## Fe – Ga Pb (Mg 1 3 Nb 2 3) O 3 – Pb Ti O 3 magnetoelectric laminate composites

Shuxiang Dong, Junyi Zhai, Naigang Wang, Feiming Bai, JieFang Li, D. Viehland, and T. A. Lograsso

Citation: Applied Physics Letters 87, 222504 (2005); doi: 10.1063/1.2137455

View online: http://dx.doi.org/10.1063/1.2137455

View Table of Contents: http://scitation.aip.org/content/aip/journal/apl/87/22?ver=pdfcov

Published by the AIP Publishing

Over **700** papers & presentations on multiphysics simulation



## Fe-Ga/Pb(Mg<sub>1/3</sub>Nb<sub>2/3</sub>)O<sub>3</sub>-PbTiO<sub>3</sub> magnetoelectric laminate composites

Shuxiang Dong, <sup>a)</sup> Junyi Zhai, Naigang Wang, Feiming Bai, JieFang Li, and D. Viehland *Materials Science and Engineering, Virginia Tech, Blacksburg, Virginia 24061* 

## T. A. Lograsso

Materials and Engineering Physics, Ames Laboratory, Ames, Iowa 50011

(Received 17 June 2005; accepted 4 October 2005; published online 22 November 2005)

We have found large magnetoelectric (ME) effects in long-type laminate composites of Fe–20%Ga magnetostrictive alloys and piezoelectric  $Pb(Mg_{1/3}Nb_{2/3})O_3-PbTiO_3$  single crystals. At lower frequencies, the ME voltage coefficient of a laminate with longitudinally magnetized and longitudinally polarized (i.e., L-L mode) layers was 1.41 V/Oe (or 1.01 V/cm Oe). Near the natural resonant frequency ( $\sim$ 91 kHz) of the laminate, the ME voltage coefficients were found to be dramatically increased to 50.7 V/Oe (36.2 V/cm Oe) for the L-L mode. In addition, the laminate can detect a minute magnetic field as low as  $\sim$ 2×10<sup>-12</sup> T at resonance frequency, and  $\sim$ 1×10<sup>-10</sup> T at lower frequencies. © 2005 American Institute of Physics.

[DOI: 10.1063/1.2137455]

Magnetostriction occurs in most ferromagnetic materials. Rare-earth systems, such as Tb<sub>0.3</sub>Dy<sub>0.7</sub>Fe<sub>2</sub>, exhibit a giant Joule magnetostriction at relatively low magnetic biases. However, these rare-earth materials are expensive. It is also commonly known that ordinary Fe has a small magnetostrictive strain  $(3/2)\lambda$ , on the order of 30 ppm. However, the introduction of Ga into crystalline solution with Fe results in a significant enhancement of its magnetostriction,  $(3/2)\lambda_{100}$  $\approx 250$  ppm, as long as the distorted AO2 phase remains stable—even though Ga reduces the spin density of the solution, it enhances its magnetostriction. Recently, a number of investigations have focused on Fe–Ga alloys (or Galfenol) due to the combination of its high mechanical strength, good ductility, relatively large  $(3/2)\lambda_{100}$  values, low saturation fields, high blocking stress, and low cost.<sup>2-7</sup> Galfenol has potential applications in acoustic projectors, acoustic sensors, and actuators.8,9

The magnetoelectric (ME) effect is a coupling effect between the magnetic field and electric field. ME effects have been found in many two-phase composites of magnetostrictive materials, such as  $Tb_{1-x}Dy_xFe_{2-y}$ , Permendur,  $Ni_{1-x}Co_xFe_2O_4$  (i.e., NFO), or  $Co_{1-x}Zn_xFe_2O_4$  (i.e., CFO), and piezoelectric materials, such as  $Pb(Zr,Ti)O_3$  or PMN-PT layers.  $^{10-21}$  In our prior article, we reported the magnetostrictive and magnetoelectric behavior of Fe-20 at.  $%Ga/Pb(Zr,Ti)O_3$  laminates.  $^{10}$  In this letter, we will show that  $Fe-Ga/Pb(Mg_{1/3}Nb_{2/3}O_3-PbTiO_3)$  (PMN-PT) ME laminates have much better ME coupling effect and higher sensitivity to a small magnetic signal.

Crystals of Fe–20 at. % Ga were grown by a Bridgman method at the Ames Laboratory. The Fe–20 at. % Ga crystals were cut into rectangular plates of dimensions  $12.7 \times 6 \times 1$  mm, and were oriented along the  $\langle 001 \rangle_c$ . The crystals were annealed at 1100 °C for 168 h, using heating and cooling rates of 10 deg/min, after which the crystals were considered to be in the "slow-cooled" state. We measured the lattice parameter of one of our Fe–20%Ga crystals (in the free condition) as a function of magnetic bias  $H_{\rm dc}$ , as shown

in Fig. 1. The lattice parameter increased with increasing H, reaching a maximum near  $H_{\rm dc}$ =300 Oe. Using these data, we then estimated the value of the magnetostrictive strain to be 262 ppm, consistent with prior results. Maximum strain was found when  $H_{dc}$  was applied along the longitudinal axis, or  $\langle 010 \rangle_c$ , of the plates. Accordingly, this longitudinal axis is here designated as the principal magnetization direction of our ME laminate design. Piezoelectric single crystals were (001)-oriented PMN-PT layers with (i) longitudinal polarization and (ii) transverse polarization. Two long-type Fe-20%Ga/PMN-PT three-layer laminates were prototyped by (i) sandwiching one longitudinally poled piezoelectric single crystal PMN-PT layer (dimensions of  $14.0 \times 6 \times 1 \text{ mm}^3$ ) between two longitudinally magnetized Fe-20%Ga layers, i.e., a (L-L) mode laminate; and, as a comparison, (ii) sandwiching one transversely poled piezoelectric single crystal PMN-PT layer (dimensions of  $14.0 \times 6 \times 0.5$  mm) between two longitudinally magnetized Fe-20%Ga layers, i.e., a (L-T) mode laminate. These configurations are similar to prior ME modes in Terfenol-D/PZT (Ref. 20) and CFO-NFO/

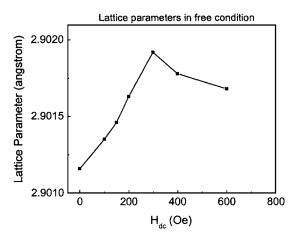


FIG. 1. Lattice parameter of a  $\langle 001 \rangle_c$  oriented Fe–20%Ga crystal as a function of dc magnetic bias. From the calculation of  $(\Delta L/L)$  using the magnetostriction, the maximum strain of a Fe–20%Ga crystal at H=300 Oe should be 262 ppm.

a) Electronic mail: sdong@mse.vt.edu

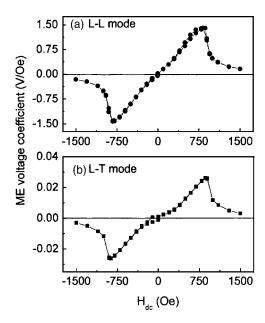


FIG. 2. ME voltage coefficient as a function of dc magnetic bias  $H_{\rm dc}$  for (a) an L-L laminate and (b) L-T laminate. The static capacitance (after assembles) of the longitudinally poled PMN–PT layers was  $\sim$ 7.8 pF.

ferroelectric lead zirconate titanate (PZT) (Refs. 16 and 19) laminates.

The voltages induced across the two electrodes of the PMN-PT layer in the Fe-20%Ga/PMN-PT laminate composite were measured for various  $H_{\rm dc}$  and ac magnetic drives  $(H_{\rm ac})$  over the frequency range of f=0.1-100 kHz, using a charge amplifier combined with a phase-referencing (i.e., lock-in) method. An electromagnet was used to apply a dc magnetic bias  $H_{\rm dc}$  and a pair of Helmholtz coils was used to generate a small  $H_{\rm ac}$ , via an input current  $I_{\rm coil}$ . For a ME composite with very low static capacitance, it necessary to use a charge amplifier to obtain correct induced ME voltages.

Figure 2(a) shows the ME voltage coefficients for the L-L mode,  $\partial V_{\text{ME}}^{L,L}/\partial H$ , as a function of magnetic field bias,  $H_{\rm dc}$ . These data were taken at frequency of f=1 kHz and a drive of  $H_{\rm ac}=1$  Oe. The value of  $\partial V_{\rm ME}^{L,L}/\partial H$  can be seen to be strongly dependent on  $H_{\rm dc}$ . In the dc magnetic bias range  $0 < H_{dc} < 700$  Oe, the ME voltages of the L-L mode of Fe-20%Ga/PMN-PT composites increased with increasing  $H_{\rm dc}$ , reaching a maximum ME effect of  $\partial V_{\rm ME}^{L,L}/\partial H$ ~1.41 V/Oe at  $H_{\rm dc}$ =750 Oe. (or, correspondingly,  $\partial E_{\rm ME}^{L,L}/\partial H$ =~1 V/cm Oe). For  $H_{\rm dc}$ >750 Oe,  $\partial V_{\rm ME}^{L,L}/\partial H$  decreased dramatically with increasing  $H_{dc}$ , as the Fe-Ga layers of the laminate approached saturation of its magnetostriction. It is relevant to note that the maximum value of  $\partial V_{\rm MF}^{L,L}/\partial H$  that we report here for Fe–20%Ga/PMN–PT laminates is  $\sim 4 \times$  higher than that for L-L configurations of Fe-20%Ga/PZT. 10 It is also comparable to that of L-L or C-C configurations of Terfenol-D/PMN-PT, 20,22 and  $\sim$ 10× higher than that for *L-T* configurations of Terfenol-D/ PZT or Terfenol-D/PMN-PT ones. 21 A long-type laminate favors the optimum combination of magnetostrictive and piezoelectric effects; in particular, the longitudinal magnetostrictive strain of Fe-20%Ga, and the longitudinal piezoelectric strain of PMN-PT are higher, than the corresponding transverse ones.

As a comparison, Fig. 2(b) shows the ME voltage coefficients,  $\partial V_{\text{ME}}^{L,T}/\partial H$ , for the *L-T* mode of Fe–20%Ga/PMN–PT laminates as a function of  $H_{\text{dc}}$ . Again, these data were also

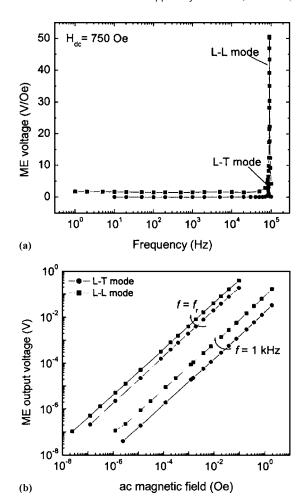


FIG. 3. ME responses of Fe–Ga/PMN–PT laminates, (a) induced ME voltage as a function of frequency of the ac magnetic field for both L-L and L-T modes. These data were taken under a constant dc magnetic bias of  $H_{\rm dc}$ =750 Oe, and ac magnetic field of  $H_{\rm ac}$ =1 Oe; and (b) induced ME voltages for L-L and L-T modes at low-frequency (1 kHz) and resonance frequency ( $f_r$ ) under a  $H_{\rm dc}$ =750 Oe as a function of ac magnetic field over the range of  $10^{-8} < H_{\rm ac} < 10$  Oe.

taken at a frequency of f=1 kHz and a drive of  $H_{\rm ac}=1$  Oe. The measured value of  $\partial V_{\rm ME}^{L,T}/\partial H$  can be seen to be  $\sim 30$  mV/Oe, which is a factor of  $\sim 40\times$  smaller than that of the L-L mode. However, its ME field coefficient,  $\partial E_{\rm ME}^{L,T}/\partial H=0.60$  V/cm Oe, is not much less than that of L-L mode.

We noted that ME voltage coefficients for both the L-L and L-T mode are antisymmetric about  $H_{\rm dc}$ . In our previous reports,  $^{10}$  we only illustrated the ME voltage amplitude as a function of  $H_{\rm dc}$ .

Magnetic field responses of the Fe–Ga alloy/PMN–PT laminates were then measured over a wider frequency range of  $10^2$ – $10^5$  Hz, as given in Fig. 3(a). The results show that the Fe–Ga/PMN–PT laminates have a much enhanced ME response when operated near its resonance frequency of  $f_r = \sim 86$  kHz. The maximum ME voltage coefficient at resonance for the L-L mode was  $\sim 50.7$  V/Oe (or 36.2 V/cm Oe for the field coefficient), which is  $\sim 36 \times$  higher than that in the low-frequency range; whereas the maximum ME voltage coefficient at resonance for the L-T mode was only 3.6 V/Oe (or  $\sim 70$  V/cm Oe).

Low-level magnetic field responses of the Fe–Ga/PMN–PT laminates operated in the L-T and L-L modes are given in Fig. 3(b). It can be seen that the induced ME voltages are near linear functions of  $H_{\rm ac}$ . In this figure, the in-

duced ME voltage can be seen to have a good linear response to  $H_{\rm ac}$  over a wide field range from  $10^{-12}$  T (or  $10^{-8}$  Oe) to  $\sim 10^{-3}$  T (or 10 Oe). The limit of magnetic field sensitivity for the L-L mode operated at resonance ( $f_r$ =8.6×10<sup>4</sup> Hz) was as low as  $2.3\times 10^{-8}$  (Oe), whereas its sensitivity at a low frequency of f=1×10<sup>3</sup> Hz was  $\sim 1.2\times 10^{-6}$  (Oe). As a comparison, Fig. 3(b) also illustrates the limit of magnetic field sensitivities for the L-T mode operated both at resonance and low frequencies. It can be seen that the limit of magnetic field sensitivities for L-T mode are 1.2  $\times 10^{-7}$  (Oe) at resonance, and  $\sim 2.5\times 10^{-6}$  (Oe) at low frequencies. These results demonstrate that our Fe–Ga/PMN–PT laminates were exceptional sensitive to minute magnetic field variations.

In summary, laminate composites of Fe–20%Ga/PMN–PT have been found to have (i) a large L-L ME voltage coefficient of  $\partial V_{\rm ME}^{L,L}/\partial H$ > 1.4 V/Oe, or  $\partial E_{\rm ME}^{L,L}/\partial H$ =  $\sim$  1.0 V/cm Oe, under modest dc magnetic biases; (ii) a large L-T ME field coefficient of  $\partial E_{\rm ME}^{L,T}/\partial H$ = 0.6 V/cm Oe, or  $\partial V_{\rm ME}^{L,T}/\partial H$ = 0.03 V/Oe, again under modest bias; (iii) a dramatic enhancement in the ME response near the resonance frequency; and (iv) a high sensitivity to minute magnetic field variations of <10<sup>-12</sup> T (10<sup>-8</sup> Oe) at resonance.

This research was supported by the Office of Naval Research under Grant Nos. N000140210340, N000140210126, and MURI N000140110761 and by the Office of Science, Materials Science Division, USDOE under Contract No. W-7405-ENG-82.

- <sup>3</sup>R. Q. Wu, J. Appl. Phys. **91**, 7358 (2002).
- <sup>4</sup>N. Srisukhumbowornchai and S. Guruswamy, J. Appl. Phys. **90**, 5680 (2001).
- <sup>5</sup>A. E. Clark, K. B. Hathaway, M. Wun-Fogle, J. B. Restorff, T. A. Lograsso, V. M. Keppens, G. Petculescu, and R. A. Taylor, J. Appl. Phys. **93**, 8621 (2003).
- <sup>6</sup>N. Srisukhumbowornchai and S. Guruswamy, J. Appl. Phys. **92**, 5371 (2002).
- <sup>7</sup>A. E. Clark, J. B. Restorff, M. Wun-Fogle, T. A. Lograsso, and D. L. Schlagel, IEEE Trans. Magn. **36**, 3238 (2000).
- <sup>8</sup>P. Mungsantisuk, R. Corson, and S. Guruswamy, Proceedings of Advanced Materials for Energy Conversion **2**, 275 (2004).
- <sup>9</sup>A. E. Clark, J. B. Restorff, M. Wun-Fogle, T. A. Lograsso, D. L. Schlagel, A. R. Ross, and D. Schlagel, *Proceedings of the Actuator 2000 Conference*, Bremen, Germany, June 2000.
- <sup>10</sup>S. Dong, J. Zhai, F. Bai, J. F. Li, D. Viehland, and T. A. Lograsso, J. Appl. Phys. **97**, 103902 (2005).
- <sup>11</sup>S. Di. Matteo and A. G. M. Jansen, Phys. Rev. B 66, 100402 (2002).
- <sup>12</sup>L. Wiegelmann, A. A. Stepanov, I. M. Vitebsky, A. G. M. Jansen, and P. Wyder, Phys. Rev. B **49**, 10039 (1994).
- <sup>13</sup>C. W. Nan, M. Li, and J. H. Huang, Phys. Rev. B **63**, 144415 (2001).
- <sup>14</sup>J. Ryu, A. Vazquez Carazo, K. Uchino, and H. Kim, Jpn. J. Appl. Phys., Part 1 40, 4948 (2001).
- <sup>15</sup>U. Lalestin, N. Padubnaya, G. Srinivasan, and C. P. Devreugd, Appl. Phys. A: Mater. Sci. Process. 78, 33 (2004).
- <sup>16</sup>G. Srinivasan, E. Rasmussen, B. Levin, and R. Hayes, Phys. Rev. B 65, 134402 (2002).
- <sup>17</sup>J. Ryu, S. Priya, K. Uchino, H. E. Kim, and D. Viehland, J. Korean Ceramic Society 39, 813 (2002).
- <sup>18</sup>S. L. Kadam, K. K. Patanka, V. L. Mathe, M. B. Kothale, and R. B. Kale, Mater. Chem. Phys. **78**, 684 (2003).
- <sup>19</sup>G. Srinivasan, E. T. Rasmussen, and R. Hayes, Phys. Rev. B 67, 014418 (2003).
- <sup>20</sup>S. Dong, J. F. Li, and D. Viehland, IEEE Trans. Ultrason. Ferroelectr. Freq. Control **50**, 1261 (2003); **50**, 1236 (2003); **51**, 794 (2004).
- <sup>21</sup>S. Dong, J. F. Li, and D. Viehland, Appl. Phys. Lett. **83**, 2265 (2003); **85**, 5305 (2004); **85**, 2307 (2004).
- <sup>22</sup>S. Dong, J. G. Bai, J. Zhai, J.-F. Li, G.-Q. Lu, S. Zhang, T. R. Shrout, and D. Viehland, Appl. Phys. Lett. **86**, 182506 (2005).

<sup>&</sup>lt;sup>1</sup>G. Engdahl, *Magnetostrictive Materials Handbook* (Academic, New York, 2000).

<sup>&</sup>lt;sup>2</sup>L. Y. Dai, J. Cullen, M. Wuttig, T. Lograsso, and E. Quandt, J. Appl. Phys. 93, 8627 (2003).