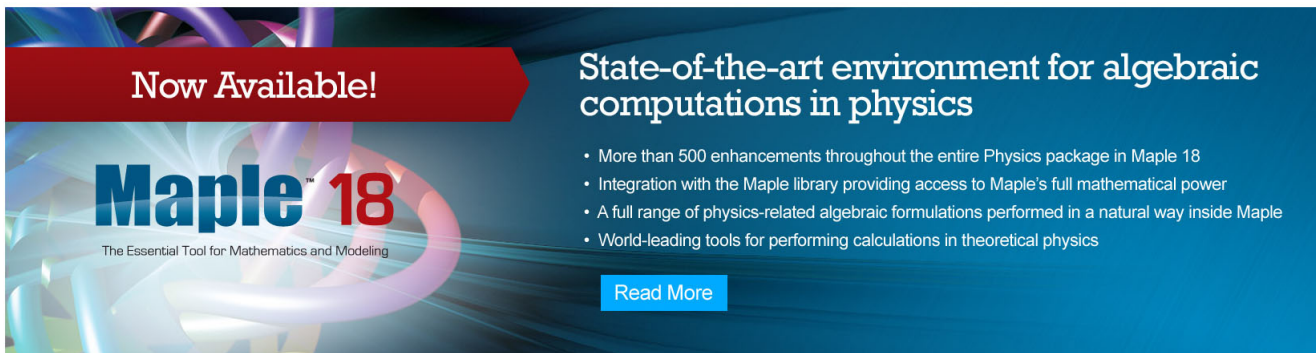
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Citation: *Journal of Mathematical Physics* **18**, 1985 (1977); doi: 10.1063/1.523156

View online: <http://dx.doi.org/10.1063/1.523156>

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The interaction function and lattice duals

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(Received 29 April 1975)

An interaction function is defined for lattice models in statistical mechanics. A correlation function expansion is derived, giving a direct proof of the duality relations for correlation functions.

A general theory of duality transformations between pairs of classical spin- $\frac{1}{2}$ lattice models has been developed by Gruber and Merlini¹ and independently by Wegner.² The theory of Gruber and Merlini is constructive, providing explicitly a family of "dual" lattices and Hamiltonians for any given spin- $\frac{1}{2}$ system. These duals are exact, all requisite boundary terms being provided for, which is necessary in considerations of correlation functions below criticality.

We define in this article the interaction functions $u_{H^*}(A, B)$ of lattice duals G and G^* , and express them in terms of correlation functions. This gives an easy derivation of the relationship between correlation functions of a lattice and its duals. The notation in this article, while somewhat different from Ref. 1 and some current usage, has the advantage, in addition to simplifying the derivations, of generalizing to higher spin lattices.³ The reader is referred to Ref. 1 for details on the construction of dual spin- $\frac{1}{2}$ lattices.

1. DUAL LATTICES

We suppose we are given a finite set Λ of lattice sites in a ν -dimensional space, along with a Hamiltonian H defined on the configuration of Λ . It is convenient to take as the configuration space the group $P_2(\Lambda)$ of functions from Λ to Z_2 , the integers modulo 2, with group multiplication

$$fg(\lambda) = f(\lambda) + g(\lambda) \pmod{2}.$$

Considering H as a function $H: P_2(\Lambda) \rightarrow \mathbb{C}$, its Fourier decomposition

$$H(g) = \sum_{\sigma \in \hat{G}} H_\sigma \sigma(g), \quad g \in P_2(\Lambda)$$

in terms of the elements of the character group \hat{G} of $G = P_2(\Lambda)$ is just the usual decomposition of H into a sum of products of spin matrices, since the characters of G are products of characters of Z_2 . Define the set of nonzero interactions

$$B = \{\sigma \in G \mid H_\sigma \neq 0\}.$$

Dual lattices are constructed with the set B . Defining $P_2(B)$ as the group of functions from B to Z_2 , let p be the group homomorphism

$$p: P_2(B) \rightarrow G \quad \text{by} \quad p(f) = \prod_{\sigma \in B} f(\sigma)$$

and denote its kernel by K_p . Suppose X is any set which generates K_p as a group. Then X defines a dual of Λ , with configuration space $G^* = P_2(X)$ and dual Hamiltonian H^* defined as follows. Let

$$q: B \rightarrow B^* \subset G^* \quad \text{by} \quad q(\sigma): h \rightarrow h(\sigma)$$

for $\sigma \in B$ and $h \in X \subset P_2(B)$. $q(\sigma)$ is indeed a character on G^* , and these $q(\sigma)$ are to be the nonzero interactions of the dual. The coefficients $H_{q(\sigma)}$ are given by

$$H_{q(\sigma)} = \frac{1}{2} \beta \log \prod_{\substack{\sigma' \in B \\ q(\sigma') = q(\sigma)}} \tanh \beta H_{\sigma'}, \quad (1)$$

and

$$H^* = \sum_{q(\sigma)} H_{q(\sigma)} q(\sigma).$$

In most models of physical interest, q is one-one, except perhaps near the boundary. Thus

$$H_{q(\sigma)} = \frac{1}{2\beta} \log \tanh \beta H_\sigma$$

except near the boundary, where (1) must be used.

The partition functions $Z(\beta H) = \sum_{g \in G} \exp(-\beta H(g))$ of G and $Z(\beta H^*) = \sum_{g^* \in G^*} \exp(-\beta H^*(g^*))$ of its duals G^* are related then by

$$Z(\beta H) = \frac{N(K_t^*)}{N(G)} \prod_{\sigma \in B} [\sinh(-\beta H_\sigma) \cosh(-\beta H_\sigma)]^{1/2} Z(\beta H^*)$$

where $N(S)$ is the cardinality of S , and K_t^* is defined after Eq. (2).

2. THE INTERACTION FUNCTION

The correlation functions $\rho(\sigma)$ of G are defined by

$$\rho(\sigma) = Z(\beta H)^{-1} \sum_{g \in G} \exp(-\beta H(g)) \sigma(g), \quad \sigma \in \hat{G}$$

with H^* replacing H for the correlation functions $\rho(\sigma^*)$ of G^* , $\sigma^* \in \hat{G}^*$. Note that $\rho(\sigma) = 0$ if σ is not a product of elements of B .¹

Define the characteristic projection $t: G^* \rightarrow P_2(B^*)$ by

$$t(g^*): \sigma \rightarrow \frac{1}{2}(1 - \sigma(g^*)), \quad \sigma \in B^*. \quad (2)$$

The support of $t(g^*)$ is precisely those characters $\sigma \in B^*$ whose value at g^* is -1 . Now if the kernel and range of t are denoted, respectively, by K_t^* and R_t^* , then the map $Q: K_p \rightarrow P_2(B^*)$ given by

$$Q(f)(q(\sigma)) = f(\sigma), \quad \sigma \in B,$$

is a group isomorphism $K_p \rightarrow R_t^*$. In particular, $f \in K_p$,

$$f(\sigma) = \begin{cases} 1, & \sigma \in S, \\ 0, & \sigma \in Y - S, \end{cases} \quad \text{if and only if } Q(f) \in R_t^*,$$

$$Q(f)(q(\sigma)) = \begin{cases} 1, & \sigma \in S, \\ 0, & \sigma \in Y - S, \end{cases}$$

and then

$$\prod_{\sigma: f^{-1}(1)} \tanh(-\beta H_\sigma) = \prod_{\sigma \in Q(f)^{-1}(1)} \exp(2\beta H_\sigma).$$

Let the symbol $\sum_f(S, T)$ with $S, T \subset B$ indicate that the summation [over $f \in P_2(B)$, say] is to be restricted to f satisfying $f(\sigma) = 0$ if $\sigma \in S$, $f(\sigma) = 1$ if $\sigma \in T$. Then the interaction function $u_{H^*}(A, C)$ is given by

$$u_{H^*}(A, C) = \sum_{f \in R_t^*} \prod_{\sigma: f^{-1}(1)} \exp(2\beta H_\sigma)$$

for $A, C \subset B^*$.

We wish to evaluate u_{H^*} in terms of the correlation functions of G^* . Thus, suppose Y and W are any disjoint subsets of B^* . Writing \bar{Y} for $\pi\sigma$, $\sigma \in Y$, etc., obtain from (2):

$$\begin{aligned} & \left(\prod_{\sigma \in B^*} e^{-\beta H_\sigma} \right)^{-1} Z^* \rho(\overline{Y \cup W}) \\ &= \sum_{g \in G^*} \prod_{\sigma \in B^*} e^{-\beta H_\sigma (g(\sigma)-1)} \prod_{\sigma' \in Y \cup W} g(\sigma') \\ &= N(K_t^*) \sum_{f \in R_t^*} \prod_{\sigma: f^{-1}(1)} e^{2\beta H_\sigma} \prod_{\sigma' \in Y} (-2f(\sigma') + 1) \\ & \quad \times \prod_{\sigma'' \in W} (2 - 2f(\sigma'') - 1). \end{aligned}$$

Now, expanding the product

$$\prod_{\sigma' \in Y} (-2f(\sigma') + 1) = \sum_{L \subset f^{-1}(1) \cap Y} (-2)^{N(L)},$$

and similarly with $\prod_{\sigma'' \in W} (2 - 2f(\sigma'') - 1)$, this becomes

$$N(K_t^*) \sum_{\substack{L \subset Y \\ M \subset W}} u_{H^*}(M, L) (-2)^{N(L)+N(M)} (-1)^{N(W)}.$$

Therefore, with a change in summation variable,

$$\begin{aligned} & \left(\prod_{\sigma \in B^*} e^{-\beta H_\sigma} \right)^{-1} Z^* \sum_{\substack{Y \subset C \\ W \subset A}} (-1)^{N(Y)} \rho(\overline{Y \cup W}) \\ &= N(K_t^*) \sum_{\substack{L \subset C \\ M \subset A}} \sum_{\substack{Z \subset C-L \\ V \subset A-M}} u_{H^*}(M, L) 2^{N(L)+N(M)} (-1)^{N(Z)+N(V)} \\ &= 2^{N(A)+N(C)} N(K_t^*) u_{H^*}(A, C), \end{aligned} \quad (3)$$

which gives the desired expression.

3. DUAL CORRELATION FUNCTIONS

The interaction functions can be used to derive directly the duality relations for correlation functions. Let $Y \subset B$. Then, using

$$\exp(-\beta H_\sigma(g)) = \cosh(-\beta H_\sigma) + \sigma(g) \sinh(-\beta H_\sigma)$$

and the orthonormality of the characters,

$$Z\rho(\bar{Y}) = N(G) \prod_{\sigma \in B} \cosh(-\beta H_\sigma) \sum_{\substack{f \in P_2(B) \\ \rho(f)=\bar{Y}}} \prod_{\sigma' \in f^{-1}(1)} \tanh(-\beta H_{\sigma'}).$$

From the one-one correspondence between $f \in P_2(B)$ with $\rho(f) = \bar{Y}$ and $f' \in K_p$ with

$$f': \sigma \rightarrow \begin{cases} f(\sigma), & \text{if } \sigma \notin Y, \\ f(\sigma) + 1, & \text{if } \sigma \in Y, \end{cases}$$

the expansion can be written as

$$\begin{aligned} & Z \left(N(G) \prod_{\sigma \in B} \cosh(-\beta H_\sigma) \right)^{-1} \rho(\bar{Y}) \\ &= \sum_{S \subset Y} \sum_{\substack{f \in P_2(B) \\ \rho(f)=\bar{Y}}} \sum_{\sigma' \in f^{-1}(1)} \prod_{\sigma' \in f^{-1}(1)} \tanh(-\beta H_{\sigma'}) \\ &= \sum_{S \subset Y} \sum_{f \in K_p} \sum_{\sigma' \in f^{-1}(1)} \prod_{\sigma' \in f^{-1}(1)} \tanh(-\beta H_{\sigma'}) \prod_{\sigma \in S} [\tanh(-\beta H_\sigma)]^{-1} \\ & \quad \times \prod_{\sigma \in Y-S} \tanh(-\beta H_\sigma) \\ &= \sum_{\substack{S \subset Y \\ S^* \cap (Y-S)^* = \emptyset}} u_{H^*}(Y^* - S^*, S^*) \prod_{\sigma \in S} (\tanh(-\beta H_\sigma))^{-1} \\ & \quad \times \prod_{\sigma \in Y-S} \tanh(-\beta H_\sigma), \end{aligned}$$

where it has been necessary to consider in the sum over S only sets $S \subset Y$ for which $S^* = \{q(\sigma) \mid \sigma \in S\}$ and $(Y-S)^*$ are disjoint. Thus, the interaction function expansion (3) gives the general relation between the correlation functions of G and the correlation functions of a dual G^* ,

$$\rho_G(\bar{Y}) = \sum_{T^* \subset W^*} \rho_{G^*}(\bar{T}^*) K(W, T^*), \quad (4)$$

where

$$\begin{aligned} K(W, T^*) &= 2^{-N(W^*)} \sum_{\substack{S \subset W \\ S^* \cap (W-S)^* = \emptyset}} (-1)^{N(S^* \cap T^*)} \\ & \quad \times \prod_{\sigma \in S} (\tanh(-\beta H_\sigma))^{-1} \prod_{\sigma \in W-S} \tanh(-\beta H_\sigma) \end{aligned}$$

for any $W \subset B$ such that $\bar{W} = \bar{Y}$.

In the event that the duality map q is one-one, Eq. (4) simplifies to the path formula of Kadanoff and Ceva.¹ Injectivity of q is equivalent to requiring that the elements of K_p separate the bonds σ of B , and is satisfied, for example, by a hexagonal Ising lattice with periodic boundary conditions, or with an external field at the boundary, but is not satisfied by this lattice with open boundary conditions.

¹D. Merlini and C. Gruber, J. Math. Phys. **13**, 1814 (1972).

²F. Wegner, J. Math. Phys. **12**, 2259 (1971).

³W. Greenberg, Commun. Math. Phys. **29**, 163 (1973).

⁴L. P. Kadanoff and H. Ceva, Phys. Rev. B **3**, 3918 (1971).