

1 **A Geodynamic Investigation of Plume-Lithosphere Interactions Beneath the East African**
2 **Rift**

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22 **Key Points:**

- 23 • Geodynamic modeling suggests observed rift parallel motions influenced by the African
24 Superplume
- 25 • Asthenospheric seismic anisotropy induced by northward flow associated with the
26 African Superplume
- 27 • Horizontal mantle tractions driving surface deformation

28

29 **Abstract**

30 Constraining the origin of forces that drive continental rifting remains a highly debated and
31 unresolved question within geodynamics. The East African Rift (EAR) provides an ideal natural
32 laboratory to examine the relative role of plate driving forces as only lithospheric buoyancy
33 forces and horizontal mantle tractions act on the system. Here, we employ high-resolution 3D
34 thermomechanical models to test whether: 1) anomalous, rift-parallel surface deformation
35 observed by GNSS data in the EAR are driven by viscous coupling to northward mantle flow
36 associated with the African superplume, and 2) the African superplume is the dominant source
37 mechanism of anomalous rift-parallel seismic anisotropy beneath the EAR. We calculate Lattice
38 Preferred Orientations (LPO) and surface deformation from two types of mantle flow: (1) a
39 scenario with multiple plumes constrained by shear wave tomography and (2) a single
40 superplume model with northward boundary condition to simulate large-scale flow. Comparison
41 of calculated LPO with observed seismic anisotropy, and surface velocities with GNSS and plate
42 kinematics reveal that there is a better fit with the superplume mantle flow model, rather than the
43 tomography-based (multiple plume) model. We also find a relatively better fit spatially between

44 observed seismic anisotropy and calculated LPO with the superplume model beneath northern
45 and central EAR, where the superplume is proposed to be shallowest. Our results suggest that the
46 viscous coupling of the lithosphere to northward mantle flow associated with the African
47 superplume drives most of the rift-parallel deformation and is the dominant source of the
48 observed LPO in the EAR. The model show that northward mantle flow associated with the
49 African superplume drives most of the rift-parallel deformation in the EAR and is the dominant
50 source of the observed LPO in the EAR.

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52 **Plain Language Summary**

53 What forces drive continental rifting remains an outstanding question of geodynamics. During
54 continental rifting, surface deformation and the underneath mantle flow are usually perpendicular
55 to the rift. In the case of the East African Rift (EAR), the largest continental rift on Earth, it has
56 been demonstrated that its E-W extension is mostly driven by forces due to its high-topography,
57 but some deformations are parallel northward to the rift. Here, we use 3-dimensional computer
58 simulation to test if these anomalous deformations are driven by hot buoyant upwelling known as
59 African Superplume. Comparison of model results with measured surface deformation, measured
60 with highly accurate Global Position System (GPS) measurements, show that mantle northward
61 associated with the African Superplume drive some of the rift parallel deformations in the EAR.
62 Our results also suggest that most of the enigmatic rift parallel seismic anisotropy, nowadays
63 observed by geophysicist beneath the EAR, can be explained by the large northward mantle flow
64 from the African Superplume.

65

66 **1. Introduction**

67 Continental rifting is an important geodynamic process during which the Earth's
68 lithosphere undergoes continuous stretching resulting in continental break-up and, ultimately, the
69 formation of new oceanic basins. In the past few decades, many geophysical studies have been
70 carried out to elucidate the driving forces of continental rifting in order to advance our
71 understanding of rift initiation and evolution (e.g. Mulugeta, 1985; Buck, 1991; Brune, 2018;
72 Glerum et al., 2020; Naliboff et al., 2020). The origin of extensional stresses responsible for
73 continental rifting can be classified into two categories: 1) horizontal tractions at the base of the
74 lithosphere arising from mantle convection (Froidevaux and Nataf, 1981; Lithgow-Bertelloni and
75 Guynn, 2004; Bird et al., 2008; Forte et al., 2010; Ghosh and Holt, 2012) and 2) variations in
76 lithospheric buoyancy forces arising from topography gradients and density variations (e.g. Bott
77 and Kuszniir, 1979; Lithgow-Bertelloni and Silver, 1998; Coblentz and Sandiford, 1994;
78 Lithgow-Bertelloni and Guynn, 2004; Ghosh et al., 2009; Naliboff et al., 2012; Stamps et al.,
79 2014, 2015; Rajaonarison et al., 2021a).

80 The EAR is region of continental extension within the diverging the Nubia-Somalia plate
81 system, with kinematic models derived from GPS data suggesting an ~E-W extension direction
82 (Calais et al., 2006; Saria et al., 2014; Stamps et al. 2021; Figure 1.1A). Previous studies suggest
83 E-W extension of the EAR is driven largely by lithospheric buoyancy forces attributed mainly to
84 the unusually high topography (higher than 1000 m above sea level) known as the African
85 Superswell (e.g., Nyblade and Robinson, 1994) of the EAR (Figure 1.1A) and to density
86 variations within the lithosphere (e.g. Coblentz and Sandiford, 1994; Stamps et al., 2014, 2015;
87 Rajaonarison et al., 2021a). However, other studies propose horizontal mantle tractions drive a
88 significant portion of the E-W extension in Africa (e.g. Ghosh and Holt, 2012; Kendall and

89 Lithgow-Bertelloni, 2016). Significantly, all of the aforementioned studies approximate
90 deformation within the lithosphere using a depth-independent 2D (i.e., thin shell or sheet)
91 modeling approach. In contrast, recent 3D thermomechanical modeling (Rajaonarison et al.,
92 2021a) suggests that the E-W extension of the EAR, (i.e., the rigid plates and microplates
93 rotation and their velocity magnitudes) is dominated by lithospheric buoyancy forces and that
94 additional forces, such as viscous coupling to mantle flow, are only needed to explain the
95 anomalous observed northward surface motions within the deforming zones as revealed by
96 GNSS/GPS velocities (Figure 1.1A, Stamps et al., 2018) and strike-slip focal mechanisms
97 observed along multiple branches of the rift (Figure 1.1B; Dziewonski et al., 1981; Ekström et
98 al., 2012).

99 Here, we expand on the work of Rajaonarison et al. (2021a) to test if plume-lithosphere
100 interactions are a plausible explanation for the rift-parallel deformation observed along the EAR,
101 where plume-lithosphere interactions have been suggested to be important in the system's long-
102 term evolution and dynamics (Koptev et al., 2016; Koptev et al., 2018a, b; Koptev, Calais et al.,
103 2018). Numerous seismic tomography studies have imaged low velocity anomalies beneath the
104 EAR and associated them with the presence of one or more upwelling thermal anomalies (e.g.,
105 Ebinger and Sleep, 1998; Montelli et al., 2006; Hansen et al., 2012; Bagley and Nyblade, 2013;
106 Emry et al., 2019; Chang et al., 2011, 2020). However, interpretations of these low velocity
107 anomalies in terms of plume structure remain controversial, with two main end-member plume
108 models (Figure 1C, D) often invoked. The first model attributes the low velocity anomalies
109 beneath the EAR as originating from multiple plumes (Figure 1C) beneath the EAR (e.g., Camp
110 and Roobol, 1992; Chang and Van der Lee, 2011; Ebinger and Sleep, 1998; George et al., 1998;
111 Montelli et al., 2006). The second model invokes the African superplume (Figure 1D), which

112 consists of a large single low velocity anomaly that penetrates the transition zone (Hansen et al.,
113 2012), originating from the core-mantle boundary beneath South Africa with upward and
114 northeastward flow beneath east Africa to Arabia (e.g. Bastow et al., 2008; Benoit et al., 2006;
115 Forte et al., 2010; Furman et al., 2006; Ritsema et al., 1999; Simmons et al., 2007, 2009). The
116 hypothesized northeastward direction of the mantle flow associated with the African superplume
117 is supported by consistent rift-parallel (NE-oriented) shear wave splitting measurements beneath
118 the EAR (Figure 1B; Bagley and Nyblade, 2013; Andriampemanana et al., 2021). Bagley and
119 Nyblade (2013) suggest that the NE-oriented azimuthal anisotropy is sourced from
120 asthenospheric Lattice Preferred Orientation (LPO) of olivine induced by northeastward mantle
121 flow associated with the African superplume.

122 In this work, we test the hypotheses that: 1) the observed northward component of
123 surface motion is driven by viscous coupling between the lithosphere and northward mantle flow
124 beneath the EAR, and 2) NE oriented azimuthal anisotropy beneath the EAR reflects
125 northeastward mantle flow associated with the African superplume. To test these hypotheses, we
126 use the open-source finite element code Advanced Solver for Problems in the Earth's
127 ConvecTion (ASPECT) (Bangerth et al., 2015, 2020; Kronbichler et al., 2012; Heister et al.,
128 2017) to simulate 3D convection and deformation of the lithosphere-asthenosphere system for
129 the EAR and surroundings. Our numerical experiments build upon the model from Rajaonarison
130 et al. (2021a), in which the deformation is driven by lithospheric buoyancy forces.

131 We incorporate sublithospheric mantle flow using thermal anomalies derived from the
132 conversion of shear wave tomography from Emry et al. (2019) to account for plume structures in
133 the upper 660 km. This shear wave model is used to represent a “multiple plumes” scenario, as it
134 contains multiple distinct thermal anomalies rather than a single large thermal anomaly (i.e.,

135 superplume) . A second numerical experiment builds on this model by adding an imposed
136 (boundary-drive) northward mantle flow beneath 200 km depth to simulate the proposed effects
137 of large-scale mantle flow induced by a superplume in this region. The mantle flow velocities
138 from each simulation are used to generate synthetic seismic anisotropy, which is compared
139 with observed azimuthal anisotropy from the EAR region. The synthetic seismic anisotropy is
140 calculated using D-Rex (Kaminski et al., 2004), which follows the kinematic model for
141 plastic deformation and dynamic recrystallization (Kaminski and Ribe, 2001, 2002; Ribe and Yu,
142 1991). To assess the influence of mantle flow on surface deformation, we quantitatively
143 compare surface motion from both models with GNSS observations from Stamps et al. (2018)
144 and block kinematic models from Stamps et al. (2021).

145 We find that the presence of multiple plumes or the superplume does not affect the rigid
146 block rotation of the Victoria and Rovuma Blocks, but produces faster velocities (up to twice in
147 velocity magnitude) for the Somalian Plate. We also find that the northward component of
148 motion along the Main Ethiopian Rift, the Western Branch, and the central Eastern branch can be
149 explained by viscous coupling of the lithosphere to northward mantle flow associated with the
150 African Superplume. The velocity angular misfit between predicted velocities and GNSS/GPS
151 observations in these regions improve from 66° , 114° , and 74° , respectively from a
152 model with deformation driven solely by lithospheric buoyancy forces (Rajaonarison et al.,
153 2021a), to 50° , 40° , and 30° , respectively. This improvement of the angular misfit
154 suggests that viscous coupling to northward mantle flow is favored within the deforming zones.
155 Our results also suggest that additional mechanisms, such as anisotropic viscosity in the
156 asthenosphere is needed to explain the rotation rate of the Somalian Plate. Although lithospheric
157 buoyancy forces dominate the force balance driving E-W extension across the EAR, this work

158 suggests horizontal tractions from northward mantle flow associated with the African
159 Superplume is needed to explain observations of rift-parallel surface motions in deforming zones
160 from GNSS/GPS data and northward oriented seismic anisotropy beneath the EAR.

161

162 **2 Methods**

163 **2.1 Model Design and Experiments**

164 We use the finite element code ASPECT (Advanced Solver for Problems in Earth's
165 ConvecTion; Heister et al., 2017; Kronbichler et al., 2012), which has been widely employed to
166 investigate melt/mantle dynamics (Dannberg et al., 2019; Dannberg & Heister, 2016; Njinju et
167 al., 2021), lithospheric deformation and continental extension (Glerum et al., 2020; Naliboboff et
168 al., 2020), and global/regional mantle convection (Rajonarison et al., 2020; Zhang & Li, 2018),
169 to simulate lithospheric deformation and mantle flow within the EAR region. Surface and
170 internal velocities are calculated with Stokes system of equations, which follows the Boussinesq
171 approximation for an incompressible Newtonian fluid. Although the aim of this study is to assess
172 the contributions of horizontal tractions from mantle flow beneath the EAR, it is important that
173 the primary driving forces acting in this region, which are lithospheric buoyancy forces and
174 horizontal mantle tractions, are incorporated in the model. Thus, our model is derived from
175 Rajonarison et al. (2021a,b) in which the lithospheric buoyancy forces are implemented using
176 ETOPO1 (Amante and Eakins, 2009) for the surface topography, CRUST1.0 for the laterally
177 varying crustal structures and densities (lower, middle, and upper crust; Laske et al., 2013), and
178 isostatically compensated to 100 km depth for the mantle lithosphere density.
179 Here, mantle flow fields driven by density heterogeneities are added in the sublithospheric
180 regions. The sublithospheric mantle densities in our simulations vary linearly with temperature:

$$181 \quad \rho = \rho_0(1 - \alpha(T - T_0)) \quad (1.1)$$

182 where α is the thermal expansion coefficient. The Stokes and temperature system of equations is
183 solved within a 3D spherical chunk geometry model domain, covering $5300 \times 3300 \times 660$ km in
184 the East, North, and radial directions, respectively.

185 The initial temperature structure throughout the lithosphere is calculated following the
186 approach of Chapman (1986), which uses an analytical solution for a conductive geothermal
187 profile to obtain temperatures throughout a layered lithosphere (upper crust, middle crust, lower
188 crust, and mantle lithosphere). A key assumption to this approach is that each layer contains a
189 constant thermal conductivity and radiogenic heat production, unique for each layer (Table B.1).
190 While the lithosphere-asthenosphere boundary is defined by the 1673 K isotherm, the
191 lithospheric geothermal gradient varies as a function of lithospheric thickness, the surface heat
192 flow, and the crustal thicknesses (CRUST1.0). We use lithospheric thicknesses that capture the
193 key tectonic regions obtained from averages of the lithospheric models LITHO1.0 Pasyanos
194 (2013), Fishwick (2010, updated), and Emry (per comm.). For the cratonic regions, including the
195 Tanzania Craton, the Congo Craton, the Bangweulu Block, and the Masai Block, the lithospheric
196 thicknesses are 150 km thick. The Eastern Branch and the Western Branch of the EARS have 70
197 km and 90 km thick lithosphere, respectively. We impose 50 km lithosphere thickness for
198 oceanic ridges. We also assume a 100 km thick lithosphere for mobile melts, and for the oceanic
199 lithosphere, and other regions not previously defined.

200 In sub-lithospheric regions, the temperature is the sum of an approximate adiabatic
201 temperature profile (with 0.3 K/km increase with depth; Schubert et al., 2009) and a
202 temperature anomaly obtained using a conversion of shear wave anomalies from Emry et al.
203 (2019), which is shown in Figure 3A,B. We choose a simple approach for the conversion of

204 shear wave velocities to temperature to minimize complexity in the model. The shear wave
205 anomalies are first converted to density perturbations $\delta\rho/\rho$ using a velocity-density factor of 0.15
206 following numerous previous studies (e.g., Becker, 2006; Conrad and Lithgow-Bertelloni, 2006;
207 Conrad and Behn, 2010). The density perturbation is further translated into a temperature
208 anomaly ΔT (Figure 2C, D) by multiplying it by the negative inverse of thermal expansion (e.g.,
209 Austermann et al., 2017; Karato, 1993). The shear wave tomography from Emry et al. (2019)
210 shows low velocities up to 6-12% beneath the Kenya Rift and the Main Ethiopian Rift (Figure
211 2A, B), which could result in a temperature excess of 500-1000 K following the conversion
212 outlined method above. However, these temperatures are higher than typical excess plume
213 temperatures ranging from 200-300 K (e.g., Schuberth et al., 2009). We therefore apply a
214 temperature excess cut-off of 300 K during the shear wave anomaly to temperature conversion.
215 This assumption is supported by geochemical evidence suggesting that low seismic shear wave
216 velocities beneath the EAR are a combined contribution from mantle temperature, melt (at
217 shallow depth), and composition beneath 200 km (Rooney et al., 2012). Details of the calculation
218 of our lithospheric geothermal gradient is described in the Supplementary material.

219

220 The numerical experiments incorporate non-linear (non-Newtonian) rheological flow
221 laws, which are needed to match first-order observations of solid Earth deformation and
222 constraints from rock deformation experiments (Figure 3). In brittle regions (i.e., in the upper
223 crust and uppermost mantle), plasticity limits the stress and reduces viscosity, which we assume
224 is in accordance with a Drucker-Prager yield criterion. The viscous flow law for the crust is
225 visco-plastic with dry quartzite and dislocation creep for dry olivine in the mantle lithosphere.
226 We assume that the main composition of the mantle is dry olivine since the area is far from a

227 subduction zone. The flow in the sublithospheric mantle lithosphere is governed by composite
228 rheology (Jadamec and Billen, 2010). A plastic strain weakening factor for cohesion and friction
229 is applied to the lithospheric viscosity in the deforming zones (Stamps et al., 2018) in order to
230 promote strain localization.

231

232 The velocity boundary condition at the model surface is free slip in order to minimize
233 unrealistic vertical motions and restrict near-surface motion to be horizontal. For the multiple
234 plumes model, the mechanical boundary conditions are free slip on all faces, except at the
235 bottom boundary where a no-slip condition is used to simulate the relatively higher viscosity and
236 slow mantle flow beneath the transition zone. For the superplume model, a 2 cm/yr northward
237 velocity boundary condition is applied at the northern and southern faces below 200 km and zero
238 velocity above this depth (See Figure 3).

239

240 **2.2 Estimating Mantle Flow Induced LPO**

241 Using the program D-Rex (Kaminski et al., 2004), we predict asthenospheric crystal
242 aggregates LPO that develop beneath the EAR from both the multiple plume and superplume
243 models. As D-Rex incorporates deformation mechanisms of olivine-enstatite aggregates
244 in flow such as plastic deformation, grain boundary sliding, and dynamic recrystallization
245 (Kaminski and Ribe, 2001, 2002), the code has been successfully used to estimate synthetic
246 LPO that develop in response to given global or regional mantle flow fields (e.g. Faccenda
247 and Capitanio, 2012, 2013; Hu et al., 2017; Rajaonarison et al., 2020). Since most seismic
248 anisotropy is often interpreted as the fast axis of olivine crystals, the transverse isotropic
249 approximation (TI) output by D-Rex is used here as a proxy for LPO (e.g., Becker et al.,

250 2006; Faccenda and Capitanio, 2012; Hu et al., 2017). The workflow for calculating synthetic
251 LPO that develops along pathlines in steady state flow is outlined below .

252 We consider particles that consist of 2000 virtual grains of olivine (70%) and enstatite
253 (30%) crystals (Kaminski et al., 2004), which represents a harzburgite composition. The particles
254 are located beneath the seismic stations where observations are available, and distributed
255 between 100 km and 400 km depths with a vertical interval of 25 km. For the purpose of
256 calculating their trajectories, first, the particles are advected backward in time, using the fourth
257 order Runge-Kutta advection scheme described in Becker et al. (2003), during which they
258 accumulate strain and stopped until a critical strain of 0.75 is reached (e.g., Rajaonarison et al.,
259 2020) or until the particles reach the 410 km discontinuity. At the location where the backward
260 advection is stopped, the particle is set with randomly oriented crystals to form an isotropic
261 mantle and advected forward along its trajectory. During this forward advection, the mineral
262 aggregates accumulate finite strain, LPO develops and evolves along the pathlines in response to
263 shear. Then, D-Rex calculates the global elastic tensor of the grain assemblage using the Voigt
264 averaging scheme for single-crystal elastic tensors and extracts the TI axis orientation that we
265 subsequently compare with observations.

266

267 **3 Results**

268 **3.1 Mantle Flow Beneath the EAR**

269 Mantle flow from the multiple plume model is mainly characterized by localized
270 upwelling beneath multiple rift segments of the EAR (see Figure 4). Figures 4A and 4B show
271 map views of the multiple plumes at depth slices 150 km and 300 km, with background colors
272 illustrating vertical velocities and the vectors indicating horizontal motions at that depth. The

273 positive vertical velocities at 150 km (Figure 4A) indicate localized upwelling occurs beneath the
274 Main Ethiopian Rift, the Eastern Branch, the Western Branch, and the northern Malawi Rift. The
275 upwelling rates across the entire EAR range from 6-10 cm/yr, whereas the upwelling beneath
276 northernmost EAR, the Main Ethiopian Rift, and Kenya Rift are more vigorous than beneath the
277 Western Branch. These patterns of upwelling are consistent with relatively low shear wave
278 anomalies beneath the Main Ethiopian Rift, the Eastern Branch, and northern Malawi Rift, which
279 result in relatively higher temperature, higher buoyancy, and lower viscosity. The associated
280 horizontal velocities show radiating patterns that result mostly in rift perpendicular flow with a
281 dominantly E-W orientation.

282 At 300 km depth (Figure 4B) the horizontal velocities associated with the localized
283 upwelling converge inward, indicating that slow seismic velocity anomalies are not connected to
284 a deeper source that passes through the transition zone. A profile view of the mantle flow (Figure
285 4C) shows that the localized upwelling extends to the 410 km discontinuity and is mostly limited
286 to that depth due to the relatively high viscosity between 410-660 km. Figure 3C also shows that
287 a downwelling beneath the Tanzania Craton separates the upwellings beneath the Western
288 Branch and Eastern Branch.

289 Asthenospheric flow from the superplume model beneath the EAR is characterized by
290 northward horizontal flow driven by the imposed boundary velocities (Figure 5A,B). The
291 associated vertical velocities show that the upwellings still occur beneath multiple segments
292 of the rifts. At 150 km (Figure 5A), the northward mantle flow imposed from the southern model
293 boundary is deflected to both the west and to the east by the rheologically

294 strong lithosphere of the Tanzania Craton. To the west, the mantle flow is channeled northward
295 between the Tanzania Craton and the Bangweulu Block and then accommodates the curvature of
296 the Western
297 Branch to the north. To the East, the mantle flow is slightly deflected eastward by the Tanzania
298 Craton lithosphere. At 300 km depth (Figure 5B), the horizontal mantle flow mostly exhibits
299 northward flow directions except beneath the Eastern Branch and Main Ethiopian Rift, where the
300 horizontal flow tends to align in the ~NE direction. Figure 5C shows that beneath the thick
301 lithosphere in the southern EAR mantle flow pattern is mostly horizontal trending northward. In
302 contrast, beneath the thin lithosphere of northern EAR, localized upwellings still occur at
303 shallow depths.

304

305 **3.2 Comparison of Mantle Flow Induced LPO**

306 In this section, we present the comparison of synthetic LPO, or TI-axis orientation, and observed
307 splitting at locations where each of the observations exist (Figure 6 and Figure 7).
308 Due to the high variability of the anisotropy orientations, we present the comparisons regionally,
309 partitioned into northern (Region A), central (Region B), and southern (Region
310 C) EAR. Also, given the changes of horizontal mantle flow pattern with depth from both the
311 multiple plumes and superplume models, we illustrate and present the comparison with
312 observations at relatively shallow asthenosphere (150 km) and at relatively deeper depths (300
313 km). Table 1 shows a summary of the comparison of predicted TI-axis with observed shear wave
314 splitting. For the multiple plumes mantle flow model (Figure 6) the overall observed (N27°E)
315 fast direction pattern (~S39°E) is poorly reproduced by the calculated synthetic LPO. At 150

316 km (Figure 6A), the synthetic LPO mostly exhibits a rift perpendicular trend with an average
317 orientation of S20°E in the northern EAR (Region A), the orientation is S15°E
318 in the central EAR (Region B), and the trend is N20°E at the southern EAR (Region
319 C). Comparisons yield a poor angular misfit with the observations with average angular
320 misfits of 51°, 52°, and 40°, respectively. Similarly, poor misfit are found at 300 km depth
321 because the synthetic LPO also exhibits rift perpendicular trends with average orientations
322 of S20°E, S15°E, and N20°E for the northern, central, and southern EAR, respectively. The
323 corresponding mean angular misfits are 51°, 51°, and 46°, respectively. The rift perpendicular
324 trend of the synthetic LPO is due to radiating pattern of the horizontal flow from the multiple
325 plumes model (see Figure 4).

326 For the superplume mantle flow model (Figure 7), the overall observed ~NE fast
327 direction patterns are well reproduced by the calculated synthetic LPO at relatively deeper depths
328 (300 km; Figure 7B) than at shallow depths (150 km; Figure 7A). At 150 km (Figure 7A), the
329 synthetic LPO mostly exhibits rift perpendicular trends (S30°E, S1°E, and N20°E, respectively,
330 for the northern, central, and southern EAR), yielding poor mean angular misfit of 60°, 44° ,
331 and 46°, respectively. The rift perpendicular trend of the synthetic LPO at 150 km is consistent
332 with the radiating pattern of horizontal flow pattern (see Figure 4A). At 300 km (Figure 7B), we
333 find a good fit between the observations and the synthetic LPO with average directions of N23°
334 E, N13°E, and N50°E for the northern, central, and southern EAR, respectively. The best fit is
335 found in the northern EAR (Region A, Figure 7B) with an average misfit of 20°. The angular
336 misfits are 27° for central EAR, and 40° in the southern EAR (Region C; Figure 7B).

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342 **3.3 Surface Velocity Comparison**

343 To test how the presence of multiple plumes or a superplume beneath the EAR drives surface

344 deformation, we quantitatively compare our dynamic velocity estimates with kinematic

345 predictions of surface motions within zones of rigidity (Stamps et al., 2021; Figure 1A) and

346 GNSS/GPS velocities in intra-rift zones (Stamps et al., 2018; Figure 1A). For velocity

347 magnitude comparison, we use the Root Mean Square velocity statistic, whereas regional

348 mean angular misfit ($[0-180]$ from aligned to opposite velocity direction) is used to quantify

349 the fit between predicted and observed velocity directions. The comparison of dynamic

350 velocities with the kinematic model and GPS observations are summarized in Table 2.

351 Dynamic velocities driven by mantle tractions from multiple plumes mantle flow and

352 lithospheric buoyancy forces predict well the rigid plate motions of the Victoria and the Rovuma

353 Block with RMS misfit of 0.5 mm/yr and 0.5 mm/yr, respectively, and an angular misfit of

354 6° and 10° , respectively (Figure 8B). For the Somalian Plate, the dynamic velocities align

355 well with the predicted clockwise rotation with an angular misfit of 5° , but the rotation

356 rate is 13 mm/yr, which is three times larger than the kinematic model (4.41 mm/yr; see

357 Table 2), resulting in a large RMS misfit of 12 mm/yr (Figure 8A). Within the deforming

358 zones (Figure 8B), the dynamic velocities are driven by mantle tractions from the multiple

359 plumes and lithospheric buoyancy forces resulting in a poor fit to the GNSS/GPS observations.

360 In the MER (Region A; Figure 8), although the dynamic velocities align well with

361 the GNSS/GPS observations at some stations within the rift, the velocities at the western

362 border of the rift and in the southern region do not align well, yielding a mean angular

363 misfit of 60° . The dynamic velocities are also faster, with a mean velocity of 7.1 mm/yr
364 compared to the GNSS/GPS velocities (2.59 mm/yr; see Table 2). In the Eastern Branch
365 (Region B,C,D; Figure 8B), we find good fits between the dynamic velocities and the GPS
366 observations in Region B and C, with RMS misfit of 2.1 mm/yr and 1.5 mm/yr, respectively, and
367 angular misfits of 5° mm/yr and 4° mm/yr, respectively. The fit is worse in the central Eastern
368 Branch (Region C; Figure 8B) where the RMS misfit is 2.5 mm/yr. In Region C, the fit is
369 consistent with some of the dynamic velocities exhibiting a \sim SE trend whereas the GPS
370 velocities trend \sim NE, yielding poor angular misfit of 50° . For the Western Branch (Region E;
371 Figure 8), the mean dynamic velocity magnitude is relatively small
372 (0.2 mm/yr) compared with the 2.55 mm/yr observed by GNSS/GPS, with an RMS misfit
373 of 1.36 mm/yr and the velocities do not align well with a poor angular misfit of 99° . For
374 the southern EAR (Region F; Figure 8) or the northern Malawi Rift, we find a good fit
375 of the dynamic velocities and the GNSS/GPS velocities with an angular misfit of 20° and
376 an RMS misfit of 0.8 mm/yr. Overall, the lithospheric buoyancy forces combined with
377 sublithospheric multiple plumes drive similar surface deformation to the lithospheric buoyancy
378 forces alone estimated by Rajaonarison et al. (2021a,b), indicating that lithospheric deformation
379 is decoupled from the mantle flow, except for the Somalian Plate, which appears to be coupled
380 with the underlying eastward mantle flow.

381 For the model in which deformation is driven by mantle tractions from the superplume
382 and lithospheric buoyancy forces, the block kinematics are also predicted well with angular
383 misfits of 5° , 7° , and 6° , respectively for the Somalian Plate, Victoria Block, and Rovuma Block.
384 Their rotation rates are slightly faster (7.5 mm/yr, 2.8 mm/yr, and 3.1 mm/yr, respectively;
385 Figure 9A) than the kinematic model (4.41 mm/yr, 2.13 mm/yr, and 2.14 mm/yr, respectively;

386 Figure 9A). The Somalian Plate rotation exhibits the largest RMS misfit (5 mm/yr), indicating
387 some degree of coupling between the plate and the underlying mantle flow. In the deforming
388 zones, we find some improvement of the fit with the lithospheric buoyancy forces and
389 superplume combined model. In the MER (Region A; Figure 9B), although the RMS misfit is
390 relatively large (6 mm/yr), the mean angular misfit decreases from 66° (from the model with
391 only GPE only, see Table 2) to 30° , with predicted velocities and observations having opposite
392 directions at a station located in the southern part of Region A. We also find in the mean angular
393 misfit at the central Eastern Branch (Region C; Figure 9B), where the dynamic velocities at some
394 stations have shifted from \sim SE trend to \sim E-W and \sim NE decreasing the mean angular misfit
395 from 73° to 30° . In the Western Branch (Region E; Figure 9B), the mean angular misfit also
396 significantly improves from 114° to 43° respectively. However, the dynamic velocities remain
397 small 0.7 mm/yr compared to the GNSS/GPS velocities (2.55 mm/yr). In the southern EAR
398 (Region F; Figure 9B), the dynamic velocities exhibit \sim E-W and \sim NE trends whereas the GPS
399 velocities are oriented \sim E-W and \sim SE, which slightly worsens the mean angular misfit from 20°
400 to 24° , but may be within the uncertainties of the model. Overall, the comparison of dynamic
401 velocities from the combined lithospheric buoyancy forces and the superplume model indicates
402 there is some degree of lithosphere coupling to northward flow, notably within the deforming
403 zones and beneath the Somalian Plate.

404

405 **4 Discussion**

406 **4.1 Source of Observed Seismic Anisotropy Beneath EAR**

407 Overall observed seismic anisotropy in the EAR exhibits N-S or NNE trends with delay
408 times ranging from 1-1.5 s. Most anisotropy with such relatively large delay time are often

409 attributed to mantle flow induced LPO (Helffrich, 1995). However, in divergent settings,
410 mantle flow induced LPO are expected to be perpendicular to the rifts. For example, along
411 mid oceanic ridges mantle flow induced LPO are perpendicular to the plate boundaries
412 (Becker et al., 2003, 2006, 2014). Thus, interpretations of the observed rift parallel anisotropy
413 along the EAR remains controversial. Several mechanisms such as fossilized anisotropy due
414 to past orogenic events (Homuth et al., 2016; Walker et al., 2004), and parallel dykes or magma
415 filled lenses (Albaric et al., 2014; Gao et al., 1997) have been invoked as explanations for
416 the rift parallel anisotropy. Bagley and Nyblade (2013) suggest that the overall consistent
417 N-S or NNE trend of the anisotropy can be attributed one single mechanism that is the
418 northward flow of the African superplume.

419 The results presented here corroborate with the African superplume hypothesis suggested
420 by Bagley and Nyblade (2013) because we find good alignment with calculated LPO and
421 observations at deeper depths (300 km) than shallower depths (150 km) for the superplume
422 mantle flow model. Moreover, the multiple plume model does not produce a good fit with the
423 observations indicating that mantle upwelling beneath the Main Ethiopian Rift, the Kenya Rift,
424 the Western Branch are influenced by northward flow from the African superplume.
425 Our results also highlight the importance of the Tanzania Craton, Congo Craton, and the
426 Bangweulu Block in channeling the northward mantle wind resulting in rift parallel mantle
427 flow mainly in the Western Branch. This rift parallel mantle flow results in rift parallel
428 anisotropy explaining most of the seismic anisotropy including those in the northernmost
429 Western Branch. It is also worth noting that sub-lithosphere mantle viscosity plays an important
430 role in favoring rift parallel mantle flow. For example, the mantle flow and the LPO
431 exhibit NNE trends at the Main Ethiopian Rift as a result of low viscosity along the strike

432 of the rift.

433

434 **4.2 Plume-Lithosphere Interactions Beneath EAR**

435

436 Our comparison of surface velocities with predicted rigid block motion and GPS velocities

437 highlights the importance of the interactions between the lithosphere and mantle flow associated

438 with the African superplume beneath the EAR. The significance of plume-lithosphere interactions

439 has been discussed by Forte et al. (2010), Ghosh and Holt (2012), and Kendall and Lithgow-

440 Bertelloni (2016), who suggest that mantle tractions from divergent flow associated with the

441 African superplume alone can cause the opening of East Africa. However, previous geodynamic

442 modeling results suggest that the E-W extension of the EAR is dominated by lithospheric

443 buoyancy forces and adding sub-lithospheric mantle flow over-predicts the observed deformation

444 (Stamps et al., 2014, 2015). Our numerical experiments of lithospheric deformation combined with

445 asthenospheric flow suggest that the multiple plumes model beneath the EAR does not

446 significantly affect the surface deformation of the EAR, except for the Somalian Plate where the

447 rotation rate is over-predicted. Moreover, the multiple plumes mantle flow cannot explain the

448 northward component of motion with the deforming zones. The inconsistency between the

449 multiple plume model and surface velocities in the deforming zones is likely because beneath the

450 rifts is dominated by upwelling and diverging E-W flow beneath the adjacent rigid blocks. In

451 contrast, the superplume model results in mantle flow beneath the rifts that is mostly oriented

452 northward and rift parallel. This northward horizontal flow generates horizontal tractions at the

453 base of the lithosphere and produces northward surface motions.

454 Our results also suggest that the counter-clockwise rotation of the Victoria Block is

455 independent of the underlying mantle flow from the multiple plumes model, but a minor

456 contribution might occur from the northward mantle flow associated with the African Superplume.
457 This interpretation is consistent with the findings of Glerum et al. (2020). Due to the presence of
458 Tanzania Craton within the Victoria Block, the lithosphere-asthenosphere coupling is expected to
459 be significant because previous studies have shown that viscous coupling between asthenosphere
460 and lithosphere is favorable beneath cratonic roots (Conrad and Lithgow-Bertelloni, 2006;
461 Stoddard and Abbott, 1996; Zhong, 2001). This strong coupling is not the case for the Victoria
462 Block possibly because of the continuation of the Congo Craton into the Tanzania Craton through
463 the northern Western Branch (Link et al., 2010) provides resisting forces (drag) to mantle tractions.
464 The resisting forces can be caused by the stability of the Congo Craton within the stable Nubian
465 Plate. The role of a craton in favoring viscous coupling of mantle flow and lithosphere can,
466 however, be observed in the central Eastern Branch, located to the north of the Masai Block. Our
467 results indicate that due to the presence of the Masai Block, the surface deformation in the central
468 Eastern Branch is influenced by the underlying northward mantle flow associated with the African
469 Superplume.

470

471 **4.3 Anisotropic Viscosity**

472 Our results indicate the rotation rate of the Somalian Plate is over-predicted by our
473 lithospheric deformation models driven by lithospheric buoyancy forces and mantle flow (multiple
474 plumes and superplume model). Previous geodynamic modeling studies have also reported this
475 over-prediction when coupling lithospheric deformation with basal shear from global mantle flow
476 models (e.g., Stamps et al., 2014, 2015). These previous studies suggest that lithospheric buoyancy
477 forces alone are capable of driving present-day plate motion of the EAR and that viscous coupling
478 between mantle flow and lithosphere should be inefficient. Stamps et al. (2015) attribute the

479 inefficiency of coupling with lower viscosity of the asthenosphere or slower mantle flow than is
480 often estimated. An alternative mechanism of decoupling between mantle flow and lithosphere
481 beneath the EAR is anisotropic viscosity in the asthenosphere. Recent studies on viscous
482 anisotropy in the asthenosphere reveal that the development of olivine aggregates' LPO can
483 weaken the viscosity in the direction parallel to the fabrics and strengthen the viscosity in the
484 perpendicular direction (Hansen et al., 2016; Király et al., 2020). The anisotropic viscosity has
485 important implications for plate motion because it can increase plate velocity in the direction of
486 the LPOs and slow down in the perpendicular direction over a period of up to ~10 Myr (Király et
487 al., 2020). Moreover, Perry-Houts and Karlstrom (2019) found that in the lithosphere, the
488 magnitude of anisotropic viscosity is highly dependent on the ratio of intrusion to host rock
489 viscosity, highlighting the importance of anisotropic viscosity in lithospheric deformation
490 modeling. Here, we found that the observed rift parallel NE seismic anisotropy beneath the EAR
491 can be associated with the northward mantle flow from the African Superplume. The presence of
492 such long-term northward flow and the induced LPO could lead to anisotropic viscosity beneath
493 the EAR. This would result in weak viscosity parallel to the rifts and thus faster northward flow
494 than estimated in this study. Conversely, the viscosity in the E-W direction would be stronger
495 causing more resistance to plate motion and would result in slower Somalian plate motion. Since
496 in this study, we have posteriorly calculated the synthetic LPO after the mantle flow simulation,
497 the anisotropic viscosity is not captured in our models. The implementation of anisotropic viscosity
498 would require using a viscosity tensor (e.g., Hansen et al., 2016; Király et al., 2020; Perry-Houts
499 and Karlstrom, 2019) rather than the scalar composite viscosity used here. The concept of
500 anisotropic viscosity beneath the EAR provides new insight into the lithosphere and mantle
501 dynamics beneath East Africa and a future direction for studies of continental rifting.

502

503 **5 Conclusions**

504 We use 3D thermomechanical modeling of the lithosphere-asthenosphere system to investigate:

505 1) the role plume-lithosphere interactions beneath the EAR in driving observed surface

506 deformation and 2) the sources of observed seismic anisotropy beneath the EAR oriented

507 approximately ~N-S. We test two types of mantle flow models including a multiple plume

508 model constrained by shear wave tomography from Emry et al. (2019) and an African superplume

509 model simulating a northward mantle wind on the multiple plume model. Our results

510 indicate that interactions of the northward African superplume with the rigid Tanzania

511 Craton, Congo Craton, and Bangweulu and Masai Blocks may be controlling mantle flow pattern

512 beneath the EAR and may be the source of most of the observed seismic anisotropy. We also

513 found that the multiple plume mantle flow model should be ruled out as a dominant source

514 mechanism of sub-lithospheric horizontal mantle flow beneath the EAR. Comparison of dynamic

515 velocities and block kinematics suggests that the Victoria and Rovuma Blocks rotate

516 independently from mantle flow driven by multiple plumes or a superplume. In contrast, we

517 found that for the multiple plumes and superplume mantle flow models, the Somalian plate is

518 over-predicted up to twice or three times its expected rotation rate. While the multiple plumes

519 mantle flow model does not significantly affect the surface deformation driven by the

520 lithospheric buoyancy forces, the northward mantle flow field from the superplume model drives

521 northward component of motions in the deforming zones that improves the fit at the Main

522 Ethiopian Rift, Central Eastern Branch, and northern Western Branch. Those results suggest that

523 sublithospheric viscosity beneath the EAR might be anisotropic due to developed olivine

524 aggregates's LPO induced by the African superplume. The anisotropic viscosity might control
525 the rotation rate of the Somalian Plate.

526

527

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537

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847 **Table 1:** A comparison of TI-axis, or synthetic LPO, from the two models (multiple plumes and
 848 superplume) produced in this work with observed anisotropy. $\langle\phi_{obs}\rangle$ and $\langle\delta t_{obs}\rangle$ are the observed
 849 regional mean fast direction, respectively. TI-axis represents the predicted synthetic LPO. $\langle\Delta\alpha\rangle$
 850 is the regional mean angular misfit between the observed fast direction and predicted LPO.

Region	Observations		Multiple Plumes Model				Superplume Model			
	$\langle\phi_{obs}\rangle^\circ$	$\langle\delta t_{obs}\rangle$ (s)	TI-axis		$\langle\Delta\alpha\rangle^\circ$		TI-axis		$\langle\Delta\alpha\rangle^\circ$	
Depth	-	-	150 km	300 km	150 km	300 km	150 km	300 km	150 km	300 km
Northern	N27°E	1.3	S20°E	S20°E	51°	51°	S30°E	N23°E	60°	20°
Center	N26°E	1.	S15°E	S20°E	52°	51°	S1°E	N13°E	44°	27°
South	N35°E	0.8	N20°E	N1°E	40°	46°	N20°E	N56°E	46°	40°

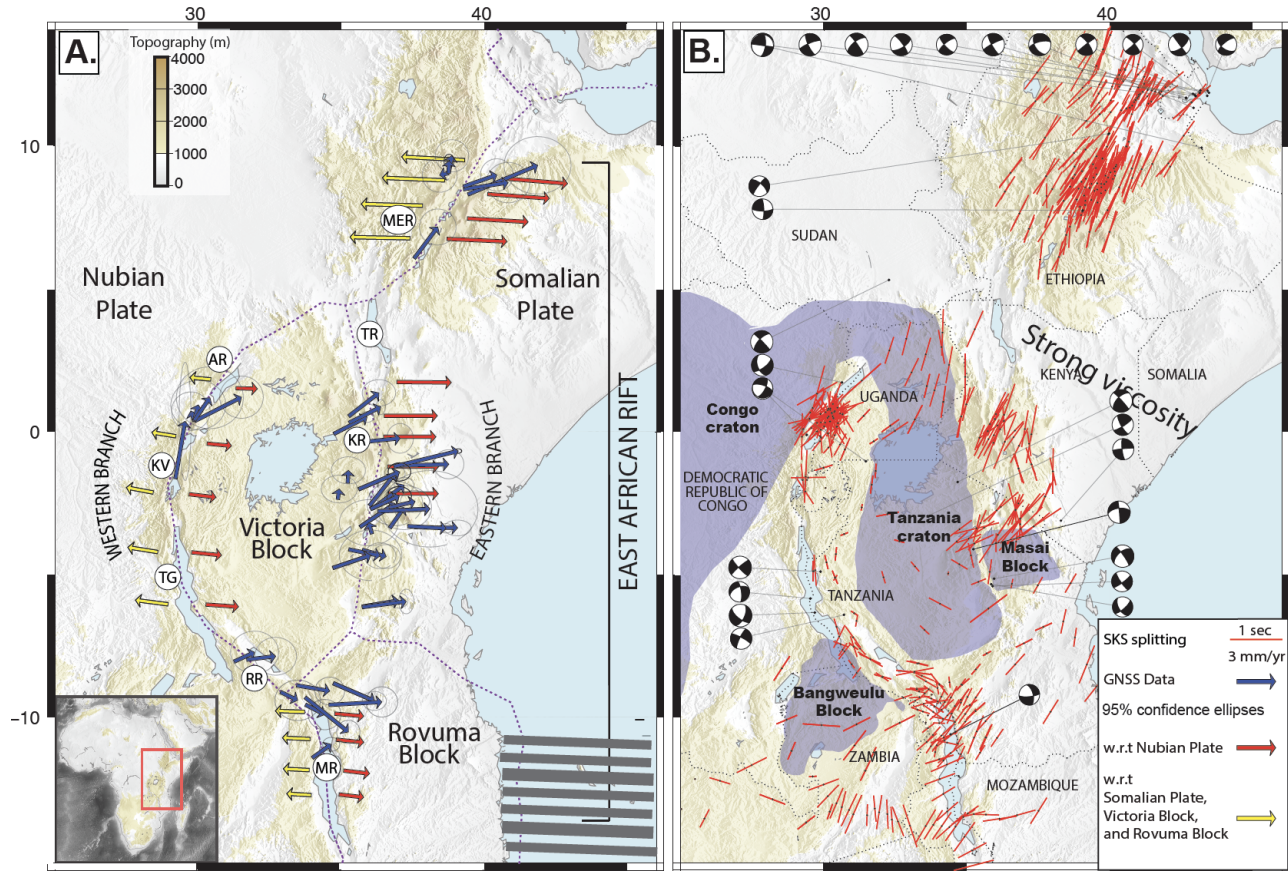
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852 **Table 2:** Summary of the comparison of dynamic velocities from this study (driven by GPE +
853 Multiple Plumes Model and GPE + Superplume Model) and from Rajaonarison et al. (under
854 review, GRL) (driven by GPE only) with kinematic model from [Stamps et al. \(2020\)](#) and
855 GNSS/GPS velocities from Stamps et al. (2018) in deforming zones defined as regions A-F. \bar{V}
856 (mm/yr) represents mean velocity, α represents mean angular misfit ($[0 - 180^\circ]$ from good to poor
857 fit), and RMS is the root mean square velocity.

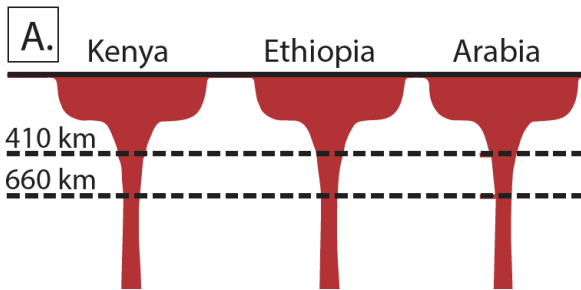
Region	Kinematic model/GPS		GPE only		GPE+Multiple Plumes Model			GPE+Superplume Model		
	\bar{V} (mm/yr)	\bar{V} (mm/yr)	RMS (mm/yr)	α	\bar{V} (mm/yr)	RMS (mm/yr)	α	\bar{V}	RMS (mm/yr)	α
Somalian Plate	4.41	6	1.8	10°	13	12.	5°	7.5	5.	5°
Victoria Block	2.13	1.39	0.5	5°	2.5	0.5	6°	2.8	0.6	7°
Rovuma Block	2.14	2.3	0.2	10°	2.	0.5	10°	3.1	1.	6°
Region A	2.59	5.9	2.6	66°	7.1	6	60°	8.1	6.	50°
Region B	2.7	4.3	1.6	12°	2.7	2.1	5°	2.7	0.7	40°
Region C	3.2	5.1	1.3	74°	2.9	2.5	50°	2.7	1.8	3°
Region D	2.8	3.2	1.1	8°	2.	1.5	4°	2.	1.8	30°
Region E	2.55	0.54	1.9	114 °	0.2	1.36	99°	0.3	1.3	4°
Region F	2.6	1.9	0.8	20°	2.4	0.8	20°	2.9	0.7	24°

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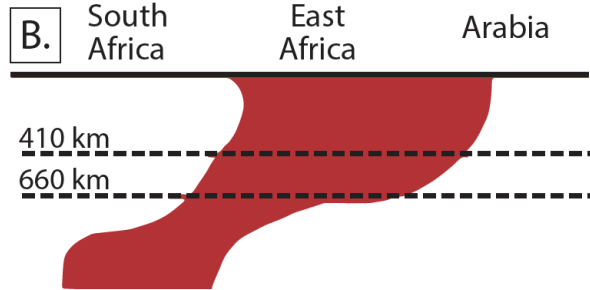
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Multiple Plume Model



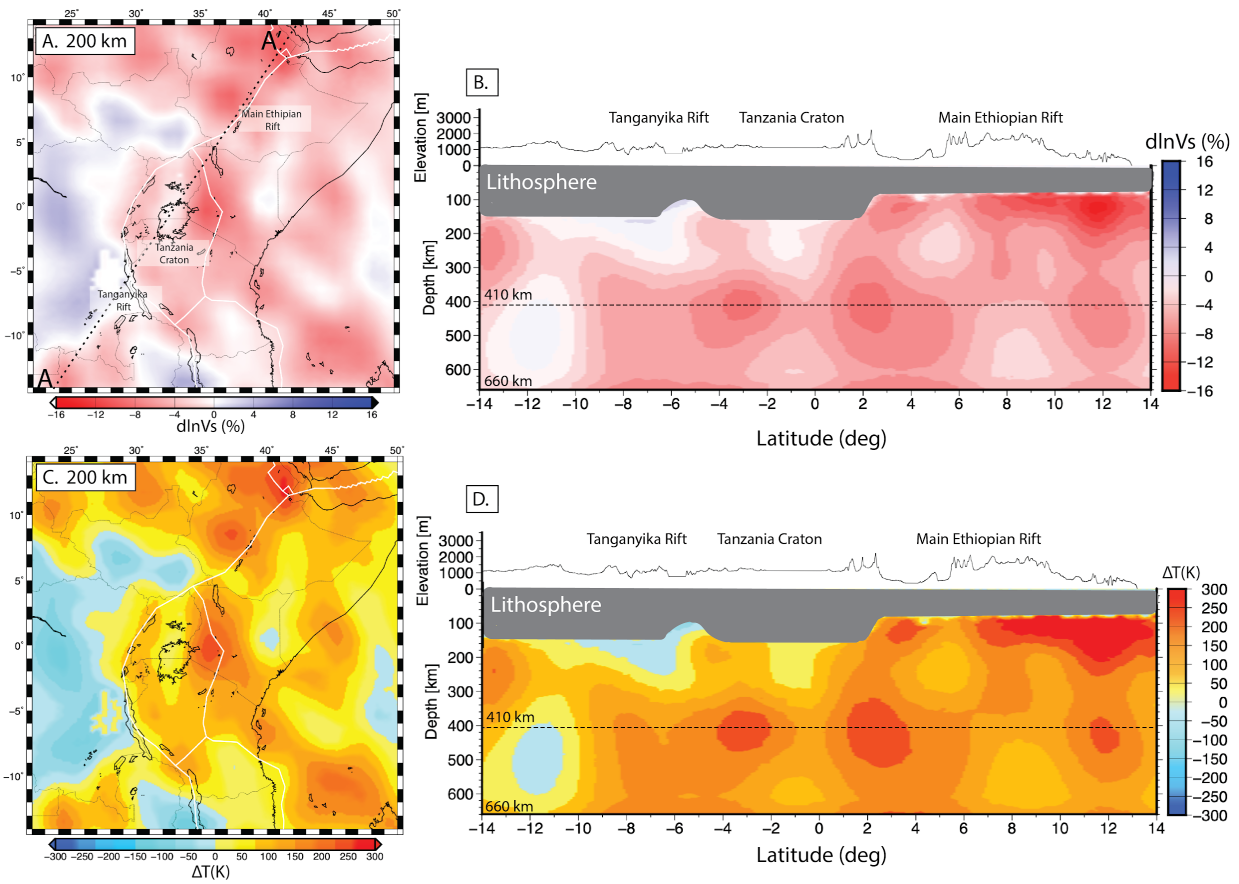
Superplume Model



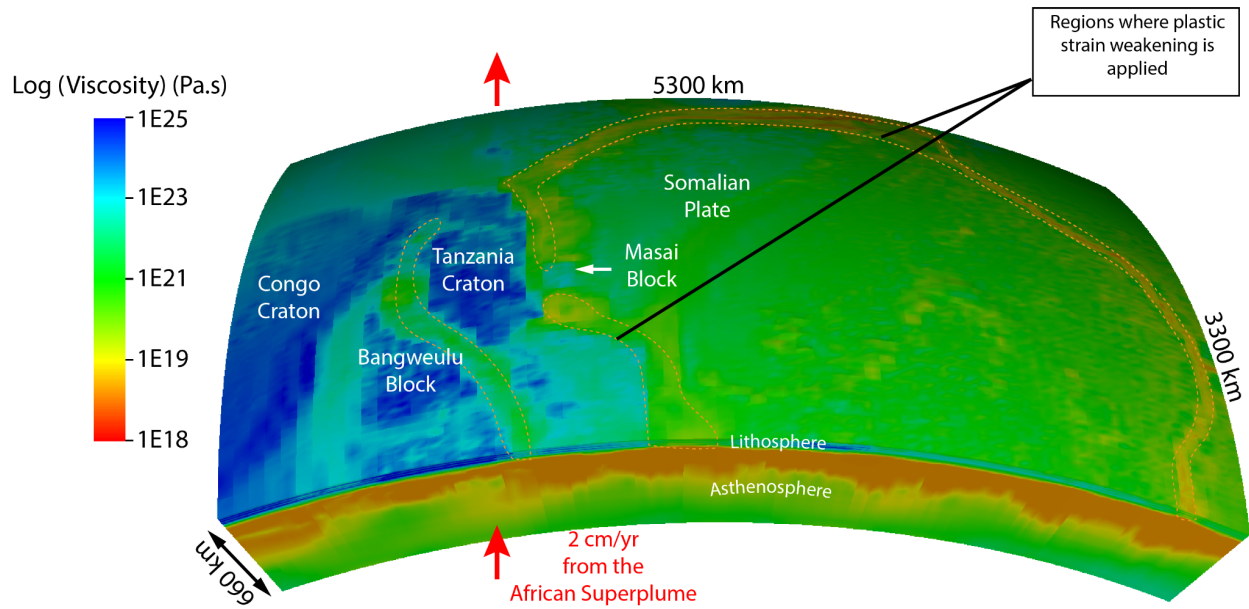
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861 **Figure 1:** A) GPS velocities along the deforming zone of the East African Rift (EAR; Stamps et
 862 al., 2018; blue vectors with 95% confidence ellipses). Red and yellow vectors are kinematic
 863 models from Saria et al. (2014) for the Somalian Plate, the Victoria Block, and the Rovuma Block,
 864 respectively. Purple dashed lines are plate boundaries. B) Observed SKS splitting (blue bars) and
 865 strike-slip focal mechanisms along the EAR. Several rifts are defined: TG = Tanganyika Rift, RR
 866 = Rukwa Rift, AR = Albertine Rift, MER = Main Ethiopian Rift, TR = Turkana Rift, KR = Kenya

867 Rift, MR = Malawi Rift, and KV = Kivu Rift. Black dots indicate earthquake locations where
 868 strike-slip focal mechanisms are observed (GCMT; Dziewonski et al., 1981; Ekström et al., 2012).
 869 Dashed lines are political boundaries. Hashed lines indicate deforming zones from Stamps et al.
 870 (2020). (C and D) Plume models for the East African Rift System invoked to explain upper mantle
 871 seismic velocity models (adapted from Hansen et al., 2012). C) The multiple plume model. D) The
 872 superplume model that is referred to as African superplume.
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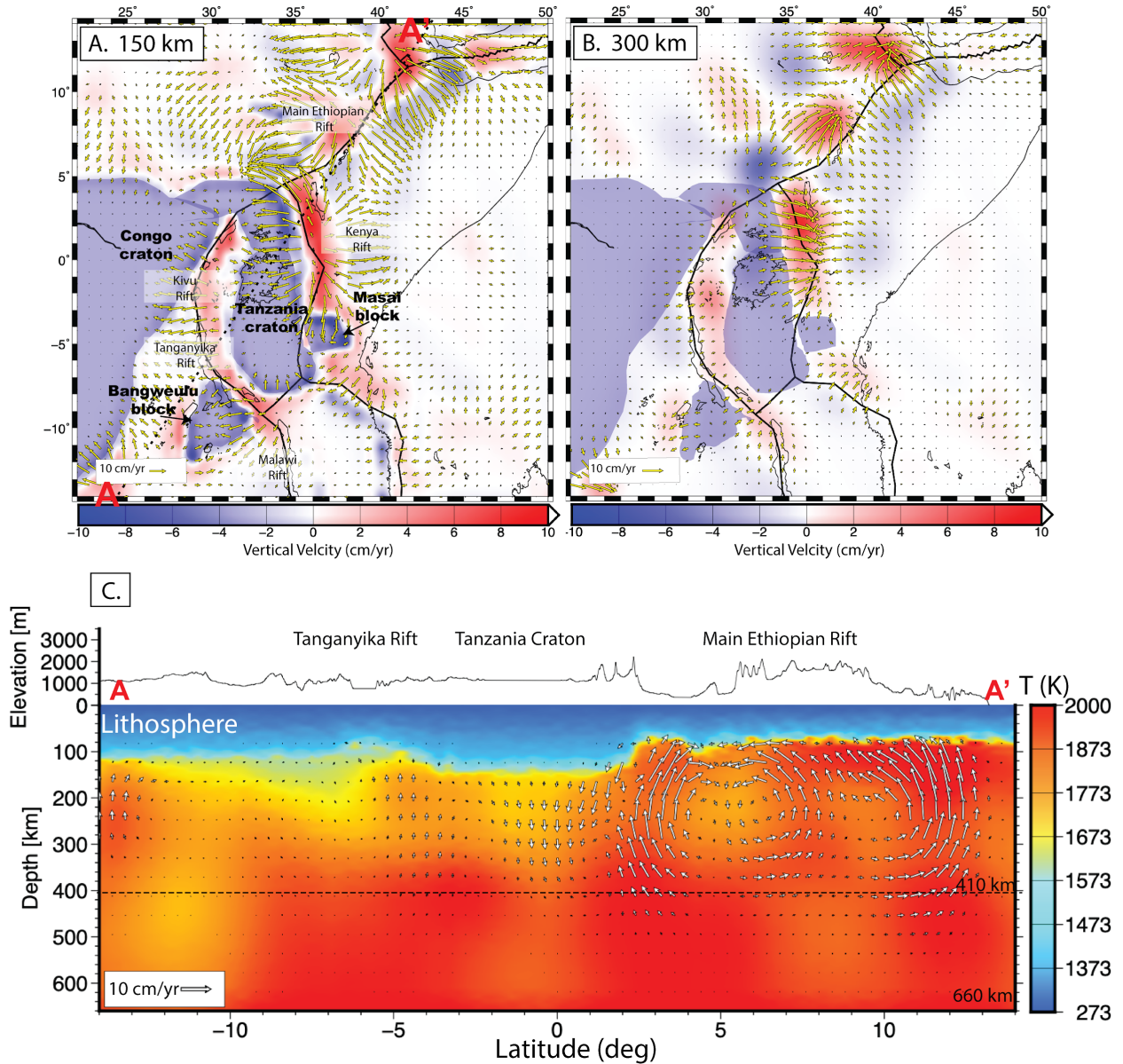


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 875 **Figure 2:** (A,B) Shear wave tomography anomaly from Emry et al. (2019): A) at 200 km depth
 876 and B) along profile AA'. (C,D) Temperature anomaly obtained by converting the shear wave
 877 tomography in A and B. C) at 200 km depth and D) along profile AA').
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880 **Figure 3:** 3D viscosity model used in this study. Dashed orange lines indicate regions where a
881 plastic weakening mechanism is applied to localize deformation. Red vectors indicate the
882 velocity boundary condition used to simulate the African Superplume.



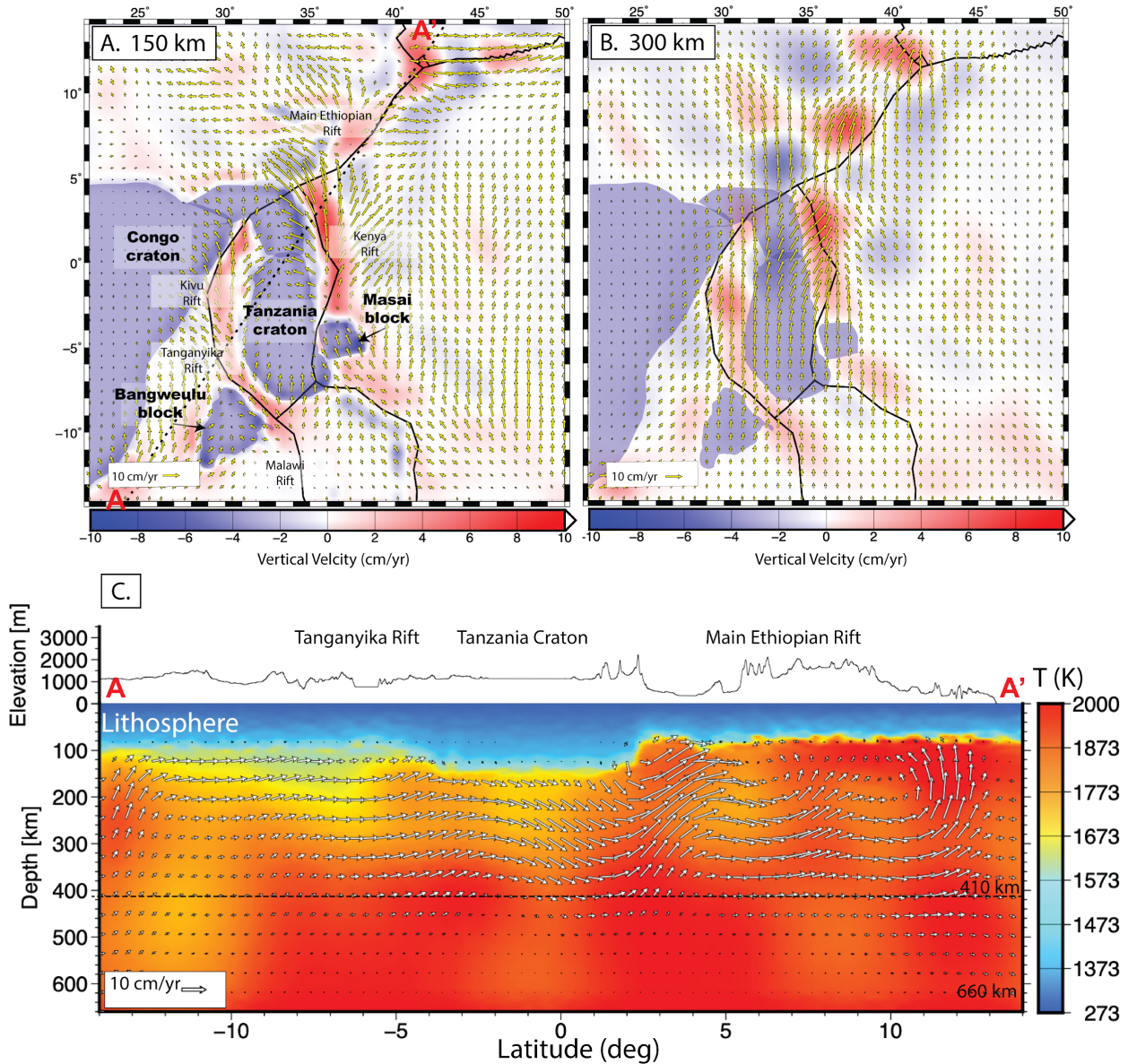
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885 **Figure 4:** Mantle flow field from the multiple plumes model: A) at 150 km depth, B) at 300 km

886 depth. Background color indicates vertical flow. Yellow vectors portray horizontal. C) along

887 profile AA' (Figure 4A). The background color indicates the temperature field.



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889 **Figure 5:** Mantle flow field from the superplume model: A) at 150 km depth, B) at 300 km

890 depth. Background color indicates vertical flow. Yellow vectors portray horizontal. C) along

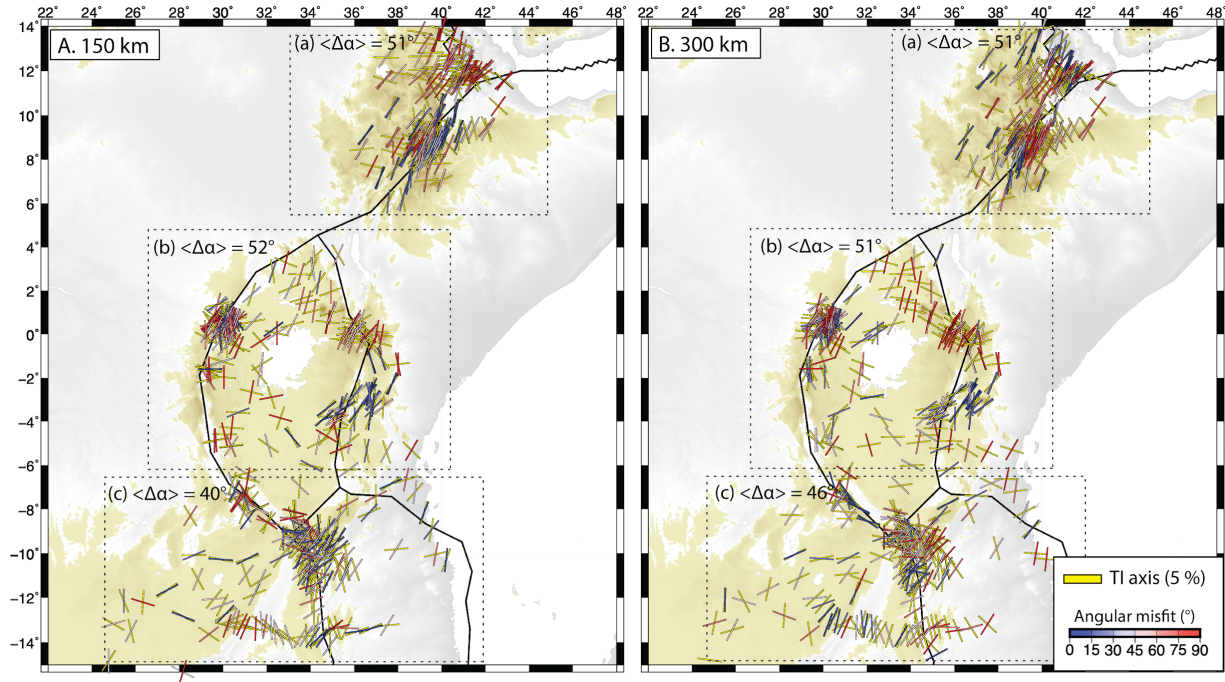
891 profile AA' (Figure 5A). The background color indicates the temperature field.

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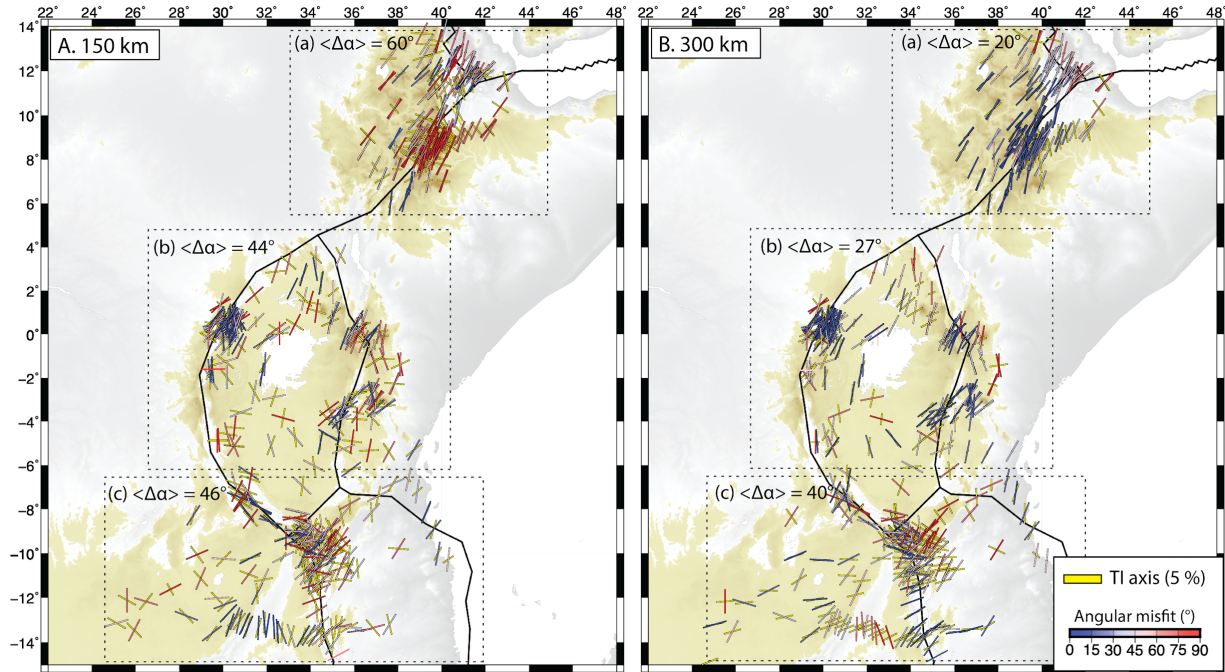
898 **Figure 6:** Comparison of calculated TI axes with observations from the multiple plume model:

899 A) at 150 km and B) at 300 km. The yellow bars indicate TI axis orientation. The SKS splitting

900 measurement bars are colored according to angular misfit [0° - 90°]. The background

901 shows topography.

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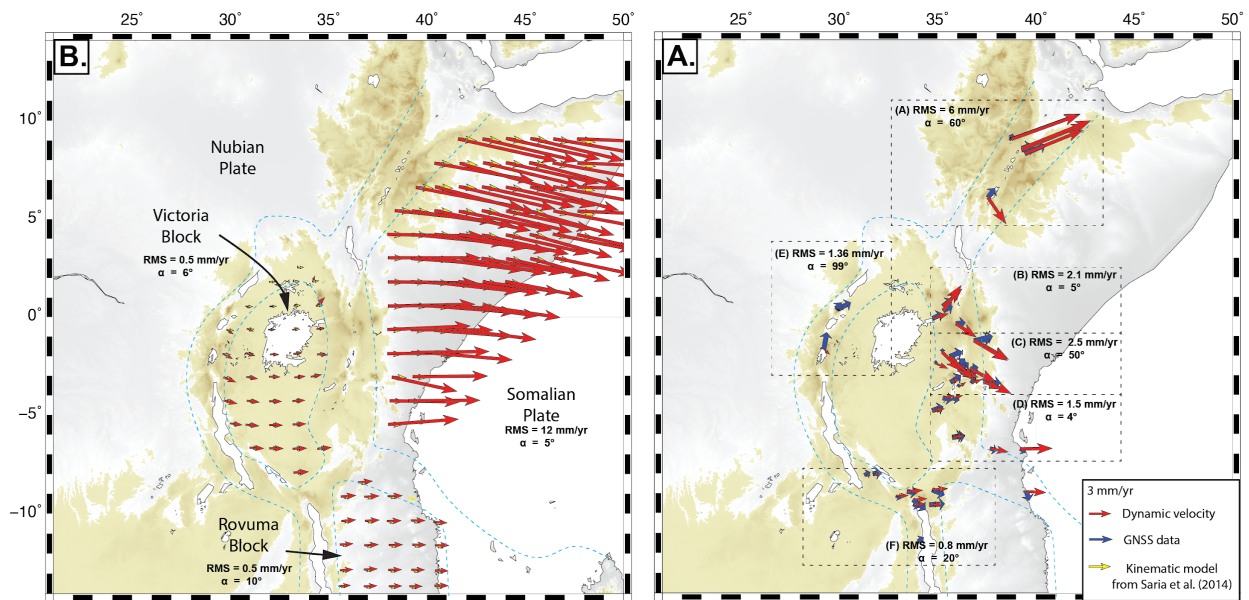


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904 **Figure 7:** Comparison of calculated TI axes with observations from the superplume model:

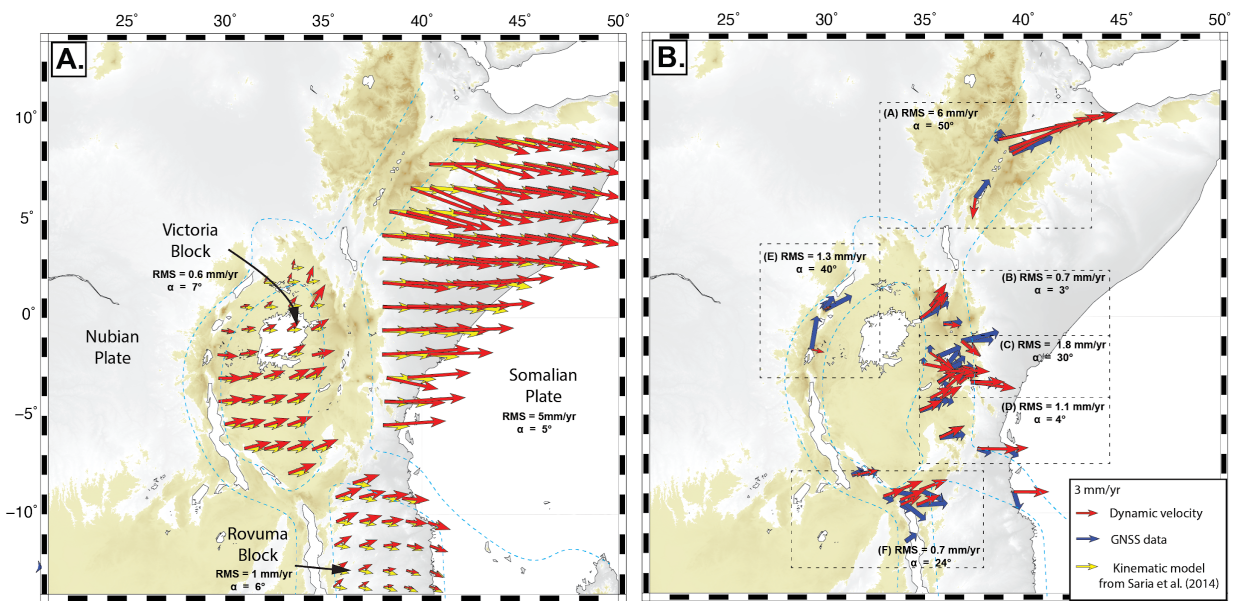
905 A) at 150 km and B) at 300 km. The yellow bars indicate TI axis orientation. The SKS splitting
 906 measurement bars are colored according to angular misfit [0°- 90°]. The background shows
 907 topography.

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910 **Figure 8:** Comparison of dynamic velocities (red vectors) driven by mantle tractions from the
 911 multiple plumes model and lithospheric buoyancy forces with: A) kinematic predictions from the
 912 Stamps et al. (2021) model (yellow vectors) within the Somalian Plate, the Victoria and the
 913 Rovuma Blocks. B) GNSS/GPS data from (Stamps et al., 2018, blue vectors) within the deforming
 914 zones (dashed blue line) and comparisons statistics (RMS and mean angular misfit) are shown
 915 inside a dashed box for each region.
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 918 **Figure 9:** Comparison of dynamic velocities (red vectors) driven by mantle tractions from
 919 the superplume model and lithospheric buoyancy forces with: A) kinematic predictions from the
 920 Stamps et al. (2021) model (yellow vectors) within the Somalian Plate, the Victoria and the
 921 Rovuma Blocks. B) GNSS/GPS data from (Stamps et al., 2018, blue vectors) within
 922 the deforming zones (dashed blue line) and Comparisons statistics (RMS and mean angular
 923 misfit) are shown inside a dashed box for each region.

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