

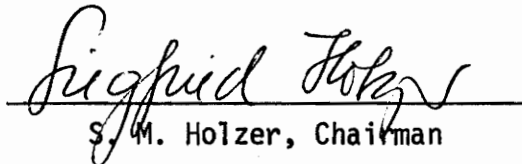
STABILITY OF RETICULATED DOMES UNDER MULTIPLE STATIC
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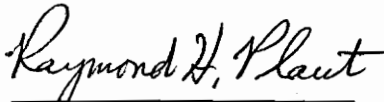
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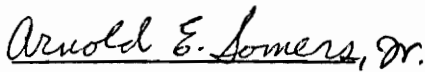
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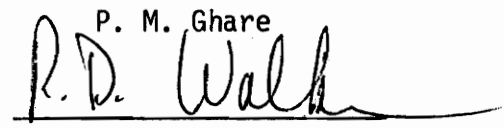
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STABILITY OF RETICULATED DOMES UNDER MULTIPLE
STATIC AND DYNAMIC LOADS

by

Ayodele Olushola Abatan

(ABSTRACT)

The primary purpose of this dissertation is to investigate the stability of reticulated domes under multiple static and dynamic loads. Two elastic geometrically nonlinear structural models of a reticulated dome with 21 and 39 degrees of freedom are considered.

The nonlinear response of the system to static loads is obtained using nonlinear programming and discrete perturbation techniques. The nonlinear programming technique is used to obtain a starting solution for the discrete perturbation technique and to optimize the choice of the perturbation parameter. Convergence criteria and error estimates to limit errors in a perturbation scheme are developed. A method for selecting a "suitable" perturbation parameter for imperfection sensitive systems is proposed.

The investigation of stability of equilibrium of the system subjected to finite disturbances is based on the concept of "degree of stability" and the associated sufficient stability condition. The stability condition is derived from a theorem on extent of asymptotic stability of Liapunov's direct method of the theory of stability of

motion. Its application requires the determination of the nonlinear fundamental path and the "nearest" unstable post-buckling path.

This is obtained via static analysis.

The perturbed motion of the system under a given set of perturbations is obtained by numerically integrating the nonlinear equations of motion. The dynamic stability tests confirm the sufficiency of the dynamic stability condition. However, they also indicate that there is a dynamic disturbance with a specific spatial distribution for which the sufficient condition of stability is also a necessary condition for each equilibrium state tested. Since in practice, the spatial distribution of the disturbances cannot be controlled, the sufficient dynamic stability condition employed is practical for the design of reticulated domes.

The stability boundaries corresponding to two independent loads on the models are presented. Limit points lie on a boundary which is convex towards the region of stability. Bifurcation points lie on a continuous but piecewise differentiable boundary. Each piece of the boundary containing bifurcation points appears to be convex towards the region of stability.

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NOTATION

A_i	cross-sectional area of i -th element
C_M	set containing the equilibrium state; a subset of the domain of asymptotic stability of the origin
E_i	modulus of elasticity of i -th element
$E^j(Q_i)$	displacements corresponding to loads Λ^j ; E^j are linear in Q_i
h	height of center of dome from supports
L_i	undeformed length of element i
L_i^*	deformed length of element i
ℓ^j	direction cosine from a loading ray in the load space to the load Λ^j
L	Liapunov function
L'	total derivative of L with respect to time
M	energy level at "nearest" unstable equilibrium state
NM	number of members (elements)
NJ	number of joints
O	Landau's error order symbol
P_s	generalized load corresponding to generalized displacement Q_s
P_j	three dimensional force vector at joint j .
P_j^T	transpose of vector P_j
Q_i	generalized displacements
\dot{Q}_i	time derivative of Q_i
$q_i(t)$	generalized velocity perturbation of i -th coordinate

q_i	generalized displacement-increment of i -th coordinate
\dot{q}_i	first derivative of q_i with respect to ϵ
R	regular matrix
s	span of the reticulated dome
T	kinetic energy function
t	time
t_d	duration of exciting force
U_j	three-dimensional displacement vector at joint j
V	total potential energy function
v	total energy; Liapunov function
v_0	value of v at time zero
W_L	limit point characteristic matrix
W_B	bifurcation point characteristic matrix
W_C	coincident critical point characteristic matrix
W_S	the s -th equilibrium equation
$W_{s,i}$	partial derivatives of the s -th equilibrium equation with respect to the i -th generalized coordinate
$W_{s,ij}$	second order partial derivatives of the s -th equilibrium equation with respect to the i -th and j -th generalized coordinates
$\det(W)$	determinant of matrix W
$y(t)$	state vector; resulting motion of system
y_0	value of y at time zero; initial state of system
y_1	generalized displacement vector of perturbed motion

y_2	generalized velocity vector of perturbed motion
$\ y\ $	vector norm; $(y^T y)^{1/2}$
y_1^u	unstable equilibrium states
Λ, λ	load parameters
Λ^j	j-th generalized independent load
Λ_i^c	critical value of Λ_i
Λ_{ok}	critical value of Λ^k when acting singly
ϵ	perturbation parameter
η	increment in load parameter λ
$\dot{\eta}$	first derivative of η with respect to ϵ
π	potential energy; strain energy of system
π_i	the strain energy of the i-th element
Ω	potential energy of external forces
δ	limit of errors in a perturbation scheme
γ	unbalanced joint forces; small positive constant
\in	is an element of

CHAPTER 1

INTRODUCTION

1.1 Purpose and Scope

The primary purpose of this thesis is to investigate the stability of reticulated domes under multiple static and dynamic loads. An elastic geometrically nonlinear structural model of a reticulated dome is considered. Geometric nonlinearities are due to admissible finite joint displacements. The pre- and post-buckling paths of the imperfection-sensitive system are obtained via nonlinear programming and discrete perturbation techniques. A method for selecting a "suitable" perturbation parameter is proposed.

Stability of equilibrium of the system subjected to instantaneous and transient disturbances is established using the concept of "degree of stability" [45]* and the associated sufficient stability condition. The sufficient stability condition is based on Liapunov's direct method [37].

The perturbed motion of the system for a given class of perturbations is obtained by numerically integrating the nonlinear differential equations of motion. The characteristics of the resulting phase-plane portraits provide two useful guides: 1. a check on the sufficient stability condition, and 2. an indication when the sufficient condition is also necessary for the stability of equilibrium.

Stability boundary curves [59] are presented for the reticulated

*Numbers in brackets refer to the references listed in the bibliography.

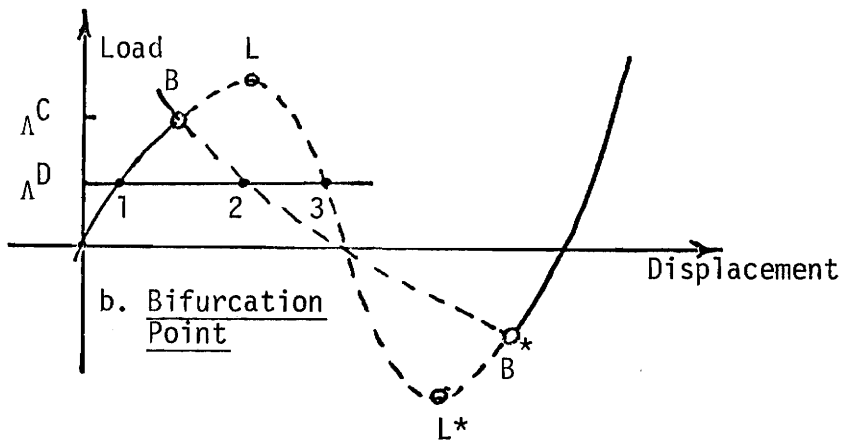
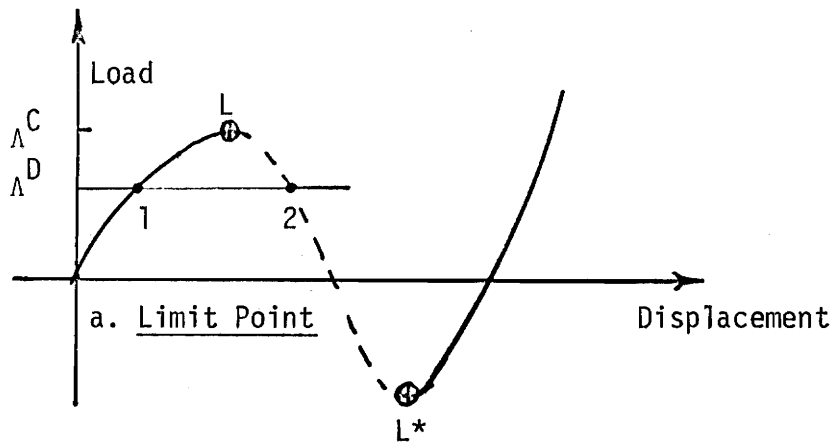
dome model to indicate the effect of several independent external loads on the stability of equilibrium.

1.2 "Nearest" Unstable Equilibrium States

In this section, the location of the "nearest" unstable equilibrium states is discussed. The unstable post-buckling path, emanating from an unstable critical point, which is closest to the fundamental path [114] contains the "nearest" unstable equilibrium points.

In this study, the potential energy difference between the stable prebuckling state and the "nearest" unstable equilibrium state is used as a measure of degree of stability of equilibrium [45]. The degree of stability is a quantitative measure of the disturbances an equilibrium state can sustain; it is a safe estimate of the domain of asymptotic stability of the equilibrium state whose stability is in question [45]. Hence the location of the "nearest" unstable equilibrium states with high precision is central to the correct usage of the stability criterion. The concept of degree of stability represents a practical stability criterion.

The nature and type of the critical point in configuration space determines the form of the post-buckling response. Post-buckling paths beyond a limit point, bifurcation point, and a coincident critical point are shown in Fig. 1.1. Definitions of these critical points are given by Thompson [109]. In all three cases shown in Fig. 1.1, point 2 represents the "nearest" unstable equilibrium state with respect to the stable equilibrium point 1. Λ^C , Λ^D represent the critical



— Stable paths

- - - Unstable paths

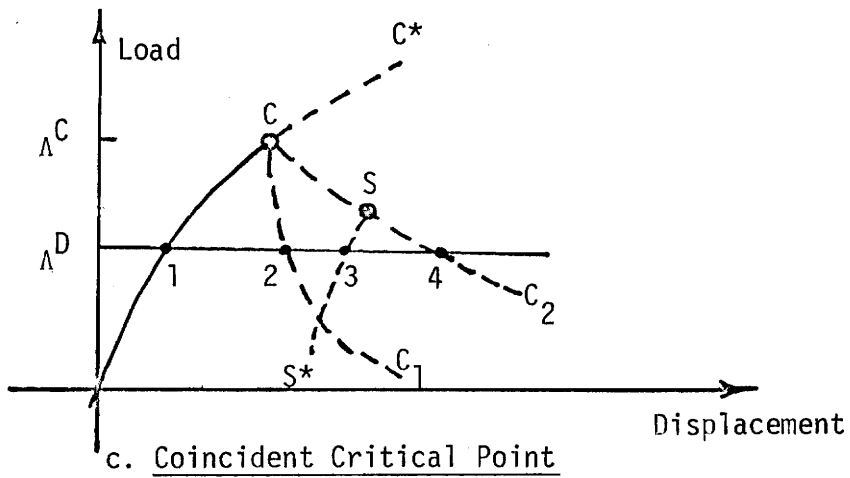


FIG. 1.1 "NEAREST" UNSTABLE EQUILIBRIUM PATHS

load, design load, on the system, respectively. In the case of a limit point, Fig. 1.1a, the unstable post-buckling path, LL^* , is unique. If a bifurcation point B (Fig. 1.1b) exists, two unstable paths, BB^* and BLL^*B^* , exist. For some loading conditions, point 2 could be to the right of point 3. In that case, point 3 represents the "nearest" unstable state. The case of a coincident critical point is shown in Fig. 1.1c. For the case of two coincident critical loads, paths CC^* , CC_1 , CC_2 are feasible. Again, the location of point 2 on CC_1 is required to use the sufficient stability criterion. Secondary bifurcation [7] from an unstable post-buckling path (e.g. path SS^*) is not considered.

The static perturbation method [92], a systematic approach, is used to locate all post-buckling paths beyond unstable critical points of the reticulated dome model used.

1.3 Literature Review

There are many factors affecting the stability of structures (Fig. 1.2). Researchers in the field of structural stability have considered several of these factors in selecting mathematical models for several classes of structures. In this section, major works relating to the present study are reviewed and classified on the basis of modelling and solution characteristics. Special attention is given to conservative, geometrically nonlinear structural models. The review contains three sub-sections: stability concepts and theories, solution techniques in stability analysis, and stability of reticulated space structures. The review concludes with a justification and need for research on stability of reticulated domes subjected to multiple

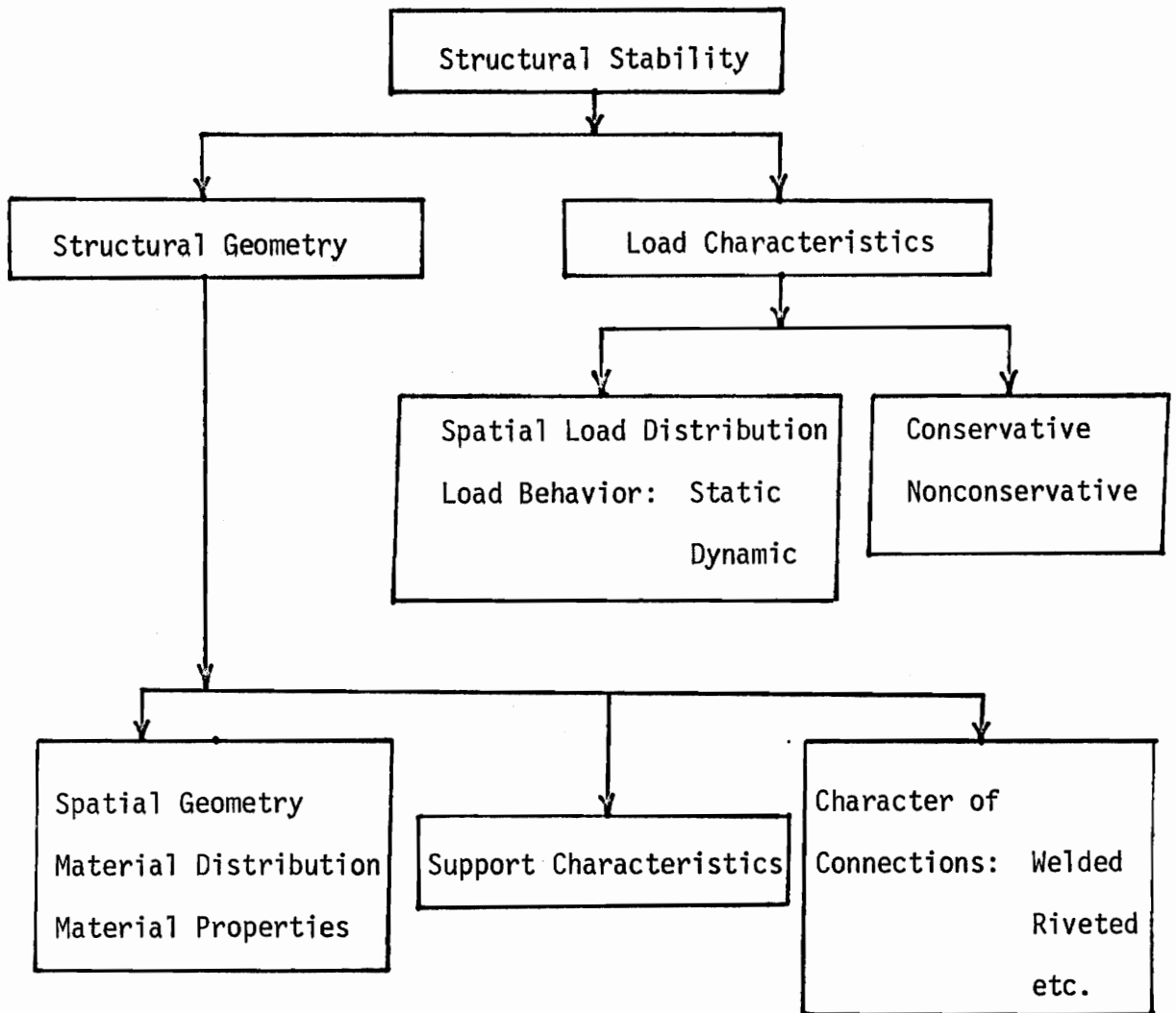


FIG 1.2 FACTORS AFFECTING STRUCTURAL STABILITY

static and dynamic loads.

a. Stability Concepts and Theories

Over two hundred years ago, Euler solved the first problems of elastic instability. He considered the lateral buckling of compressed members [116]. Euler's method permits the question of stability to be replaced by the question of bifurcation of the equilibrium modes [13]. Other eminent scientists such as Lagrange, Dirichlet, Bryan, and Poincaré clarified the connection between the Euler method and the stability concept [13]. Although the Euler method yields accurate results for the stability analysis of conservative elastic systems [116], stability analysis should be based on the theory of stability of motion [12]. The need to utilize dynamic analysis in stability investigations was stressed by Ziegler [122] and later by others [71, 58]. For conservative (divergence) type structures, a static approach to stability is permissible [13, 71] since nontrivial positions of equilibrium always determine the limit of stability. For flutter-prone structures [12, 58], the dynamic approach to stability is essential in order to detect the loss of stability by flutter. Published erroneous results (e.g. Pflüger's static stability analysis of an elastic cantilever rod under a follower load [13]) emphasize the limitations of static analyses in stability investigations.

A rigorous definition of stability of equilibrium (motion) is due to Liapunov [37], whose treatise on the theory of stability of motion (equilibrium) was published in Russia in 1892. Liapunov's direct method

[37] represents a powerful tool for the solution of stability problems and has been successfully used, e.g., [12, 43, 37, 82]. Recently, the method has been used to establish a rigorous basis for the concept of degree of stability of equilibrium of geometrically nonlinear discrete systems subjected to finite aperiodic disturbances [45]. It is the basis of a sufficient stability condition proposed by Hsu [50] and has appeared implicitly in other definitions of dynamic stability [e.g. 18, 54]. A review article [47] places the existing dynamic stability criteria for elastic imperfection-sensitive shells in the framework of Liapunov's theory of stability of motion (equilibrium).

Several other stability theories have been proposed in the literature. Extensive review of some of these theories appears in a paper by Hoff [41] and by several others [e.g. 19, 18, 36, 46]. In the particular case of shallow spherical shells [61, 87, 115], most of these theories are in disagreement with each other and with existing experimental data. Karman and Tsien [61] showed in 1939 that the large discrepancies between experiment and theory for shell structures were due to the highly-unstable post-buckling behavior of these structures. Since then, researchers in shell stability analysis have presented more accurate experimental and theoretical approaches to account for the discrepancies.

In 1945, Koiter, in a classic thesis [65], presented a general theory of stability for elastic systems by means of asymptotic expansion of the potential energy of the system. Koiter's theory received little attention until recently [66] when Thompson [106, 107], and Sewell [92, 94] developed the basic concepts of elastic stability for a discrete

system described by N generalized coordinates and a variable loading parameter. These researchers in their theories expanded the potential energy function as a power series and conducted an extensive study relating to the occurrence of limit and bifurcation critical states. This "static perturbation method" [92] became a powerful mathematical tool in a wide class of stability investigations [e.g. 93, 96, 107, 108, 110, 114, 59].

b. Solution Techniques in Stability Analysis

The solution techniques available in structural stability analysis are mostly numerical and/or approximate. This is because the analysis often involves systems of nonlinear equations with non-available exact solutions. Common solution techniques include the Newton-Raphson method [29], the Finite Element Method [68, 79, 80]. Nonlinear Programming [33, 73] and Static Perturbation Methods [92, 109]. The relative advantages and disadvantages of these methods are discussed by Sticklin, et. al. [99] and Oden [80].

The static perturbation method has received considerable attention lately in obtaining post-buckling response and measures of imperfection-sensitivity of structural models. Application of the method is reviewed by Thompson [109]. Walker [117] proposed the technique as a general method of solution for nonlinear simultaneous equations. Various perturbation methods are available in text books [10, 78]. Perturbation approaches have been used in various stability analyses both in the continuum context [2, 77, 85, 98, 108] and in the discrete context [114,

59]. A combined approach using the static perturbation method with other nonlinear analysis techniques appears in recent papers. Techniques used recently in conjunction with the perturbation method include the Newton-Raphson method [117], the Finite Element Method [79], and Nonlinear Programming Techniques [118].

For dynamic stability analysis, Liapunov's direct method [37] represents a consistent approach. This method provides a safe estimate of the domain of asymptotic stability of equilibrium. It can be used to study motion of the system "in the large" [75]. Liapunov's direct method does not require the solution of the differential equations for an investigation of system stability [75]. Its main drawback is the lack of general methods for the construction of Liapunov functions. A Liapunov-type stability analysis provides criteria which are lower bounds on the regions of stability [45, 75, 37]. For conservative systems, the total energy represents a suitable Liapunov function [45].

c. Stability of Reticulated Space Structures

A reticulated space structure is an assemblage of one-dimensional elements interconnected at a finite number of points [46]. The points of connections are structural joints. Because of their three-dimensional rigidity, flexibility and economy of construction, erection, and fabrication, reticulated space structures are widely accepted in the building industry [24]. However, because of their relative lightness in comparison with continuous solid shells, reticulated space structures pose fundamental stability problems under the influence of

static and dynamic loads. The stability problems which arise in reticulated space structures are of three types: bar stability (Euler critical load), snap-through buckling, and general buckling [46, 104]. Hence, under increasing external load, the change in the geometry of these structures becomes significant. Therefore, the structure exhibits a nonlinear load-deflection relationship [22].

In reticulated shells, local buckling leads to general buckling because the shell cannot resist shear loads caused by local buckling, and hence the area of buckling increases [24]. An extensive review of the stability of reticulated space structures is given by Holzer and Buchert [46].

In 1963, a large reticulated dome of 93.5 meters span, constructed in Bucharest as the roof structure of the National Economy Pavilion, suffered catastrophic failure after a fall of approximately one meter of fresh snow. This alerted practicing engineers and researchers to the importance of buckling problems of reticulated shell structures. Beles, et. al. [9] and Wright [119] examined the failure of the Bucharest dome and presented criteria for the buckling of reticulated shells using continuous, solid shell analogies. Other related papers appeared in a book on Space Structures [5].

The post-buckling behavior of reticulated shells is often studied by means of Koiter's general post-buckling theory [66]. This theory provides the possibility of determining the sensitivity of a structure to initial imperfections which may cause a great reduction in the

buckling load [19]. Application of Koiter's theory to reticulated shells leads to a system of linear equations for the pre-buckling state [104]. At the buckling state, a linear eigenvalue problem is solved. Recent analyses have shown [63] that the assumption of a linear pre-buckling state may lead to inaccurate results. Perturbation techniques have shown excellent results in solving eigenvalue problems with non-linear pre-buckling deformations [117, 118].

In spite of many studies on linear eigenvalue problems of discrete systems [e.g., 104, 90], there are relatively few publications on non-linear eigenvalue analyses [e.g., 6]. The basic difficulty lies in solving nonlinear simultaneous algebraic equations with large numbers of unknowns. Although the 'equivalent' continuous shell analogy for reticulated shell analysis has been developed and utilized in many practical designs [24, 46], one question still remains unanswered: to what extent can these analogies represent the true nature of general reticulated shells of arbitrary composition? Also, it is doubtful if these analogical concepts simulate the behavior of the original reticulated shell for localized large deflections common in the buckling of shells [38].

The present study addresses itself to the assessment of degree of stability of equilibrium [45] of reticulated domes under multiple static and dynamic loads. None of the studies cited above included the effect of several independent loads. Multiple load considerations are necessary for consideration of symmetrical as well as unsymmetrical load distribution on the dome. Unsymmetrical loads account for wind, snow,

and special cases.

In Chapter 2, the basic theorems and definitions used in the sequel are discussed. In Chapter 3, the solution techniques applied are discussed and analyzed in detail with special emphasis on the static perturbation method. Chapter 4 presents the error estimates and convergence criteria in a static perturbation scheme. In Chapter 5, a detailed account of the application of the theory to the stability analysis of geodesic domes is given. Chapter 6 presents the conclusions of the present work.

CHAPTER 2

STABILITY DEFINITIONS, THEOREMS, AND PRINCIPLES

The basic stability definitions, theorems, and principles pertinent to the present study are stated. Their implications and limitations with respect to the model used are discussed where necessary. The discussions on the concept of degree of stability follows those of references 44, 45, and 48. The practical implication of prior knowledge of convexity or concavity of stability boundaries in the case of multiple loads is stated.

2.1 Basic Principles and Theorems

Static and dynamic stability principles and theorems applicable to discrete systems are stated in this section. The discrete system is defined by a total potential function, $V(Q_i, \Lambda^j)$, and kinetic energy function, $T(Q_i, \dot{Q}_i, \Lambda^j)$; $i = 1, 2, \dots, N$, $j = 1, 2, \dots, M$. Q_i, Λ^j are the generalized coordinates (displacements), generalized independent external loads, respectively; N, M are the number of system degrees of freedom, number of generalized independent loads, respectively. At least in the region of interest, V and T are assumed to be single valued and continuous functions of the Q_i ; also $\partial V/\partial Q_i, \partial T/\partial Q_i$ and $\partial T/\partial \dot{Q}_i$ are assumed continuous in the same region. The system is holonomic and scleronomic, since there are no constraints on the variations of the coordinates, and V and T do not contain time explicitly [114].

The energy criterion and the principles of virtual work form the basis of the static stability analysis. In this regard, the following

principle is stated [75].

Virtual Work Principle: A mechanical system is in a state of equilibrium if, and only if, the virtual work of all forces (external and internal) vanishes for any virtual displacement.

This is a necessary and sufficient condition. It implies that V is stationary with respect to every kinematically admissible displacement, Q_i , from a state of equilibrium.

A Theorem, due to Lagrange, relating Liapunov's dynamical definition of stability [37] to the extremum properties of V is now stated.

Theorem 1: A complete relative minimum of the total potential energy with respect to the generalized coordinates is both necessary and sufficient for stability of equilibrium.

Lagrange's Theorem has not been proved in complete generality for the system under study [114, 59]. By introducing damping with a positive definite energy dissipation, Koiter established a general theorem [59].

Theorem 2: An equilibrium configuration in which the potential energy has no proper relative minimum is always unstable in the presence of damping with a positive definite energy dissipation.

In view of the presence of damping in actual physical systems, Theorem

2 is adequate for practical purposes [59]. Koiter [65, 66] in his studies on post-buckling behavior provided another theorem.

Theorem 3: The initial stage of post-critical behavior is completely specified by the stability or instability at the critical point itself.

Theorem 3 is used in determining the sensitivity of a system to imperfections. If the critical point is unstable, the associated system is imperfection-sensitive [47, 25]; i.e., the critical load of the perfect system provides an upper bound to that of the imperfect system.

Two fundamental theorems, due to Thompson [112], concerning the stability of equilibrium beyond bifurcation points are now stated.

Theorem 4: An initially stable (fundamental) equilibrium path rising monotonically with the loading parameter cannot become (thoroughly) unstable without intersecting a second distinct (post-buckling) equilibrium path.

Theorem 5: An initially stable (fundamental) equilibrium path rising monotonically with the loading parameter cannot approach an unstable critical equilibrium state (from which the system would snap dynamically) without the approach of a second distinct (post-buckling) equilibrium path at sub-critical values of the

loading parameter.

Topological and analytical proofs of Theorems 4 and 5 for a one degree-of-freedom system are given in [112]. Coincident critical points are excluded in the proofs.

For the discrete system under dynamic loads, Liapunov's definitions of stability, asymptotic stability, and instability of equilibrium are applicable. Let a state vector, $y(t)$, be defined as

$$y(t) = \begin{bmatrix} y_1(t) \\ \text{---} \\ y_2(t) \end{bmatrix} \quad (2.1)$$

where

$$y_1(t) = [q_i(t)], \quad y_2(t) = [\dot{q}_i(t)]; \quad (2.2)$$

i.e., $y_1(t)$, $y_2(t)$ are generalized displacement, velocity vectors, respectively, of the perturbed motion. The differential equation of perturbed motion [37] can be expressed in the $2N$ -dimensional phase-space as

$$y' = f(y) \quad (2.3)$$

Primes in Eqs. 2.2, 2.3 denote differentiation with respect to time.

The vector f is assumed to be continuous in a domain D of the $2N$ -dimensional Euclidean-space, E_{2N} . Let the trivial equilibrium state be an isolated equilibrium point. Consider an initial disturbance y_0 applied

at time $t = 0$. The Euclidean norm

$$||y|| = \sqrt{y^T y} = \sqrt{y_1^T y_1 + y_2^T y_2} \quad (2.5)$$

provides a measure of the deviation of the initial state, y_0 , and the resulting motion, $y(t)$, from the equilibrium state. The following definitions [37, 75] apply to the resulting motion $y(t)$.

Definition 1: The equilibrium state $y = 0$ is **STABLE** if for any $\epsilon^* > 0$, there exists a $\delta^*(\epsilon^*) > 0$ such that if $||y(0)|| < \delta^*$, then $||y(t)|| < \epsilon^*$ for all $t \geq 0$.

Definition 2: The equilibrium state $y = 0$ is **ASYMPTOTICALLY STABLE** if it is stable and if $\lim_{t \rightarrow \infty} ||y(t)|| = 0$.

Definition 3: The equilibrium state $y = 0$ is **UNSTABLE** if there exists an $\epsilon_1^* > 0$ such that for arbitrarily small $\delta^* > 0$, there is a motion $y(t)$ for which $||y(0)|| < \delta^*$ and $||y(t)|| > \epsilon_1^*$ at some time t_1 .

Definition 1 implies that the motion can be confined to the domain $||y|| < \epsilon^*$ if the initial disturbance is restricted to the domain $||y|| < \delta^*$. Definition 2 is a stronger stability definition, and it implies that sufficiently small disturbances die out. While the concept of stability refers to all disturbances which are sufficiently

small, in the case of instability, Definition 3, it is only necessary that there exist one particular disturbance which leads to an unstable behavior of the system [75].

Of the two distinct methods proposed by Liapunov for investigating stability of equilibrium (motion), Liapunov's direct method [37] is the most powerful. The direct method requires no approximations and the equations of motion need not be solved. However, it involves Liapunov functions, $L(y)$. For conservative systems, the total energy, E , represents a suitable Liapunov function. This follows from Definition 4 and Theorems 6-8 stated below.

Definition 4: In the region D of E_{2N} , the Liapunov function $L(y)$ is said to be

1. Positive Definite if

$$L(y) > 0 \text{ for all } y \neq 0$$

and

$$L(y) = 0 \text{ at } y = 0$$

2. Negative Definite if

$$L(y) < 0 \text{ for all } y \neq 0$$

and

$$L(y) = 0 \text{ at } y = 0$$

3. Positive Semi-definite if

$$L(y) \geq 0 \text{ for all } y$$

4. Negative Semi-definite if

$$L(y) \leq 0 \text{ for all } y.$$

The rate of change of $L(y)$ during motion is denoted $L'(y)$; i.e., $L'(y) = dL(y)/dt$. For conservative systems, $L = E$ and $L'(y) = 0$. The following theorems, stated without proof, form the basis of Liapunov's direct method [75].

Theorem 6: If there is a function $L(y)$ such that $L(y)$ is positive definite and $L'(y)$ is negative semi-definite, then the equilibrium state $y = 0$ is STABLE

Theorem 7: If there is a function $L(y)$ such that $L(y)$ is positive definite and $L'(y)$ is negative definite, then the equilibrium state $y = 0$ is ASYMPTOTICALLY STABLE

Theorem 8: If there is a function $L(y)$ such that $L(y_0) > 0$ for some y_0 with arbitrarily small norm and $L'(y)$ is positive definite, then the equilibrium state $y = 0$ is UNSTABLE

To apply these theorems necessitates finding a suitable Liapunov function, $L(y)$. The lack of general methods of constructing functions $L(y)$ and the lack of uniqueness of such functions are the major drawbacks of the direct method. However, the solutions, $y(t)$, are not

required. What is needed is $L'(y)$ which comes directly from the equations of motion. Liapunov's theorems provide estimates for the region of asymptotic stability, the domain of attraction of the equilibrium state whose stability is in question. In [45], the theorem on extent of asymptotic stability of Liapunov's direct method was used to derive a subdomain of asymptotic stability for a discrete system. This region is identical to that obtained by Hsu [50]. Geometric interpretations [75, 37] of Theorems 6-8 are given in Fig. 2.1.

Liapunov's theorems permit a study of stability of equilibrium relative to finite disturbances. Hence, if the disturbances which are likely to occur during the life of a system do not cause the perturbed motion to escape the domain of asymptotic stability, equilibrium is stable relative to these disturbances [44, 45]. This is the notion of PRACTICAL STABILITY. This notion is different from that of INFINITESIMAL STABILITY where only stability of equilibrium relative to infinitesimal disturbances is considered. Classical stability analyses [116, 122] ensure only infinitesimal stability.

2.2 Degree of Stability

The discussion of the concept of degree of stability in this section follows that of [44, 45] directly. The degree of stability of equilibrium [45] is a measure of the disturbances an equilibrium state can sustain; i.e., the system continues to perform its intended function. The concept of degree of stability is particularly important in the design of imperfection-sensitive systems, e.g., reticulated domes.

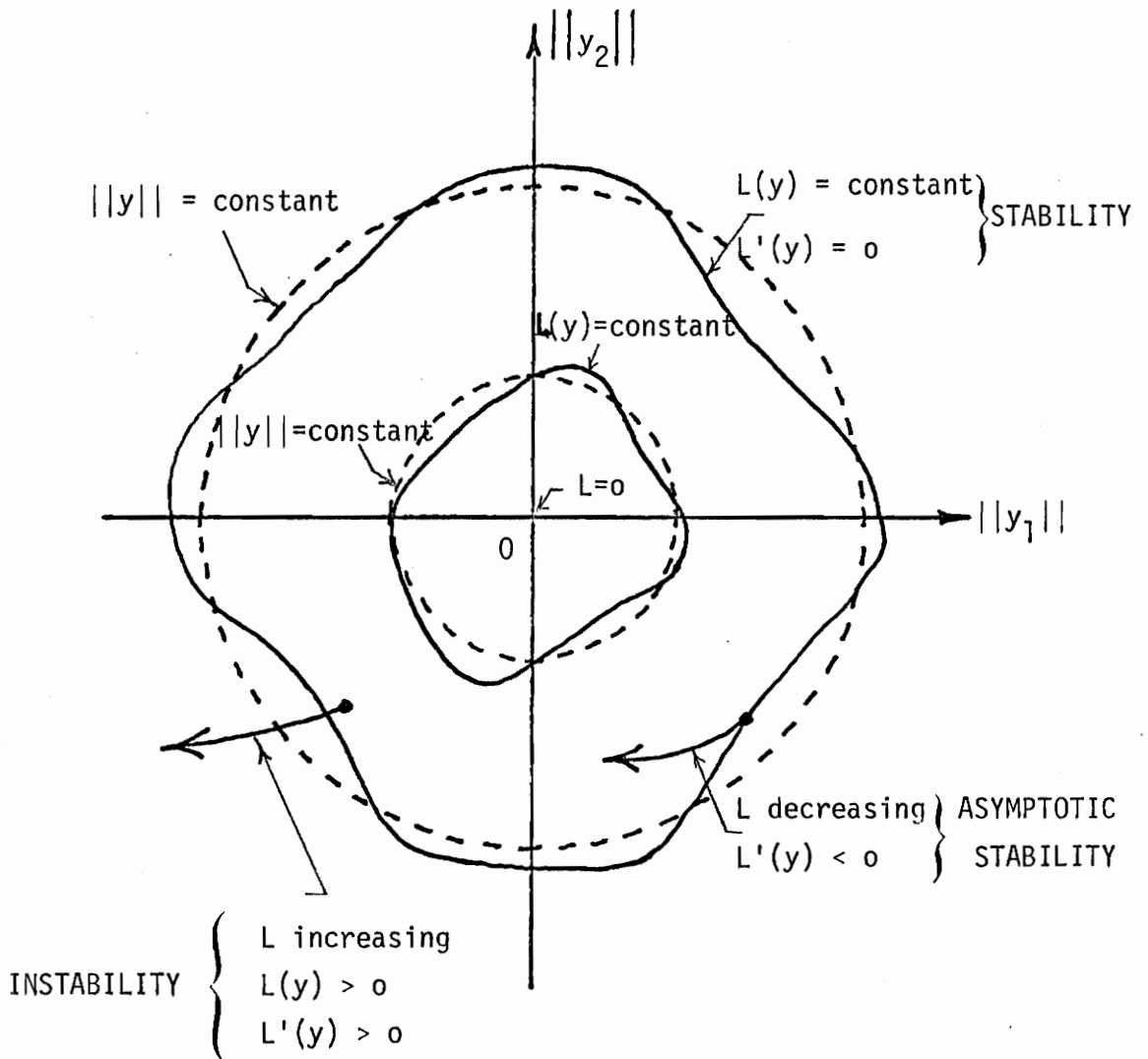


FIG. 2.1 GEOMETRIC INTERPRETATION OF LIAPUNOV'S THEOREMS

This is because imperfection-sensitive systems exhibit multi-equilibrium states below the buckling load, and sufficiently large disturbances can force the system from one equilibrium state to another [45]. It is therefore necessary to know (or at least obtain a safe estimate of) the degree of stability of equilibrium for imperfection-sensitive systems. With such knowledge, imperfection-sensitive systems can be designed for practical stability. In [45], the size of the region of asymptotic stability is identified as a direct measure of degree of stability.

Suppose the origin of the phase-space defined by Eq. 2.3 is asymptotically stable. The finite stability of the origin is in question. Let C_M be a subset of the region of asymptotic stability of the origin [45] having the property

$$C_M = \{y | v(y) < M\}. \quad (2.6)$$

In Eq. 2.6, $v(y)$ is the total energy of the system which is positive definite in C_M ; and M is the change in the potential energy between the stable origin and the "nearest" unstable equilibrium state, y_1^u . Mathematically

$$M = v(y_1^u, 0). \quad (2.7)$$

Hence, the origin is stable with respect to a specific disturbance if the resulting motion initiates in C_M and the total energy along this motion does not reach the value M . Analytically, the conditions

$$y_0 \in C_M \quad (2.8)$$

and

$$v(y) < M; 0 < t \leq t_d \quad (2.9)$$

are sufficient for the stability of the origin. In Eqs. 2.8, 2.9, y_0 and y are the initial state and the resulting motion, respectively; t_d denotes the duration of the disturbance. Instantaneous disturbances are defined in terms of initial conditions, and transient disturbances are expressed in terms of exciting forces of finite duration. Eqs. 2.8-2.9 form the basis of a sufficient stability condition proposed by Hsu [50] and it was shown [44] that other definitions of dynamic stability [e.g. 54, 18] are related to it. An application of this sufficient stability condition to the stability of equilibrium of reticulated domes under multiple static and dynamic loads is presented in Chapter 5.

2.3 Stability Boundaries - Multiple Loads

The concept of stability boundaries for discrete systems under several independent external loads is discussed in this section. The practical implications of prior knowledge of the shapes of the stability boundaries for such systems are mentioned.

In classical mechanics of discrete finite systems [114], the external loading on a structural model is represented by a single variable parameter. Thus, the primary target of the stability analysis is the determination of the critical value of this parameter. Such a representation occasionally breaks down [56]. The non-classical

studies of Huseyin on problems of multiple loading, which appeared in his recent book [59], emphasize the need to consider several independent loading parameters. The following discussion follows that in reference 59.

Consider the system defined in section 2.1 with the potential function $V(Q_i, \Lambda^j)$. For the system under the independent loads Λ^j , the location of the "stability boundary" corresponds to the location of the critical value of the single variable parameter used in the classical theory. The "stability boundary" is the locus of all critical points associated with an initial loss of stability in the load-space [56]. The load-space is the M-dimensional subspace of the (N+M) dimensional configuration space. The points of the load-space which correspond to states of critical equilibrium constitute the critical surface [107, 93]. The stability boundary refers to portions of the critical surface associated with an initial loss of stability. Hence, the projections of the critical zone (of the configuration space) onto the load space constitute the stability boundary. This is illustrated in Fig. 2.2.

The analysis of a discrete system subjected to independent external loads Λ^j is simplified by the consideration of paths on the equilibrium surface (Fig. 2.2) emerging from an unloaded state and corresponding to the rays [59]

$$\Lambda^j = \ell^j \lambda. \quad (2.10)$$

In Eq. 2.10, ℓ^j are direction cosines and λ is a loading parameter as

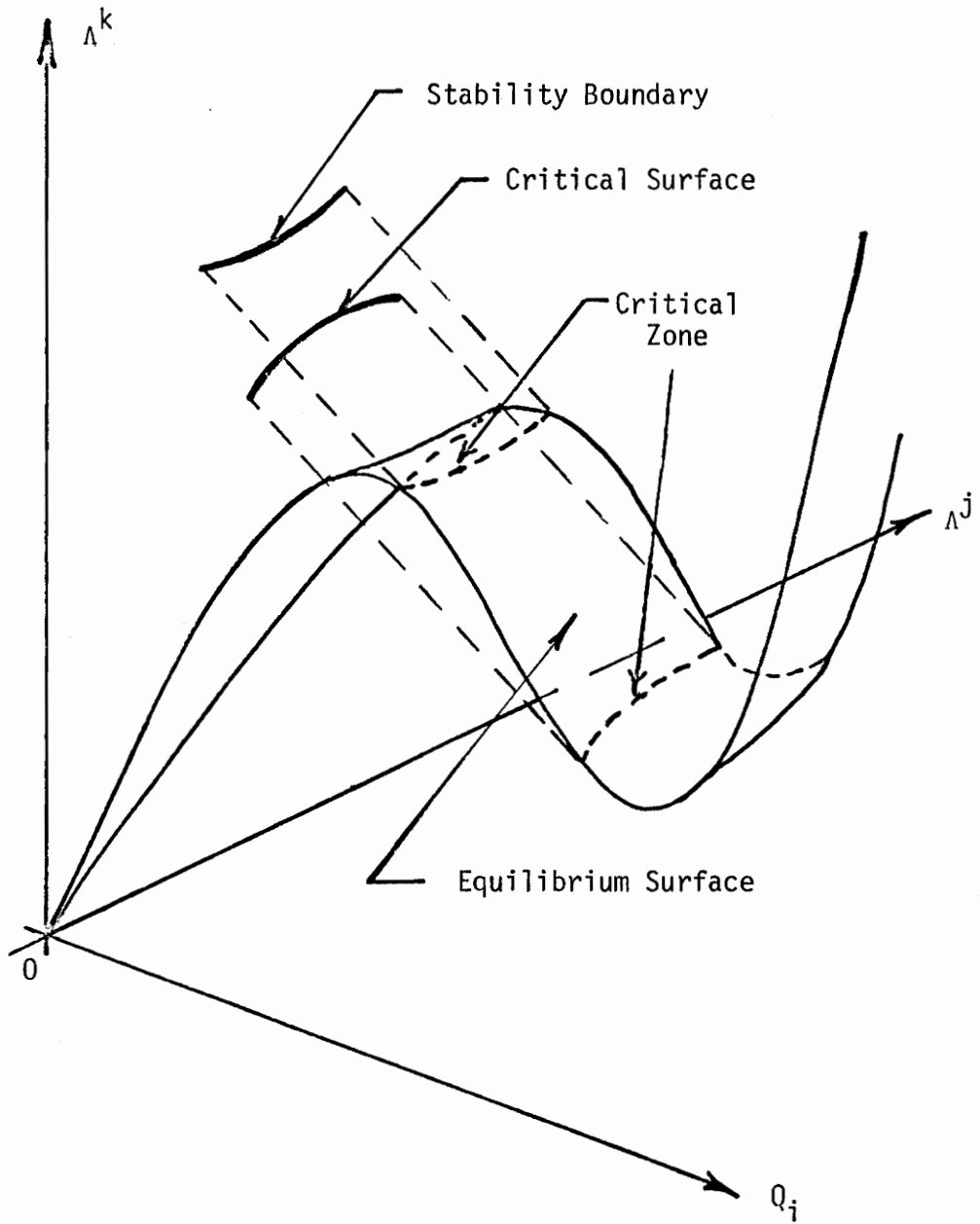
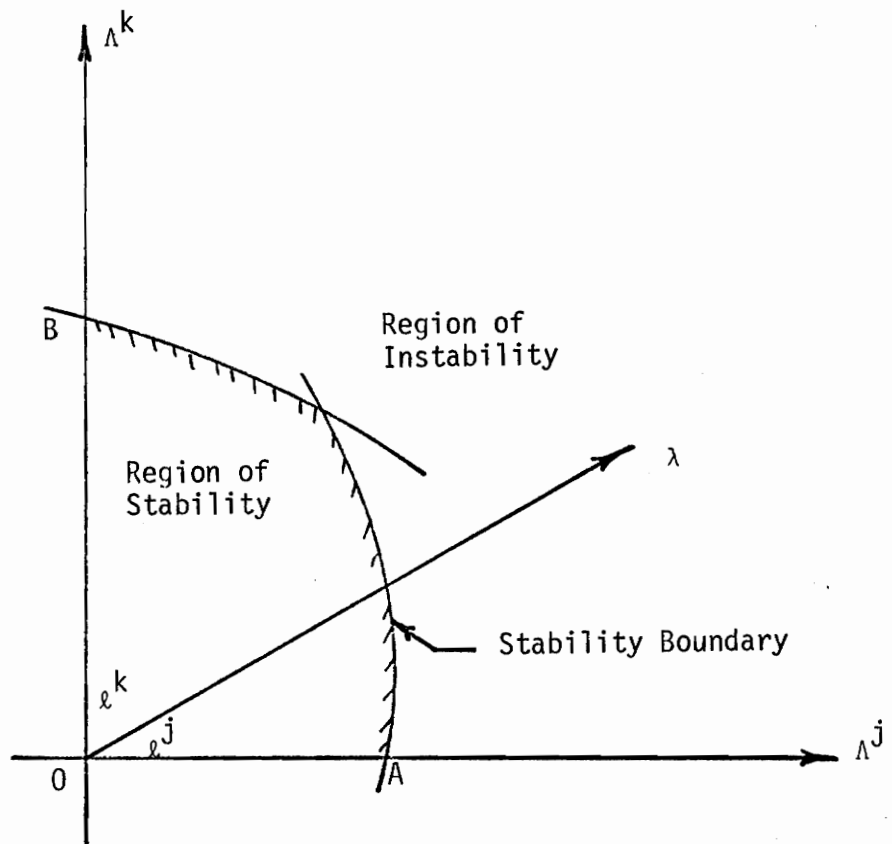


FIG. 2.2 STABILITY BOUNDARY AND EQUILIBRIUM SURFACE

shown in Fig. 2.3.

The concavity or convexity (i.e., the shapes) of the stability boundaries have practical implications. For example, if the stability boundary is concave (as in Fig. 2.4a) for two independent loads, Λ^j and Λ^k , then BA represents a lower bound, while lines BF* and AF* represent upper bounds on the region of stability. BF* and AF* are tangents to the boundary curve BA at B and A, respectively. The critical loads (points A and B) on the axes can be obtained without much computational effort since their location does not involve load interaction. Hence, at least a lower bound can be established immediately if the stability boundary for the system is known a priori to be concave towards the origin, 0. Similarly, if the stability boundary is convex towards 0, Fig. 2.4b, the line BA represents an upper bound on the region of stability. Again, point F* defining the lower bounds AF* and BF* can be obtained as in Fig. 2.4a. Fig. 2.4c depicts the case of a singular critical point F [59]. Singular points are associated with unstable symmetric points of bifurcation occurring along the loading ray of Fig. 2.3. In this case, the stability boundary takes the form of a cusp (point F, Fig. 2.4c). By joining F to the critical points A and B on the axes, the upper bounds AF and BF on the region of stability are obtained. The line BA represents a lower bound. A safe upper bound is defined by BF* and AF*, where F* connects tangents at B, A to the stability boundary curves BF, AF.

The location of the stability boundary poses serious fundamental difficulties for geometrically nonlinear systems [59]. However, for

FIG. 2.3 LOADING RAY λ IN LOAD-SPACE

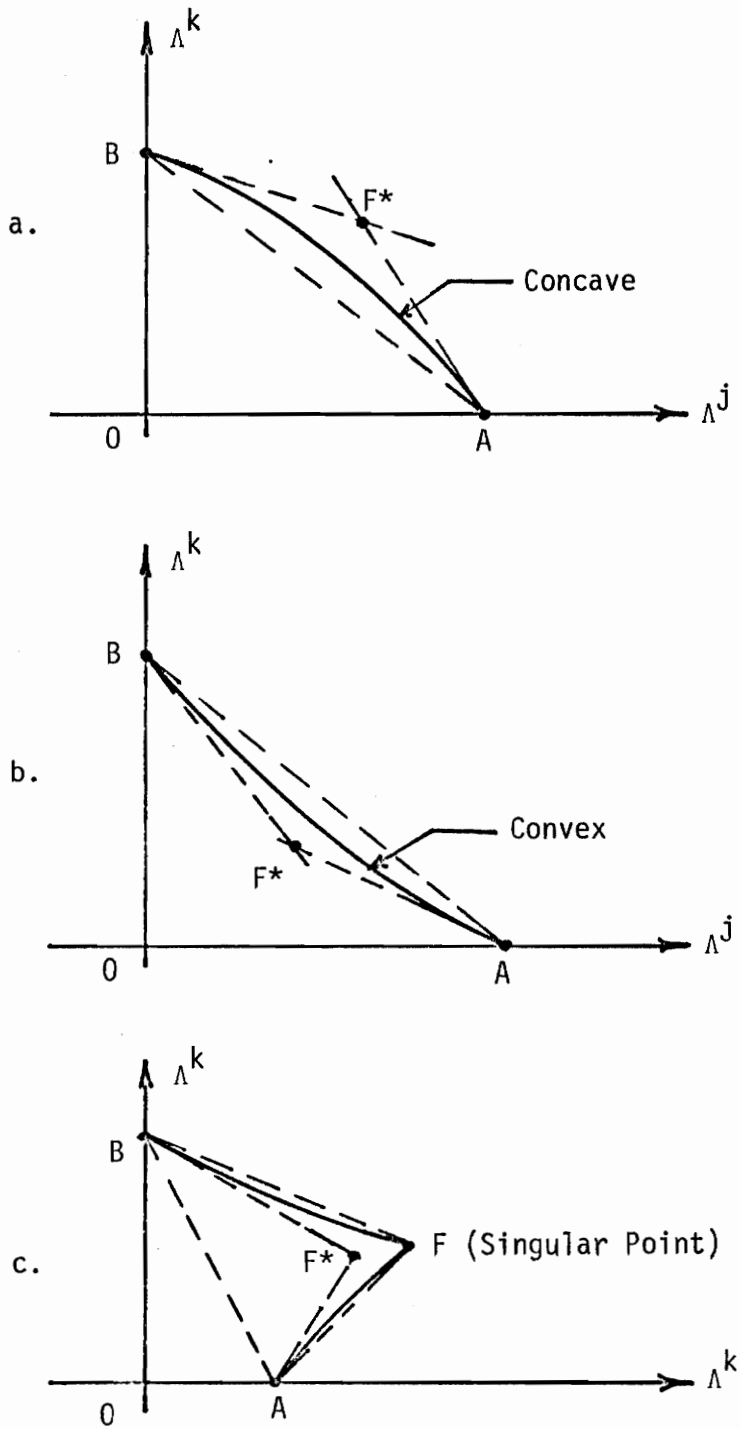


FIG. 2.4 CONVEXITY AND CONCAVITY OF STABILITY BOUNDARIES

some systems, estimates of stability boundaries exist [57], and some theorems have been proposed [59] for these systems.

Stability boundaries for reticulated dome models under two independent external loads are presented in Chapter 5.

CHAPTER 3

SOLUTION TECHNIQUES

The discrete perturbation, nonlinear programming, and dynamic analysis techniques used for obtaining the response of the reticulated dome model to static and dynamic loads are described in this chapter. Special attention is given to the discrete (static) perturbation technique, the key solution process for the determination of the nonlinear response to static loads. A non-diagonalized perturbation scheme [114] is used, in which the generalized displacements are selected as generalized coordinates. This scheme is compared with Thompson's [110, 114], where a set of "sliding" coordinates is introduced as generalized coordinates.

3.1 Discrete Perturbation Method

In this section, the concepts underlying the discrete perturbation technique are introduced. A general analysis is given, and the modifications necessary to obtain response paths beyond critical equilibrium points are described in detail.

a. Introduction

The discrete (static) perturbation method [92] is a parameter perturbation technique. In this technique, the solution is represented by the first few terms of an asymptotic expansion. It is analogous to the continuum perturbation technique [78, 65]. For the elastica problem, it was shown [108] that the contraction of the equilibrium

equations in the discrete perturbation method is equivalent to the removal of the secular terms (sources of nonuniformity) in the corresponding continuum analysis. In the discrete perturbation method (Fig. 3.1), the continuum potential function is approximated (discretized) in terms of "appropriate" generalized coordinates. The generalized coordinates reflect response characteristics of the system. Stationarity conditions applied to the approximate total potential function yield a series of nonlinear algebraic equations. A perturbation analysis in terms of a "suitable" parameter is then conducted to obtain the solution to the set of nonlinear algebraic equations. The result is a series of linear, ordered perturbation equations which can be solved recursively. Such an analysis provides the locations of equilibrium points satisfying the nonlinear differential equations of equilibrium. Hence, no information is obtained regarding the stability of the equilibrium states directly. However, the properties of the obtained equilibrium states (e.g., generalized displacements, etc.) can be used in a stability analysis based on higher variations (usually the second variation) of the total potential energy of the system. When the total potential energy assumes a relative minimum, Lagrange's theorem [69] states that the equilibrium state is stable.

b. General Analysis

A general discrete perturbation analysis is presented in this section. The presentation follows the works of Sewell [92, 94], Thompson [114], and Huseyin [56].

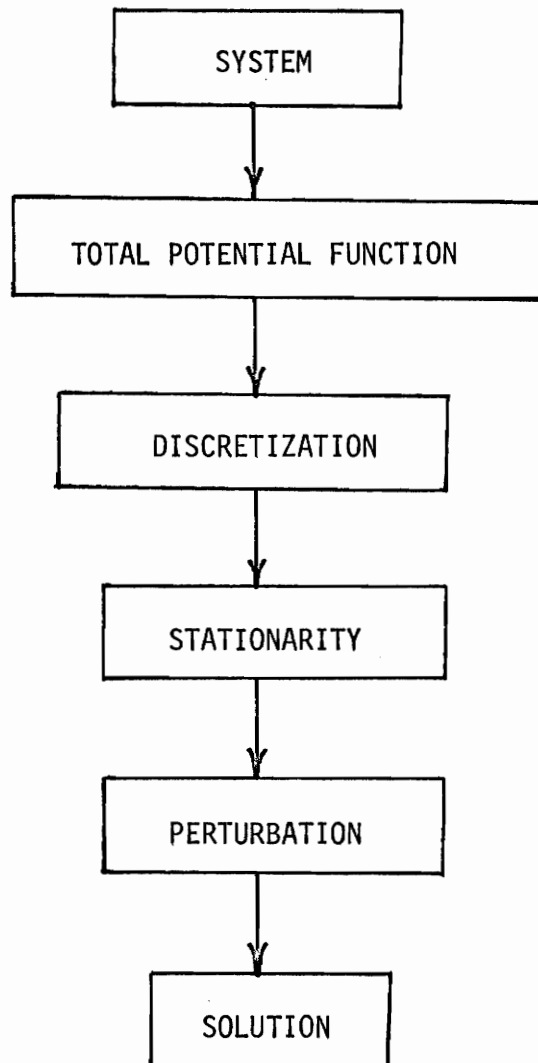


FIG. 3.1 DISCRETE PERTURBATION METHOD

Consider the discrete, conservative system defined in section 2.1. It is assumed that the equilibrium point (Q_i, Λ^j) in configuration space (Fig. 3.2) is known. Using the static perturbation technique, the location of a neighboring equilibrium point $(Q_i+q_i, \Lambda^j+\lambda^j)$ is sought. The search is conducted by considering incremental perturbations q_i in the generalized displacements, Q_i , and λ^j in the independent loads, Λ^j . The equilibrium equations for the known point are

$$\left. \begin{aligned} W_s(Q_i, \Lambda^j) &= 0; \quad i = 1, 2, \dots, N \\ & \quad j = 1, 2, \dots, M \\ & \quad s = 1, 2, \dots, N \end{aligned} \right\} \quad (3.1)$$

where

$$W_s = \frac{\partial V}{\partial Q_s} \quad (3.2)$$

Correspondingly, at $(Q_i+q_i, \Lambda^j+\lambda^j)$, the equilibrium equations

$$W_s(Q_i+q_i, \Lambda^j+\lambda^j) = 0 \quad (3.3)$$

are satisfied. In Eqs. 3.1-3.3, N and M are the system degrees of freedom and number of independent loads, respectively, and V is the total potential energy function. V satisfies all properties defined in section 2.1.

The concept of a loading ray (section 2.3) in the M -dimensional load space with n representing the distance along the ray from Λ^j to $\Lambda^j + \lambda^j$ leads to the relation

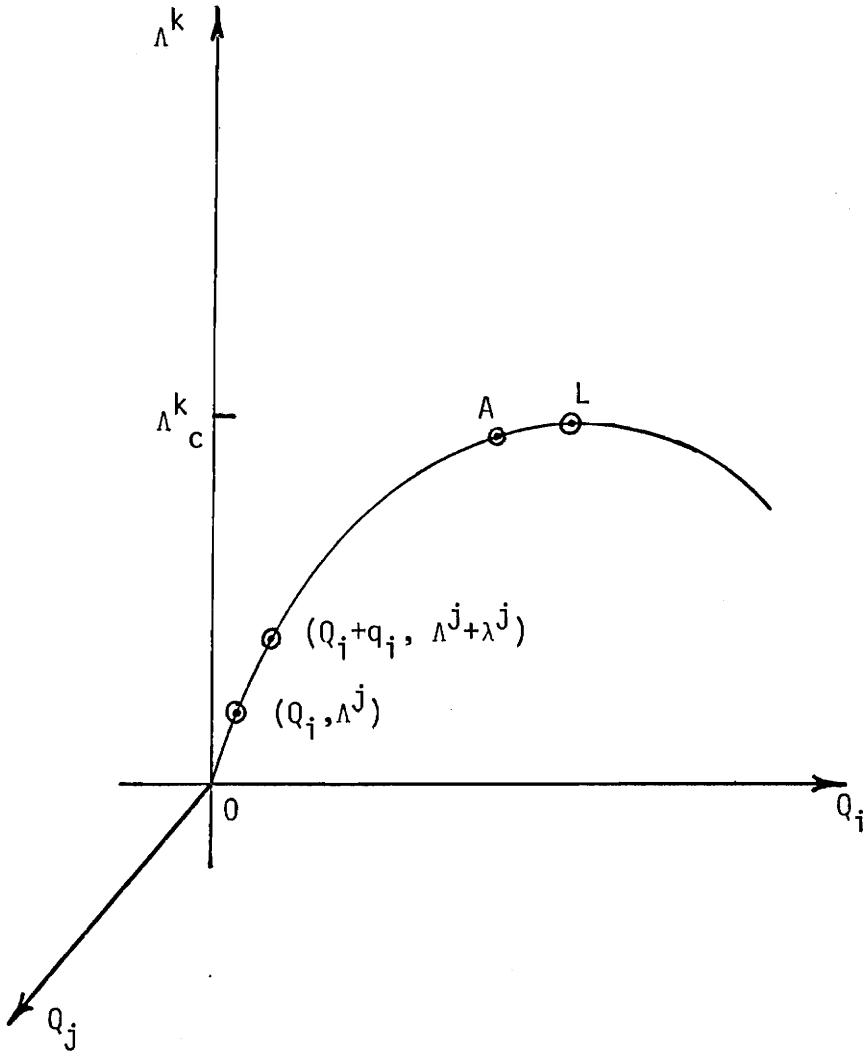


FIG. 3.2 EQUILIBRIUM PATH

$$\lambda^j = \lambda^j \eta \quad (3.4a)$$

for a path with proportional loading (Fig. 2.3). From Eqs. 2.10 and 3.4a, one obtains

$$\Lambda^j + \lambda^j = \lambda^j (\lambda + \eta) \quad (3.4b)$$

Eq. 3.4b transforms Eqs. 3.3 into the form

$$W_S [Q_i + q_i, \lambda^j (\lambda + \eta)] = 0 \quad (3.5)$$

which is the equilibrium equation for the desired equilibrium state defined by parameters q_i and η .

Expanding Eqs. 3.3 about (Q_i, Λ^j) in a Taylor series expansion with increments (q_i, η) gives

$$\begin{aligned} W_S(Q_i + q_i, \Lambda^j + \lambda^j) &= W_S(Q_i, \Lambda^j) + (W_{S,i} q_i + W_{S,\lambda} \lambda^n) \\ &+ \frac{1}{2!} (W_{S,ij} q_i q_j + 2W_{S,i\lambda} q_i \eta + W_{S,\lambda\lambda} \eta^2) \\ &+ \frac{1}{3!} (W_{S,ijk} q_i q_j q_k + 3W_{S,ij\lambda} q_i q_j \eta \\ &+ 3W_{S,i\lambda\lambda} q_i \eta^2 + W_{S,\lambda\lambda\lambda} \eta^3) \\ &+ O[q_i^4, \eta^4]. \end{aligned} \quad (3.6a)$$

where

$$\left. \begin{aligned} W_{s,i} &= \frac{\partial W_s}{\partial Q_i} \\ W_{s,\lambda} &= \frac{\partial W_s}{\partial \lambda} \\ W_{s,i\lambda} &= \frac{\partial^2 W_s}{\partial Q_i \partial \lambda} ; \dots \end{aligned} \right\} \quad (3.6b)$$

All derivatives of W_s are evaluated at $q_i = 0$, $\lambda = 0$ in Eqs. 3.6 and thereafter; summation convention is utilized. The symbol O in Eqs. 3.6a represents an error term of order q_i^4 or η^4 . It is Landau's error symbol [78]. Eqs. 3.1 and 3.3 together with Eqs. 3.6 lead to the following set of nonlinear differential equations

$$\begin{aligned} & (W_{s,i} q_i + W_{s,\lambda} \eta) \\ & + \frac{1}{2!} (W_{s,ij} q_i q_j + 2W_{s,i\lambda} q_i \eta + W_{s,\lambda\lambda} \eta^2) \\ & + \frac{1}{3!} (W_{s,ijk} q_i q_j q_k + 3W_{s,ij\lambda} q_i q_j \eta \\ & + 3W_{s,i\lambda\lambda} q_i \eta^2 + W_{s,\lambda\lambda\lambda} \eta^3) + O[q_i^4, \eta^4] = 0. \end{aligned} \quad (3.7)$$

The simultaneous solution of Eqs. 3.7 can be described in parametric form by the single-valued functions

$$q_i(\epsilon) = \epsilon \dot{q}_i + \frac{1}{2!} \epsilon^2 \ddot{q}_i + \frac{1}{3!} \epsilon^3 \dddot{q}_i + O[\epsilon^4] \quad (3.8)$$

$$\eta(\epsilon) = \epsilon \dot{\eta} + \frac{1}{2!} \epsilon^2 \ddot{\eta} + \frac{1}{3!} \epsilon^3 \dddot{\eta} + O[\epsilon^4]. \quad (3.9)$$

In Eqs. 3.8 and 3.9, ϵ is a "suitable" perturbation parameter, and dots denote differentiation with respect to ϵ ; i.e.

$$\dot{q}_i = \frac{\partial q_i}{\partial \epsilon}, \quad \ddot{\eta} = \frac{\partial^2 \eta}{\partial \epsilon^2}, \quad \text{etc.} \quad (3.10)$$

The selection of a "suitable" perturbation parameter is discussed in section 4.3. The known state (Q_i, Λ^j) corresponds to $\epsilon = 0$.

There are two fundamental assumptions [92] in the Taylor series representations defined by Eqs. 3.8 and 3.9. First, a non-zero radius of convergence is assumed (for $\epsilon > 0$) to rule out the possibility of converging to the known state (Q_i, Λ^j) . Secondly, the parametric representations $q_i = q_i(\epsilon)$, $\eta = \eta(\epsilon)$ are assumed to be such that not all $\dot{q}_i, \dot{\eta}$ vanish simultaneously at $\epsilon = 0$.

Eqs. 3.8 and 3.9 transform Eqs. 3.7 into

$$\epsilon \alpha_1 + \epsilon^2 \alpha_2 + \epsilon^3 \alpha_3 + O[\epsilon^4] = 0 \quad (3.11)$$

where

$$\alpha_i = \alpha_i(\dot{q}_k, \dot{n}, \dots); i = 1, 2, 3. \quad (3.12)$$

These equations must hold for all values of ϵ . Since sequences of ϵ are linearly independent, each coefficient $\alpha_i \neq \alpha_i(\epsilon)$ must vanish independently. Hence, the following linear, ordered perturbation equations are obtained:

$$\epsilon^1: \quad W_{s,i} \dot{q}_i + W_{s,\lambda} \dot{n} = 0 \quad (3.13)$$

$$\epsilon^2: \quad W_{s,i} \ddot{q}_i + W_{s,\lambda} \ddot{n} + W_{s,ij} \dot{q}_i \dot{q}_j + 2W_{s,i\lambda} \dot{q}_i \dot{n} + W_{s,\lambda\lambda} \dot{n}^2 = 0 \quad (3.14)$$

$$\begin{aligned} \epsilon^3: \quad & W_{s,i} \dddot{q}_i + W_{s,\lambda} \dddot{n} + \frac{3}{2} W_{s,ij} (\dot{q}_i \ddot{q}_j + \ddot{q}_i \dot{q}_j) \\ & + 3W_{s,i\lambda} (\dot{q}_i \ddot{n} + \ddot{q}_i \dot{n}) + 3W_{s,\lambda\lambda} \dot{n} \ddot{n} \\ & + W_{s,ijk} \dot{q}_i \dot{q}_j \dot{q}_k + 3W_{s,ij\lambda} \dot{q}_i \dot{q}_j \dot{n} \\ & + 3W_{s,i\lambda\lambda} \dot{q}_i \dot{n}^2 + W_{s,\lambda\lambda\lambda} \dot{n}^3 = 0 \end{aligned} \quad (3.15)$$

where ϵ^k ($k = 1, 2, 3$) is the k -th order perturbation equation.

In many problems of practical interest, the total potential energy is a linear function of the loading parameters [114]*. In such a case,

* This is particularly true when Λ^j represent external loads only. If the structural model is acted upon by a nonlinearly elastic loading device, Eq. 3.16 takes the form [93]

$$V = \pi(Q_i) - \int_0^e \bar{\Lambda}^j(e, \Lambda^j) de \quad \text{where } e = e(Q_i) \text{ and } \bar{\Lambda}^j = \bar{\Lambda}^j(\Lambda^j)$$

V can be expressed as [114]

$$V = \pi(Q_i) - \Lambda^j E^j(Q_i). \quad (3.16)$$

In Eq. 3.16, π is the system strain energy and E^j are the deflections corresponding to the loads Λ^j , and E^j are assumed linear in displacements, Q_i . Eqs. 3.13-3.15 are therefore reduced to the following form:

$$\epsilon^1: \quad W_{s,i} \dot{q}_i + W_{s,\lambda} \dot{\eta} = 0 \quad (3.17)$$

$$\epsilon^2: \quad W_{s,i} \ddot{q}_i + W_{s,\lambda} \ddot{\eta} + W_{s,ij} \dot{q}_i \dot{q}_j = 0 \quad (3.18)$$

$$\epsilon^3: \quad W_{s,i} \dddot{q}_i + W_{s,\lambda} \dddot{\eta} + \frac{3}{2} W_{s,ij} (\dot{q}_i \ddot{q}_j + \ddot{q}_i \dot{q}_j) + W_{s,ijk} \dot{q}_i \dot{q}_j \dot{q}_k = 0 \quad (3.19)$$

It follows from Eqs. 3.17-3.19 that the key solution is the first order perturbation solution. For example, \ddot{q}_i and $\ddot{\eta}$ of Eqs. 3.18 can be defined completely once \dot{q}_i and $\dot{\eta}$ are obtained from Eqs. 3.17. In general, the solution to the k -th order perturbation equation, ϵ^k , depends on solutions to the first $(k-1)$ perturbation equations. Hence the new equilibrium state $(Q_i + q_i, \Lambda^j + \lambda^j)$ is completely defined by the recursive solutions of Eqs. 3.17-3.19.

Eqs. 3.17 contain $(N+1)$ unknowns, q_i and η . Since only N equations exist, one of the q_i 's or η must be specified a priori. Such a specification reduces Eqs. 3.17 to N equations in N unknowns which can be solved simultaneously.

If one of the q_i 's is selected as ϵ , the scheme is called displace-

ment-increment perturbation procedure [92]. This is characterized by the equations

$$\left. \begin{aligned} q_i &= \epsilon, \dot{q}_i = 1, \ddot{q}_i = \ddot{\ddot{q}}_i = \dots = 0 \\ q_j &= \epsilon q_j + \frac{\epsilon^2}{2!} \ddot{q}_j + \frac{\epsilon^3}{3!} \ddot{\ddot{q}}_j + O[\epsilon^4]; j \neq i \\ \eta &= \epsilon \eta + \frac{\epsilon^2}{2!} \ddot{\eta} + \frac{\epsilon^3}{3!} \ddot{\ddot{\eta}} + O[\epsilon^4]. \end{aligned} \right\} \quad (3.20)$$

Similarly if η is selected as ϵ , the scheme is called load-increment perturbation procedure [92]. This is characterized by the equations

$$\left. \begin{aligned} \eta &= \epsilon, \dot{\eta} = 1, \ddot{\eta} = \ddot{\ddot{\eta}} = \dots = 0 \\ q_i &= \epsilon q_i + \frac{\epsilon^2}{2!} \ddot{q}_i + \frac{\epsilon^3}{3!} \ddot{\ddot{q}}_i + O[\epsilon^4] \end{aligned} \right\} \quad (3.21)$$

The above general analysis assumes that only one equilibrium path leaves any equilibrium point. The analysis requires modifications in the neighborhood of critical points. The required perturbation analysis through various critical points is discussed in the next three subsections.

c. Perturbation Analysis Through a Limit Point

In this section, the required modification to Eqs. 3.13 to obtain the equilibrium path through a limit point is discussed.

At a limit point, L, the fundamental path [114] OA of Fig. 3.2a reaches a local extremum. The initial post-buckling path associated

with the limit point is unstable and falling. Also the coefficient matrix of ϵ^1 in Eq. 3.13 vanishes [114]; i.e.,

$$\det (W_{S,i}) = 0 \quad (3.22)$$

The coefficient matrix $W_{S,i}$ is referred to as W_L in later discussions since it is associated with the limit point, L.

In addition to Eq. 3.22, the relation (using Eqs. 3.20)

$$\dot{\eta} = 0 \quad (3.23)$$

holds at the limit point. Hence, Eqs. 3.13 require modification to obtain a defined set of \dot{q}_i .

The perturbation analysis for the equilibrium path through a limit point can be established by using the displacement-increment perturbation scheme defined by Eqs. 3.20. The solutions to Eqs. 3.13 are sought by replacing the column corresponding to the "suitable" q_i with that of η . Such a transformation leads to Eqs. 3.24, i.e.,

$$\begin{bmatrix} W_{1,\lambda} & W_{1,2} \cdots W_{1,n} \\ W_{2,\lambda} & W_{2,2} \cdots W_{2,n} \\ \cdot & \cdot \quad \cdot \\ \cdot & \cdot \quad \cdot \\ \cdot & \cdot \quad \cdot \\ W_{n,\lambda} & W_{n,2} \quad W_{n,n} \end{bmatrix} \begin{bmatrix} \dot{\eta} \\ \dot{q}_2 \\ \cdot \\ \cdot \\ \cdot \\ \dot{q}_n \end{bmatrix} + \begin{bmatrix} W_{1,1} \\ W_{2,1} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,1} \end{bmatrix} \dot{q}_1 = \begin{bmatrix} 0 \\ 0 \\ \cdot \\ \cdot \\ \cdot \\ 0 \end{bmatrix} \quad (3.24)$$

if q_1 is assumed to be the "suitable" q_i .

The coefficient matrix of Eqs. 3.24 is denoted W_B for later reference. Since W_B^{-1} exists at equilibrium points excluding a bifurcation point, a unique post-buckling path through the limit point is obtained using Eqs. 3.24. High order approximations are then obtainable from Eqs. 3.14, 3.15.

d. Perturbation Analysis Through a Bifurcation Point

The construction of equilibrium paths through a bifurcation point is presented in this section.

At a point of bifurcation, the fundamental path intersects a second, distinct equilibrium path (Theorems 4 and 5, section 2.1). At bifurcation points, the coefficient matrix of Eqs. 3.24, W_B , becomes singular. Hence, the perturbation scheme of section 3.1c breaks down at bifurcation points and cannot be used for obtaining post-buckling response beyond such points. Consequently, a non-diagonalized perturbation scheme [114] is used, in which the generalized displacements are selected as generalized coordinates. This is presented below.

Let the singularity expressed in Eqs. 3.22 be characterized by the vanishing of the smallest eigenvalue of W_L . This implies the existence of a non-zero minor determinant, $\det (W_C)$, where

$$W_C = \begin{bmatrix} W_{2,2} \cdots W_{2,n} \\ \cdot & \cdot \\ \cdot & \cdot \\ & & \cdot \\ W_{n,2} \cdots W_{n,n} \end{bmatrix} \quad (3.25)$$

Also, let q_1 be a "suitable" perturbation parameter ϵ . Hence, Eqs. 3.20 take the form:

$$\left. \begin{aligned} q_1 &= \epsilon, \dot{q}_1 = 1, \ddot{q}_1 = \ddot{\ddot{q}}_1 = \dots = 0 \\ q_i(q_1) &= q_1 \dot{q}_i + \frac{1}{2!} q_1^2 \ddot{q}_i + \frac{1}{3!} q_1^3 \ddot{\ddot{q}}_i + O[q_1^4]; \quad i \neq 1 \\ n(q_1) &= q_1 \dot{n} + \frac{1}{2!} q_1^2 \ddot{n} + \frac{1}{3!} q_1^3 \ddot{\ddot{n}} + O[q_1^4] \end{aligned} \right\} \quad (3.26)$$

Inserting Eqs. 3.26 into Eqs. 3.24, one obtains the following two equations:

$$W_{1,\lambda} \dot{n} + [W_{1,2} \dots W_{1,n}] \begin{bmatrix} \dot{q}_2 \\ \vdots \\ \dot{q}_n \end{bmatrix} + W_{1,1} = 0 \quad (3.27)$$

and

$$\begin{bmatrix} W_{2,\lambda} \\ \vdots \\ W_{n,\lambda} \end{bmatrix} \dot{n} + \begin{bmatrix} W_{2,2} \dots W_{2,n} \\ \vdots \\ W_{n,2} \dots W_{n,n} \end{bmatrix} \begin{bmatrix} \dot{q}_2 \\ \vdots \\ \dot{q}_n \end{bmatrix} + \begin{bmatrix} W_{2,1} \\ \vdots \\ W_{n,1} \end{bmatrix} = \begin{bmatrix} 0 \\ \vdots \\ 0 \end{bmatrix} \quad (3.28)$$

To simplify algebraic operations, the following symbolic representations are introduced:

$$a_{11} = W_{1,\lambda} \quad (3.29)$$

$$a_{12} = [W_{1,2} \cdots W_{1,n}] \quad (3.30)$$

$$a_{21} = \begin{bmatrix} W_{2,\lambda} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,\lambda} \end{bmatrix} \quad (3.31)$$

$$b_{11} = W_{1,1} \quad (3.32)$$

$$b_{21} = \begin{bmatrix} W_{2,1} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,1} \end{bmatrix} \quad (3.33)$$

These representations (Eqs. 3.29-3.33) and Eq. 3.25 transform Eqs. 3.28 into

$$\begin{bmatrix} \cdot \\ q_2 \\ \cdot \\ \cdot \\ \cdot \\ \cdot \\ q_n \end{bmatrix} = -W_c^{-1} (a_{21} \dot{n} + b_{21}) \quad (3.34)$$

and together with Eqs. 3.27 lead to the relation:

$$\alpha \dot{\eta} + \beta = 0 \quad (3.35)$$

In Eqs. 3.35,

$$\alpha = a_{11} - a_{12} W_C^{-1} a_{21} \quad (3.36)$$

and

$$\beta = b_{11} - a_{12} W_C^{-1} b_{21} \quad (3.37)$$

But

$$\det (W_B) = 0 \quad (3.38)$$

at a bifurcation point. Hence, $\dot{\eta}$ is indefinite in both Eqs. 3.24 and 3.35. This indicates that in order for the bifurcation point to satisfy the first order perturbation equations, Eqs. 3.27 and 3.28, the relation

$$\alpha = \beta = 0 \quad (3.39)$$

must hold in Eqs. 3.35.

The proof of Eq. 3.39 can be established if a regular matrix [34] is introduced as follows:

$$R = \begin{bmatrix} 1 & | & -a_{12} W_C^{-1} \\ \hline 0 & | & I \end{bmatrix} \quad (3.40)$$

such that

$$\det (R) = 1 \quad (3.41)$$

Pre-multiplying W_B in Eqs. 3.24 by R, one obtains

$$\left[\begin{array}{c|c} 1 & -a_{12} W_C^{-1} \\ \hline 0 & I \end{array} \right] \left[\begin{array}{c|c} a_{11} & a_{12} \\ \hline a_{21} & W_C \end{array} \right] = \left[\begin{array}{c|c} a_{11} - a_{12} W_C^{-1} a_{21} & 0 \\ \hline a_{21} & W_C \end{array} \right] \quad (3.42)$$

Taking determinants on both sides of Eq. 3.42 and using conditions expressed in Eqs. 3.38 and 3.41, one finds that

$$(a_{11} - a_{12} W_C^{-1} a_{21}) \cdot \det (W_C) = 0 \quad (3.43)$$

But

$$\det (W_C) \neq 0 \quad (3.44)$$

Hence

$$(a_{11} - a_{12} W_C^{-1} a_{21}) = \alpha = 0 \quad (3.45)$$

Similarly, if W_L is pre-multiplied by R, the resulting equation is

$$\left[\begin{array}{c|c} 1 & -a_{12} W_C^{-1} \\ \hline 0 & I \end{array} \right] \left[\begin{array}{c|c} b_{11} & a_{12} \\ \hline b_{21} & W_C \end{array} \right] = \left[\begin{array}{c|c} b_{11} - a_{12} W_C^{-1} b_{21} & 0 \\ \hline b_{21} & W_C \end{array} \right] \quad (3.46)$$

Taking determinants on both sides of Eq. 3.46 leads to

$$(b_{11} - a_{12} W_C^{-1} b_{21}) = \beta = 0 \quad (3.47)$$

since $\det(W_L)$ vanishes at bifurcation points [114].

The conditions expressed in Eqs. 3.45 and 3.47 imply that the first order perturbation equations do not yield a unique solution for \dot{n} at a point of bifurcation. In order to obtain definite values for \dot{n} , the second order perturbation equations, ϵ^2 , must be used. Rewriting ϵ^2 , defined by Eqs. 3.14, in matrix form, one obtains

$$\begin{bmatrix} W_{1,1} & W_{1,2} \cdots W_{1,n} \\ W_{2,1} & W_{2,2} \cdots W_{2,n} \\ \cdot & \cdot \quad \cdot \\ \cdot & \cdot \quad \cdot \\ \cdot & \cdot \quad \cdot \\ W_{n,1} & W_{n,2} \cdots W_{n,n} \end{bmatrix} \begin{bmatrix} \ddots \\ q_1 \\ \ddots \\ q_2 \\ \cdot \\ \cdot \\ \cdot \\ \ddots \\ q_n \end{bmatrix} + \begin{bmatrix} W_{1,\lambda} \\ W_{2,\lambda} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,\lambda} \end{bmatrix} \ddots n$$

$$+ \begin{bmatrix} W_{1,ij} \\ W_{2,ij} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,ij} \end{bmatrix} \ddots q_i q_j + 2 \begin{bmatrix} W_{1,i\lambda} \\ W_{2,i\lambda} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,i\lambda} \end{bmatrix} \ddots q_i n + \begin{bmatrix} W_{1,\lambda\lambda} \\ W_{2,\lambda\lambda} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,\lambda\lambda} \end{bmatrix} n^2 = \begin{bmatrix} 0 \\ 0 \\ \cdot \\ \cdot \\ \cdot \\ 0 \end{bmatrix} \quad (3.48)$$

All partial derivatives in Eqs. 3.48 are evaluated at the known equilibrium state $\epsilon = 0$; e.g.

$$W_{i,j} = \left. \frac{\partial W_i}{\partial q_j} \right|_{\epsilon = 0} ; W_{i,\lambda} = \left. \frac{\partial W_i}{\partial \eta} \right|_{\epsilon = 0} \quad (3.49)$$

Using Eqs. 3.25 and 3.29-3.31, Eqs. 3.48 can be written as:

$$a_{11} \ddot{\eta} + a_{12} \begin{bmatrix} \ddot{q}_2 \\ \cdot \\ \cdot \\ \cdot \\ \cdot \\ \ddot{q}_n \end{bmatrix} + W_{1,ij} \dot{q}_i \dot{q}_j + 2W_{1,i\lambda} \dot{q}_i \dot{\eta} + W_{1,\lambda\lambda} \dot{\eta}^2 = 0 \quad (3.50)$$

and

$$a_{21} \ddot{\eta} + W_C \begin{bmatrix} \ddot{q}_2 \\ \cdot \\ \cdot \\ \cdot \\ \cdot \\ \ddot{q}_n \end{bmatrix} + \begin{bmatrix} W_{2,ij} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,ij} \end{bmatrix} \dot{q}_i \dot{q}_j + 2 \begin{bmatrix} W_{2,i\lambda} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,i\lambda} \end{bmatrix} \dot{q}_i \dot{\eta} + \begin{bmatrix} W_{2,\lambda\lambda} \\ \cdot \\ \cdot \\ \cdot \\ W_{n,\lambda\lambda} \end{bmatrix} \dot{\eta}^2 = \begin{bmatrix} 0 \\ \cdot \\ \cdot \\ \cdot \\ 0 \end{bmatrix} \quad (3.51)$$

Eqs. 3.51 imply that

$$\begin{aligned}
 \begin{bmatrix} \ddots \\ \ddot{q}_2 \\ \vdots \\ \vdots \\ \ddots \\ \ddot{q}_n \end{bmatrix} &= -W_C^{-1} (a_{21} \ddot{n} + \begin{bmatrix} W_{2,ij} \\ \vdots \\ \vdots \\ W_{n,ij} \end{bmatrix} \dot{q}_i \dot{q}_j \\
 &+ 2 \begin{bmatrix} W_{2,i\lambda} \\ \vdots \\ \vdots \\ W_{n,i\lambda} \end{bmatrix} \dot{q}_i \dot{n} + \begin{bmatrix} W_{2,\lambda\lambda} \\ \vdots \\ \vdots \\ W_{n,\lambda\lambda} \end{bmatrix} \dot{n}^2) \quad (3.52)
 \end{aligned}$$

Inserting Eqs. 3.52 into Eqs. 3.50 yields

$$\begin{aligned}
 -a_{12} W_C^{-1} & \left(\begin{bmatrix} W_{2,ij} \\ \vdots \\ \vdots \\ W_{n,ij} \end{bmatrix} \dot{q}_i \dot{q}_j + 2 \begin{bmatrix} W_{2,i\lambda} \\ \vdots \\ \vdots \\ W_{n,i\lambda} \end{bmatrix} \dot{q}_i \dot{n} + \begin{bmatrix} W_{2,\lambda\lambda} \\ \vdots \\ \vdots \\ W_{n,\lambda\lambda} \end{bmatrix} \dot{n}^2 \right) \\
 & + W_{1,ij} \dot{q}_i \dot{q}_j + 2W_{1,i\lambda} \dot{q}_i \dot{n} + W_{1,\lambda\lambda} \dot{n}^2 \\
 & + \ddot{n} (W_{1,\lambda} - a_{12} W_C^{-1} a_{21}) = 0 \quad (3.53)
 \end{aligned}$$

The last term of Eqs. 3.53 is equivalent to $\alpha \ddot{n}$ which vanishes

since α is zero at a bifurcation point. Hence, Eqs. 3.53 in view of Eqs. 3.25 and 3.34 lead to the following quadratic equation in \dot{n} :

$$a \dot{n}^2 + b \dot{n} + c = 0 \quad (3.54)$$

where

$$\begin{aligned} a = & G_{1,k} W_{k,ij} G_{i,\lambda} G_{j,\lambda} + 2W_{1,i\lambda} G_{i,\lambda} \\ & + 2G_{1,k} W_{k,i\lambda} G_{i,\lambda} + G_{i,k} W_{k,\lambda\lambda} \\ & + W_{1,\lambda\lambda} + W_{1,ij} G_{i,\lambda} G_{j,\lambda}; \end{aligned} \quad (3.55)$$

$$\begin{aligned} b = & G_{1,k} (W_{k,li} + W_{k,il}) G_{i,\lambda} \\ & + G_{1,k} W_{k,ij} (G_{i,\lambda} G_{1,j} + G_{1,j} G_{j,\lambda}) \\ & + 2G_{1,k} W_{k,i\lambda} + 2G_{1,k} W_{k,i\lambda} G_{1,i} \\ & + (W_{1,il} + W_{1,li}) G_{i,\lambda} + 2W_{1,i\lambda} G_{1,i} \\ & + W_{1,ij} (G_{1,i} G_{j,\lambda} + G_{i,\lambda} G_{1,j}) + 2W_{1,1\lambda}; \end{aligned} \quad (3.56)$$

$$\begin{aligned} c = & G_{1,k} W_{k,11} + G_{1,k} (W_{k,li} + W_{k,il}) G_{1,i} \\ & + G_{1,k} W_{k,ij} G_{1,i} G_{1,j} + G_{1,11} \\ & + (W_{1,il} + W_{1,li}) G_{1,i} + W_{1,ij} G_{1,i} G_{1,j}; \end{aligned} \quad (3.57)$$

$$i, j, k = 2, 3, \dots, N.$$

and

$$\begin{bmatrix} G_{1,2} \\ \vdots \\ G_{1,k} \\ \vdots \\ G_{1,n} \end{bmatrix} = - \begin{bmatrix} W_{2,2} & \dots & W_{2,n} \\ \vdots & & \vdots \\ W_{k,2} & \dots & W_{k,n} \\ \vdots & & \vdots \\ W_{n,2} & & W_{n,n} \end{bmatrix}^{-1} \begin{bmatrix} W_{2,1} \\ \vdots \\ W_{k,1} \\ \vdots \\ W_{n,1} \end{bmatrix} \quad (3.58)$$

$$\begin{bmatrix} G_{2,\lambda} \\ \vdots \\ G_{i,\lambda} \\ \vdots \\ G_{n,\lambda} \end{bmatrix} = - \begin{bmatrix} W_{2,2} & \dots & W_{2,n} \\ \vdots & & \vdots \\ W_{i,2} & \dots & W_{i,n} \\ \vdots & & \vdots \\ W_{n,2} & \dots & W_{n,n} \end{bmatrix}^{-1} \begin{bmatrix} W_{2,\lambda} \\ \vdots \\ W_{i,\lambda} \\ \vdots \\ W_{n,\lambda} \end{bmatrix} \quad (3.59)$$

The roots of Eq. 3.54 are

$$\left. \begin{array}{l} \dot{\eta}_1 \\ \dot{\eta}_2 \end{array} \right\} = \frac{-b \pm \sqrt{b^2 - 4ac}}{2a} \quad (3.60)$$

Using $\dot{\eta}_1, \dot{\eta}_2$ from Eq. 3.60 in conjunction with Eqs. 3.34, two sets of

$[\dot{q}_2 \ \dot{q}_3 \ \dots \ \dot{q}_n]$ are obtained. These together with Eqs. 3.14 and 3.15 yield \dot{q}_i , $\ddot{\eta}$, etc. With such information, together with Eqs. 3.26, the construction of the post-buckling paths beyond a bifurcation point becomes possible. For N large, computer programming is required to carry out the analysis described.

If Eqs. 3.17-3.19 instead of Eqs. 3.13-3.15 are used, the resulting quadratic equation in $\dot{\eta}$ is given by

$$a^* \dot{\eta}^2 + b^* \dot{\eta} + c^* = 0 \quad (3.61)$$

where

$$a^* = a - (2G_{1,k} W_{k,i\lambda} G_{i,\lambda} + G_{1,k} W_{k,\lambda\lambda} + W_{1,\lambda\lambda}) \quad (3.62)$$

$$b^* = b - (2G_{1,k} W_{k,i\lambda} + 2G_{1,k} W_{k,i\lambda} G_{1,i} + 2W_{1,i\lambda} G_{1,i} + 2W_{1,1\lambda}) \quad (3.63)$$

$$c^* = c \quad (3.64)$$

and

$$\left. \begin{array}{l} \dot{\eta}_1 \\ \dot{\eta}_2 \end{array} \right\} = \frac{-b^* \pm \sqrt{b^{*2} - 4a^* c^*}}{2a^*} \quad (3.65)$$

Eqs. 3.62 - 3.64 reflect the savings in time and effort in computer analysis when the total potential energy is a linear function of the loading parameters. These equations (i.e., Eqs. 3.62-3.64) form the

basis for comparing the non-diagonalized scheme above with that of Thompson [110, 114].

For the purpose of comparison, consider the potential energy function in two generalized coordinates Q_1 , Q_2 and a loading parameter λ [114, p. 142]

$$V = \frac{1}{2} Q_1^2 + \frac{1}{2} Q_2^2 - \frac{1}{2} Q_1 Q_2^2 - \lambda Q_1 \quad (3.66)$$

The equilibrium equations defined by $\partial V / \partial Q_i = 0$, $i = 1, 2$, yield the fundamental path solutions

$$Q_1 = \lambda \quad (3.67)$$

$$Q_2 = 0$$

Thompson [114] defined a set of incremental coordinates X_1 , X_2 emerging from the fundamental path as follows

$$Q_1 = \lambda + X_1 \quad (3.68)$$

$$Q_2 = 0 + X_2$$

From Eqs. 3.68 and 3.66, one obtains a new function

$$W = \frac{1}{2} X_1^2 + \frac{1}{2} X_2^2 - \frac{1}{2} X_1 X_2^2 - \frac{1}{2} \lambda X_2^2 - \frac{1}{2} \lambda^2 \quad (3.69)$$

It follows from Eqs. 3.66 and 3.69 that the transformation from V to W destroys the linearity in λ when X_1 , X_2 are treated as generalized coordinates. Hence Thompson's 'sliding' transformation forces the

generalized coordinates to be measured along curvilinear coordinates. This observation was made by Sewell [94]. However, Sewell [94] obtained branching paths beyond bifurcation points by expressing the paths explicitly in terms of orthonormal eigenvectors of an algebraic eigenvalue problem. Such an approach necessitates finding N eigenvalues and eigenvectors of the characteristic stability matrix for paths emanating from the critical point. For N large, this approach is disadvantageous in terms of computing time. It could lead to numerical instability problems for large N [38].

The non-diagonalized perturbation scheme resulting in Eq. 3.61 does not require eigenvalue and eigenvector computations. The generalized displacements used as generalized coordinates prevent any loss of the linearity of V in Λ^j (or V in λ). The conclusion can therefore be drawn that the scheme proposed is computationally advantageous for a large system.

e. Coincident Critical Point

Critical points (limit, bifurcation points) normally lie isolated on the fundamental path. In that case, unique buckling modes are associated with the critical loads. If simultaneous buckling occurs, two or more critical points coincide. At such coincident critical points, buckling modes are mixed [114, 59]. Depending on the number of post-buckling paths leaving a coincident critical point, post-buckling response for systems exhibiting it can be complex and undefined [60, 21].

Coincident critical points have been observed in the analysis of rigid link models with two degrees of freedom [21] and in cylindrical shells [65]. Some general results for N-degree of freedom systems, satisfying certain symmetry and pre-buckling deformation restrictions, have been proposed; e.g., [102, 101, 60, 59].

In the present study of reticulated domes, interest is focused on unstable coincident critical points (if they exist) and the associated unstable post buckling paths. Accordingly, the presence or absence of a coincident critical point is monitored on the basis of the theories of sections 3.1 b-d. For at least two critical points to be coincident, the relations

$$\det (W_L) = \det (W_B) = \det (W_C) = 0 \quad (3.70)$$

must hold. The relation $\det (W_C) = 0$ indicates the presence of at least one zero eigenvalue for the matrix W_C . In this case, the analysis of section 3.1d is extended to include the third order perturbation equation since ϵ^1 and ϵ^2 will not yield unique values of η .

In the current investigation, coincident critical points are absent for all loading distributions considered. This fact is established from the non-vanishing of $\det (W_C)$ for all cases considered.

3.2 Nonlinear Programming Technique

The basic ideas in nonlinear structural analysis using nonlinear programming technique is discussed in this section. The various

nonlinear programming algorithms employed in the nonlinear analysis of geodesic domes (Chapter 5) are referenced. The role of the nonlinear programming technique in obtaining the nonlinear response of reticulated domes to static loads is stated.

The basic law (Theorem 1, section 2.1) underlying the nonlinear programming technique is the law of minimum potential energy. This law states [69] that "a conservative holonomic system is in a configuration of stable equilibrium if, and only if, the value of the total potential energy is a relative minimum."

At the equilibrium point Q_i^* , the equilibrium equations

$$W_S(Q_i, \Lambda^j) \Big|_{Q_i = Q_i^*} = 0; \quad \begin{array}{l} i = 1, 2, \dots, N \\ j = 1, 2, \dots, M \end{array} \quad (3.71)$$

are satisfied. The equilibrium point Q_i^* is stable if

$$V(Q_i^*, \Lambda^j) < V(Q_i, \Lambda^j) \quad (3.72)$$

for all Q_i in some neighborhood of Q_i^* . Eqs. 3.72 imply that the total potential energy surface is convex in the neighborhood of stable equilibrium states. Hence, the minimum potential energy principle provides a useful device for casting the structural analysis problem into a mathematical programming problem [33]. One can prove [80] that such a minimization of a scalar valued function of N variables is mathematically equivalent to solving N simultaneous nonlinear equations (if V is a quadratic form).

For geometrically nonlinear, reticulated dome models, the potential energies of the line elements can be summed to obtain the total potential energy of the assemblage. Hence, Eqs. 3.71 and 3.16 imply that at equilibrium,

$$\sum_{k=1}^{NM} \frac{\partial \pi_k}{\partial Q_S} - P_S = 0 \quad (3.73)$$

where

$$P_S = P_S(\Lambda^j) \quad (3.74)$$

In Eqs. 3.73 and 3.74, π_k is the strain energy of element k ; NM is the total number of elements, and P_S is the generalized load corresponding to the generalized displacement Q_S . A convergence criterion based on Eq. 3.71 or Eq. 3.73 is [73, 80]

$$|W_S(Q_i, \Lambda^j)| < \gamma \quad (3.75)$$

where γ is a small positive constant ($\gamma \ll 1$). Physically, γ represents the unbalanced joint forces and determines the precision involved in gradient computations.

Three unconstrained nonlinear optimization algorithms are employed in the present study: Davidon-Fletcher-Powell algorithm [31]; Conjugate Gradient technique [32]; Powell's direct search method [83]. The first two techniques are gradient methods. These are efficient for large systems when analytical derivatives of the total potential energy

(objective function) are available [39]. The third technique is a direct-search method which is unreliable for large systems because of its slow rate of convergence.

The Conjugate Gradient technique requires less storage space than the Davidon-Fletcher-Powell method. However, for precise location of the minimum, the Davidon-Fletcher-Powell method is more reliable [39]. The main advantage of Powell's direct-search technique is that it requires no derivative expressions. Hence for the direct-search method, the total potential energy need not satisfy regularity and continuity conditions; neither is a proof for the existence of derivatives required [39]. For certain simple functions, Powell's method is quite efficient and requires less function evaluations for convergence compared to gradient methods [39].

In the search for a minimum, all algorithms described above perform a series of one-dimensional searches in specified directions. They are all quadratically convergent algorithms in that they will minimize a quadratic function of N variables in N steps. In highly nonlinear problems, this property is quite desirable since a Taylor series expansion of the total potential energy about stable equilibrium states shows that the quadratic and lower terms dominate in the vicinity of the minimum. Hence, it is sufficient to pass a quadratic polynomial through three points in the neighborhood of the minimum.

Employing the three algorithms stated above provides flexibility since the solution procedure most suitable to a particular test problem

can be selected.

The role of the nonlinear programming technique in the present study is to construct the initial portions of the fundamental path [114] at the onset of loading. Such a role provides a starting solution for the static perturbation technique, and optimizes the choice of a "suitable" perturbation parameter. The choice of a "suitable" perturbation parameter is presented in section 4.3.

Detailed descriptions of all nonlinear programming techniques employed are in reference 39.

3.3 Dynamic Analysis

The perturbed motion of the discrete system modeled, under dynamic loads, is obtained using numerical integration. Numerical integration is necessary since the governing differential equations of perturbed motion are nonlinear.

The differential equations of perturbed motion can be expressed as (section 2.1, Eq. 2.3)

$$y' = f(y) \quad (3.76)$$

in the $2N$ dimensional phase space. The properties of the vector f are given in section 2.1. Consider an initial disturbance y_0 applied at time $t = 0$. Hamming's modified Predictor-Corrector numerical integration scheme [49, 84] applied to Eqs. 3.76 with the given initial disturbance provides the desired solutions y of Eq. 3.76 during motion.

Phase-plane portraits of perturbed motions provide a valuable

check on the sufficient stability criterion [45, 50] used in the present study. Closed trajectories surrounding only the equilibrium point (an isolated singularity) under investigation indicate perturbed motions of stable origin. This corresponds to a stable node (center) in the phase plane. If the perturbed motion encompasses more than one singular (equilibrium) point in the phase plane, the origin is unstable. The perturbed motion surrounding more than one stable node and passing through saddle points is termed a separatrix. A separatrix indicates the domains of attraction, the regions of asymptotic stability, of stable equilibrium states along the real y_1 axis. In general, for conservative systems, closed trajectories enclose odd numbers of singularities (equilibrium points) with the number of centers exceeding the number of saddle points by one [75].

CHAPTER 4

CONVERGENCE CHARACTERISTICS AND ERROR ESTIMATES

In this chapter, the convergence characteristics and error estimates in a perturbation scheme are considered. The role of the right choice of the perturbation parameter in the construction of the equilibrium surface is discussed, and a method for selecting a "suitable" perturbation parameter is proposed.

4.1 Convergence Characteristics of the Static Perturbation Technique

When a series representation for the increments in the generalized displacements, q_i , and the generalized external loads, η , is used in a static perturbation method, the result is a series of successive approximations in the location of a nearby equilibrium state. This section considers the convergent characteristics of such series representations.

Consider the parametric representation of any prospective equilibrium path leaving a known equilibrium point [92] in the form (Chapter 3)

$$\left. \begin{aligned} q_i(\epsilon) &= \sum_{k=1}^3 \frac{\epsilon^k}{k!} \frac{\partial^k q_i}{\partial \epsilon^k} + O(\epsilon^4) \\ \eta(\epsilon) &= \sum_{k=1}^3 \frac{\epsilon^k}{k!} \frac{\partial^k \eta}{\partial \epsilon^k} + O(\epsilon^4). \end{aligned} \right\} \quad (4.1)$$

In the neighborhood of the equilibrium point (Q_i^j, Λ^j) , information about prospective paths emanating from it is obtained from the following

linear, ordered, perturbation equations [92]

$$\left. \begin{aligned} \frac{d^r}{d\varepsilon^r} [W_s\{Q_i(\varepsilon), \Lambda^j(\varepsilon)\}] \Big|_{\varepsilon=0} &= 0; \\ r &= 1, 2, 3, \dots \\ s, i &= 1, 2, \dots, N \\ j &= 1, 2, \dots, M. \end{aligned} \right\} \quad (4.2)$$

The resulting potential energy, V , for the desired equilibrium state is a series containing terms of homogeneous order,

$$\left. \begin{aligned} V &= V_0 + V_1 q_i + V_2 q_i q_j + V_3 q_i q_j q_k + O(q_i q_j q_k q_\ell); \\ i, j, k, \ell &= 1, 2, \dots, N. \end{aligned} \right\} \quad (4.3)$$

In Eqs. 4.3, V_0 is the potential energy of the known equilibrium state; V_1 , V_2 , V_3 are the associated coefficients that characterize the first, second, third order perturbation solutions, respectively, of the physical problem. Also

$$V_m \neq V_m(\varepsilon); \quad m = 1, 2, 3. \quad (4.4)$$

Eqs. 4.3 can be rewritten as

$$(V - V_0) = V_m \delta_m(\varepsilon) + O(\varepsilon^4) \quad (4.5)$$

where

$$\delta_m(\varepsilon) = O(\varepsilon^m); \quad m = 1, 2, 3. \quad (4.6)$$

A 'Ratio Test' [78] applied to the right hand side of Eq. 4.5 establishes the limiting process

$$\lim_{\epsilon \rightarrow 0} \left[\frac{n^{\text{th}} \text{ term}}{(n-1)^{\text{th}} \text{ term}} \right] = \lim_{\epsilon \rightarrow 0} [O(\epsilon)] = 0. \quad (4.7)$$

Eq. 4.7 implies that Eq. 4.5 represents an asymptotic series which is uniformly convergent as $\epsilon \rightarrow 0$. This series belongs to a general class of asymptotic expansions [78] represented as

$$V = \sum_{m=0}^{n-1} a_m \delta_m(\epsilon) + O[\delta_n(\epsilon)] \quad (4.8)$$

where

$$a_m \neq a_m(\epsilon); \epsilon \rightarrow 0. \quad (4.9)$$

Asymptotic expansions of this type are uniformly valid (in the absence of singularities in the region of interest) and uniformly convergent. For a given asymptotic sequence, $\delta_m(\epsilon)$, the asymptotic representation of V is unique [78]. The conclusion is that in a static perturbation method, there exists a power series solution in terms of a path parameter, ϵ , which is uniformly convergent provided $\epsilon \rightarrow 0$.

A proof of convergence of the power series solution, in terms of a parameter, ϵ , for a system of first order differential equations for "sufficiently" small ϵ is given by Minorsky [76].

4.2 Error Estimates in a Perturbation Scheme

This section is concerned with the estimation of errors in a static

perturbation scheme. An upper bound on the value of the perturbation parameter, ϵ , for a specified error bound, δ , is given.

Referring to Eqs. 4.1, the pertinent error equations for third order perturbation solutions can be written as follows:

$$\Delta q_i(\epsilon) = \frac{1}{4!} \epsilon^4 \frac{\partial^4 q_i}{\partial \epsilon^4} + \frac{1}{5} \epsilon^5 \frac{\partial^5 q_i}{\partial \epsilon^5} + \dots; \quad i = 1, 2, \dots, N \quad (4.10)$$

$$\Delta \eta(\epsilon) = \frac{1}{4!} \epsilon^4 \frac{\partial^4 \eta}{\partial \epsilon^4} + \frac{1}{5} \epsilon^5 \frac{\partial^5 \eta}{\partial \epsilon^5} + \dots \quad (4.11)$$

$\Delta q_i(\epsilon)$, $\Delta \eta(\epsilon)$, in Eqs. 4.10, 4.11 represent errors in the generalized displacements, errors in the generalized independent loads, respectively. Hence an attempt to limit errors to a value δ can be made by selecting ϵ to satisfy the inequality

$$\epsilon \leq \max \left\{ \left(\frac{4! \delta}{q_i^{(4)}} \right)^{\frac{1}{4}}, \left(\frac{4! \delta}{\eta^{(4)}} \right)^{\frac{1}{4}} \right\} \quad (4.12)$$

where

$$\left. \begin{aligned} q_i^{(4)} &= \frac{\partial^4 q_i}{\partial \epsilon^4} \\ \eta^{(4)} &= \frac{\partial^4 \eta}{\partial \epsilon^4} \end{aligned} \right\} \quad (4.13)$$

Eq. 4.12 represents an upper bound on the value of ϵ ; it also implies that errors in a perturbation scheme can be controlled by limiting the value of the chosen perturbation parameter, ϵ .

4.3 Choice, and Role of a Perturbation Parameter

Sections 4.1 and 4.2 are developed on the assumption that a "suitable" perturbation parameter, ϵ , exists. Here, a possible way to choose a "suitable" ϵ is proposed. Also, the role played by ϵ in the construction of the equilibrium surface is discussed.

The perturbation parameter, ϵ , represents an increment in length along the equilibrium path [92]. It is a measure of progress in the response of the system. Geometrically, ϵ indicates possible expansions or contractions of the equilibrium surface to changing actions. It depicts changes in the physical variables (generalized external loads, generalized displacements) of the system. If the wrong ϵ is chosen, differentiation with respect to ϵ is not justified [78]. This leads to nonuniform expansions and, hence, to numerical instability of the solution process. This is illustrated by the response of a 21-degree-of freedom space truss (Chapter 5, Fig. 5.35) for different values of the perturbation parameter.

For some systems exhibiting stable critical states, Fig. 4.1, the load parameter, η , represents a "suitable" choice for ϵ [114]. As shown in Fig. 4.1, the post-buckling paths of such systems are rising. Hence, at pre-buckling loads, there is a unique equilibrium path. The associated pre-buckling deformations are very small and can be neglected [114,59]. Axially loaded columns, frames, and plates are typical examples in this class [114].

For imperfection-sensitive systems, Fig. 4.2, the critical equilibrium states are unstable. This implies the existence of at least one

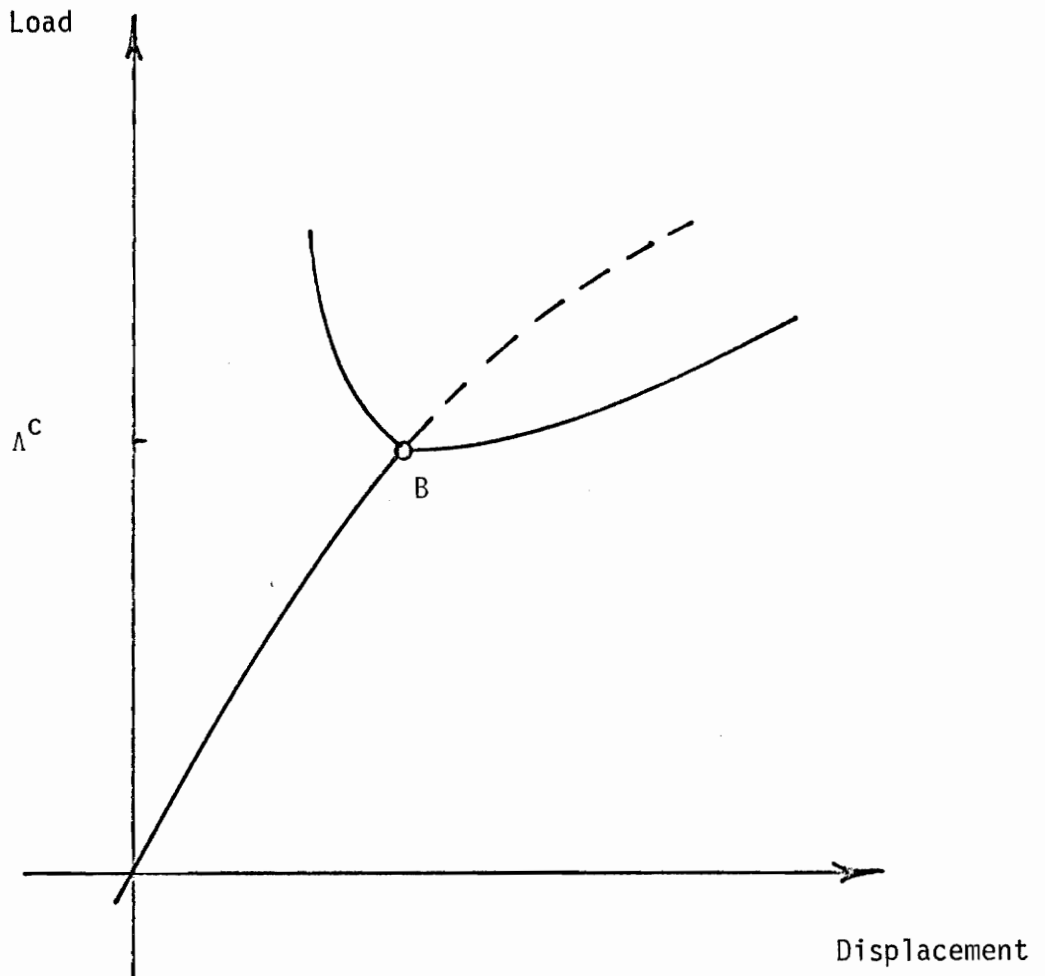


FIG. 4.1 EQUILIBRIUM PATHS THROUGH A STABLE-SYMMETRIC POINT OF BIFURCATION
(Imperfection - Insensitive System)

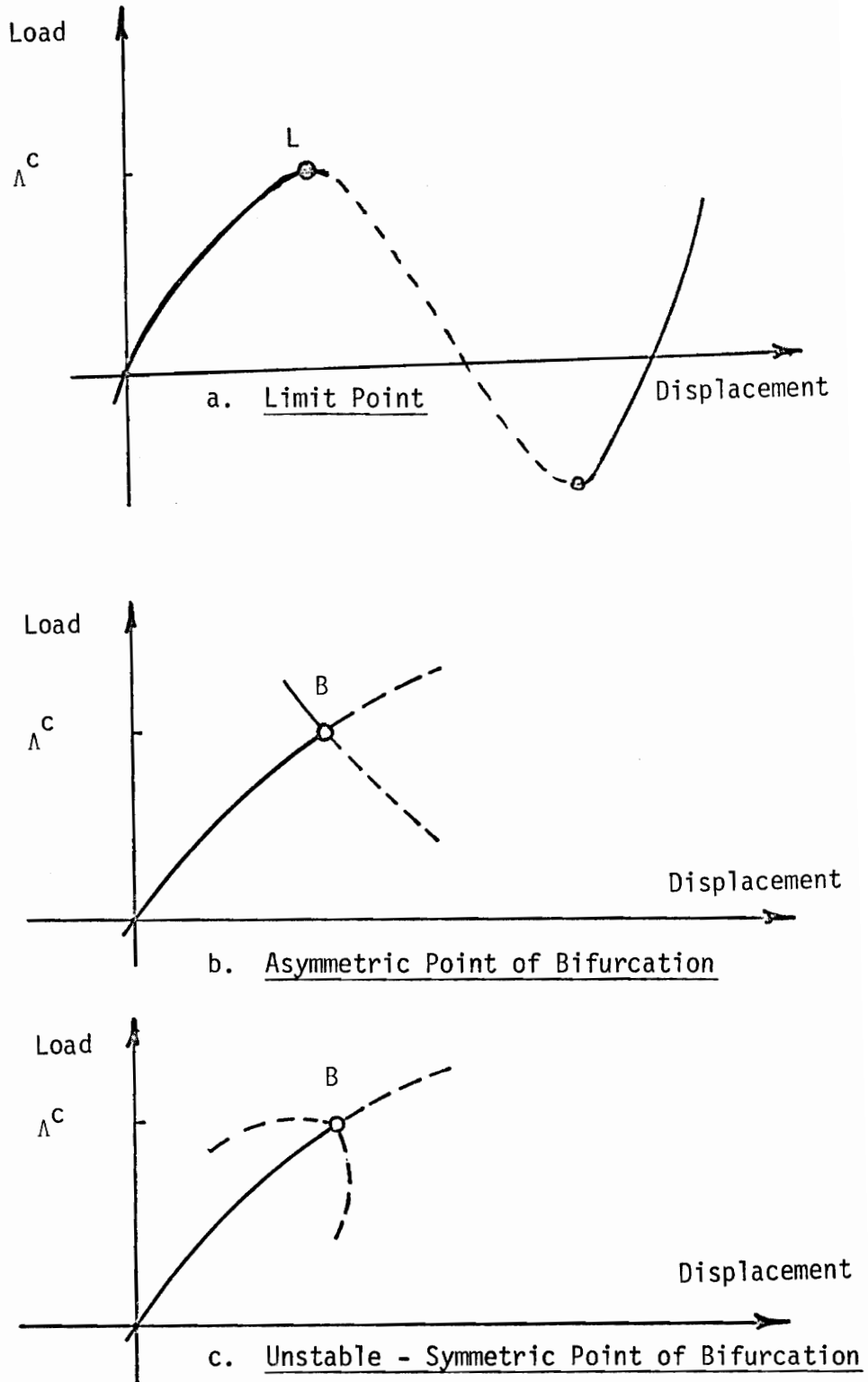


FIG. 4.2 EQUILIBRIUM PATHS THROUGH UNSTABLE CRITICAL POINTS

(Imperfection - Sensitive Systems)

falling post-buckling path. Hence at pre-buckling loads, the system could exhibit multi-equilibrium states [114] as in Fig. 4.2a. Also the pre-buckling deformations are not negligible as they play significant roles in the buckling behavior of the system [57]. In this case, η is an inappropriate choice for post-buckling response, although, estimates of the buckling load can be made. A generalized displacement-increment represents a proper choice for ϵ . This choice is necessary to obtain complete post-buckling (stable and unstable equilibrium paths) response beyond unstable critical points. Laterally loaded shallow arches, spherical shells, and reticulated domes are typical examples in this class [114].

If the imperfection-sensitive system can be adequately represented by one or two degrees of freedom, (i.e. $N \leq 2$), the choice of ϵ can be obtained by trial and error between the two generalized displacements. However, for multi-degree of freedom systems with N large, subjected to several independent external loads, the choice of a "suitable" ϵ is not obvious and cannot be arbitrary. A study in N -space topology can probably provide desirable guidelines [75]. Without topological studies, however, it is essential that the chosen ϵ ensures the expected shape for the final construction of the equilibrium surface. Also, the condition that all remaining variables are single-valued functions of ϵ must be satisfied.

The "largest" displacement-increment, q_1 , at the initial stage of loading, from the zero-state in configuration-space, is proposed as the "suitable" choice for ϵ in the case of imperfection-sensitive

systems. This choice imposes the necessary control, and restriction, on the generalized displacement that drastically distorts the equilibrium surface if other generalized coordinates are slightly perturbed. Such a control leads to an accurate construction of the equilibrium surface which is desirable in stability investigations [56].

A nonlinear programming technique is used in this study to obtain the generalized displacements corresponding to a prescribed set of external loads at the onset of loading. The associated increments of all the N generalized displacements are stored and compared. This operation is repeated for a few load steps and the "suitable" ϵ is then chosen as proposed above. A static-perturbation procedure is then initiated using the chosen ϵ as the "suitable" perturbation parameter.

CHAPTER 5
DEMONSTRATION PROBLEMS

In this chapter, the theorems, principles, solution techniques, and error measures discussed and developed in earlier chapters are applied to the stability analysis of geodesic domes. Two geodesic dome models (Figs. 5.1, 5.2), with 21 and 39 degrees-of-freedom, are studied in detail. Responses of the dome models to multiple static and dynamic loads are documented. The associated stability boundaries of the models under two independent external loads are obtained. For brevity, the term degree-of-freedom is abbreviated as DOF in subsequent discussions.

The reticulated dome models are formed by an assemblage of one-dimensional elements interconnected at the joints. Depending on the joint coordinates, the domes could be regarded as shallow or deep [87, 88]. Independent loads and finite displacements are admitted at the joints. The resulting geometrically nonlinear discrete system is described by its potential function, $V(Q_i, \Lambda^j)$, where V has all the properties outlined in section 2.1.

The potential energy of the dome model can be expressed as

$$V = \pi + \Omega \tag{5.1}$$

where

$$\pi = \sum_{i=1}^{NM} \frac{A_i E_i}{2L_i} (L_i^* - L_i)^2 \tag{5.2}$$

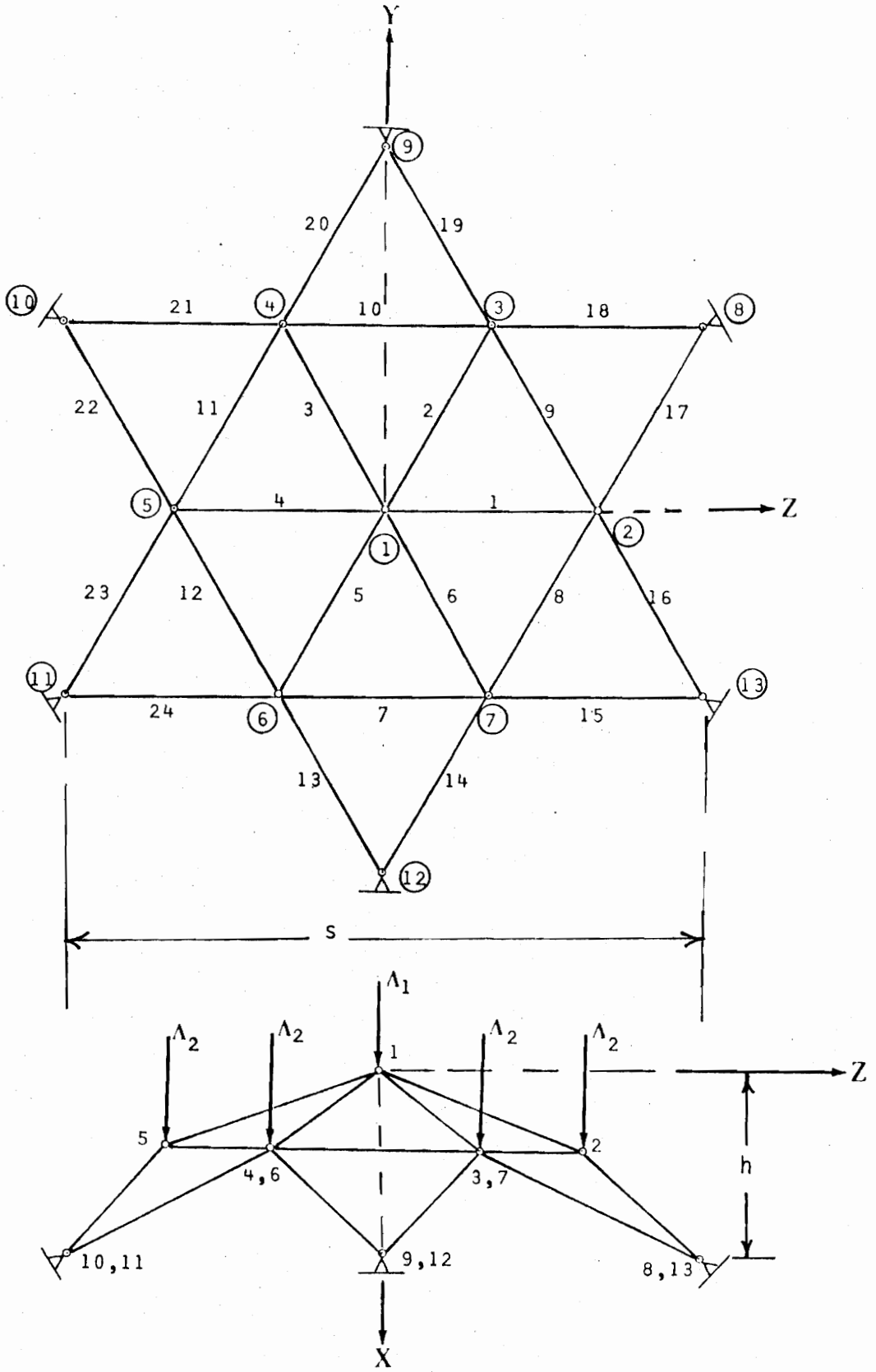


FIG. 5.1 GEODESIC DOME: 21 DOF MODEL

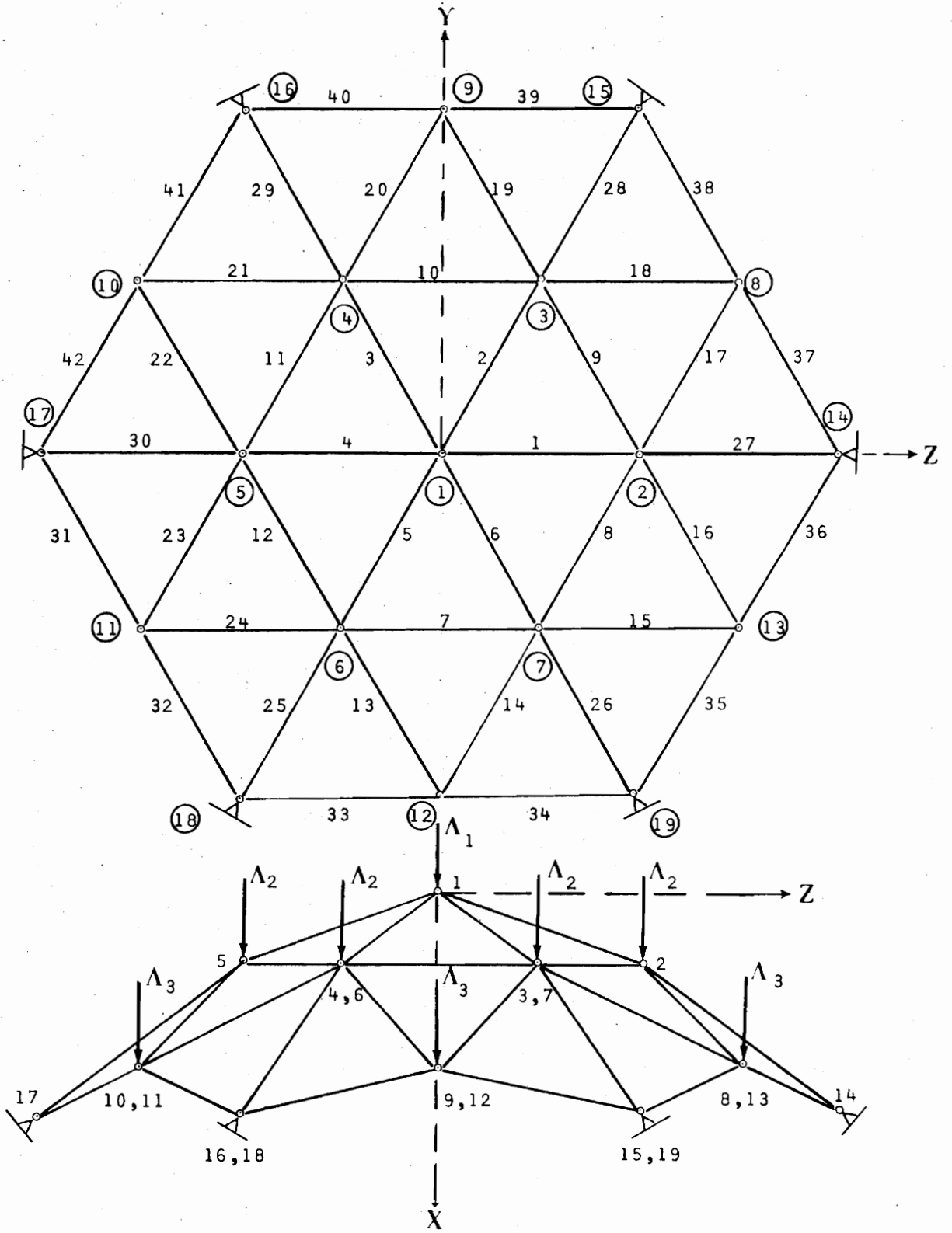


FIG. 5.2 GEODESIC DOME: 39 DOF MODEL

and

$$\Omega = - \sum_{j=1}^{NJ} P_j^T U_j \quad (5.3)$$

In Eqs. 5.2 - 5.3, A_i , E_i , L_i denote the area, the modulus of elasticity, and the initial length of element i , respectively; L_i^* defines the deformed length of the element; NM , NJ are the number of elements, the number of joints, respectively; P_j , U_j are three-dimensional force, displacement vectors at joint j , respectively; the superscript T denotes matrix transposition. The relative axial deflection is the difference between L_i^* and L_i . The initial coordinates and the displacements U_j of the joints completely define L_i and L_i^* [48]; π , Ω are the strain energy, potential energy of external forces, respectively. Eq. 5.1 is identical to Eq. 3.16. Hence, the ordered perturbation equations, Eqs. 3.17 - 3.19, can be used to obtain the nonlinear response of the dome models to static loads.

Eq. 5.1 together with the kinetic energy function $T(Q_i, \dot{Q}_i, \Lambda^j)$ and Rayleigh's dissipation function [75] forms the basis of the dynamic analysis described in section 3.3.

The 21 DOF dome model is subjected to two independent loads Λ_1 , Λ_2 while the 39 DOF model is under three independent loads Λ_1 , Λ_2 and Λ_3 .

5.1 Geodesic Dome: 21 DOF Model

The twenty-one DOF geodesic dome model is shown in Fig. 5.1. This model is analyzed for seven different loading conditions:

1. $\Lambda_1 \neq 0, \Lambda_2 = 0$
2. $\Lambda_1 = \Lambda_2 = \Lambda$
3. $\Lambda_1 > \Lambda_2$
4. $\Lambda_1 < \Lambda_2$
5. Asymmetric Loading; i.e., joints 1,2,3,7 loaded
6. $\Lambda_2 \neq 0, \Lambda_1 = 0$
7. Single concentrated load at joint 2

The variations of Λ_1 and Λ_2 are arbitrary since they are independent loads. At failure, the critical load on the system is denoted Λ^C . If both Λ_1 and Λ_2 are present, Λ^C takes the values of Λ_1 at failure. If only Λ_2 is present, then Λ^C takes the value of Λ_2 at failure. Similarly, if only Λ_1 is present, Λ^C represents the critical value of Λ_1 at failure. This notation is used in all the static response ($\Lambda - Q_i$) curves.

The effect of shallowness for the dome model is shown in Fig. 5.3. In this illustrative figure, s denotes the span of the dome while h represents the height of joint 1 above the supports. The dome admits only limit points if $s \geq 10h$ while bifurcation states are possible if $s \leq 8h$. For $8 < s/h < 10$, the form of failure depends on the loading configuration; i.e., both limit and bifurcation states are possible in this range. The structure has the freedom to assume any bifurcation configuration. The inequality $s \geq 8h$ satisfies the conditions for shallowness in the classical shallow shell theory [87, 88]. As

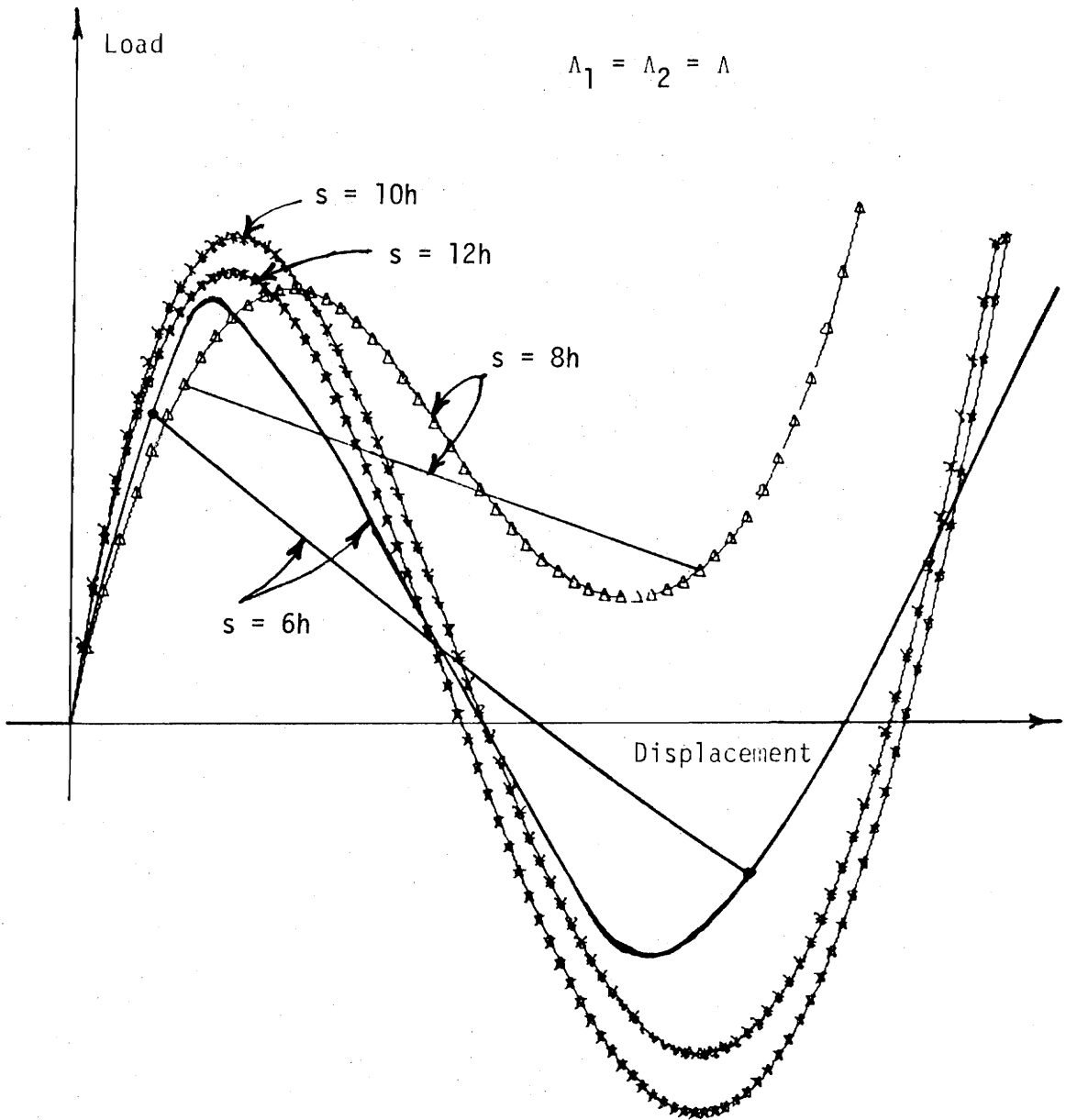


FIG. 5.3 EFFECT OF SHALLONESS

shown in Fig. 5.3, a shallow shell admits both limit and bifurcation critical states under some loading conditions. However, if the shell is sufficiently shallow, $s \geq 10h$, only limit points will be admitted as critical states regardless of loading configurations.

As indicated in Fig. 5.3, the reticulated dome exhibits unstable critical (limit and bifurcation) states. Hence, collapse is into an equilibrium configuration which is a finite distance away from the fundamental state. In this case, the buckling load of the structure is sensitive to geometric and/or loading imperfections [114]. At the initial stage of buckling, deformations of the system are symmetric in the case of a limit point. Bifurcation is characterized by initial unsymmetrical deformations in the post-buckling range prior to complete collapse.

Figs. 5.4 and 5.5 depict the process involved in the search for critical states (limit, bifurcation, and coincident critical points). In the case of a limit point (Fig. 5.4), the determinant $\det(W_L)$ vanishes and $\dot{\eta} = 0$. For bifurcation points (Fig. 5.5), $\det(W_L)$ and $\det(W_B)$ vanish but $\dot{\eta} \neq 0$. The occurrence of a critical state is detected by changes in sign of the stability determinant.

In Figs. 5.6 - 5.8, the effects of various loading conditions on the forms of instability are depicted. When the center joint is loaded at a rate equal to or higher than the ring joints, the model attains a limit point (Figs. 5.6, 5.7). However, when the ring joints take on more loads than the center joint, buckling occurs at a bifurcation

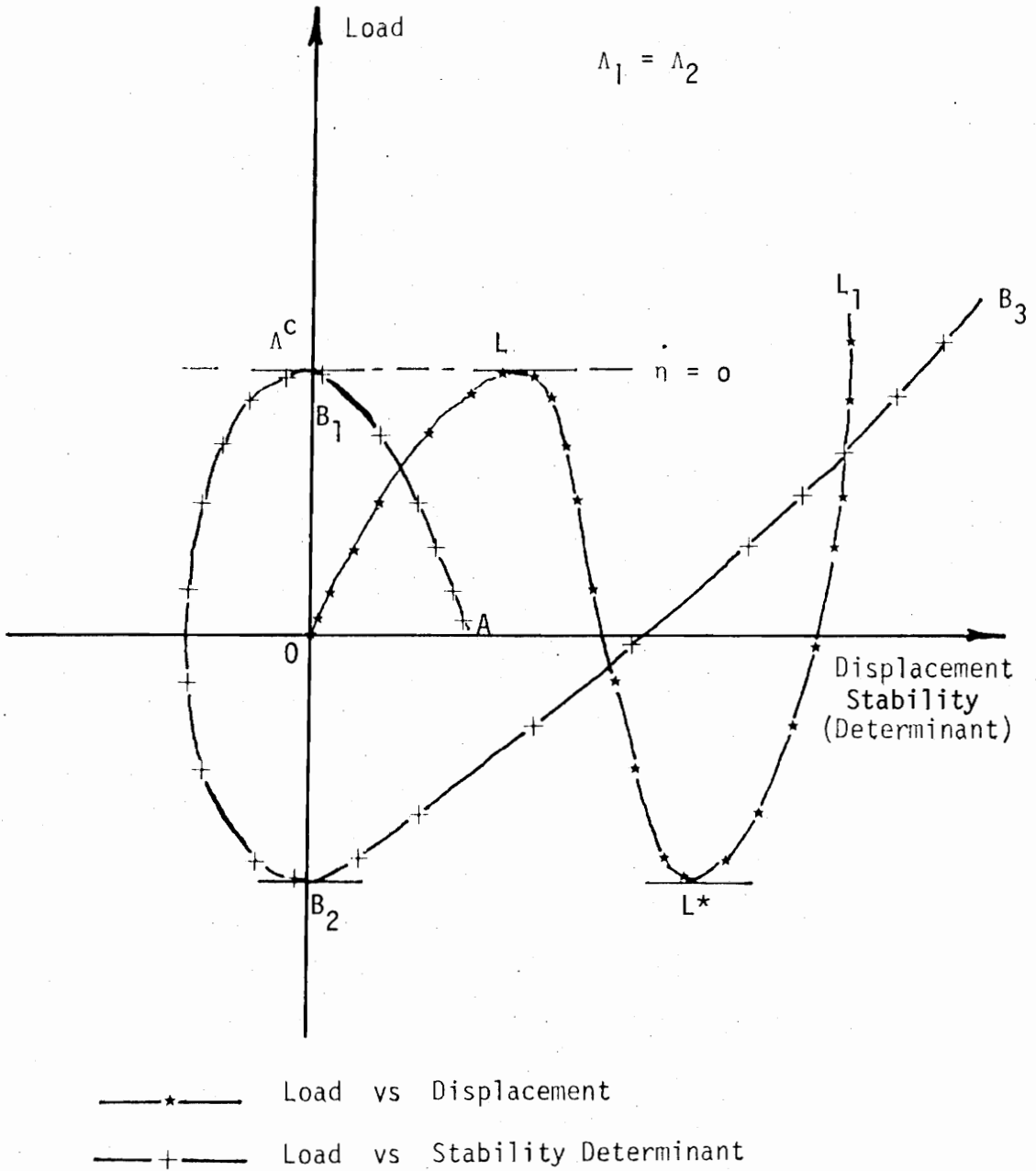


FIG. 5.4 LIMIT POINT LOCATION

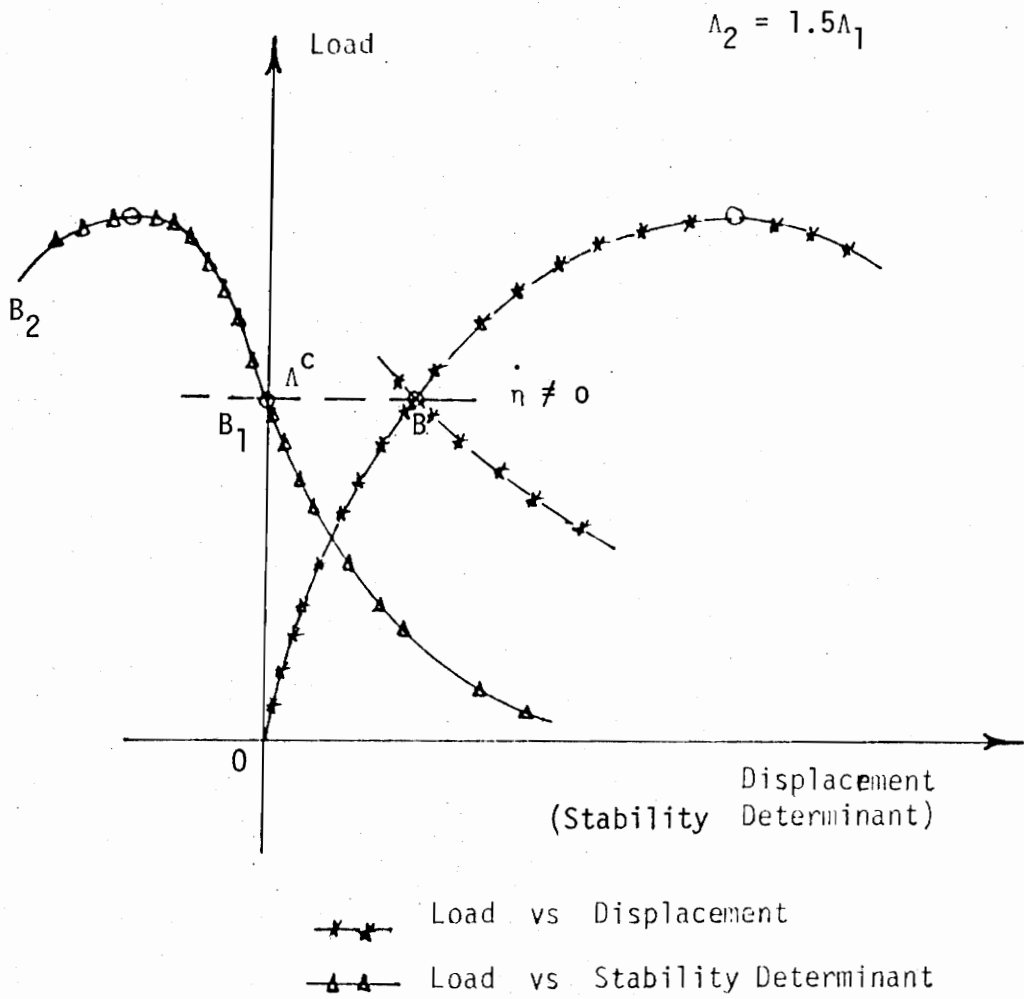


FIG. 5.5 BIFURCATION POINT LOCATION

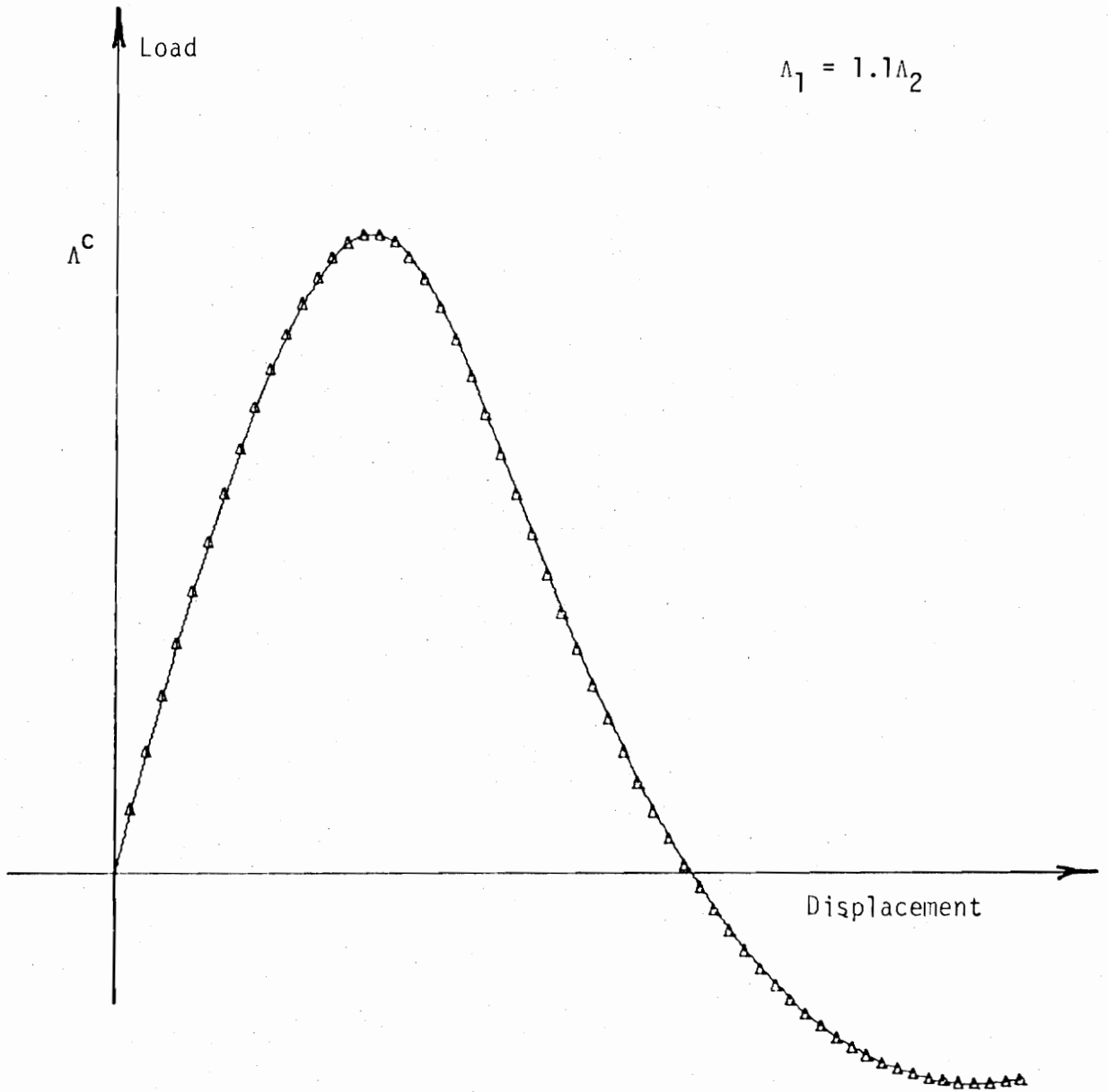


FIG. 5.6 EQUILIBRIUM PATH
(Limit-Type Failure)

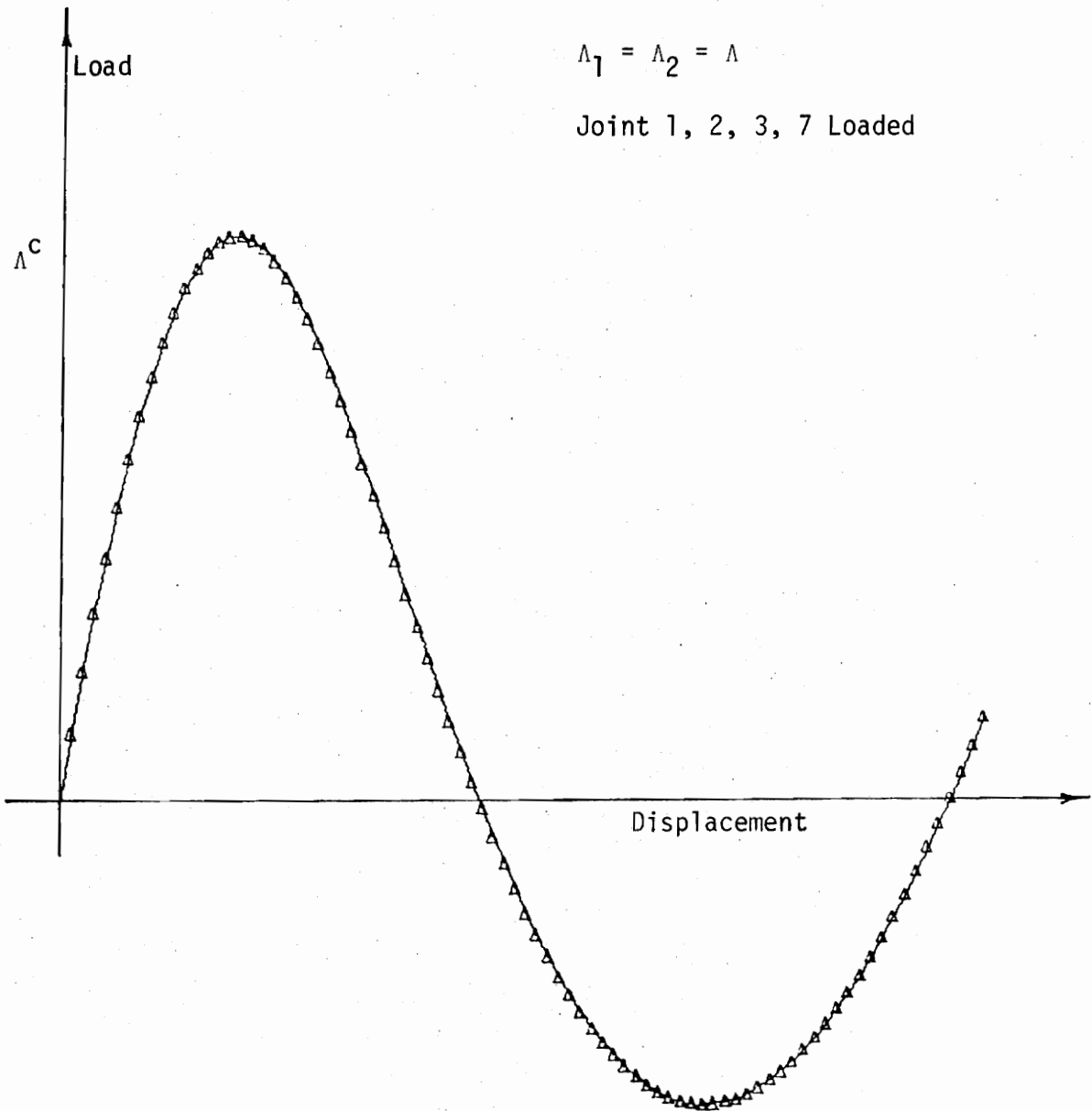


Fig. 5.7 EQUILIBRIUM PATH-ASYMMETRIC LOADING
(Limit-Type Failure)

point and the critical load Λ^C is considerably reduced (Fig. 5.8). In this case, equilibrium on the primary path becomes unstable when the load-level reaches Λ^C and thereafter.

The sensitivity of the reticulated dome model to imperfections is shown in Fig. 5.9. In the presence of loading imperfections, limit points are admitted as critical states. In Fig. 5.9, loading imperfections are introduced at joints 2 and 5, i.e. at these joints, Λ_2/Λ_1 equals 1.8. The perfect system corresponds to the ratio Λ_2/Λ_1 equals 2; i.e. all ring joints loaded at twice the rate at which joint 1 is loaded. The bifurcation load of the perfect system is an upper bound on the limit load of the imperfect system (Fig. 5.9). This figure demonstrates the need to locate all critical states of the perfect system for design. If the bifurcation load were not found, any design based on the limit load of the perfect system would be nonconservative and hence impractical.

Figs. 5.10 and 5.11 illustrate the relation between the load parameter, Λ , and the degree of stability, M . As the load increases, the degree of stability decreases. M approaches zero at the critical load, Λ^C . It appears that equilibrium is more sensitive to small (but finite) disturbances in the neighborhood of bifurcation points (characterized by $\Lambda_2 > \Lambda_1$) than around limit point (characterized by $\Lambda_1 > \Lambda_2$). In both cases, the equilibrium states in the vicinity of the critical point are quite sensitive to disturbances since the slope of the $\Lambda - M$ curves approaches infinity.

Fig. 5.12 shows the effect of geometric imperfections on the dome

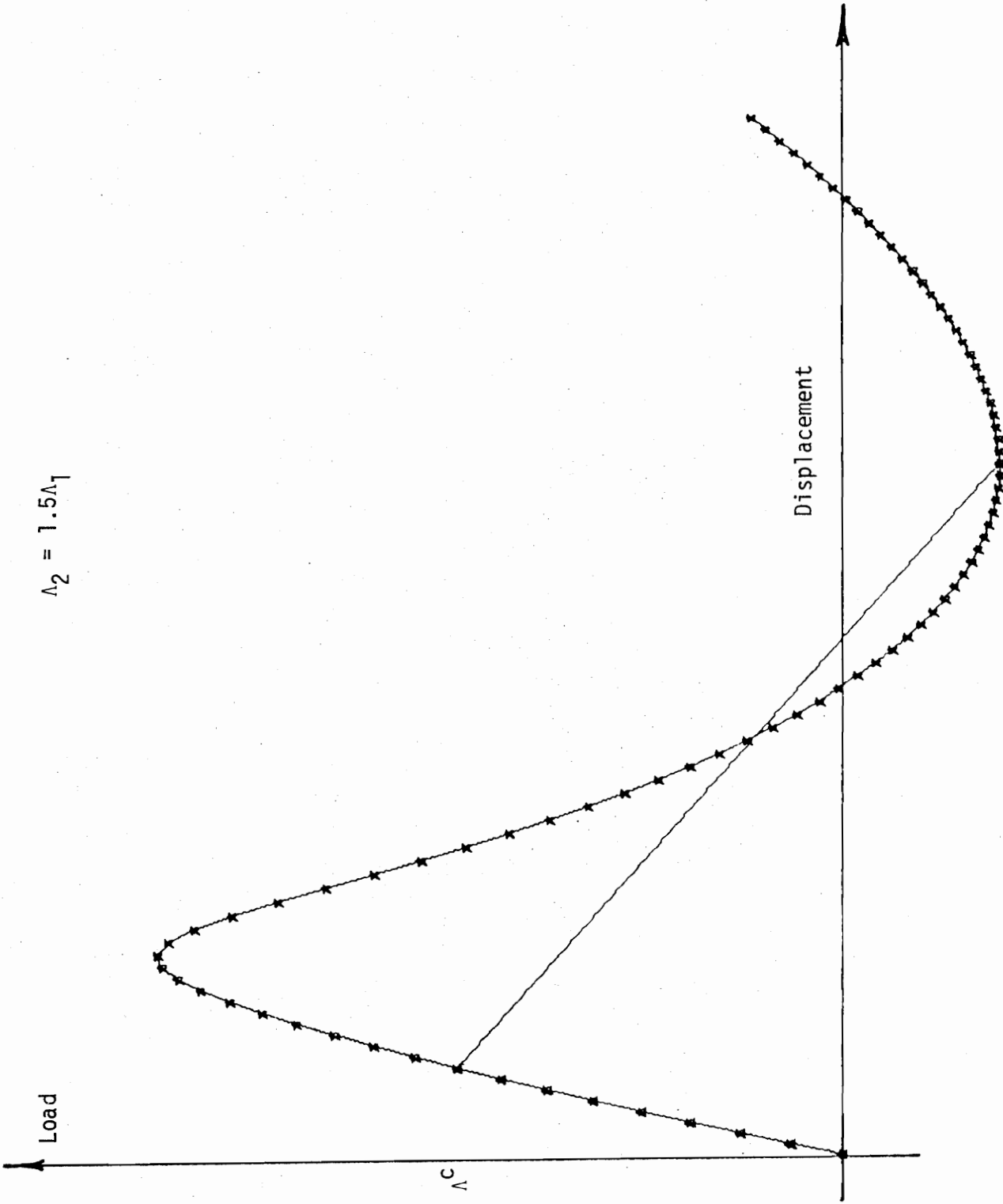


FIG. 5.8 EQUILIBRIUM PATH
(Bifurcation-Type Failure)

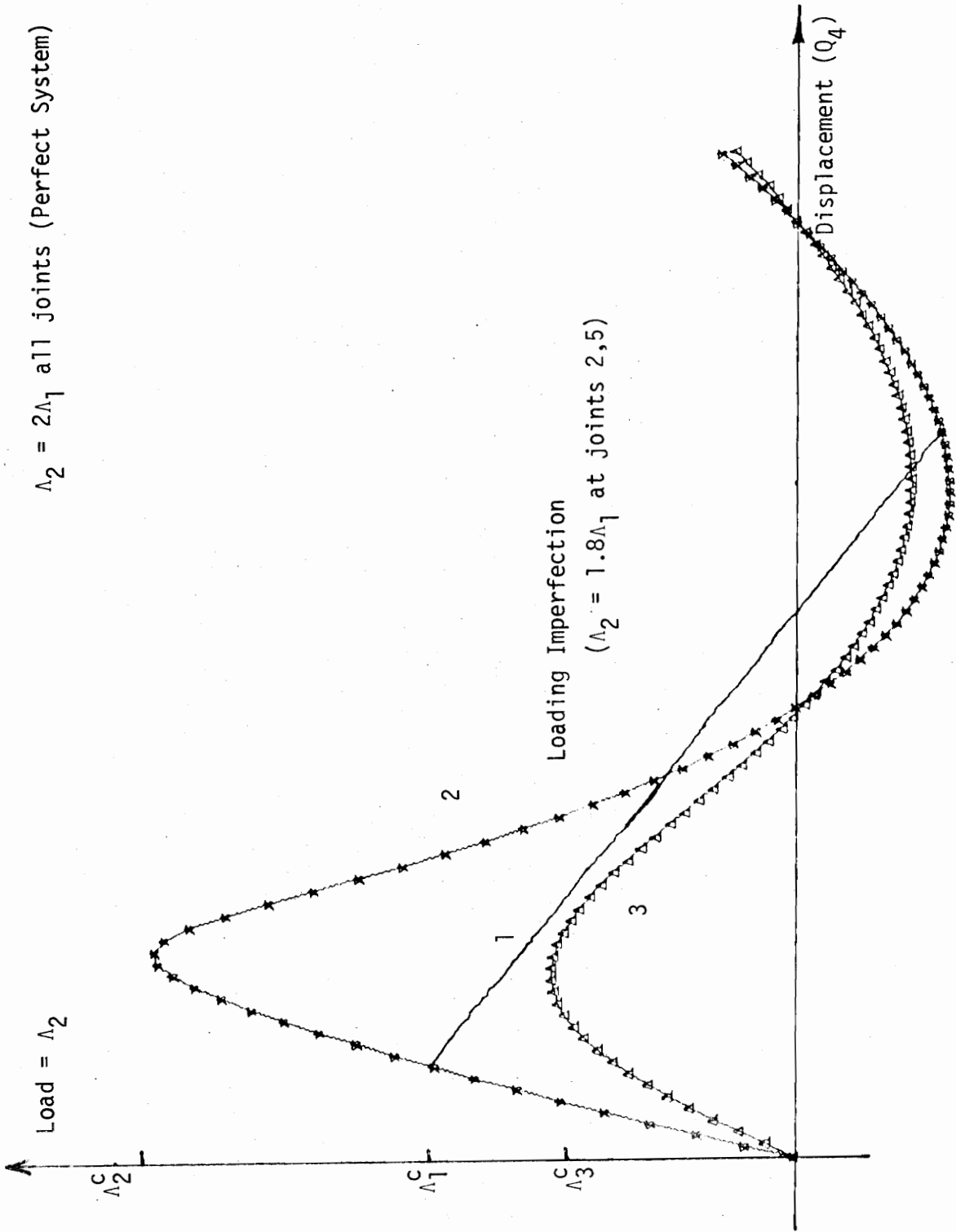


FIG. 5.9 BIFURCATION BUCKLING
(Imperfection-Sensitivity)

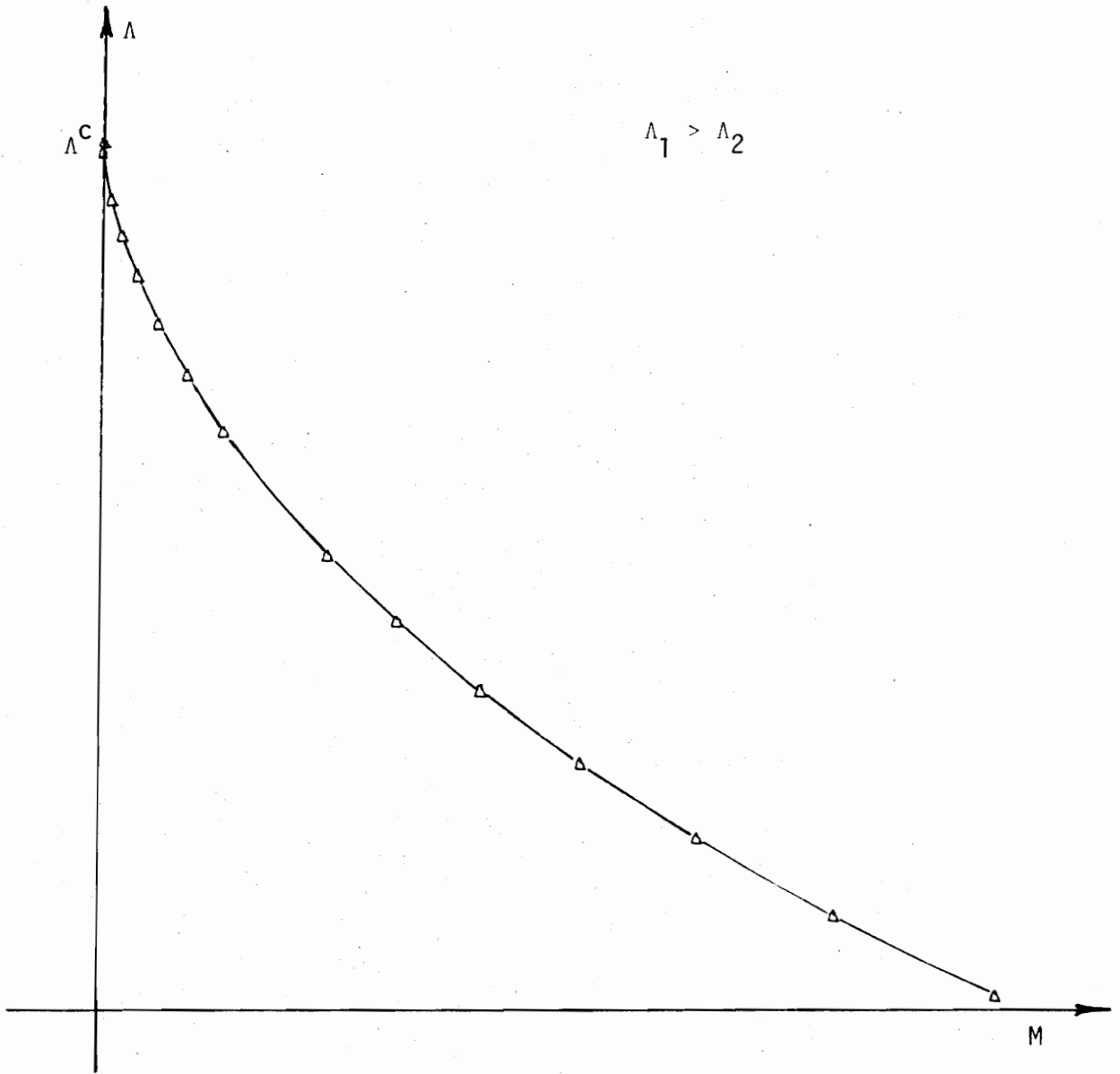


FIG. 5.10 DEGREE OF STABILITY

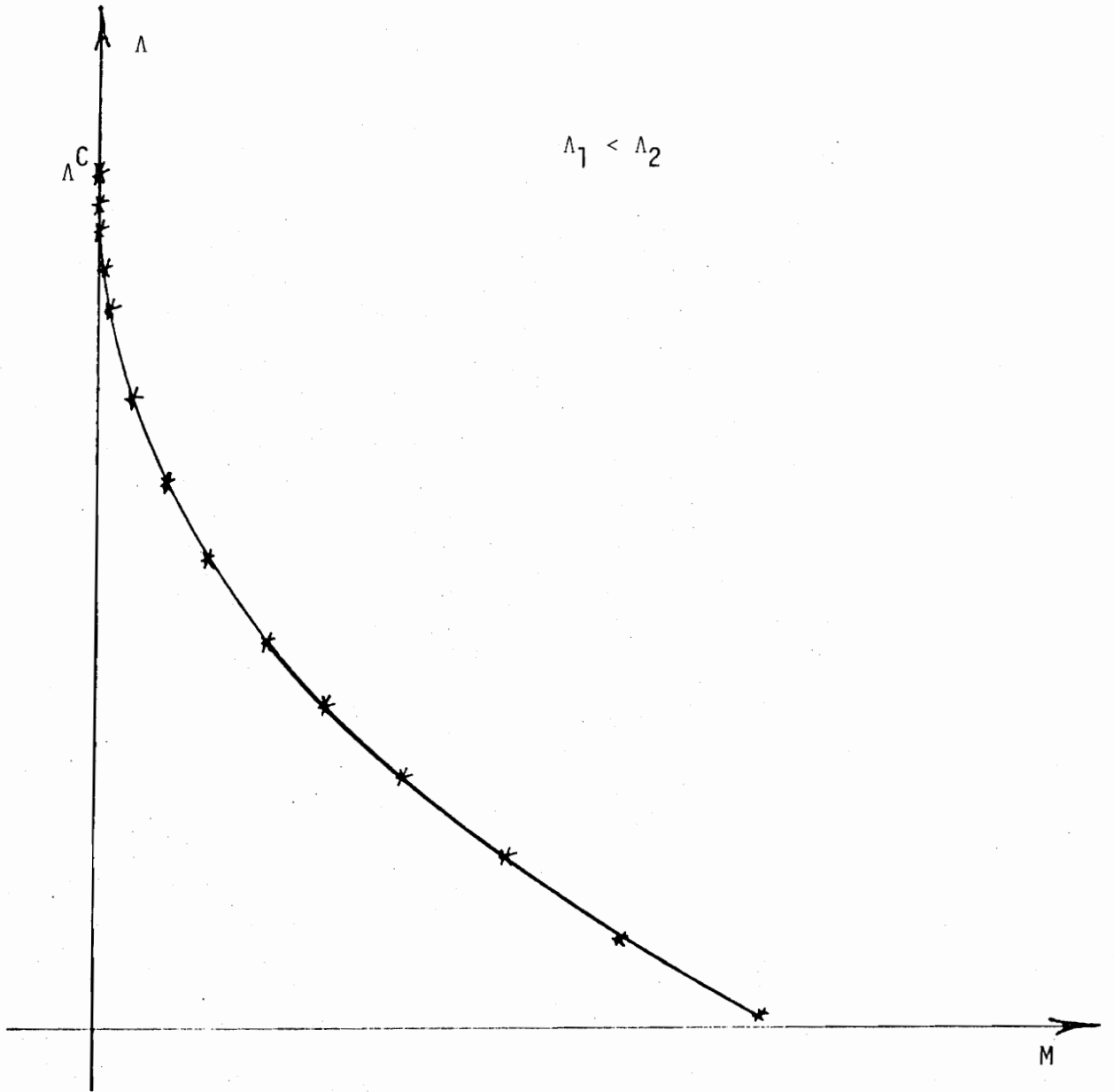


FIG. 5.11 DEGREE OF STABILITY

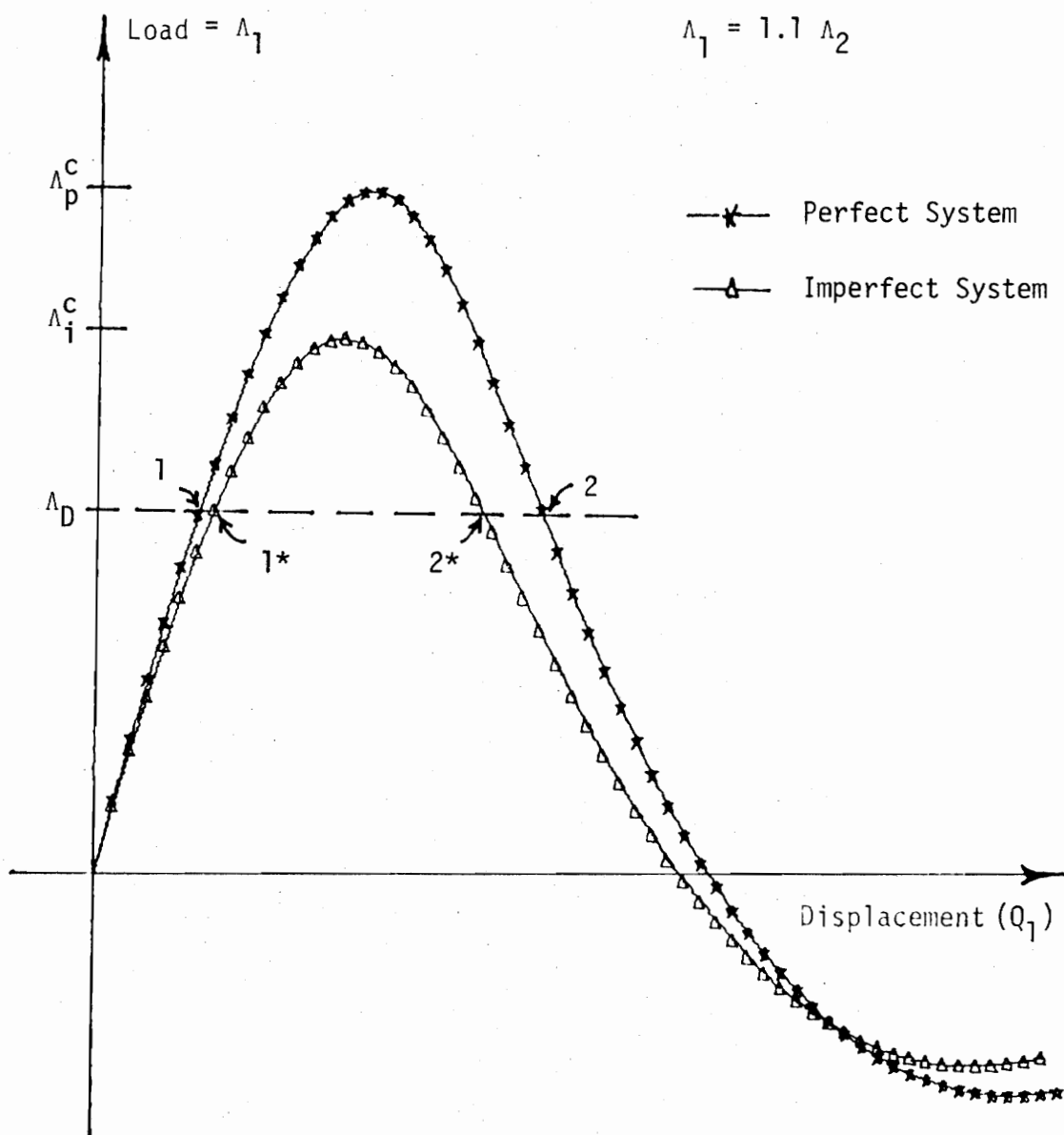


FIG. 5.12 GEOMETRIC IMPERFECTION

model. Geometric imperfections are introduced by reducing by 0.2 the coordinates of joints 2 and 5 in the X-direction (Fig. 5.1). As in the case of loading imperfections (Fig. 5.9), the limit load drops from Λ_p^C for the perfect dome to Λ_i^C for the imperfect dome. Point 2 represents the "nearest" unstable equilibrium state for the perfect dome. Correspondingly, point 2* represents the "nearest" unstable equilibrium state for the imperfect dome for a given design load, Λ_D . As shown in Fig. 5.13, the degree of stability, M_D , for the perfect system is reduced to a value M_D^* for the imperfect system. Hence imperfection sensitivity in reticulated domes are characterized by a reduction of critical load, Λ^C , and a reduction of M at the same load level.

In Figs. 5.14 - 5.16, the selection of the "nearest" unstable equilibrium state is illustrated. In the neighborhood of Λ_B^C , path 1 is "nearest" to the fundamental path. For loads less than Λ_1^* , path 2 is "nearest" to the fundamental path. As shown in Fig. 5.15, for loads Λ less than Λ_1^* , using path 1 as opposed to path 2 to estimate M will be nonconservative. This is further illustrated in Fig. 5.16 which includes the effect of loading imperfection where $\Lambda_2 = 2\Lambda_1$ defines the loading configuration for the perfect system. When the unstable equilibrium states "nearest" to the fundamental path are used, a safe estimate of the degree of stability of equilibrium is obtained. Fig. 5.16 again shows that the $\Lambda - M$ curve for an imperfect system is a lower bound for the $\Lambda - M$ curve of the corresponding perfect system.

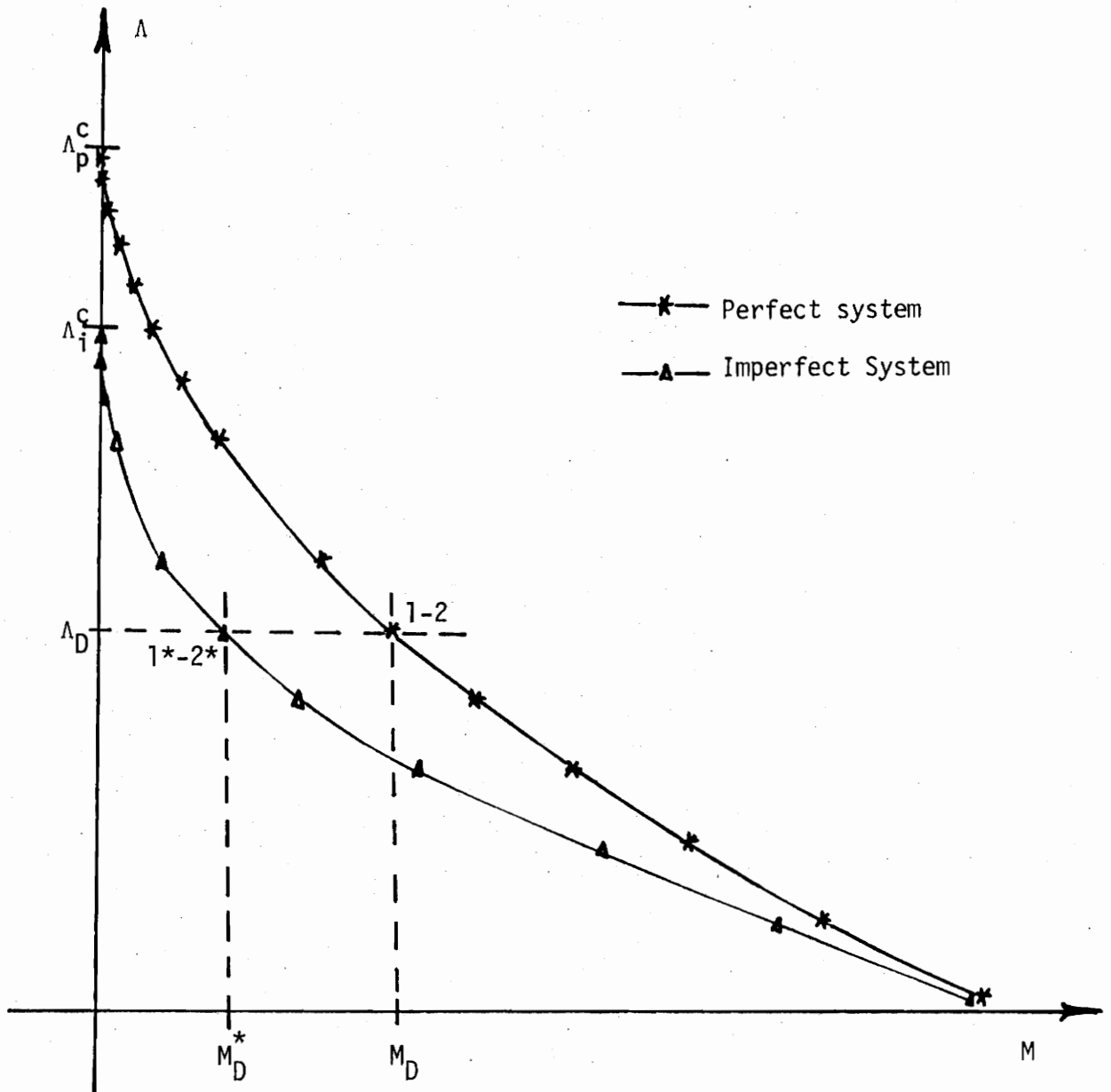


FIG. 5.13 EFFECT OF IMPERFECTIONS ON DEGREE OF STABILITY

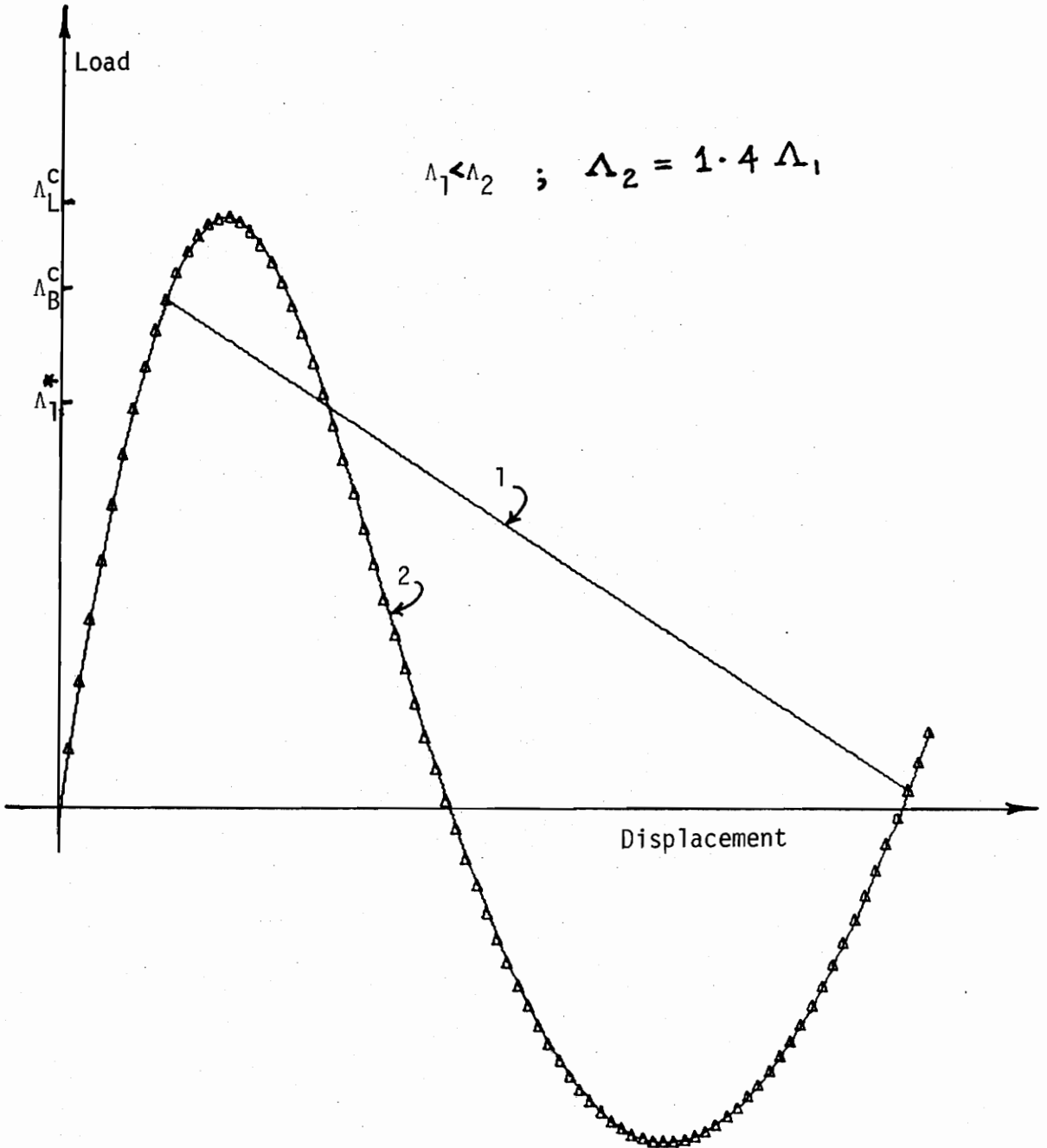


FIG. 5.14 EQUILIBRIUM PATH
("Nearest" Unstable Equilibrium States)

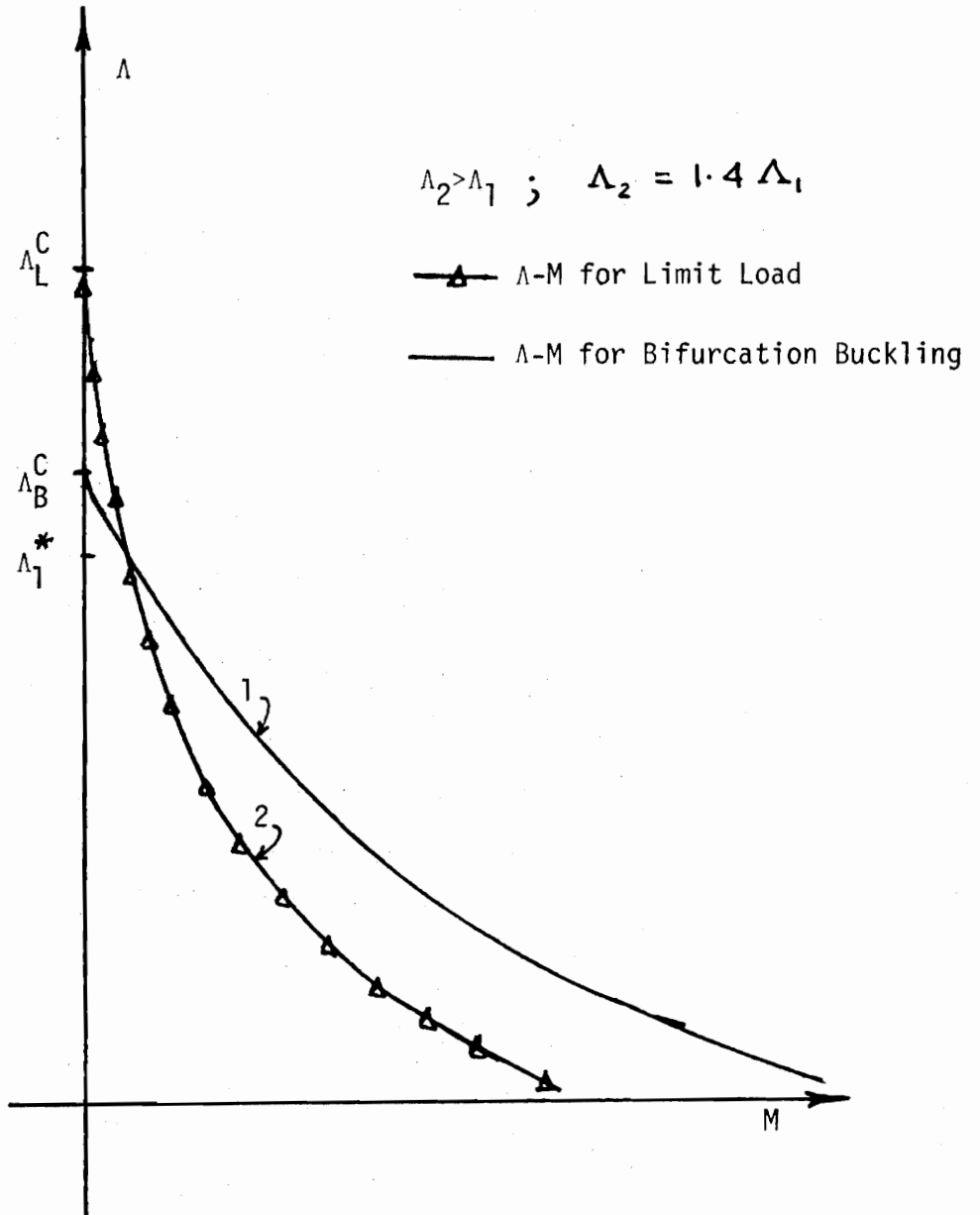


FIG. 5.15 RELATIVE DEGREE OF STABILITY

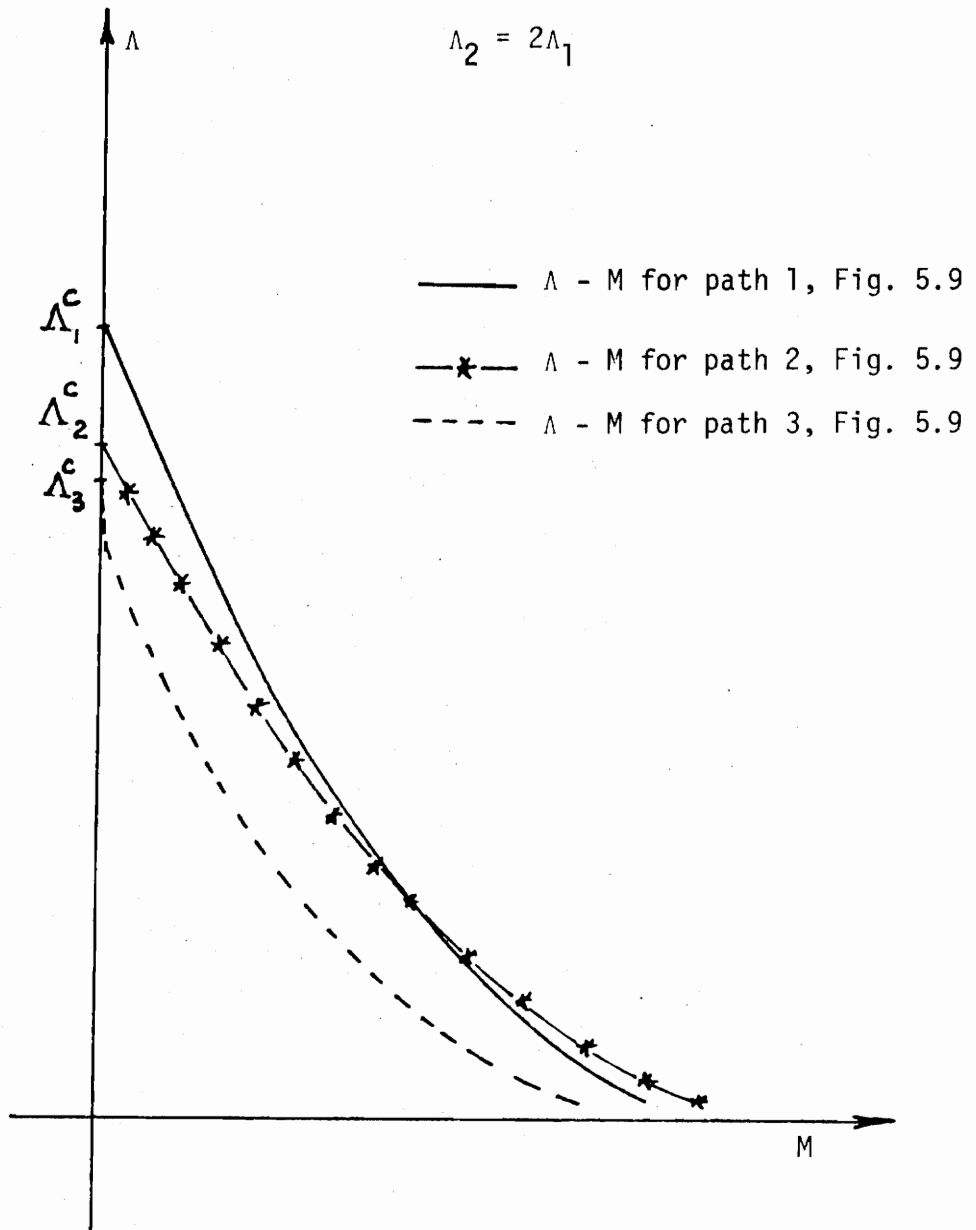


FIG. 5.16 IMPERFECTION SENSITIVITY AND DEGREE OF STABILITY

a. Stability of Equilibrium relative to Instantaneous Disturbances

The stability of equilibrium of the reticulated dome relative to instantaneous disturbances is tested using the sufficient stability condition (section 2.2) and by studying the perturbed motion. Instantaneous disturbances are defined in terms of initial conditions. The equilibrium state of the dome corresponds to a load level that is approximately 0.70 - 0.75 of the critical static load, Λ^C . The sufficient stability criterion defined by Eqs. 2.6 - 2.9 of section 2.2 are used as a basis for selecting initial conditions.

Figs. 5.17 and 5.18 depict the responses of joint 1 and joint 2 to an instantaneous disturbance applied at joint 1. The equilibrium configuration, the origin 0 in the phase plane, corresponds to two independent loads Λ_1, Λ_2 acting at all joint as shown in Fig. 5.1. The loading configuration is such that $\Lambda_1 = 1.1 \Lambda_2$. The disturbance is defined by an initial state of zero displacements and a single velocity component acting at joint 1 in the X-direction. The initial energy v_0 is such that $v_0 = 0.95M$. The resulting undamped motions are closed trajectories surrounding the isolated singularity 0. Hence the system (i.e., the origin) is stable with respect to the disturbance associated with an initial energy v_0 . This implies that the perturbed motion is confined to the domain of asymptotic stability of the origin. It further implies that the initial disturbance, y_0 , belongs to C_M (a subset of the domain of attraction of origin). This is the sufficient stability criterion [45].

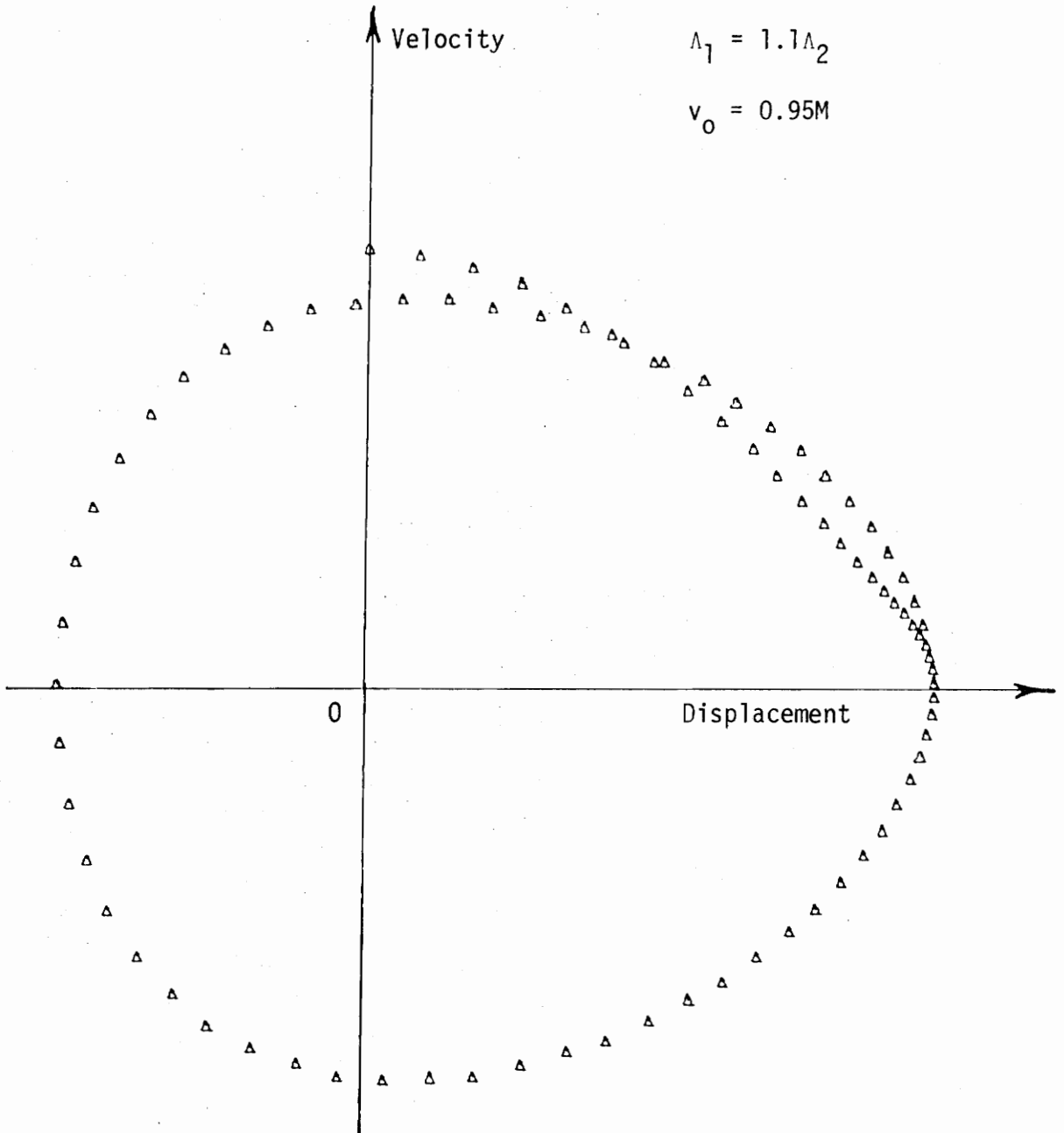


FIG. 5.17 RESPONSE OF JOINT 1
(Joint 1 Loaded with Instantaneous Disturbance)

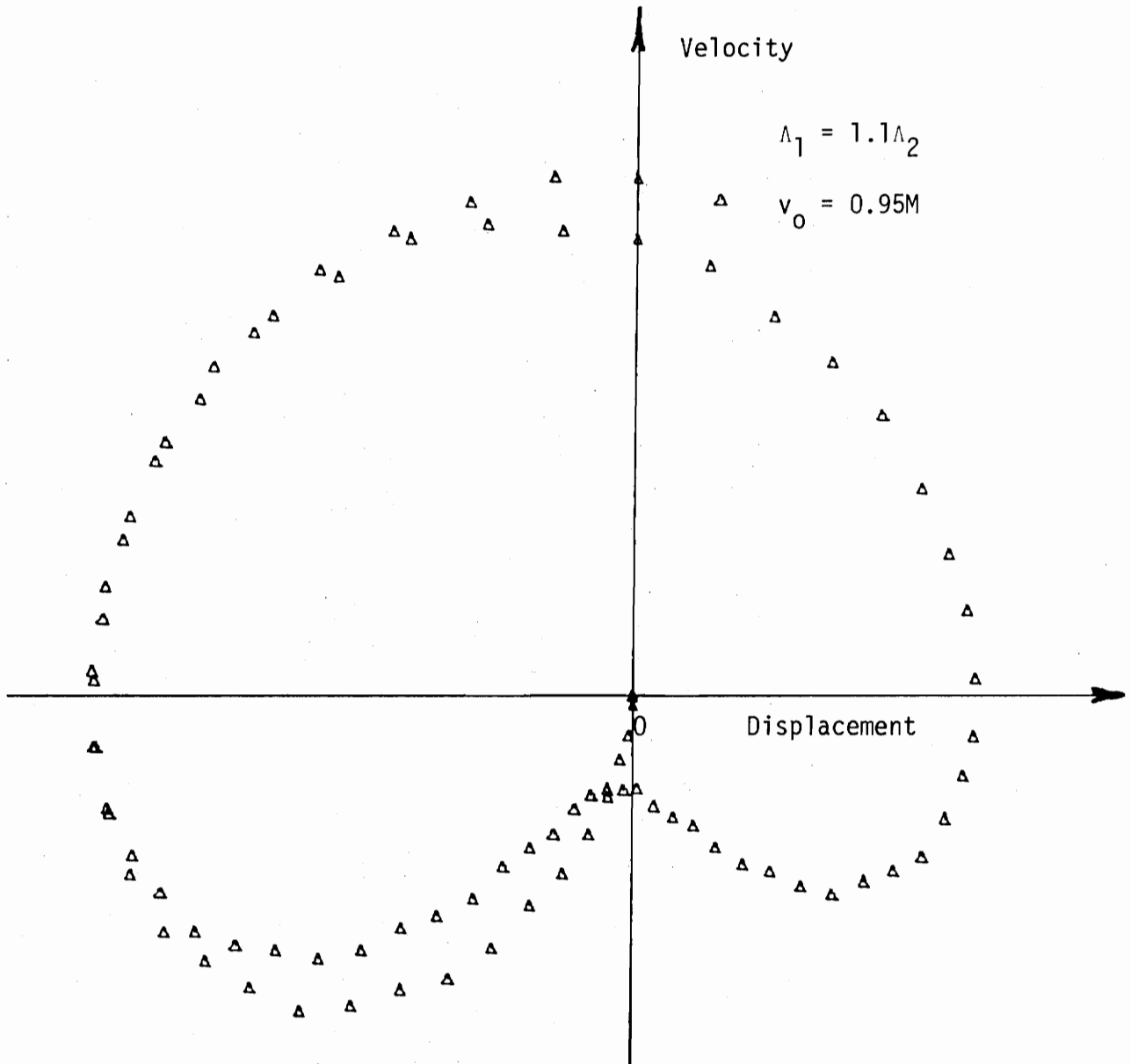


FIG. 5.18 RESPONSE OF JOINT 2
(Joint 1 Loaded with Instantaneous Disturbance)

With damping, the perturbed motion of the system relative to the initial disturbance described above converges to the origin of the phase plane. In this case, equilibrium (origin) is asymptotically stable. This is illustrated in Fig. 5.19.

Fig. 5.20 depicts the same equilibrium configuration subjected to an initial disturbance with an initial energy greater than M . In this case, with $v_0 = 1.04M$, the motion indicates instability of the origin. The effect of damping is shown in Fig. 5.21. The damped system converges to a new state, the buckled state of the reticulated dome. Since an infinitesimal change in the disturbance causes a finite change in the maximum deviation of the perturbed motion, the sufficient condition is also necessary for stability of equilibrium for this particular case.

With $\Lambda_1 = 1.1 \Lambda_2$, the responses of joint 2 to initial disturbances at joint 2 are shown in Figs. 5.22 and 5.23. In Fig. 5.22, $v_0 = 0.95M$ and the origin is stable as expected from the sufficient stability condition (section 2.3). However, with $v_0 = 1.15M$, the response of joint 2 indicates that the origin is a stable equilibrium point. This inference is made since the perturbed motion encloses only the origin in the phase plane. This case confirms that the stability criterion employed is only sufficient and need not be necessary for stability of equilibrium. This is illustrated in Fig. 5.23.

Fig. 5.24 illustrates the responses of joint 2 to instantaneous disturbance applied at joint 2. In this case, $\Lambda_2 = 2 \Lambda_1$, and the resulting perturbed motions indicate stability, asymptotic stability

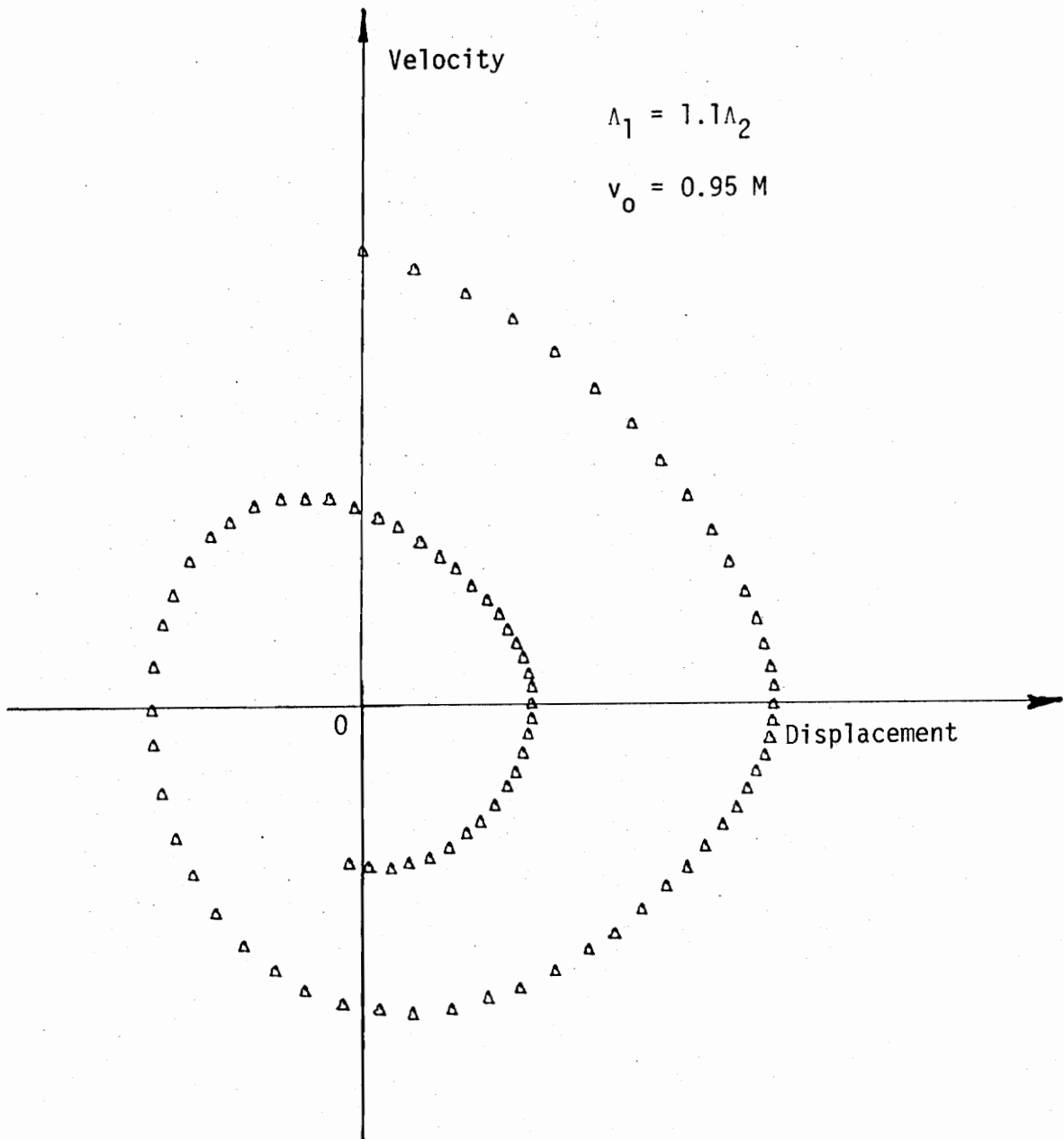


FIG. 5.19 DAMPED RESPONSE OF JOINT 1
(Instantaneous Disturbance at Joint 1)

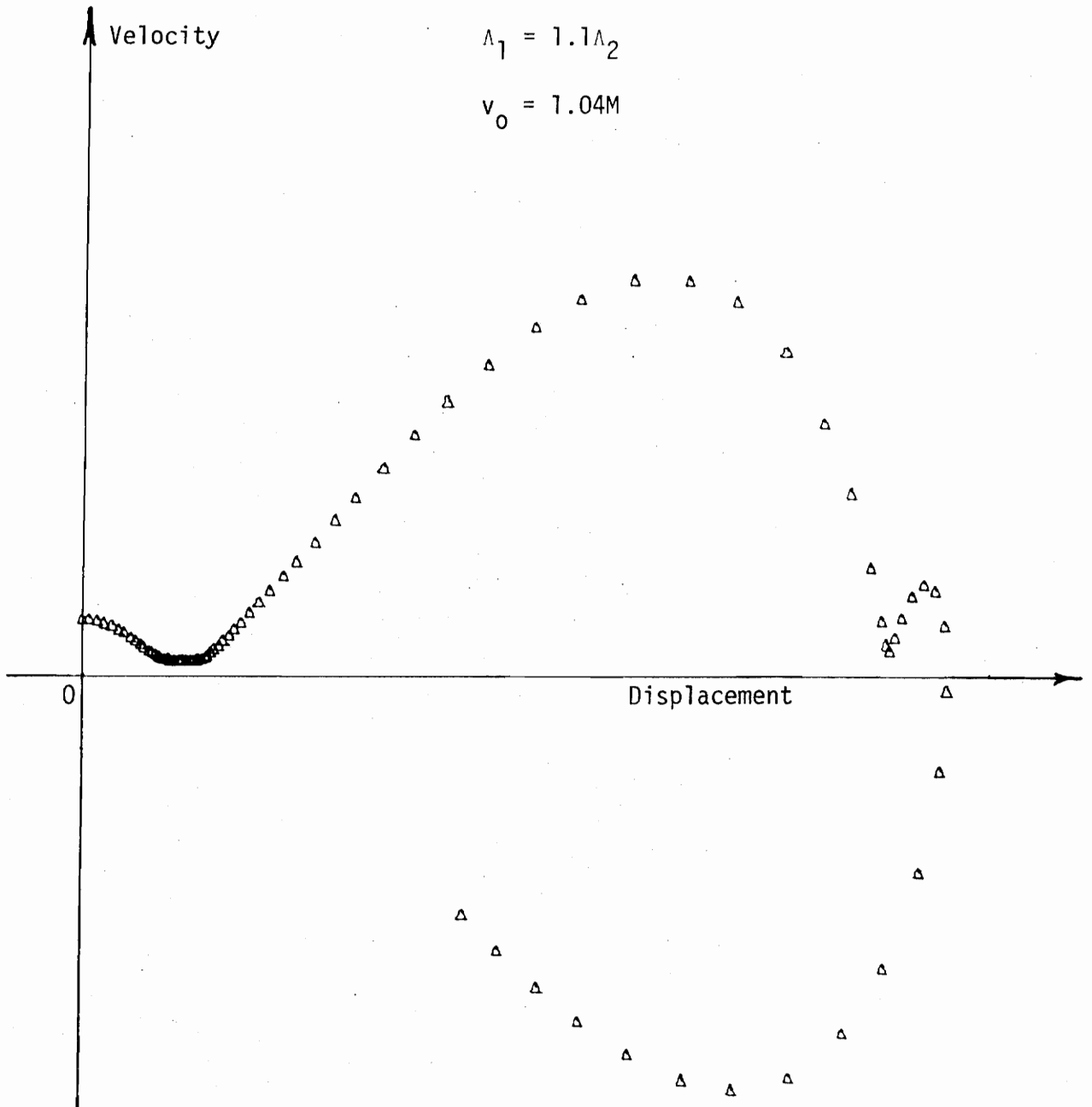


FIG. 5.20 UNDAMPED RESPONSE OF JOINT 1
(Instantaneous Disturbance at Joint 1)

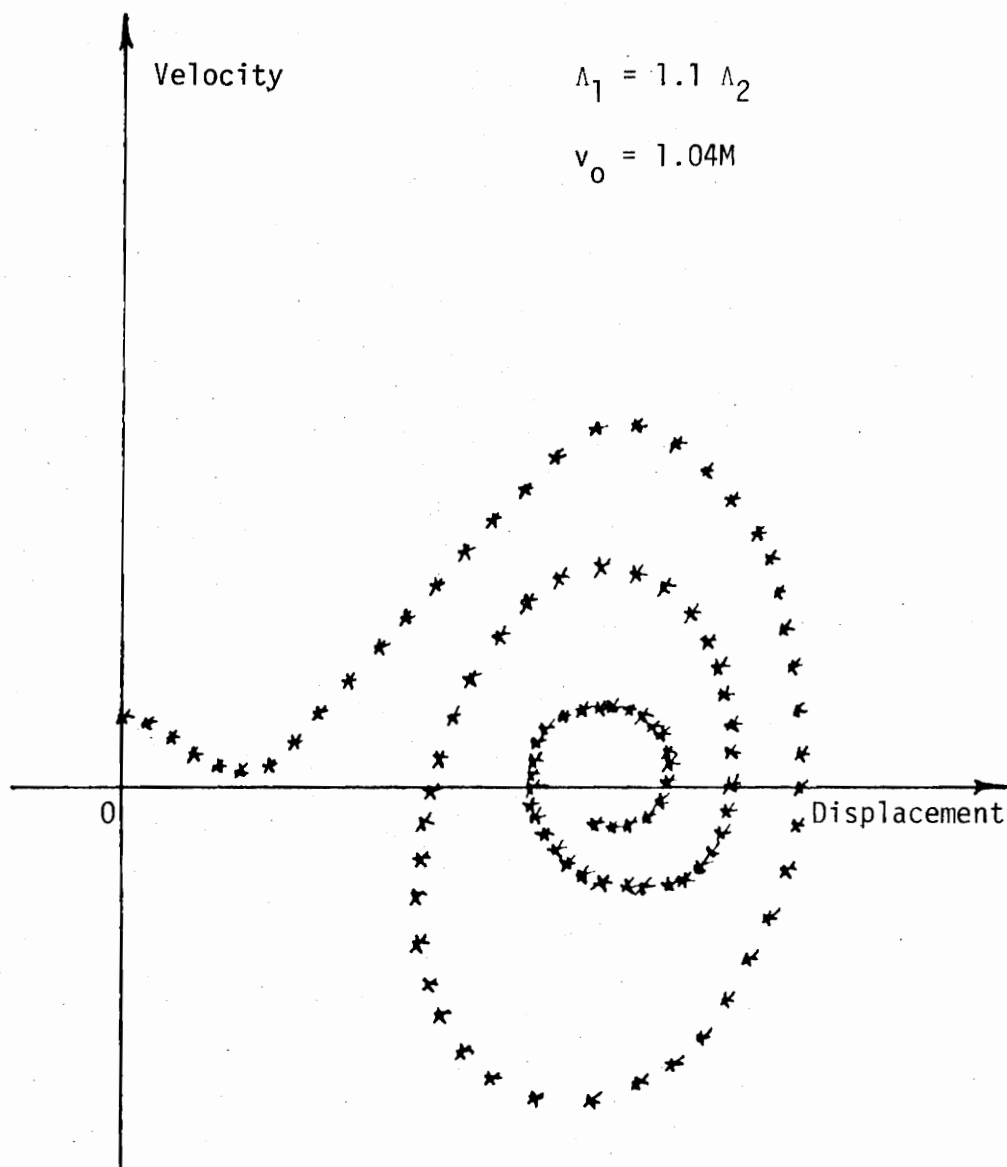


FIG. 5.21 DAMPED RESPONSE OF JOINT 1
(Instantaneous Disturbance at Joint 1)

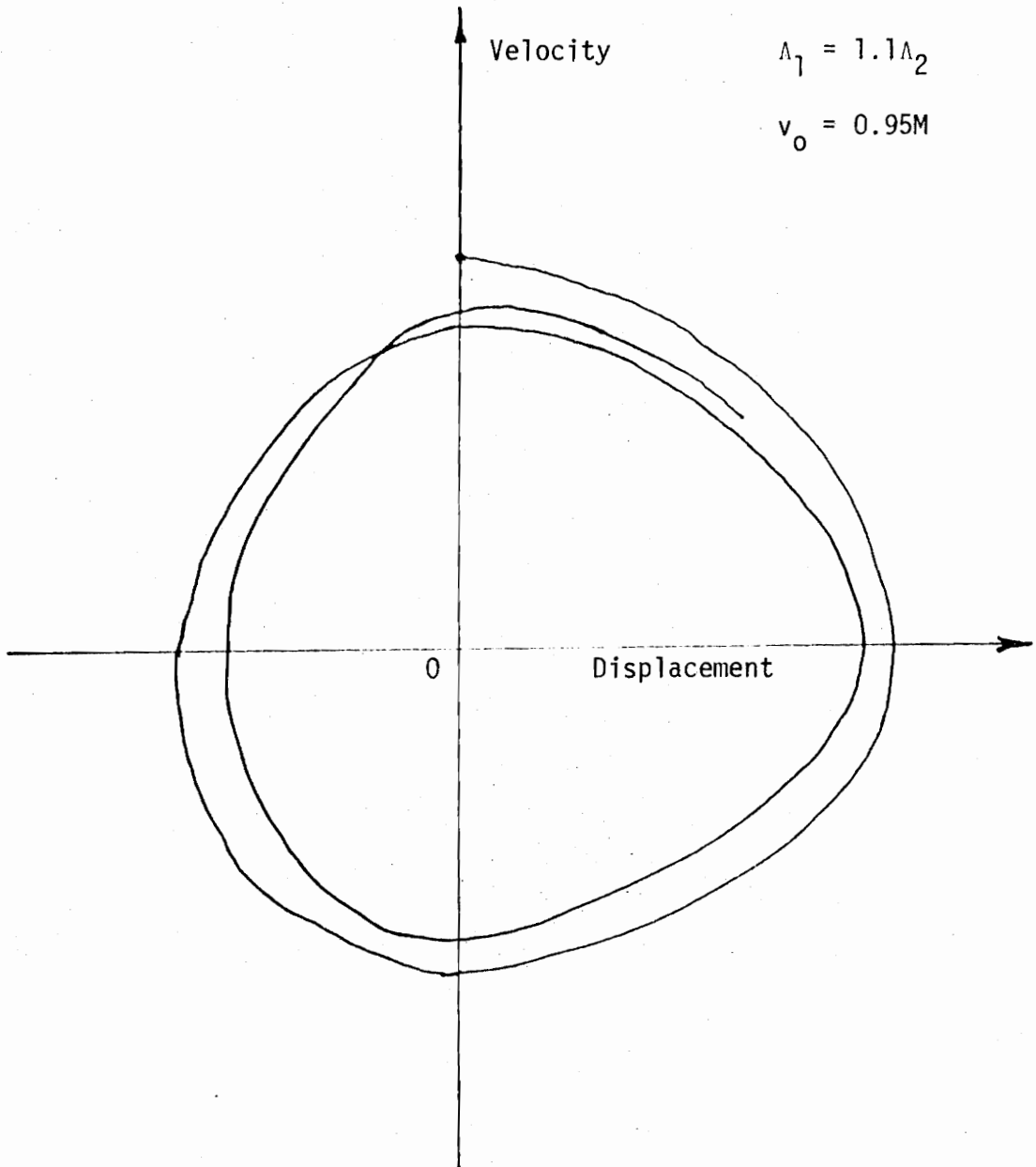


FIG. 5.22 RESPONSE OF JOINT 2
(Instantaneous Disturbance at Joint 2)

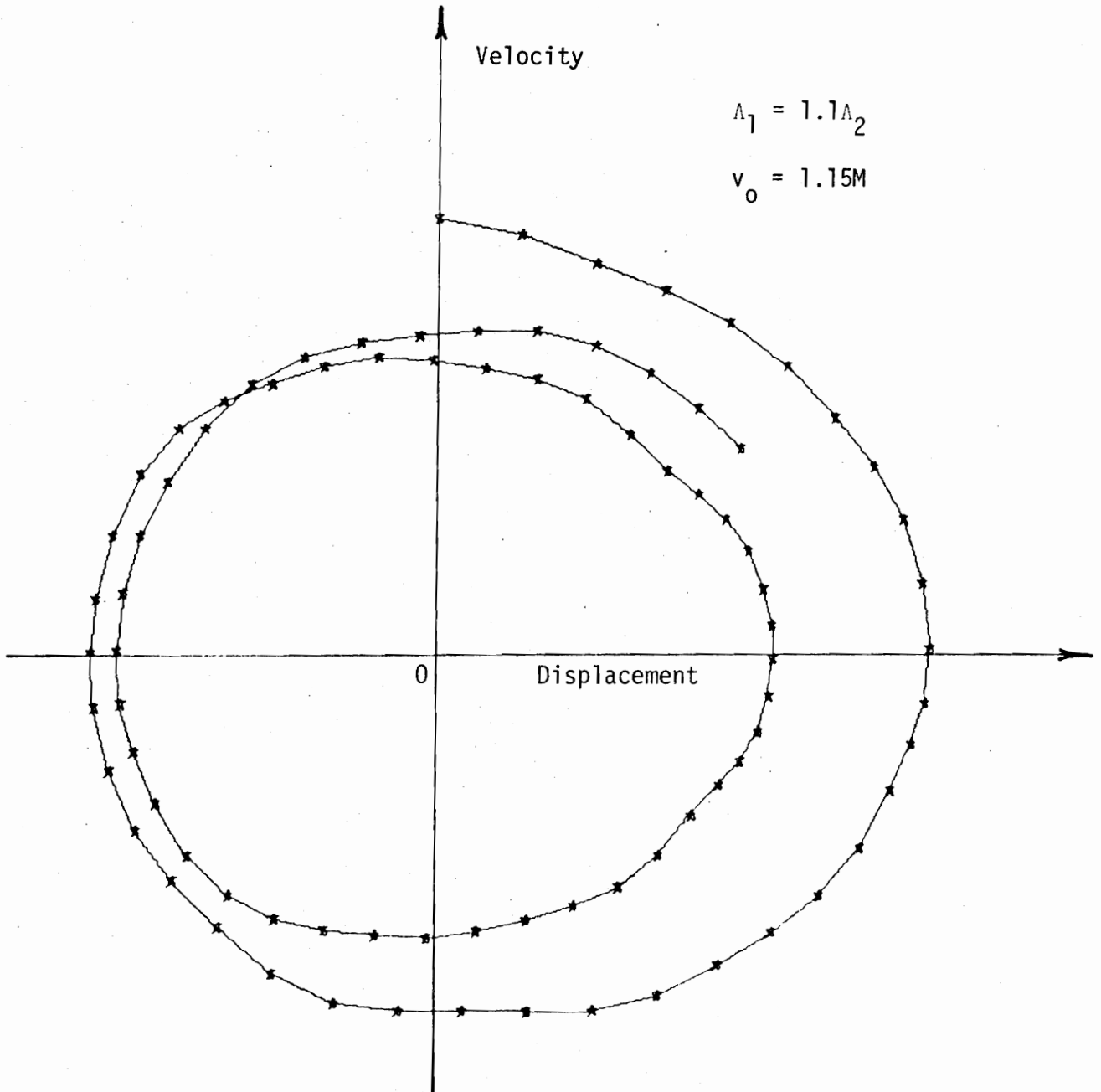


FIG. 5.23 RESPONSE OF JOINT 2
(Instantaneous Disturbance at Joint 2)

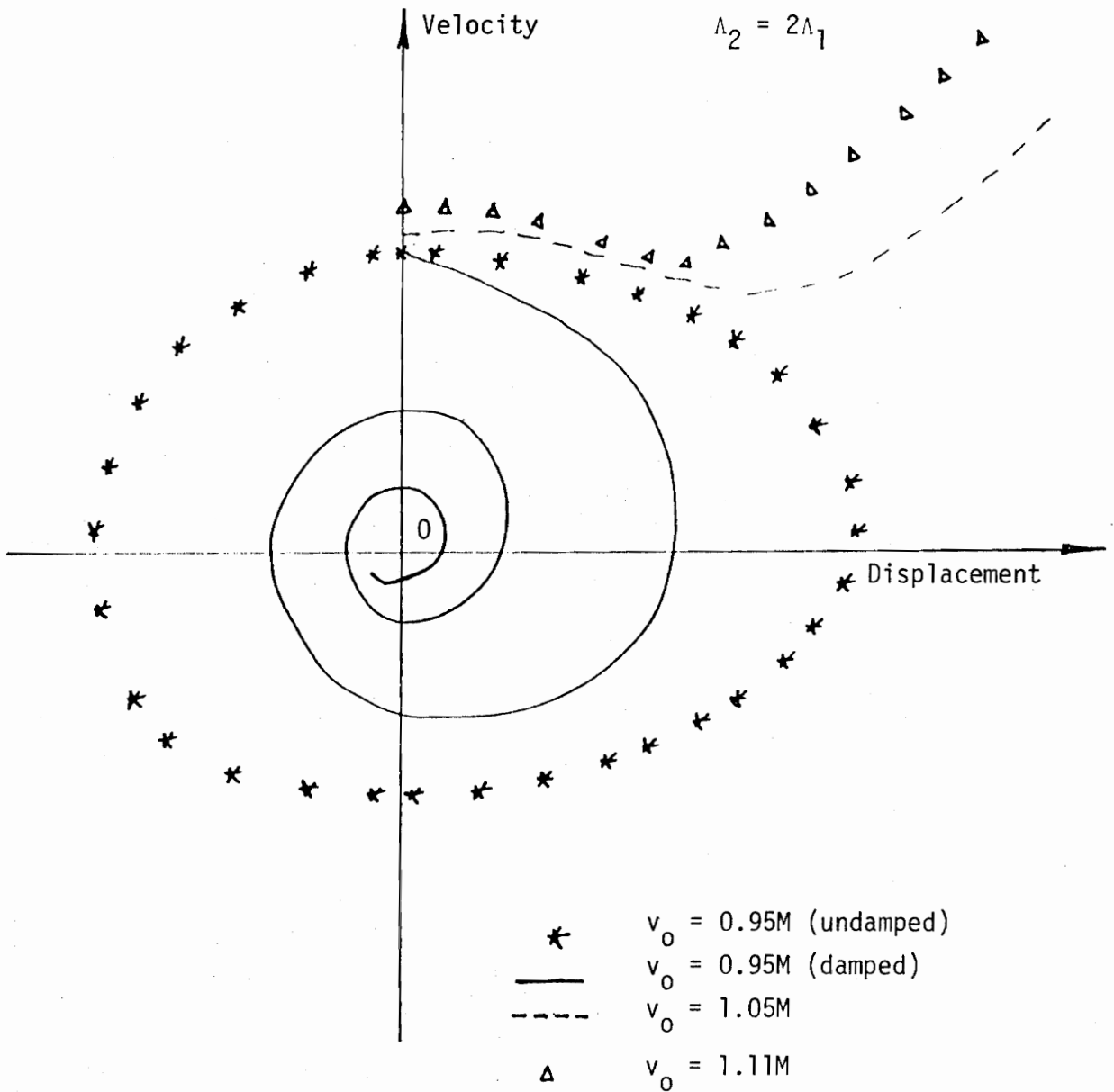


FIG. 5.24 RESPONSE OF JOINT 2
(Instantaneous Disturbance at Joint 2)

and instability of the origin depending on the magnitude of v_0 and the dissipative nature of the system. The sufficient condition is also necessary, in this case, for stability of equilibrium. A qualitatively similar response to that shown in Fig. 5.24 was observed for the case $\Lambda_1 = 0$, $\Lambda_2 \neq 0$ with joints 2, 5 instantaneously disturbed.

Under the equilibrium configuration corresponding to the load distribution $\Lambda_2 = 2 \Lambda_1$, the responses of joint 1 to an instantaneous disturbance applied at joint 1 is shown in Fig. 5.25. The perturbed motions indicate that the stability criterion is only sufficient but not necessary for stability of the equilibrium state 0.

With $\Lambda_1 \neq 0$ and $\Lambda_2 = 0$, the responses of joint 1 to an instantaneous disturbance applied at joint 1 are qualitatively similar to those of Fig. 5.24.

If the dome model is dissipative, the total energy along the perturbed motion might drop below M even when the initial energy, v_0 , is greater than M . This is illustrated in Fig. 5.26 for the loading configuration $\Lambda_1 = 1.1 \Lambda_2$. For such a case, the origin is asymptotically stable and the stability criterion remains only strictly a sufficient stability condition.

The above examples show that the geodesic dome model is more sensitive to disturbances that appear to magnify the equilibrium configuration than to disturbances that tend to excite a different configuration. When the equilibrium configuration is magnified by the dynamic disturbance, the sufficient stability criterion is also necessary for stability of equilibrium. For all cases tested and

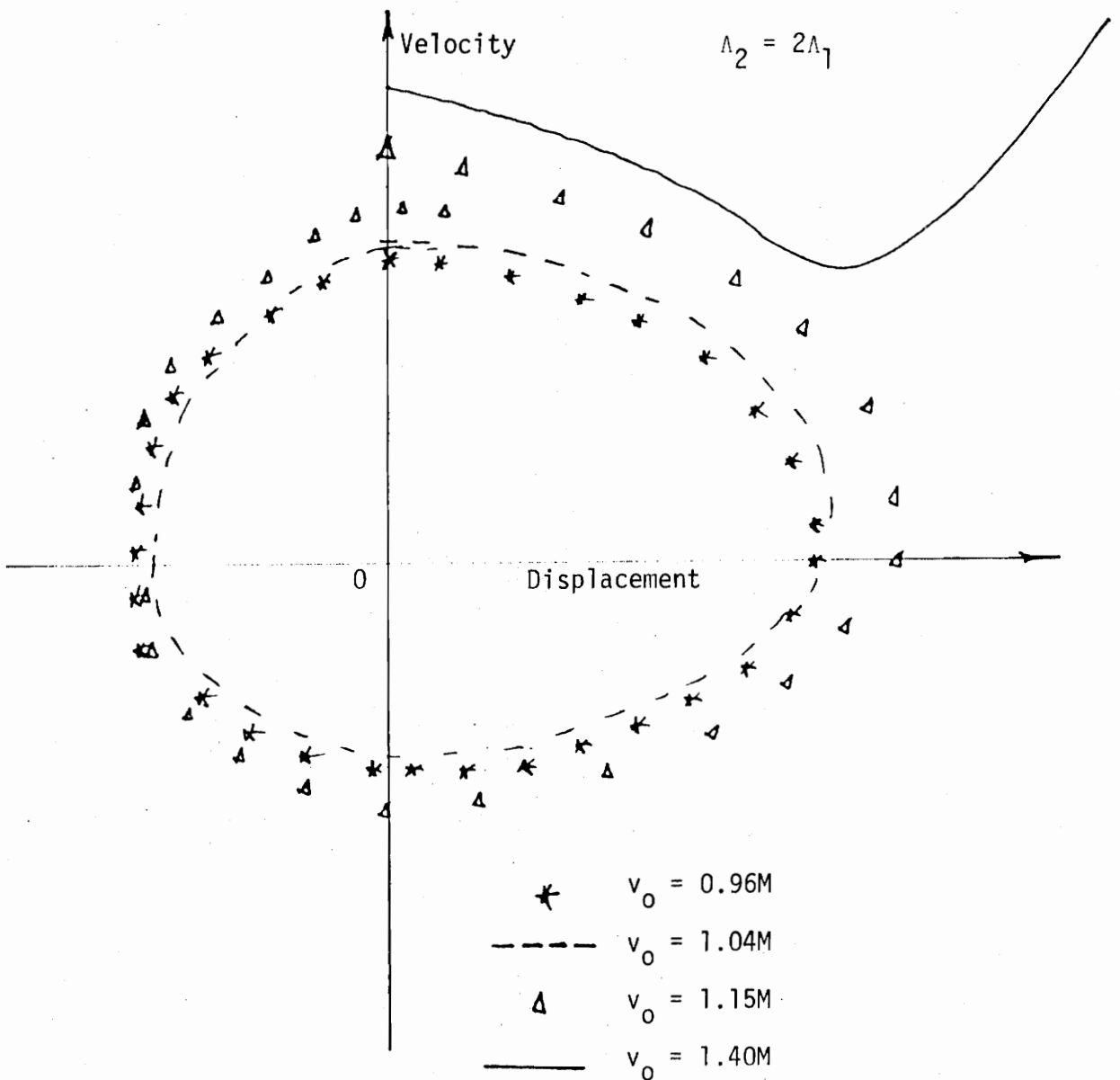


FIG. 5.25 RESPONSE OF JOINT 1
(Instantaneous Disturbance at Joint 1)

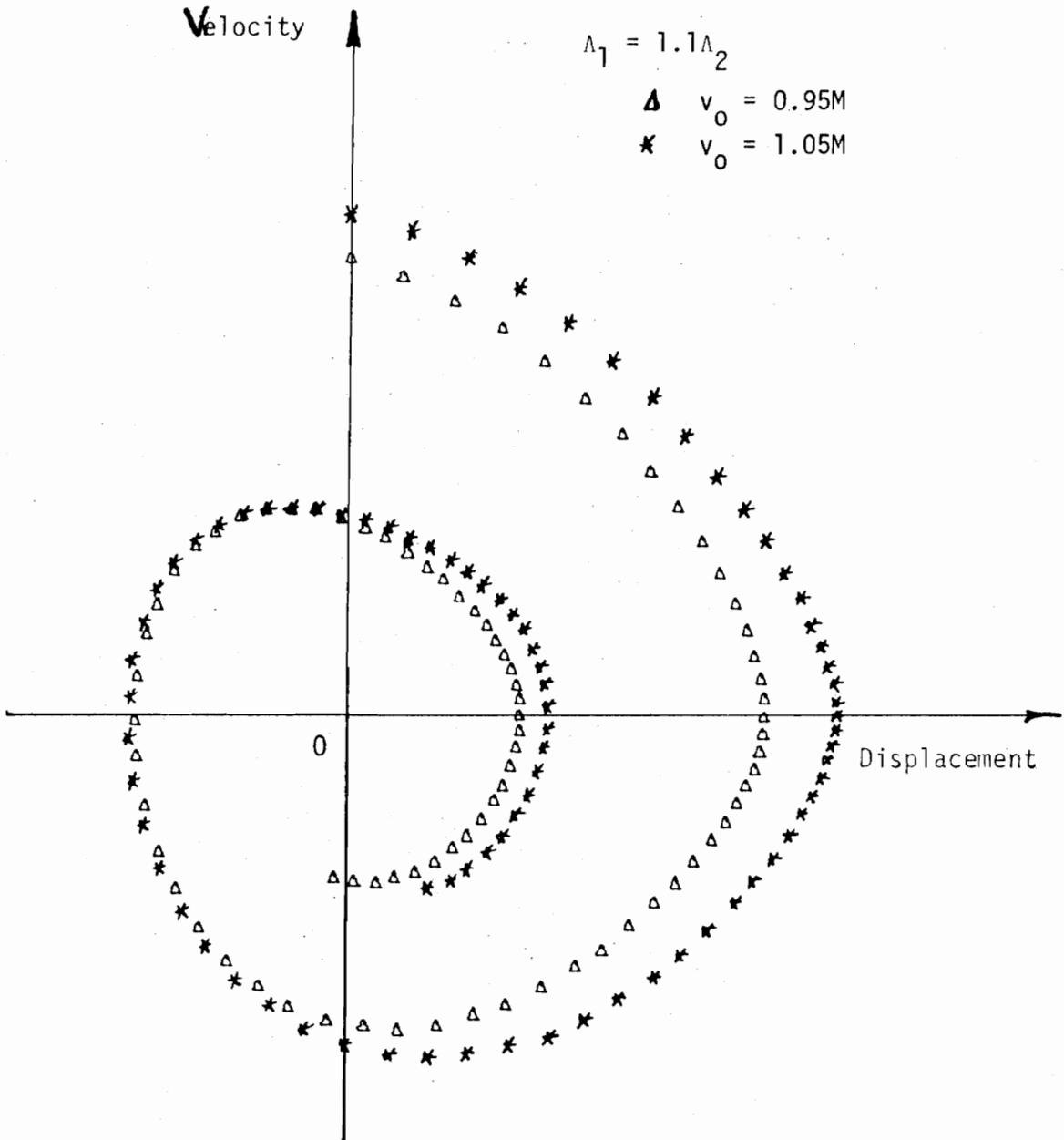


FIG. 5.26 DAMPED RESPONSE OF JOINT 1
(Instantaneous Disturbance at Joint 1)

discussed above, the stability criterion is necessary and sufficient for stability of equilibrium if the following condition is satisfied: the disturbance consisting of a single velocity component is applied in the direction of the generalized displacement whose increment was proposed as a "suitable" choice for the perturbation parameter, ϵ , in the static analysis. In all cases, the disturbance was applied at a loaded joint. Hence, for each static loading condition, there exists a dynamic disturbance (satisfying condition stated above) with a specific distribution for which the sufficient condition of dynamic stability is also a necessary condition. The sufficient stability condition of section 2.2 is, therefore, practical for the design of reticulated domes.

b. Stability of Equilibrium Relative to Transient Disturbances

Under two independent loads Λ_1, Λ_2 with $\Lambda_1 > \Lambda_2$, Fig. 5.27 shows the response of joint 1 to a transient disturbance applied at joint 1. The equilibrium state corresponds to $0.75 \Lambda_i^C$; where Λ_i^C are the critical values for Λ_1 and Λ_2 . A single exciting force with constant amplitude and finite duration (step function) is applied at joint 1. The degree of the disturbance is controlled via the duration of the exciting force. At point 1 along the perturbed motion, the total energy $v_1 = 0.94M$. According to the stability criterion, equilibrium is stable relative to the prescribed disturbance. This is illustrated by the resulting motion (Fig. 5.27), which initiates at the origin, enters a closed trajectory at the end of the excitation, and remains on this

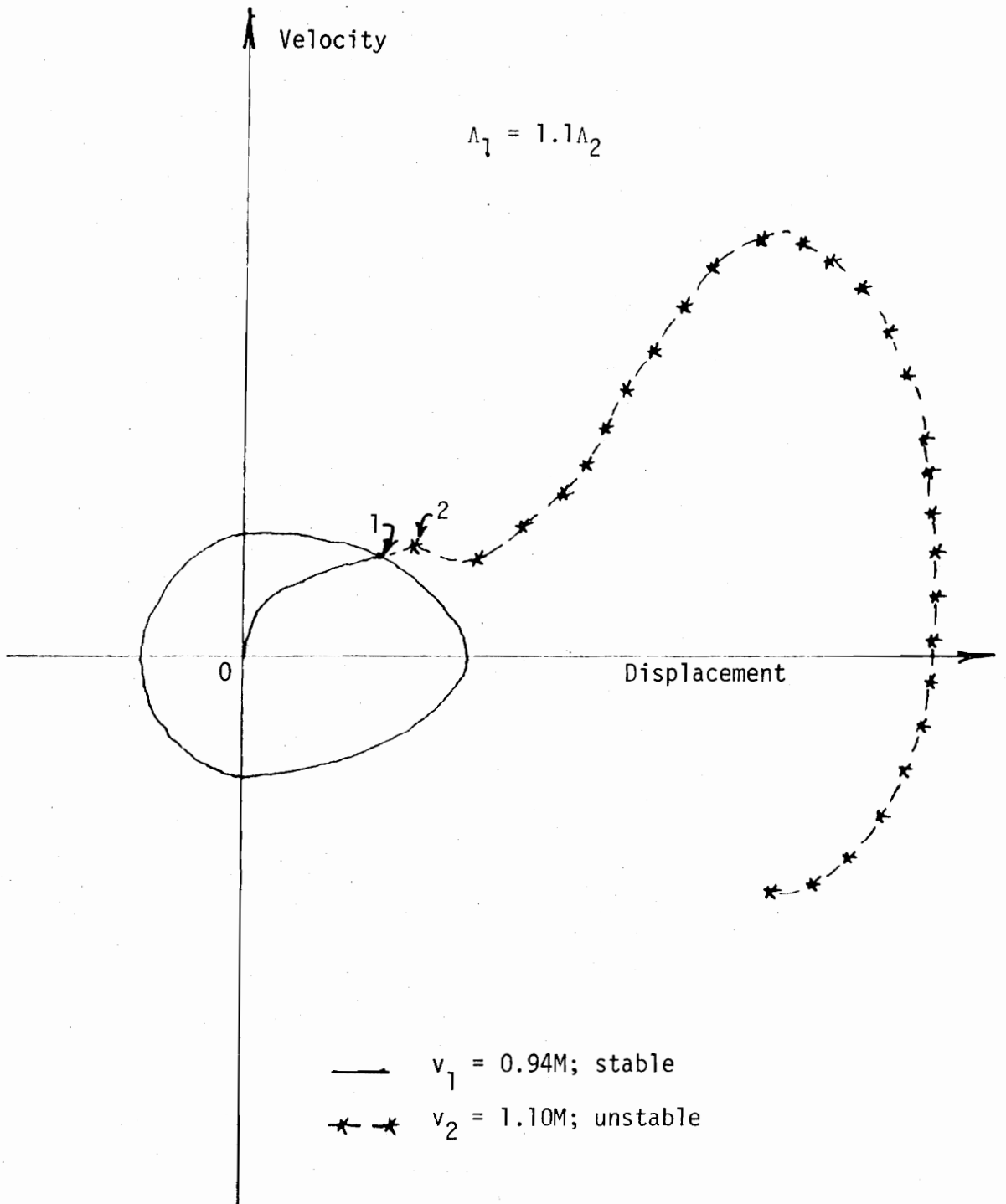


FIG. 5.27 RESPONSE OF JOINT 1
(Transient Disturbance at Join 1)

trajectory. The second excitation causes the motion to escape the domain of attraction of the origin, 0. The maximum energy along the perturbed motion v_2 (point 2, Fig. 5.27) is such that $v_2 = 1.1M$. In this case, the resulting motion shows instability of the equilibrium state 0.

Fig. 5.28 shows the effect of damping on the responses discussed in relation to Fig. 5.27. At point 1, $v_1 = 0.95M$ and the origin is asymptotically stable. At point 2, $v_2 = 1.1M$ and the perturbed motion converges to the buckled state showing instability of the origin.

5.2 Geodesic Dome: 39 DOF Model

The thirty-nine DOF geodesic dome model is shown in Fig. 5.2. The model is analyzed for different loading conditions:

1. $\Lambda_1 \neq 0, \Lambda_2 = \Lambda_3 = 0$
2. $\Lambda_1 = \Lambda_2 > \Lambda_3$
3. $\Lambda_1 = 0, \Lambda_2 \neq 0, \Lambda_3 = 0$
4. $\Lambda_1 = \Lambda_2, \Lambda_3 = 0$
5. $\Lambda_1 > \Lambda_2, \Lambda_3 = 0$
6. $\Lambda_1 < \Lambda_2, \Lambda_3 = 0$

The 39 DOF geodesic dome model exhibits features that are qualitatively similar to those of the 21 DOF model under comparable loading conditions.

For illustrative purposes, two loading conditions are shown in Figs. 5.29a and 3.29b. Both loading cases display limit points as

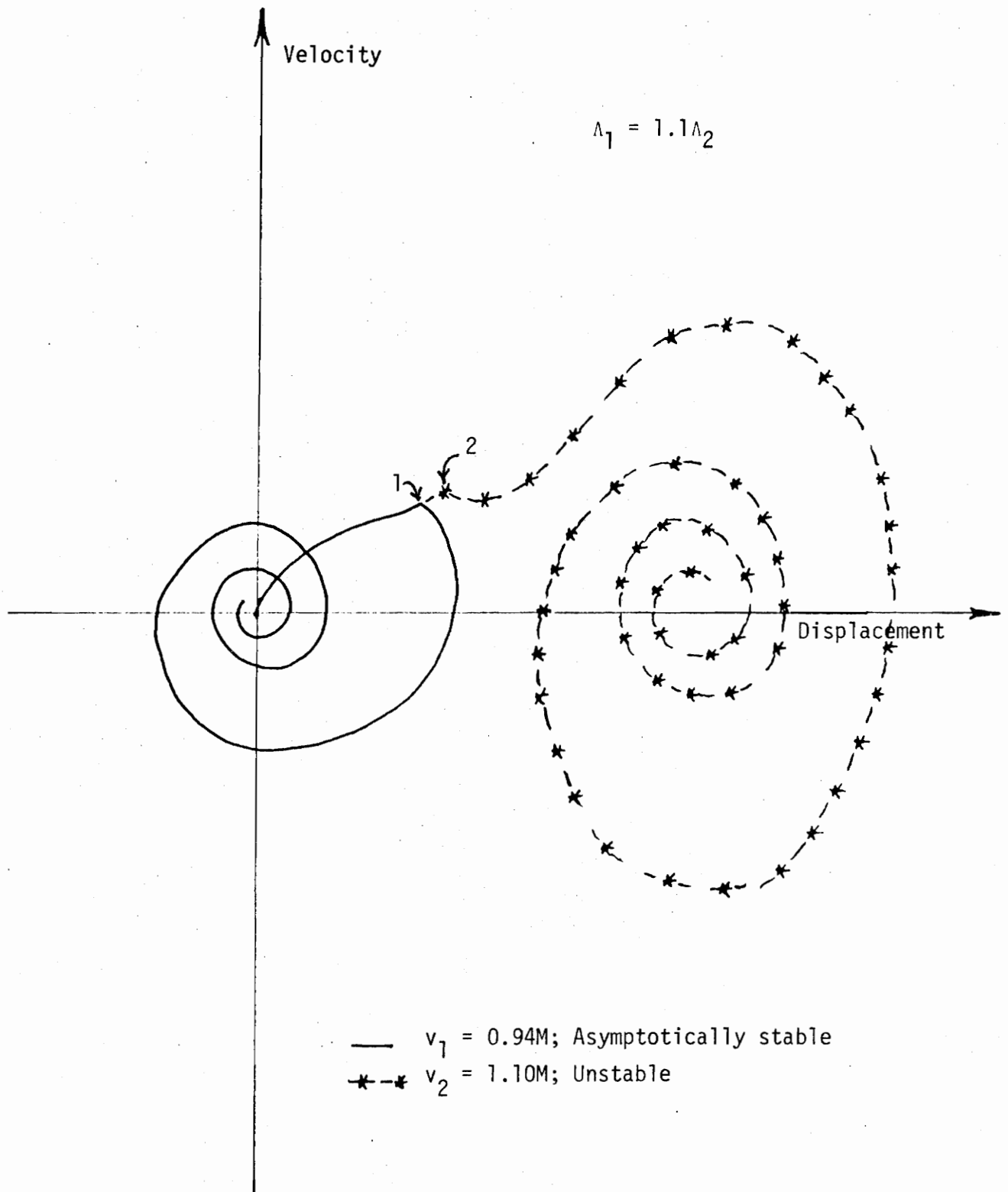


FIG. 5.28 DAMPED RESPONSE OF JOINT 1
(Transient Disturbance at Joint 1)

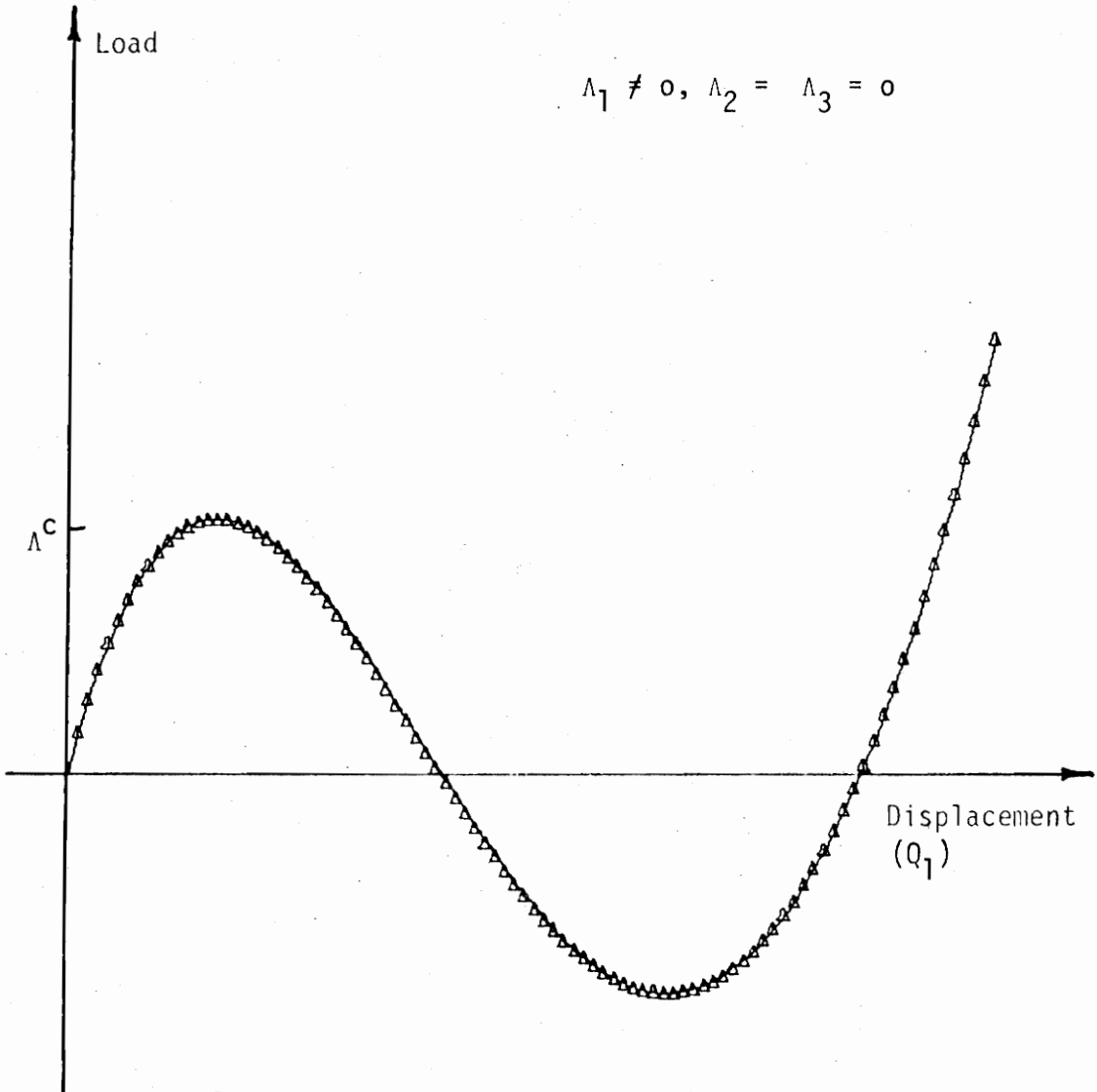


FIG. 5.29a STATIC RESPONSE OF 39 DOF GEODESIC DOME MODEL
(Only Joint 1 Loaded)

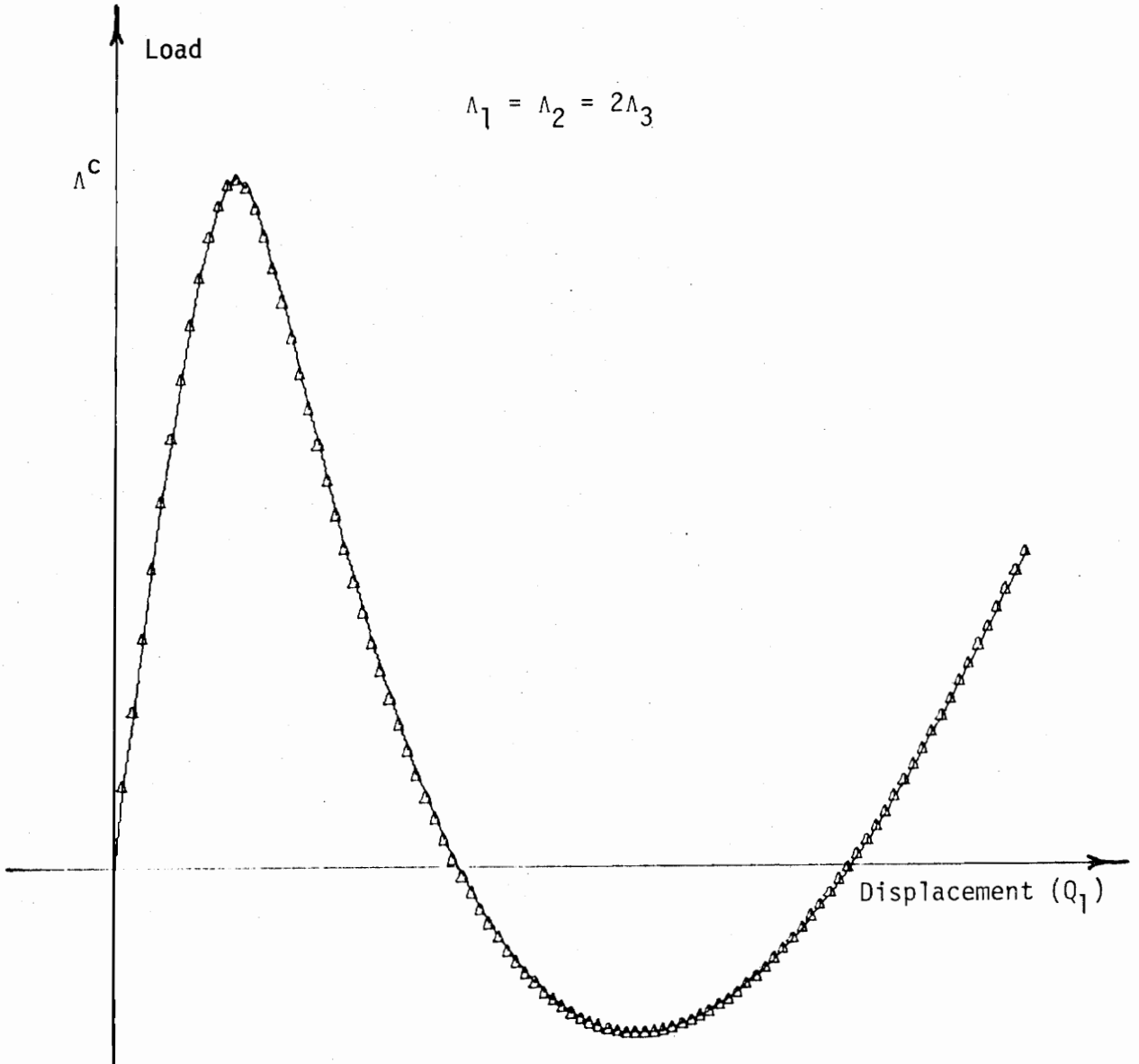


FIG. 5.29b STATIC RESPONSE OF 39 DOF GEODESIC DOME MODEL
(All Joints Loaded)

critical points. Under a single load Λ_1 , the dome fails at a much lower value of Λ^C than for the case when all joints are loaded. This is due to the effect of load concentration at joint 1 (Fig. 5.2). This concentration causes local instability of joint 1. However, since the dome cannot resist the shear loads resulting from local instability, general instability cannot be prevented.

The degrees of stability for the loading cases of Figs. 5.29a and 5.29b are illustrated in Fig. 5.30. The dome with all joints loaded is more sensitive to small (but finite) disturbances in the vicinity of the limit load than the dome with the single load. These examples indicate that all static loads that are likely to act on the dome during its lifetime should be included in the model to determine its degree of stability.

With $\Lambda_1 = \Lambda_2 = \Lambda$ and $\Lambda_3 = 0$, the response of the dome to an instantaneous disturbance at joint 1 is shown in Fig. 5.31. As with the 21 degree-of-freedom model, the equilibrium state (origin of the phase-plane) is stable or unstable depending on the initial energy. The effect of damping is shown in Fig. 5.32.

When $\Lambda_1 = 2\Lambda_2$, $\Lambda_3 = 0$ and joint 2 is instantaneously disturbed, the responses of the ring joints (e.g. joint 2) are shown in Fig. 5.33. When $v_0 = 0.96M$, the undamped system is stable while the damped system is asymptotically stable. When $v_0 = 1.3M$, the undamped system is still stable. This inference is made from the perturbed motion, a closed trajectory surrounding only the origin. The sufficient dynamic stability condition is not necessary for stability of equilibrium (the

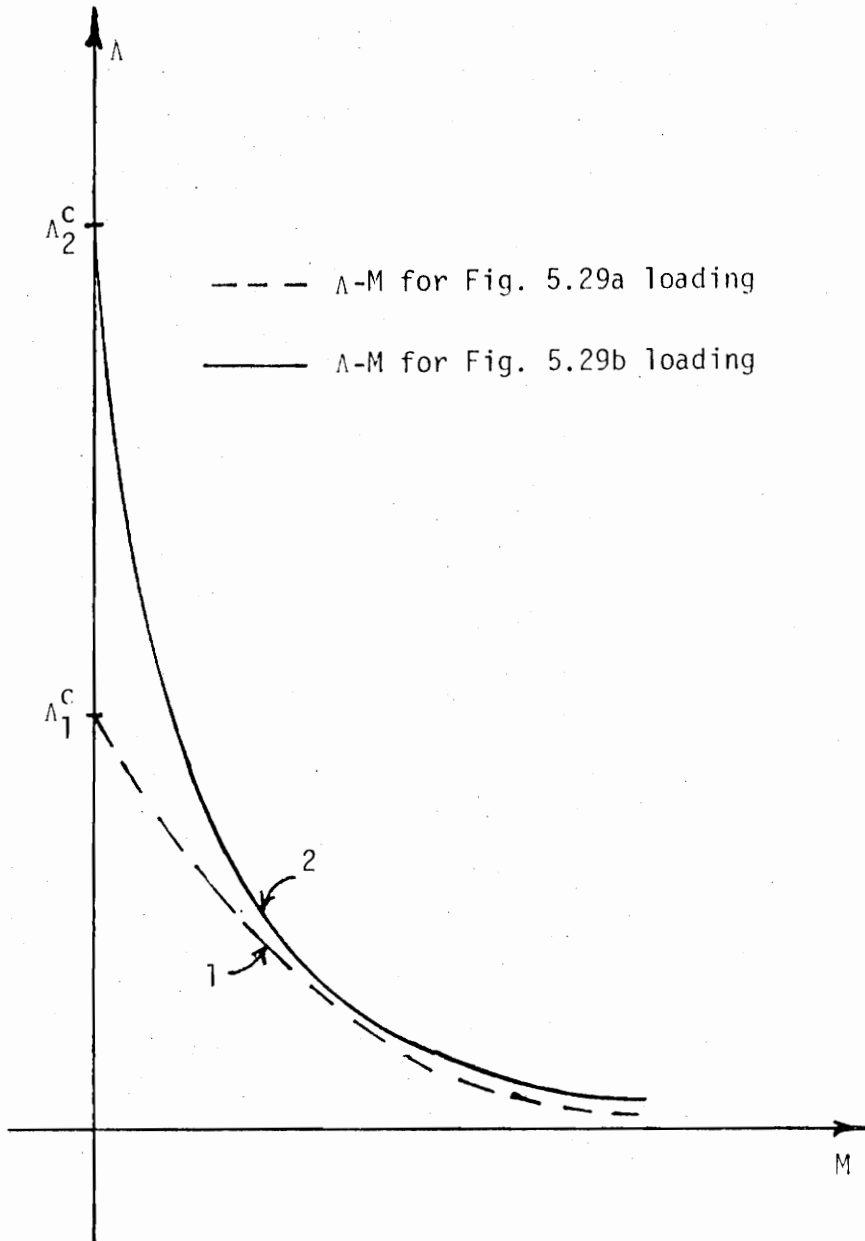


FIG. 5.30 DEGREE OF STABILITY
(Geodesic Dome: 39 DOF Model)

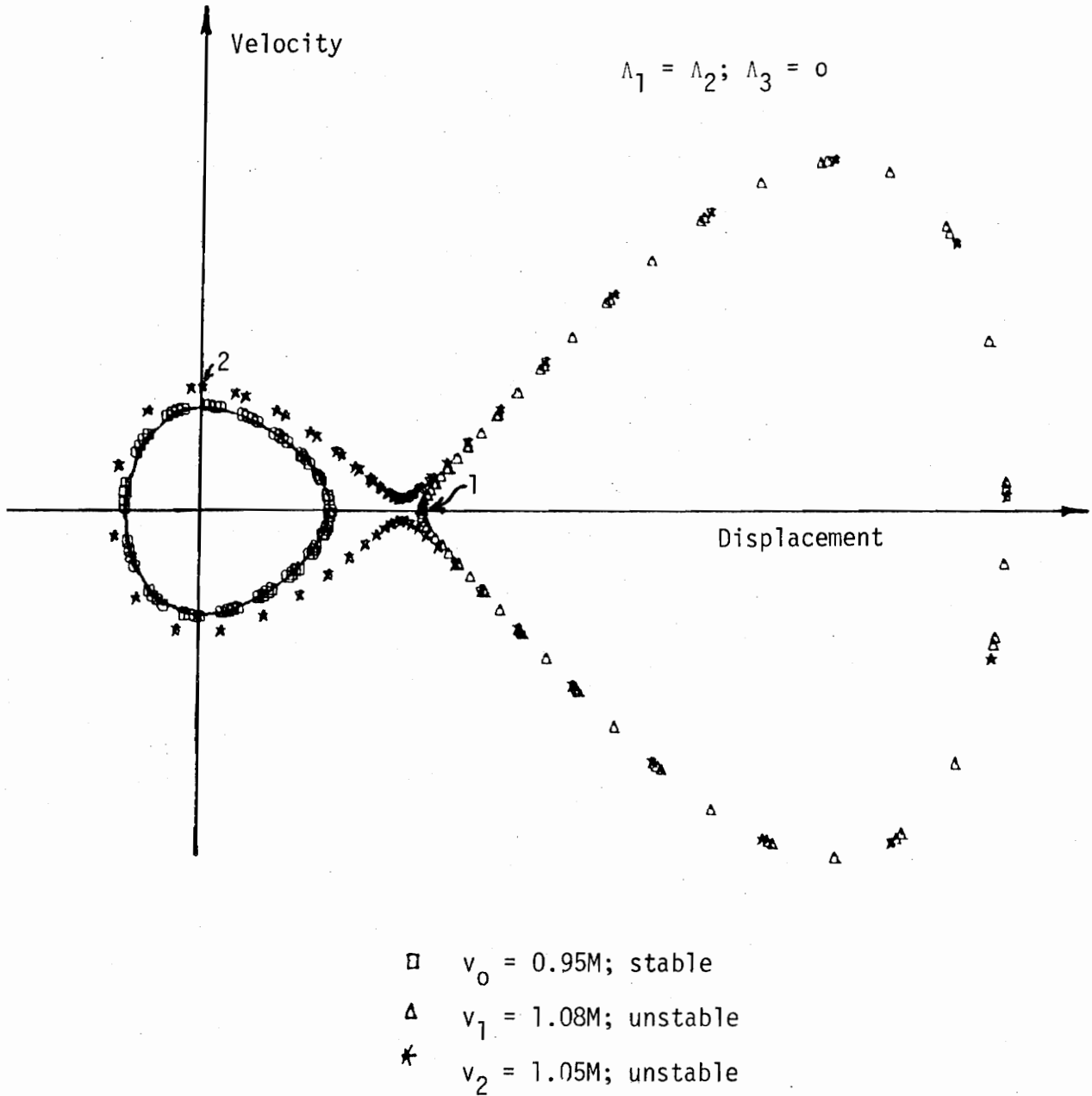


FIG. 5.31 RESPONSE OF JOINT 1
(Instantaneous Disturbance at Joint 1)

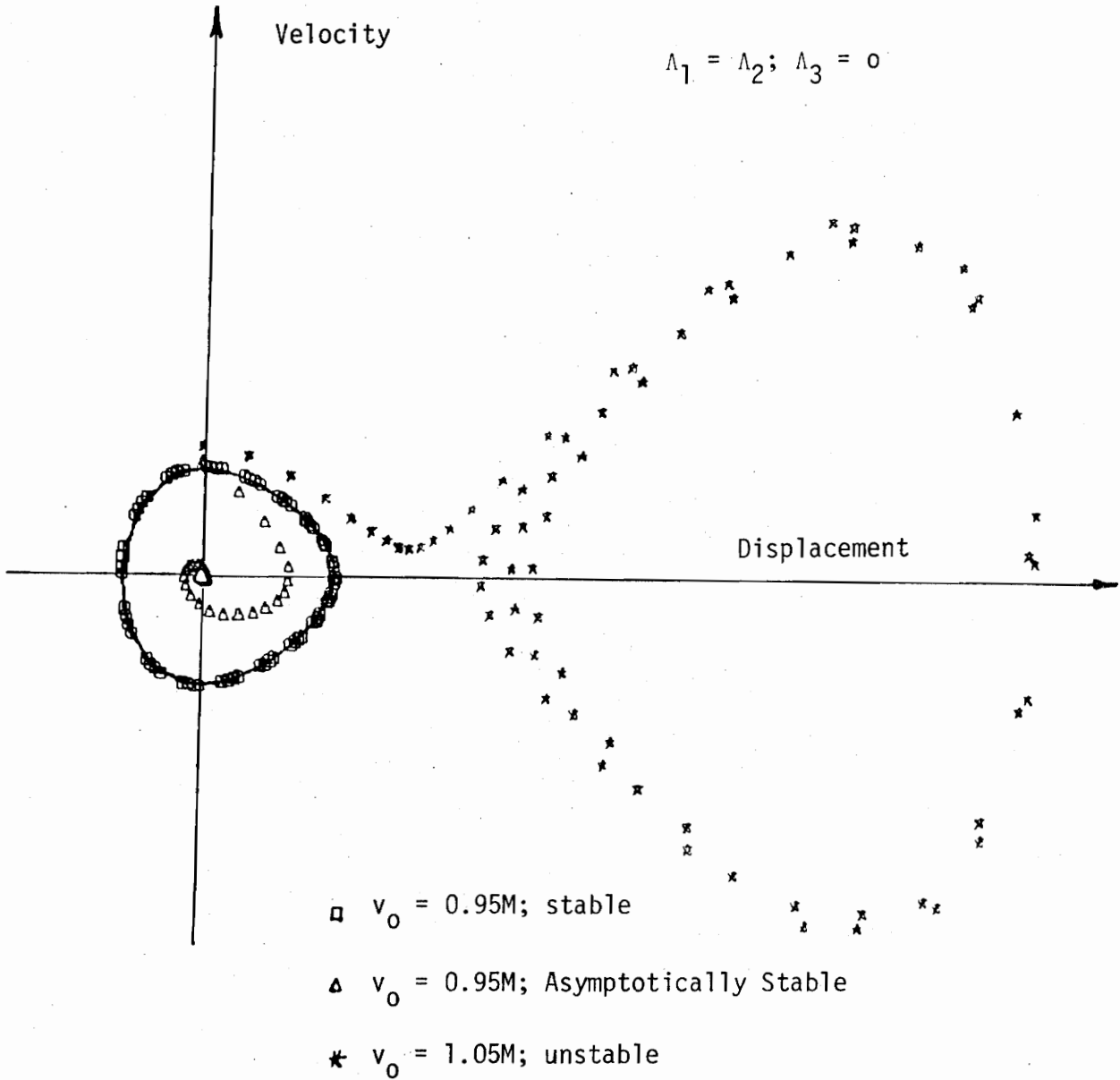


FIG. 5.32 UNDAMPED AND DAMPED RESPONSES OF JOINT 1
(Instantaneous Disturbance at Joint 1)

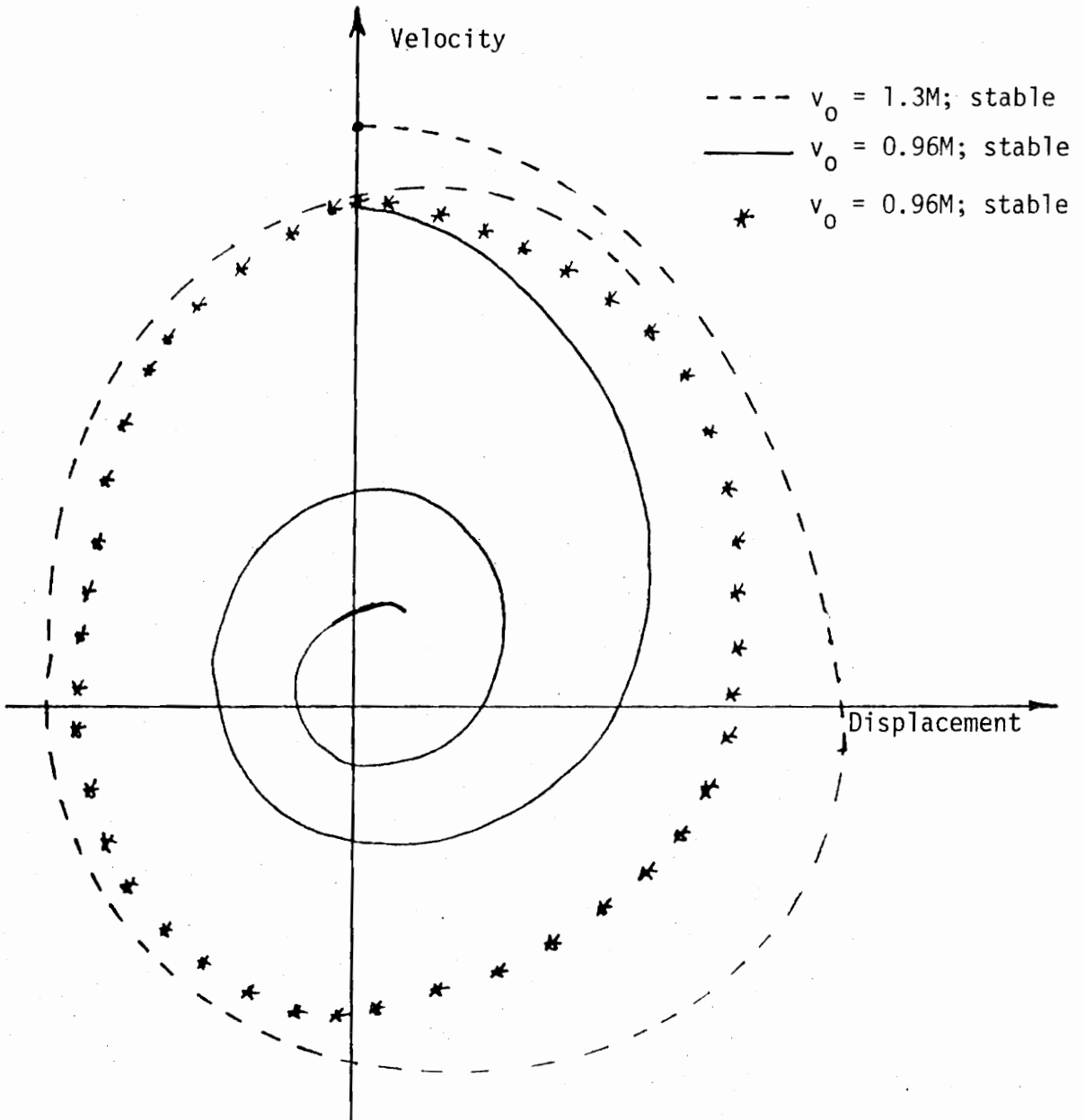


FIG. 5.33 RESPONSE OF RING JOINTS
(Instantaneous Disturbance at Joint 2)

origin in the phase plane). This is in agreement with observations on the 21 DOF model since the disturbance is not applied in the direction corresponding to that of the displacement-increment selected as the "suitable" perturbation parameter, ϵ .

The response of joint 1 to a transient disturbance at joint 1 is shown in Fig. 5.34. The equilibrium configuration corresponds to a single concentrated load at joint 1. At point 1, the total energy is less than M and the origin is stable. At point 2, the total energy is greater than M and the perturbed motion escapes the domain of attraction of the origin. Later convergence is to the buckled state to which the system is attracted. Hence stability can be decided by integrating the nonlinear equations of perturbed motion for the duration of the disturbance and checking if the total energy at the end of the excitation is less than M .

5.3 Convergence Characteristics

In all figures discussed above, the static analysis was carried out to third order, ϵ^3 , approximation. Also, a proper parameter, ϵ , was selected for locating equilibrium paths as described in Chapter 4. The convergence characteristics of the perturbation technique are illustrated in Figs. 5.35-5.37.

Fig. 5.35 shows the effect of choosing a poor perturbation parameter. Λ_1^C, Λ_2^C are fictitious critical loads which are results of numerical instability of the perturbation scheme. Point 1 corresponds to choosing q_2 as ϵ ; q_2 is the generalized displacement increment at

△ $v_2 = 1.1M$; unstable

★ $v_1 = 0.94M$; stable

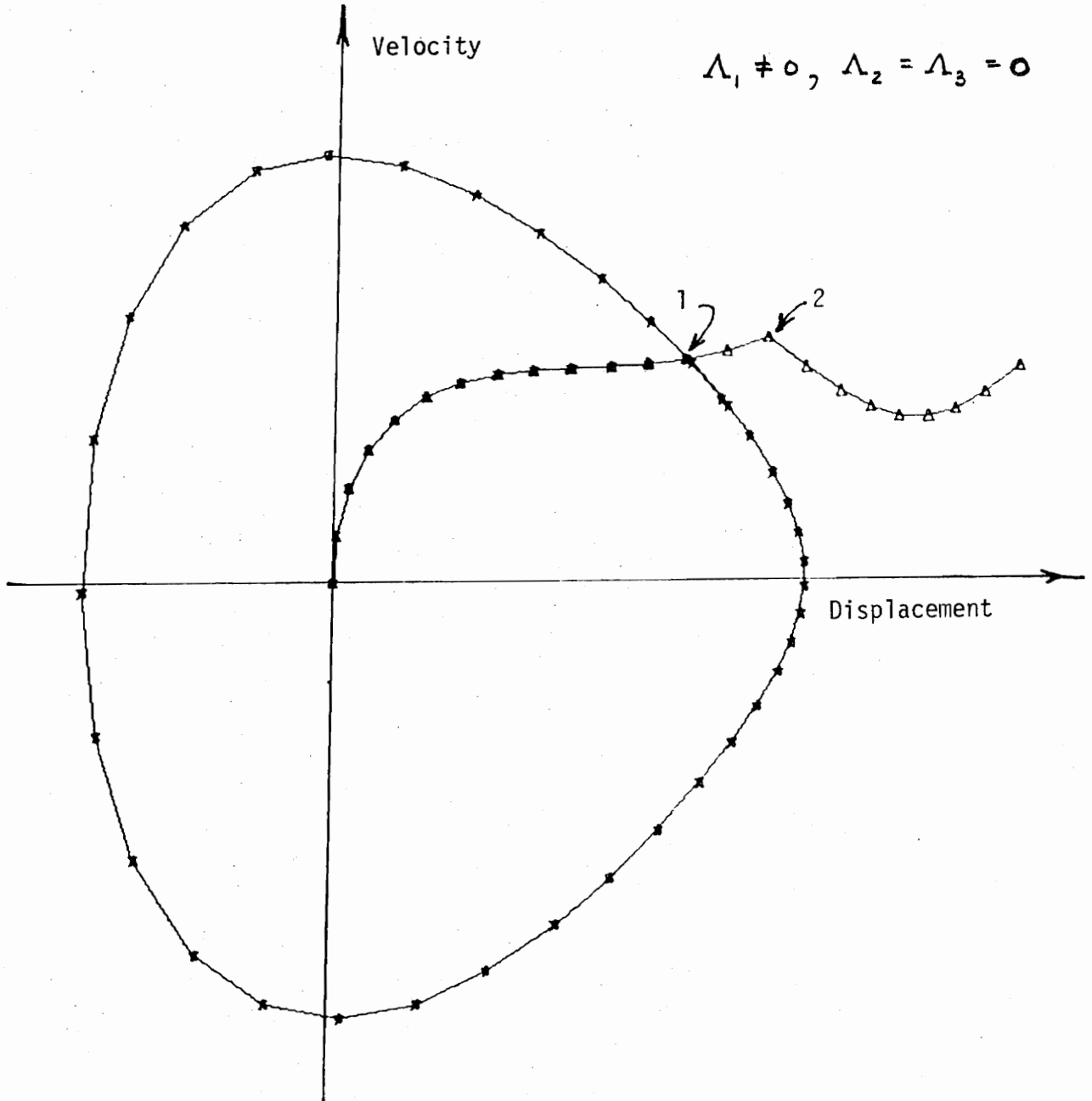


FIG. 5.34 RESPONSE OF JOINT 1 TO TRANSIENT DISTURBANCE

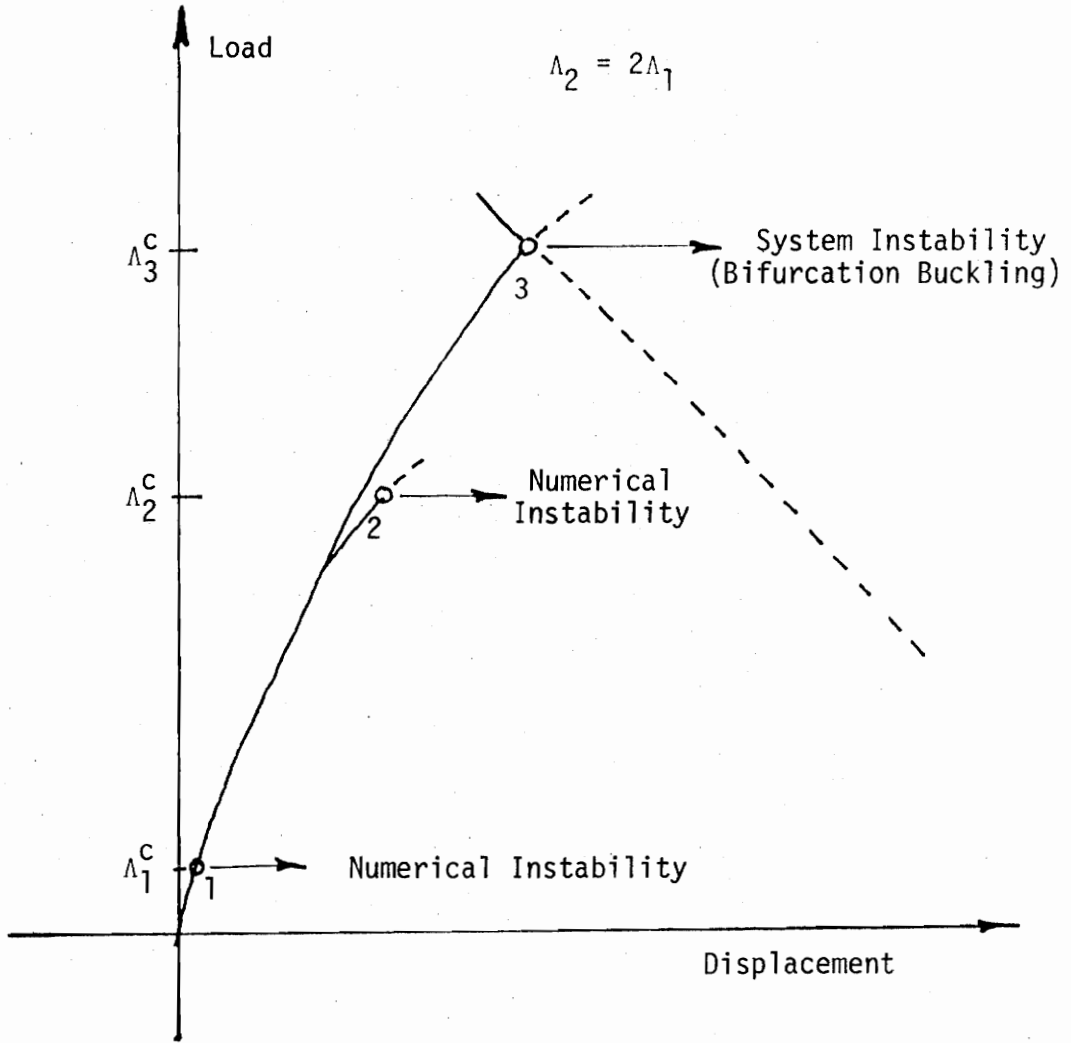


FIG. 5.35 CHOICE OF A PERTURBATION PARAMETER
(Geodesic Dome: 21 DOF Model)

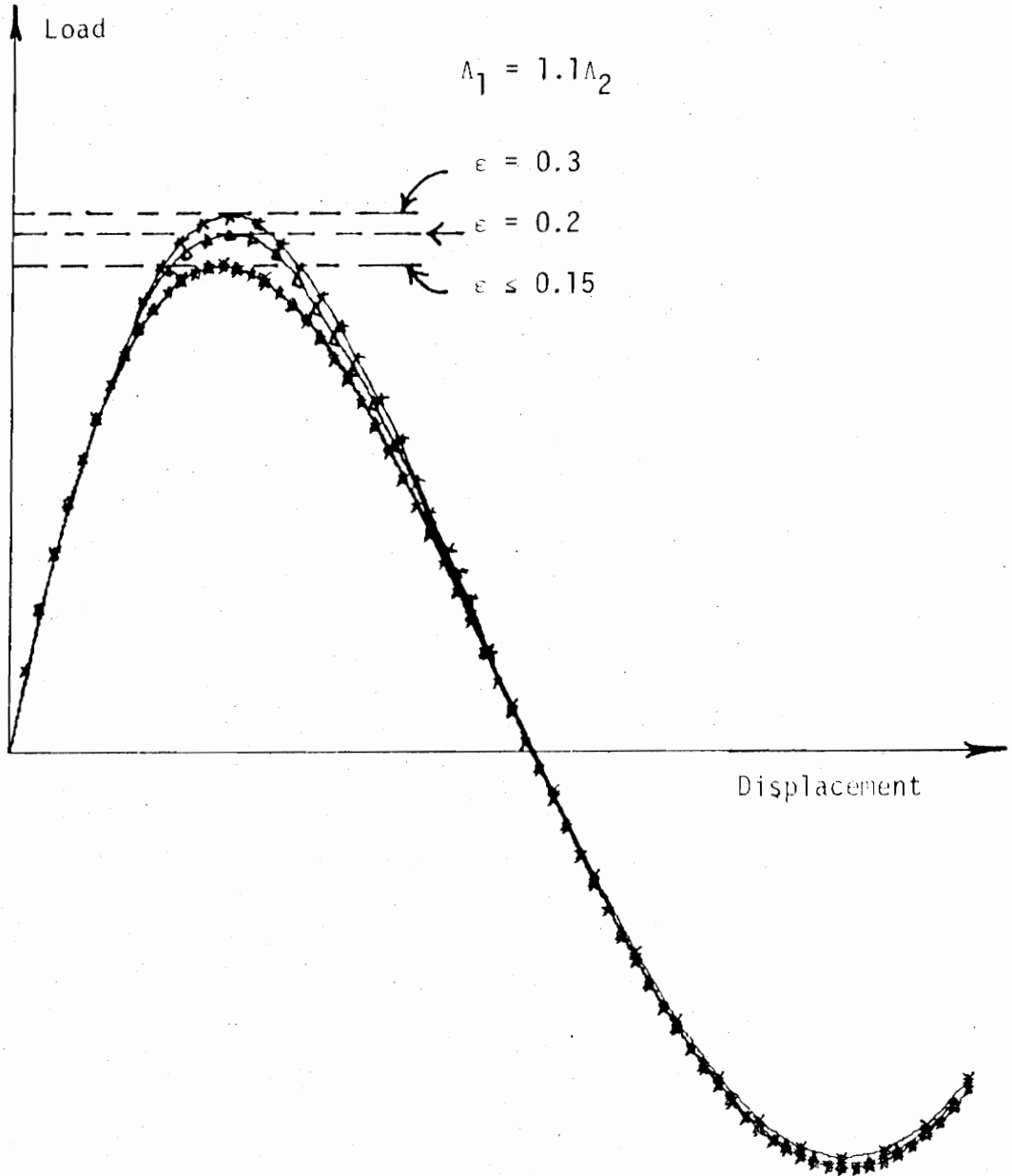


FIG. 5.36 BOUND ON PERTURBATION PARAMETER
(Geodesic Dome: 21 DOF Model)

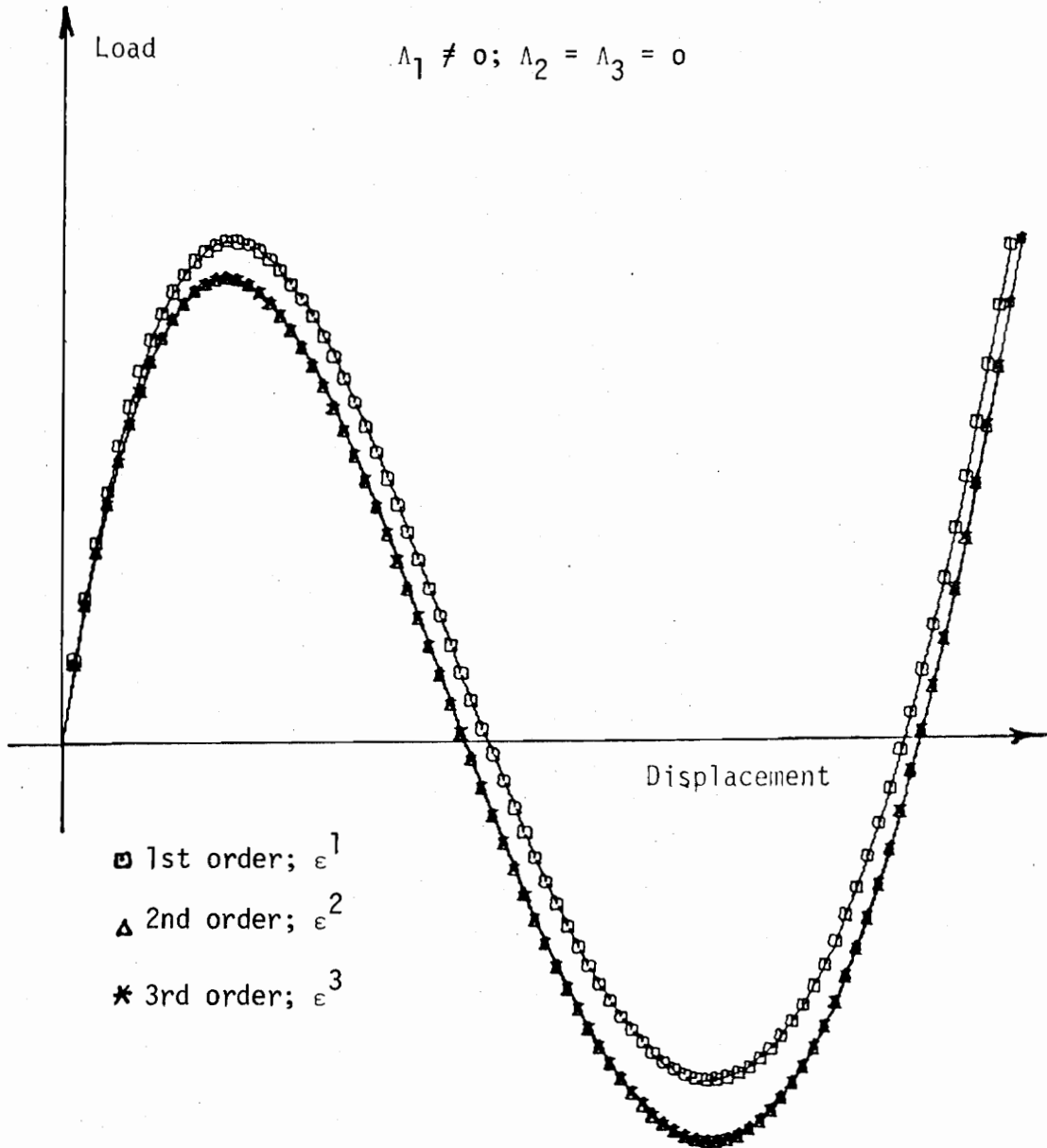


FIG. 5.37 UNIFORM CONVERGENCE OF PERTURBATION SCHEME
(Geodesic Dome: 39 DOF Model)

joint 1 in the Y-direction. Point 2 corresponds to choosing q_1 as ϵ where q_1 is the displacement increment at joint 1 in the X-direction. The correct choice for ϵ when $\Lambda_2 > \Lambda_1$ is q_4 , the generalized displacement increment at joint 2 in the X-direction. This is the proposed choice for ϵ (section 4.3). In this case, the true critical load Λ_3^C is obtained and the post-buckling path beyond the bifurcation point, 3, is obtained without any numerical instability problems.

Fig. 5.36 depicts the characteristics of the nonlinear response to static loads as functions of the value of ϵ . For an error limit of 10^{-7} and a 0.5 percent allowance for \ddot{q}_i , the value of ϵ according to Eq. 4.12 is equal to or lower than 0.15. Convergence of third order solutions for $\epsilon > 0.15$ show a nonconservative estimate of the critical load Λ . For $\epsilon \leq 0.15$, a unique critical load is obtained.

Fig. 5.37 shows the uniform convergence of the perturbation solutions for a properly chosen parameter, ϵ , and a limiting value (based on Eq. 4.12) placed on ϵ . As in Fig. 5.36, convergence is from above; i.e. the first order perturbation solution provides a nonconservative estimate for the critical load. The solutions for the second and third order perturbation approximations are practically identical (Fig. 5.37).

5.4 Stability Boundaries

The stability boundaries for the two geodesic dome models are depicted in Figs. 5.38-5.40. In both cases, Λ_{01} and Λ_{02} correspond to the critical loads of Λ_1 , Λ_2 , respectively when acting singly. Hence the locations of the critical states on the axes involve no load

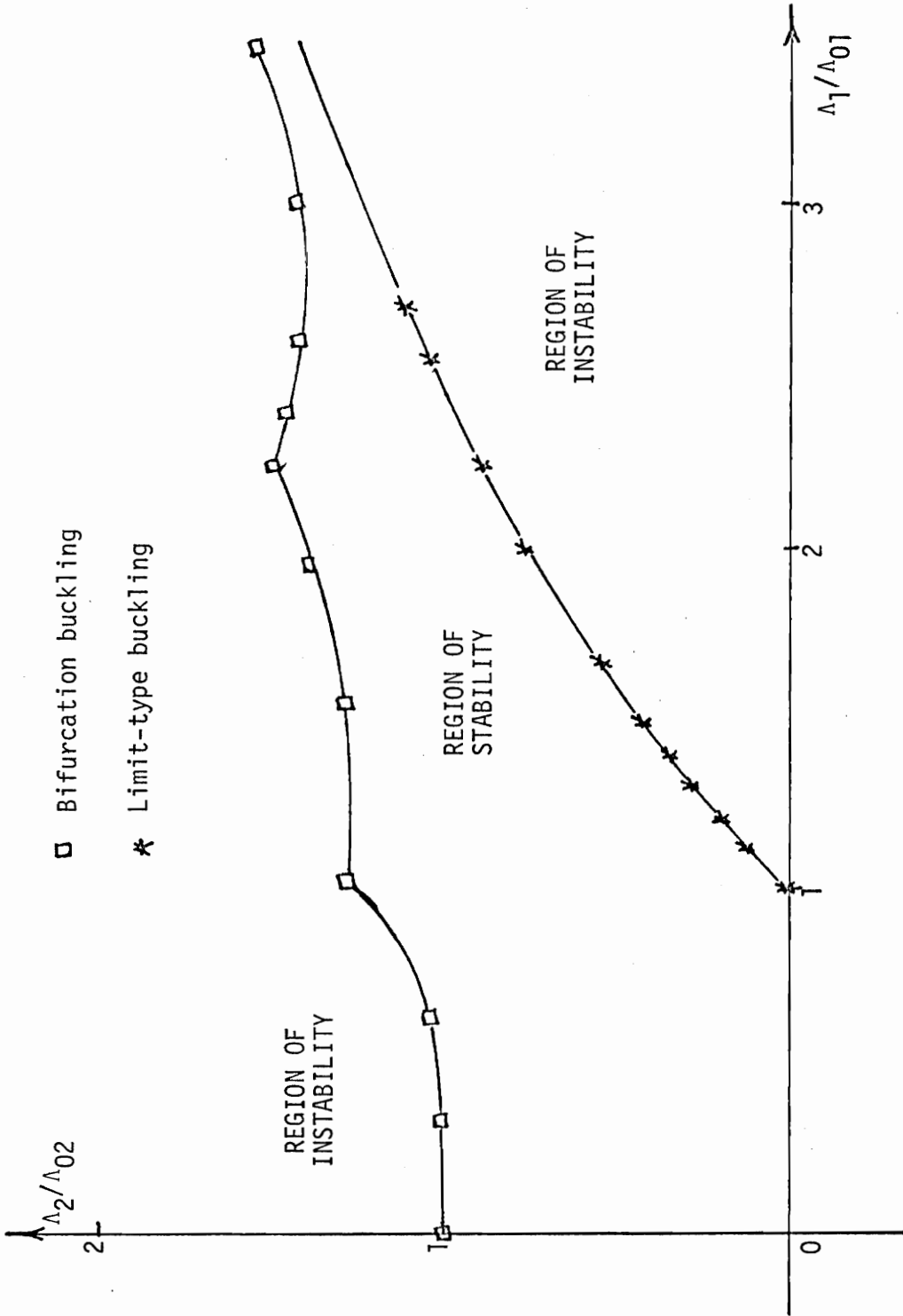


FIG. 5.38 STABILITY BOUNDARIES: 21 DOF MODEL

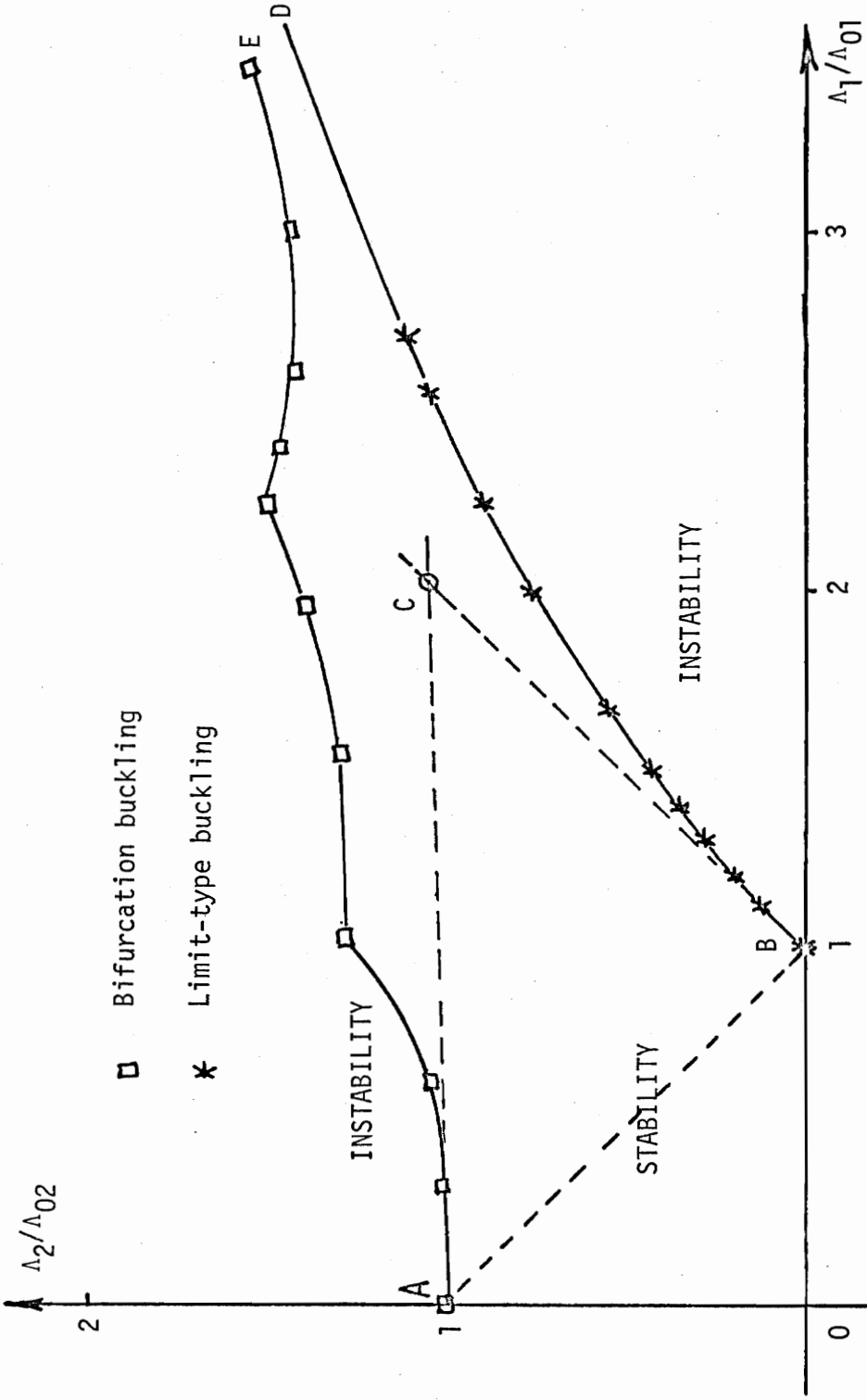


FIG. 5.39 ESTIMATES OF BOUNDS ON STABILITY BOUNDARIES
(Geodesic Dome: 21 DOF Model)

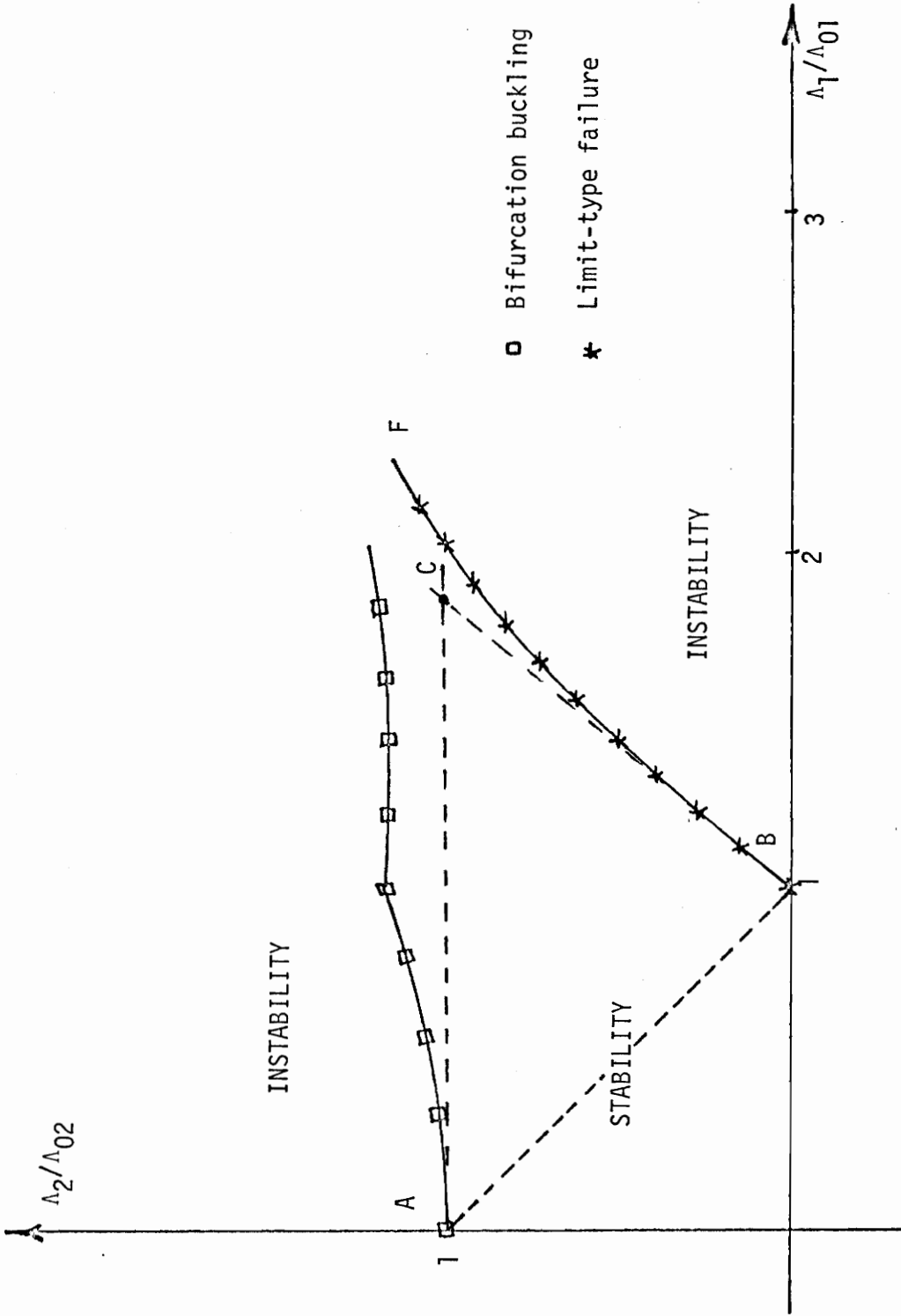


FIG. 5.40 STABILITY BOUNDARIES
(Geodesic Dome: 39 DOF Model)

interaction. The stability boundary BD (Fig. 5.39) for all points corresponding to a limit-type buckling of the 21 DOF model under independent loads Λ_1, Λ_2 is convex towards the region of stability. Bifurcation critical states lie on a stability boundary which is continuous but piecewise differentiable. This boundary (curve AE in Fig. 5.39) appears to be piecewise convex towards the region of stability. All points lying on a particular piece of AE exhibit identical buckling modes at failure.

The line AB represents the least lower bound estimate to the region of stability. The region OACBO represents a safe estimate of the region of stability. AC, BC are safe greatest lower bound estimates of the region of stability obtained from tangents, to AE at A, and to BD at B.

The stability boundaries for the 39 DOF geodesic dome model are shown in Fig. 5.40. Boundary BF is convex toward the region of stability while boundary AF is piecewise convex toward the same region. The two curves approach the singular critical point F associated with an unstable symmetric point of bifurcation. In a way similar to Fig. 5.39, AB represents the least lower bound and AC, BC are safe greatest lower bound estimates on the region of stability. Also, the points lying on each segment of AF exhibit identical buckling modes at failure.

CHAPTER 6

CONCLUSIONS

The stability of reticulated domes under multiple static and dynamic loads is investigated. Two elastic geometrically nonlinear structural models of a reticulated dome are considered .

The nonlinear response of the system to static loads is obtained using nonlinear programming and discrete perturbation techniques. The nonlinear programming technique is used to obtain a starting solution for the discrete perturbation technique and to optimize the choice of the perturbation parameter, ϵ . Convergence criteria and error estimates to limit errors in a perturbation scheme are developed.

The largest displacement-increment at the initial stage of loading from the zero-state in configuration space is proposed as a "suitable" choice for ϵ in the case of imperfection-sensitive systems. A poor choice of ϵ leads to numerical instabilities which appear as 'fictitious' system instabilities.

The investigation of stability of equilibrium of the system subjected to finite disturbances is based on the concept of "degree of stability" and the associated sufficient stability condition, which follow from Liapunov's direct method.

The perturbed motion of the system under a given set of perturbations is obtained by numerically integrating the nonlinear equations of motion . Trajectories in the phase-plane provide valuable checks on the stability condition. These trajectories confirm the sufficiency of the

dynamic stability condition. However, they also indicate that there is a dynamic disturbance with a specific spatial distribution for which the sufficient condition of dynamic stability employed is also a necessary condition. This is true for each static loading condition considered. This disturbance corresponds to that which magnifies the static equilibrium configuration. The direction of application of the disturbance corresponds to that of the generalized displacement-increment selected as the "suitable" choice of ϵ .

The stability boundaries corresponding to two independent loads on the reticulated dome models are presented. Limit points lie on a boundary which is convex towards the region of stability. Bifurcation points lie on a continuous but piecewise differentiable boundary. Each piece of the boundary containing bifurcation points appears to be convex towards the region of stability. Critical points lying on separate pieces of the stability boundary exhibit identical buckling modes. Knowledge of the general properties of such boundaries can form the basis of a rigorous mathematical study to establish the shape of the stability boundaries for reticulated domes exhibiting limit and bifurcation points. This is recommended for future study.

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APPENDIX
COMPUTER PROGRAM

COMPUTER PROGRAM

The digital computer program performs three major tasks in the complete stability analysis of reticulated domes under multiple static and dynamic loads. First, nonlinear programming is used to obtain points on the equilibrium path of the system at the onset of loading, and to optimize the choice of a "suitable" perturbation parameter. Second, a static perturbation scheme is initiated using the chosen perturbation parameter and the known solutions from nonlinear programming technique. The static perturbation scheme is used to obtain the remaining portions of the fundamental path and the entire post-buckling response path. These first two tasks complete the static analysis. Third, the stability of equilibrium (0.70-0.75 of critical static loads) relative to instantaneous and transient disturbances is investigated via the concept of "degree of stability" and the associated sufficient stability condition. A dynamic analysis is performed to check the stability criterion. Hammings modified Predictor-Corrector numerical integration scheme is used to obtain the perturbed motions of the system under a prescribed set of perturbations.

Both the static and dynamic portions of the program generate print and plot outputs. The program is coded in the FORTRAN language and has been used with the IBM 360-75. Only ELASTIC structures are considered. Independent loads and finite displacements are admitted at the joints. All free joints have hinged connections.

The main sequence of operations of the computer program are

described within the program using COMMENT cards. Details of coding are omitted . The fortran listing of the basic computer program for two independent joint loads is included for completeness.

PROGRAM LISTING

```

C===== S T A T I C   A N A L Y S I S   =====
C
C   OBJECTIVE : TO OBTAIN THE NONLINEAR RESPONSE OF RETICULATED
C             DOMES TO STATIC LOADS
C
C   METHODS : NONLINEAR PROGRAMMING AND
C             STATIC PERTURBATION TECHNIQUES
C
C             IMPLICIT REAL*8(A-H,O-Z)
C             REAL*4 P1,P2
C             EXTERNAL V
C             REAL FS,ACC
C             DIMENSION P1(150),P2(150),LABELX(1),LABELY(1)
C             DIMENSION U(39),G(39)
C             DIMENSION H(294),X(19),Y(19),Z(19)
C             DIMENSION ICON(19),LN(39),MW(39)
C             COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
C             COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
C             COMMON SUMX(42),SUMY(42),SUMZ(42)
C             COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
C             COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
C
C             RETICULATED DOMES UNDER TWO INDEPENDENT LOADS
C
C             KK=1
C             WRITE(6,10)
C             10  FORMAT('1',27X,'*****')
C             $28X,'*',44X,'*'
C             $28X,'*' A SPACE TRUSS WITH HINGED CONNECTED JOINTS *'
C             $28X,'*' SUBJECTED TO LOADS AT THE JOINTS ONLY *'

```

```

AYD 0010
AYD 0020
AYD 0030
AYD 0040
AYD 0050
AYD 0060
AYD 0070
AYD 0080
AYD 0090
AYD 0100
AYD 0110
AYD 0120
AYD 0130
AYD 0140
AYD 0150
AYD 0160
AYD 0170
AYD 0180
AYD 0190
AYD 0200
AYD 0210
AYD 0220
AYD 0230
AYD 0240
AYD 0250
AYD 0260
AYD 0270
AYD 0280
AYD 0290
AYD 0300
AYD 0310
AYD 0320

```

```

$28X,1* INDEPENDENT JOINT LOADS ARE CONSIDERED **//
$28X,1*,44X,1*1/
$28X,1*****//)
C
C INPUT DATA FOR MEMBER NUMBERS, NM AND NUMBER OF JOINTS, NJ
C
  READ(5,20) NM,NJ
  FORMAT(2I5)
  WRITE(6,30) NM,NJ
  30 FORMAT(//10X,'NUMBER OF ELEMENTS, NM =',I3//10X,'NUMBER OF JOINTS,
    $ NJ =',I3)
C
C JOINT COORDINATES AND CONSTRAINT OF THE JOINT
C ICON(J): MEANS CONSTRAINT AT JOINT 'J'
C ITS VALUES ARE: ICON(J)=1 -COMPLETE CONSTRAINED OR FIXED
C ICON(J)=0 -COMPLETELY FREE
C
C NF= NUMBER OF UNCONSTRAINED JOINTS
C
NF=NJ
WRITE(6,40)
40 FORMAT(//20X,'JOINT COORDINATES AND CONSTRAINTS,////11X,'JOINT'
  $ NO',5X,'X-COORD',10X,'Y-COORD',10X,'Z-COORD',8X,'CONSTRAINT'//)
  DO 70 J=1,NJ
  READ(5,50) K,X(K),Y(K),Z(K),ICON(K)
  50 FORMAT(I5,3F15.6,4X,I1)
  WRITE(6,60) K,X(K),Y(K),Z(K),ICON(K)
  60 FORMAT(10X,I5,6X,D12.5,6X,D12.5,5X,D12.5,10X,I1)
  NF=NF-ICON(K)
  70 CONTINUE
C

```

AYG 0330

AYG 0340

AYG 0350

AYG 0360

AYG 0370

AYG 0380

AYG 0390

AYG 0400

AYG 0410

AYG 0420

AYG 0430

AYG 0440

AYG 0450

AYG 0460

AYG 0470

AYG 0480

AYG 0490

AYG 0500

AYG 0510

AYG 0520

AYG 0530

AYG 0540

AYG 0550

AYG 0560

AYG 0570

AYG 0580

AYG 0590

AYG 0600

AYG 0610

AYG 0620

AYG 0630

AYG 0640


```

C MEMBER INCIDENCE TABLE AYO 0650
C WRITE(6,80) AYO 0660
80 FORMAT('1',///10X, 'MEMBER INCIDENCE TABLE AND ITS LENGTH'///5X, AYO 0670
$MEMBER NO',10X,'MEMBER INCIDENCE',13X,'MEMBER LENGTH'///) AYO 0680
DO 110 M=1,NM AYO 0690
READ(5,90) J,INP(J),INQ(J) AYO 0700
90 FORMAT(3I5) AYO 0710
I=INP(J) AYO 0720
K=INQ(J) AYO 0730
DX(M)=X(K)-X(I) AYO 0740
DY(M)=Y(K)-Y(I) AYO 0750
DZ(M)=Z(K)-Z(I) AYO 0760
XL(M)=DSQRT(DX(M)**2+DY(M)**2+DZ(M)**2) AYO 0770
WRITE(6,100) J,INP(J),INQ(J),XL(J) AYO 0780
100 FORMAT(4X,I6,15X,2I5,15X,D15.5) AYO 0790
110 CONTINUE AYO 0800
C AYO 0810
C AYO 0820
C AYO 0830
C AYO 0840
C AYO 0850
C AYO 0860
C AYO 0870
C AYO 0880
C AYO 0890
C AYO 0900
C AYO 0910
C AYO 0920
C AYO 0930
C AYO 0940
C AYO 0950
C AYO 0960

```

```

C      WRITE(6,130)
      130 FORMAT('1'///10X,' MINIMIZATION VIA CONJUGATE GRADIENT METHOD'////)
C
C      INITIAL GUESS OF VARIABLES
      140 WRITE(6,140)
      140 FORMAT('//20X,' INITIAL GUESS OF VARIABLES'///11X,' JOINT',8X,' X-DISPA'1040
      $,' 10X,' Y-DISP.',10X,' Z-DISP.'//)
      DO 160 I=1,NF
      I1=3*I-2
      I2=3*I-1
      I3=3*I
      U(I1)=0.000
      U(I2)=0.000
      U(I3)=0.000
      150 WRITE(6,150) I,U(I1),U(I2),U(I3)
      160 FORMAT(10X,15,9X,D12.5,6X,D12.5)
      160 CONTINUE
C
C      INITIALIZED ALL JOINT LOADS TO BE ZERO
C
C      DO 170 I=1,NF
      PX(I)=0.000
      PY(I)=0.000
      PZ(I)=0.000
      170 CONTINUE
C
C      THE ERROR INDICATOR, IER = 9 IS USED AS THE INITIAL INPUT
C      THE RETURN VALUE ARE SUCH THAT :
C      IER = 0.....MEANS CONVERGENCE IS ACHIEVED WITHIN SPECIFIED ACCURACY
C      IER = 1.....MEANS NO CONVERGENCE IN THE LIMIT ITERATION

```

AYO 0970

AYO 0980

AYO 0990

AYO 1000

AYO 1010

AYO 1020

AYO 1030

AYO 1040

AYO 1050

AYO 1060

AYO 1070

AYO 1080

AYO 1090

AYO 1100

AYO 1110

AYO 1120

AYO 1130

AYO 1140

AYO 1150

AYO 1160

AYO 1170

AYO 1180

AYO 1190

AYO 1200

AYO 1210

AYO 1220

AYO 1230

AYO 1240

AYO 1250

AYO 1260

AYO 1270

AYO 1280

```

C IER =-1.....MEANS ERRORS IN THE GRADIENT CALCULATION
C IER = 2..... MEANS IT IS LIKELY THAT THERE EXISTS NO MINIMUM
C
C IER=9
ACC=1.0-10
LIMIT=100
FS=0.0
KOUNT=0
C INPUT LOADING AT THE JOINT
C
C DO 181 I=1,NF
READ(5,180) K,PX(K)
180 FORMAT(I5,F10.6)
181 CONTINUE
C INPUT THE RATIO OF THE TWO INDEPENDENT LOADS
C (PROPORTIONAL LOADING)
READ(5,182) FLOAD
182 FORMAT(F10.5)
797 CONTINUE
CALL DEMCG(V,N,U,F,G,FS,ACC,LIMIT,IER,H)
WRITE(6,186) LIMIT,KOUNT
186 FORMAT(/////5X,'NUMBER OF ITERATION =',I4//
$ 5X,I5,5X,'FUNCTION SUBROUTINE CALLED'/////))
IW=1
CALL OUTP(U,G,F,N,0,0,IW)
C SEARCH FOR LARGEST DISPLACEMENT-INCREMENT
C
C CALL MAX(N,U,IND)
AYD 1290
AYD 1300
AYD 1310
AYD 1320
AYD 1330
AYD 1340
AYD 1350
AYD 1360
AYD 1370
AYD 1380
AYD 1390
AYD 1400
AYD 1410
AYD 1420
AYD 1430
AYD 1440
AYD 1450
AYD 1460
AYD 1470
AYD 1480
AYD 1490
AYD 1500
AYD 1510
AYD 1520
AYD 1530
AYD 1540
AYD 1550
AYD 1560
AYD 1570
AYD 1580
AYD 1590
AYD 1600

```

```

C      284 EPS=0.05
C      999 CONTINUE
C
C      C      INITIALIZE ALL INCREMENTAL SOLUTIONS TO ZERO
C
C      PARM1=0.000
C      PARM2=0.000
C      PARM3=0.000
C      DO 280 I=1,N
C      Q1(I)=0.000
C      Q2(I)=0.000
C      Q3(I)=0.000
C      Q4(I)=0.000
C      280 CONTINUE
C
C      C      USE LARGEST DISPLACEMENT-INCREMENT AS A
C      C      'SUITABLE' PERTURBATION PARAMETER
C
C      Q1(IND)=1.000
C
C      CALL FIRIN(U,Q1,PARM1,EPS,N,IND,IW)
C      CALL SECIN(U,PARM1,PARM2,EPS,N,IND,IW)
C      CALL THRIN(U,PARM1,PARM2,PARM3,EPS,N,IND,IW)
C      DO 287 I=2,NF
C      IF(PX(I).NE.0.000) PX(I)=PX(I)+PARM1+PARM2+PARM3
C      IF(PY(I).NE.0.000) PY(I)=PY(I)+PARM1+PARM2+PARM3
C      IF(PZ(I).NE.0.000) PZ(I)=PZ(I)+PARM1+PARM2+PARM3
C      287 CONTINUE
C      IF(PX(1).NE. 0.000) PX(1)=PX(1)+((PARM1+PARM2+PARM3)/FLOAD)
C      DO 285 I=1,N
C      U(I)=U(I)+EPS#Q1(I)+0.5#EPS#2#Q2(I)+EPS##3#Q3(I)/6.

```

```

AYO 1610
AYO 1620
AYO 1630
AYO 1640
AYO 1650
AYO 1660
AYO 1670
AYO 1680
AYO 1690
AYO 1700
AYO 1710
AYO 1720
AYO 1730
AYO 1740
AYO 1750
AYO 1760
AYO 1770
AYO 1780
AYO 1790
AYO 1800
AYO 1810
AYO 1820
AYO 1830
AYO 1840
AYO 1850
AYO 1860
AYO 1870
AYO 1880
AYO 1890
AYO 1900
AYO 1910
AYO 1920

```

```

285 CONTINUE
C
C   STORE DISPLACEMENT U(IND) AND THE CORRESPONDING LOAD PX(J)
C   FOR PLOTTING THE 'EQUILIBRIUM PATH'
C
      P1(KK)=U(IND)
      P2(KK)=PX(I)
C
      KK=KK+1
      KK2=KK-2
      IF(U(IND).GE.4.5) GO TO 286
      GO TO 284
286 CONTINUE
      DO 410 I=1, KK2
      WRITE(7,400) P1(I), P2(I)
      400 FORMAT(2E15.4)
      410 CONTINUE
C
C   PLOT THE 'EQUILIBRIUM PATH'
C
      CALL ABPLOT(KK2, P1, P2)
      READ(5,300) LABELX, LABELY
      300 FORMAT(A4, 5X, A4)
      CALL EZPLOT(P1, P2, LABELX, LABELY, 1, 1, KK2, -11, 0)
      CALL PLOT(0.0, 0.0, -4)
      STOP
      END
AYD 1930
AYD 1940
AYD 1950
AYD 1960
AYD 1970
AYD 1980
AYD 1990
AYD 2000
AYD 2010
AYD 2020
AYD 2030
AYD 2040
AYD 2050
AYD 2060
AYD 2070
AYD 2080
AYD 2090
AYD 2100
AYD 2110
AYD 2120
AYD 2130
AYD 2140
AYD 2150
AYD 2160
AYD 2170
AYD 2180
AYD 2190

```

```

SUBROUTINE V(N,A,VAL,GRAD)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION A(39),GRAD(39),DE(42)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
KOUNT=KOUNT+1
WORK=0.000
DO 6000 J=1,NF
  J1=3#J-2
  J2=3#J-1
  J3=3#J
  WORK=WORK+PX(J)*A(J1)+PY(J)*A(J2)+PZ(J)*A(J3)
  DE(J1)=0.000
  DE(J2)=0.000
  DE(J3)=0.000
6000 CONTINUE
STENG=0.000
DO 6001 M=1,NM
  J=INP(M)
  K=INQ(M)
  IF(J.GT.NF.AND.K.GT.NF) GO TO 6001
  M1=3#J-2
  M2=3#J-1
  M3=3#J
  M4=3#K-2
  M5=3#K-1
  M6=3#K
  IF(J.GT.NF) GO TO 700
  IF(K.GT.NF) GO TO 750

```

```

VEE 0010
VEE 0020
VEE 0030
VEE 0040
VEE 0050
VEE 0060
VEE 0070
VEE 0080
VEE 0090
VEE 0100
VEE 0110
VEE 0120
VEE 0130
VEE 0140
VEE 0150
VEE 0160
VEE 0170
VEE 0180
VEE 0190
VEE 0200
VEE 0210
VEE 0220
VEE 0230
VEE 0240
VEE 0250
VEE 0260
VEE 0270
VEE 0280
VEE 0290
VEE 0300
VEE 0310
VEE 0320

```

```

700 GO TO 760
    DU(M)=A(M4)
    DV(M)=A(M5)
    DW(M)=A(M6)
    SUMX(M)=DX(M)+DU(M)
    SUMY(M)=DY(M)+DV(M)
    SUMZ(M)=DZ(M)+DW(M)
    XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
    DL(M)=XLP(M)-XL(M)
    STENG=STENG+(0.5*DL(M)**2)/XL(M)
    DE(M4)=DE(M4)+(DL(M)*SUMX(M))/(XL(M)*XLP(M))
    DE(M5)=DE(M5)+(DL(M)*SUMY(M))/(XL(M)*XLP(M))
    DE(M6)=DE(M6)+(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
750 GO TO 6001
    DU(M)=-A(M1)
    DV(M)=-A(M2)
    DW(M)=-A(M3)
    SUMX(M)=DX(M)+DU(M)
    SUMY(M)=DY(M)+DV(M)
    SUMZ(M)=DZ(M)+DW(M)
    XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
    DL(M)=XLP(M)-XL(M)
    STENG=STENG+(0.5*DL(M)**2)/XL(M)
    DE(M1)=DE(M1)-(DL(M)*SUMX(M))/(XL(M)*XLP(M))
    DE(M2)=DE(M2)-(DL(M)*SUMY(M))/(XL(M)*XLP(M))
    DE(M3)=DE(M3)-(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
760 GO TO 6001
    DU(M)=A(M4)-A(M1)
    DV(M)=A(M5)-A(M2)
    DW(M)=A(M6)-A(M3)
    SUMX(M)=DX(M)+DU(M)
    SUMY(M)=DY(M)+DV(M)

```

```

VEE 0330
VEE 0340
VEE 0350
VEE 0360
VEE 0370
VEE 0380
VEE 0390
VEE 0400
VEE 0410
VEE 0420
VEE 0430
VEE 0440
VEE 0450
VEE 0460
VEE 0470
VEE 0480
VEE 0490
VEE 0500
VEE 0510
VEE 0520
VEE 0530
VEE 0540
VEE 0550
VEE 0560
VEE 0570
VEE 0580
VEE 0590
VEE 0600
VEE 0610
VEE 0620
VEE 0630
VEE 0640

```

```

SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
STENG=STENG+(0.5*DL(M)**2)/XL(M)
DE(M1)=DE(M1)-(DL(M)*SUMX(M))/(XL(M)*XLP(M))
DE(M2)=DE(M2)-(DL(M)*SUMY(M))/(XL(M)*XLP(M))
DE(M3)=DE(M3)-(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
DE(M4)=DE(M4)+(DL(M)*SUMX(M))/(XL(M)*XLP(M))
DE(M5)=DE(M5)+(DL(M)*SUMY(M))/(XL(M)*XLP(M))
DE(M6)=DE(M6)+(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
6001 CONTINUE
VAL=STENG-WORK
DO 6002 J=1,NF
J1=3*J-2
J2=3*J-1
J3=3*J
GRAD(J1)=DE(J1)-PX(J)
GRAD(J2)=DE(J2)-PY(J)
GRAD(J3)=DE(J3)-PZ(J)
6002 CONTINUE
RETURN
END
VEE 0650
VEE 0660
VEE 0670
VEE 0680
VEE 0690
VEE 0700
VEE 0710
VEE 0720
VEE 0730
VEE 0740
VEE 0750
VEE 0760
VEE 0770
VEE 0780
VEE 0790
VEE 0800
VEE 0810
VEE 0820
VEE 0830
VEE 0840
VEE 0850
VEE 0860

```



```

SUBROUTINE MAX(N,X,JND)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION X(39)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOOD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
XMAX=DABS(X(1))
JND=1
IF(N.EQ.1) GO TO 20
DO 10 I=2,N
IF(DABS(X(I)).LE.XMAX) GO TO 10
XMAX=DABS(X(I))
JND=I
10 CONTINUE
20 RETURN
END
MAX 0010
MAX 0020
MAX 0030
MAX 0040
MAX 0050
MAX 0060
MAX 0070
MAX 0080
MAX 0090
MAX 0100
MAX 0110
MAX 0120
MAX 0130
MAX 0140
MAX 0150
MAX 0160
MAX 0170
MAX 0180

```

```

SUBROUTINE INVE(A,N,D,L,M)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION A(1),L(1),M(1)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
D=1.0
NK=-N
DO 80 K=1,N
NK=NK+N
L(K)=K
M(K)=K
KK=NK+K
BIGA=A(KK)
DO 20 J=K,N
IZ=N*(J-1)
DO 20 I=K,N
IJ=IZ+I
10 IF(DABS(BIGA)-DABS(A(IJ))) 15,20,20
15 BIGA=A(IJ)
L(K)=I
M(K)=J
20 CONTINUE
J=L(K)
IF(J-K) 35,35,25
25 KI=K-N
DO 30 I=1,N
KI=KI+N
HOLD=-A(KI)
JI=KI-K+J
INV 0010
INV 0020
INV 0030
INV 0040
INV 0050
INV 0060
INV 0070
INV 0080
INV 0090
INV 0100
INV 0110
INV 0120
INV 0130
INV 0140
INV 0150
INV 0160
INV 0170
INV 0180
INV 0190
INV 0200
INV 0210
INV 0220
INV 0230
INV 0240
INV 0250
INV 0260
INV 0270
INV 0280
INV 0290
INV 0300
INV 0310
INV 0320

```

```

A(KI)=A(JI)
30 A(JI) =HOLD
35 I=M(K)
IF(I-K) 45,45,38
38 JP=N*(I-1)
DO 40 J=1,N
JK=NK+J
JI=JP+J
HOLD=-A(JK)
A(JK)=A(JI)
40 A(JI) =HOLD
45 IF(BIGA) 48,46,48
46 D=0.0
RETURN
48 DO 55 I=1,N
IF(I-K) 50,55,50
50 IK=NK+I
A(IK)=A(IK)/(-BIGA)
55 CONTINUE
DO 65 I=1,N
IK=NK+I
HOLD=A(IK)
IJ=I-N
DO 65 J=1,N
IJ=IJ+N
IF(I-K) 60,65,60
60 IF(J-K) 62,65,62
62 KJ=IJ-I+K
65 CONTINUE
A(IJ)=HOLD*A(KJ)+A(IJ)
KJ=K-N
DO 75 J=1,N

```

```

INV 0330
INV 0340
INV 0350
INV 0360
INV 0370
INV 0380
INV 0390
INV 0400
INV 0410
INV 0420
INV 0430
INV 0440
INV 0450
INV 0460
INV 0470
INV 0480
INV 0490
INV 0500
INV 0510
INV 0520
INV 0530
INV 0540
INV 0550
INV 0560
INV 0570
INV 0580
INV 0590
INV 0600
INV 0610
INV 0620
INV 0630
INV 0640

```

```

KJ=KJ+N
IF(J-K) 70,75,70
70 A(KJ)=A(KJ)/BIGA
75 CONTINUE
IF(DABS(D).LE.1.D-40) D=0.0D0
D=D*BIGA
A(KK)=1.0/BIGA
80 CONTINUE
K=N
100 K=(K-1)
IF(K) 150,150,105
105 I=L(K)
IF(I-K) 120,120,108
108 JQ=N*(K-1)
JR=N*(I-1)
DO 110 J=1,N
JK=JQ+J
HOLD=A(JK)
JI=JR+J
A(JK)=-A(JI)
110 A(JI) =HOLD
120 J=M(K)
IF(J-K) 100,100,125
125 KI=K-N
DO 130 I=1,N
KI=KI+N
HOLD=A(KI)
JI=KI-K+J
A(KI)=-A(JI)
130 A(JI) =HOLD
GO TO 100
150 RETURN
INV 0650
INV 0660
INV 0670
INV 0680
INV 0690
INV 0700
INV 0710
INV 0720
INV 0730
INV 0740
INV 0750
INV 0760
INV 0770
INV 0780
INV 0790
INV 0800
INV 0810
INV 0820
INV 0830
INV 0840
INV 0850
INV 0860
INV 0870
INV 0880
INV 0890
INV 0900
INV 0910
INV 0920
INV 0930
INV 0940
INV 0950
INV 0960

```

END

INV 0970

```

SUBROUTINE FIRIN(X,A,PARM,EPS,N,IND,IW)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION X(N),A(N)
DIMENSION E(21,21),Q(21),R(21),LN(21),MW(21)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
IORDER=1
IF(A(IND).EQ.1.0D0) Q(IND)=EPS
IF(A(IND).NE.1.0D0) INS=2
IF(INS.EQ.2) GO TO 5000
INS=1
CALL FIRST(E,X,N)
DO 1000 I=1,N
R(I)=-E(I,IND)
E(I,IND)=0.0D0
1000 CONTINUE
DO 2000 J=1,NF
J1=3*J-2
J2=3*J-1
J3=3*J
IF(PX(J).NE.0.0D0) E(J1,IND)=-1.0D0
IF(PY(J).NE.0.0D0) E(J2,IND)=-1.0D0
IF(PZ(J).NE.0.0D0) E(J3,IND)=-1.0D0
2000 CONTINUE
IF(PX(1).NE.0.0D0) E(1,IND)=-1.0D0/FLOAD
CALL INVE(E,N,DET,LN,MW)
DO 3000 I=1,N
PARM=PARM+E(IND,I)*R(I)
DO 3000 J=1,N
FRN 0010
FRN 0020
FRN 0030
FRN 0040
FRN 0050
FRN 0060
FRN 0070
FRN 0080
FRN 0090
FRN 0100
FRN 0110
FRN 0120
FRN 0130
FRN 0140
FRN 0150
FRN 0160
FRN 0170
FRN 0180
FRN 0190
FRN 0200
FRN 0210
FRN 0220
FRN 0230
FRN 0240
FRN 0250
FRN 0260
FRN 0270
FRN 0280
FRN 0290
FRN 0300
FRN 0310
FRN 0320

```

```

IF(I.EQ.IND) GO TO 3000
A(I)=A(I)+E(I,J)*R(J)
3000 CONTINUE
DO 4000 I=1,N
IF(I.EQ.IND) PARM=EPS*PARM
IF(I.NE.IND) Q(I)=A(I)*EPS
4000 CONTINUE
GO TO 9050
5000 CALL FIRST(E,X,N)
CALL INVE(E,N,DET,LN,MW)
DO 6000 I=1,N
R(I)=0.000
6000 CONTINUE
DO 7000 J=1,NF
J1=3*J-2
J2=3*J-1
J3=3*J
IF(PX(J).NE.0.000) R(J1)=1.000
IF(PY(J).NE.0.000) R(J2)=1.000
IF(PZ(J).NE.0.000) R(J3)=1.000
7000 CONTINUE
DO 8000 I=1,N
DO 8000 J=1,N
A(I)=A(I)+E(I,J)*R(J)
8000 CONTINUE
PARM=EPS
DO 9000 I=1,N
Q(I)=A(I)*EPS
9000 CONTINUE
9050 CONTINUE
RETURN
END
FRN 0330
FRN 0340
FRN 0350
FRN 0360
FRN 0370
FRN 0380
FRN 0390
FRN 0400
FRN 0410
FRN 0420
FRN 0430
FRN 0440
FRN 0450
FRN 0460
FRN 0470
FRN 0480
FRN 0490
FRN 0500
FRN 0510
FRN 0520
FRN 0530
FRN 0540
FRN 0550
FRN 0560
FRN 0570
FRN 0580
FRN 0590
FRN 0600
FRN 0610
FRN 0620
FRN 0630
FRN 0640

```

```

SUBROUTINE SFCIN(U, PARM1, PARM2, EPS, N, IND, IW)
IMPLICIT REAL*8(A-H, O-Z)
DIMENSION U(N), S(39, 4)
DIMENSION E(21, 21), Q(21), R(21), LN(21), MW(21), GRAD(21)
COMMON PX(13), PY(13), PZ(13), DX(42), DY(42), DZ(42)
COMMON XL(42), XLP(42), DU(42), DV(42), DW(42), DL(42)
COMMON SUMX(42), SUMY(42), SUMZ(42)
COMMON Q1(39), Q2(39), Q3(39), Q4(39), FLOAD, SDET
COMMON INP(42), INQ(42), NM, NJ, NF, KOUNT
IORDER=2
IF(Q1(IND).NE.1.000) GO TO 2000
INS=1
CALL FIRST(E, U, N)
DO 2000 I=1, N
E(I, IND)=0.000
2000 CONTINUE
DO 2050 J=1, NF
J1=3*J-2
J2=3*J-1
J3=3*J
IF(PX(J).NE.0.000) E(J1, IND)=-1.000
IF(PY(J).NE.0.000) E(J2, IND)=-1.000
IF(PZ(J).NE.0.000) E(J3, IND)=-1.000
2050 CONTINUE
IF(PX(1).NE.0.000) E(1, IND)=-1.000/FLOAD
CALL INVE(E, N, DET, LN, MW)
CALL SECOND(S, U, N, 4)
DO 2100 I=1, N
DO 2100 J=1, N
IF(I.EQ.IND) PARM2=PARM2-E(I, J)*S(J, I)
IF(I.NE.IND) Q2(I)=Q2(I)-E(I, J)*S(J, I)
2100 CONTINUE
SIN 0010
SIN 0020
SIN 0030
SIN 0040
SIN 0050
SIN 0060
SIN 0070
SIN 0080
SIN 0090
SIN 0100
SIN 0110
SIN 0120
SIN 0130
SIN 0140
SIN 0150
SIN 0160
SIN 0170
SIN 0180
SIN 0190
SIN 0200
SIN 0210
SIN 0220
SIN 0230
SIN 0240
SIN 0250
SIN 0260
SIN 0270
SIN 0280
SIN 0290
SIN 0300
SIN 0310
SIN 0320

```



```

2200 GO TO 2500
2200 INS=2
2200 CALL FIRST(E,U,N)
2200 CALL INVE(E,N,DET,LN,MW)
2200 CALL SECOND(S,U,N,4)
2200 PARM2=0.000
2200 DO 2300 I=1,N
2200 DO 2300 J=1,N
2200 Q2(I)=Q2(I)-E(I,J)*S(J,I)
2300 CONTINUE
2300 PARM=EPS
2300 DO 2400 I=1,N
2300 Q(I)=0.5*EPS**2*Q2(I)
2400 CONTINUE
2400 GO TO 2800
2500 Q(IND)=0.000
2500 DO 2700 I=1,N
2500 IF(I.EQ.IND) GO TO 2600
2500 Q(I)=0.5*EPS**2*Q2(I)
2600 GO TO 2700
2600 PARM2=0.5*EPS**2*PARM2
2700 CONTINUE
2800 CONTINUE
2800 PARM=PARM1+PARM2
2800 DO 2900 I=1,N
2800 Q(I)= EPS*Q1(I) + 0.5*EPS**2*Q2(I)
2900 CONTINUE
2900 DO 2950 I=2,NF
2900 IF(PX(I).NE.0.000) PX(I)=PX(I)+PARM
2900 IF(PY(I).NE.0.000) PY(I)=PY(I)+PARM
2900 IF(PZ(I).NE.0.000) PZ(I)=PZ(I)+PARM
2950 CONTINUE
SIN 0330
SIN 0340
SIN 0350
SIN 0360
SIN 0370
SIN 0380
SIN 0390
SIN 0400
SIN 0410
SIN 0420
SIN 0430
SIN 0440
SIN 0450
SIN 0460
SIN 0470
SIN 0480
SIN 0490
SIN 0500
SIN 0510
SIN 0520
SIN 0530
SIN 0540
SIN 0550
SIN 0560
SIN 0570
SIN 0580
SIN 0590
SIN 0600
SIN 0610
SIN 0620
SIN 0630
SIN 0640

```

```

IF(PX(I) .NE. 0.000) PX(I)=PX(I)+(PARM/FLOAD)
DO 3000 I=1,N
  Q(I)=Q(I) + U(I)
3000 CONTINUE
  CALL V(N,Q,FUNT,GRAD)
  CALL OUTP(Q,GRAD,FUNT,N,IORDER,INS,IW)
DO 3050 I=2,NF
  IF(PX(I).NE.0.000) PX(I)=PX(I)-PARM
  IF(PY(I).NE.0.000) PY(I)=PY(I)-PARM
  IF(PZ(I).NE.0.000) PZ(I)=PZ(I)-PARM
3050 CONTINUE
  IF(PX(I) .NE. 0.000) PX(I)=PX(I)-(PARM/FLOAD)
  RETURN
END

```

```

SIN 0650
SIN 0660
SIN 0670
SIN 0680
SIN 0690
SIN 0700
SIN 0710
SIN 0720
SIN 0730
SIN 0740
SIN 0750
SIN 0760
SIN 0770
SIN 0780

```

```

SUBROUTINE THRIN(U,PARM1,PARM2,PARM3,EPS,N,IND,IW)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION U(N),S(39,4),T(39,3)
DIMENSION E(21,21),Q(21),R(21),LN(21),MW(21),GRAD(21)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
IORDER=3
IF(Q1(IND).NE.1.000) GO TO 3200
INS=1
CALL FIRST(E,U,N)
CALL INVE(E,N,DET3,LN,MW)
SDET=DET3*(10.**25)
CALL FIRST(E,U,N)
DO 3000 I=1,N
E(I,IND)=0.000
3000 CONTINUE
DO 3050 J=1,NF
J1=3*J-2
J2=3*J-1
J3=3*J
IF(PX(J).NE.0.000) E(J1,IND)=-1.000
IF(PY(J).NE.0.000) E(J2,IND)=-1.000
IF(PZ(J).NE.0.000) E(J3,IND)=-1.000
3050 CONTINUE
IF(PX(1).NE.0.000) E(1,IND)=-1.000/FLOAD
CALL INVE(E,N,DET,LN,MW)
CALL SECOND(S,U,N,4)
CALL THIRD(T,U,N)
DO 3100 I=1,N
TRN 0010
TRV 0020
TRN 0030
TRN 0040
TRN 0050
TRN 0060
TRN 0070
TRN 0080
TRN 0090
TRV 0100
TRN 0110
TRN 0120
TRN 0130
TRN 0140
TRN 0150
TRN 0160
TRN 0170
TRN 0180
TRN 0190
TRN 0200
TRV 0210
TRN 0220
TRN 0230
TRN 0240
TRN 0250
TRN 0260
TRV 0270
TRN 0280
TRN 0290
TRN 0300
TRN 0310
TRN 0320

```

```

DO 3100 J=1,N
IF(I.EQ.IND) PARM3=PARM3-E(I,J)*(S(J,2)+T(J,1))
IF(I.NE.IND) Q3(I)=Q3(I)-E(I,J)*(S(J,2)+T(J,1))
3100 CONTINUE
GO TO 3500
3200 INS=2
CALL FIRST(E,U,N)
CALL INVE(E,N,DET,LN,MW)
CALL SECOND(S,U,N,4)
CALL THIRD(T,U,N)
PARM3=0.0D0
DO 3300 I=1,N
DO 3300 J=1,N
Q3(I)=Q3(I)-E(I,J)*(S(J,2)+T(J,1))
3300 CONTINUE
DO 3400 I=1,N
Q(I)=EPS**3*Q3(I)/6.
3400 CONTINUE
GO TO 3800
3500 Q(IND)=0.0D0
DO 3700 I=1,N
IF(I.EQ.IND) PARM3=EPS**3*PARM3/6.
IF(I.NE.IND) Q(I)=EPS**3*Q3(I)/6.
3700 CONTINUE
3800 CALL TAYLOR(Q,U,PARM3,N,INS,IORDER,IW)
PARM=PARM1+PARM2+PARM3
DO 3900 I=1,N
Q(I)=EPS*Q1(I)+0.5*EPS**2*Q2(I)+EPS**3*Q3(I)/6.
3900 CONTINUE
DO 3950 I=2,NF
IF(PX(I).NE.0.0D0) PX(I)=PX(I)+PARM
IF(PY(I).NE.0.0D0) PY(I)=PY(I)+PARM

```

```

TRN 0330
TRN 0340
TRN 0350
TRN 0360
TRN 0370
TRN 0380
TRN 0390
TRN 0400
TRN 0410
TRN 0420
TRN 0430
TRN 0440
TRN 0450
TRN 0460
TRN 0470
TRN 0480
TRN 0490
TRN 0500
TRN 0510
TRN 0520
TRN 0530
TRN 0540
TRN 0550
TRN 0560
TRN 0570
TRN 0580
TRN 0590
TRN 0600
TRN 0610
TRN 0620
TRN 0630
TRN 0640

```

```

IF(PZ(I).NE.0.000) PZ(I)=PX(I)+PARM
3950 CONTINUE
IF(PX(I).NE.0.000) PX(I)=PX(I)+(PARM/FLCAD)
DO 4000 I=1,N
Q(I)=Q(I)+U(I)
4000 CONTINUE
CALL V(N,Q,FUNT,GRAD)
CALL OUTP(Q,GRAD,FUNT,N,IORDER,INS,IW)
DO 4050 I=2,NF
IF(PX(I).NE.0.000) PX(I)=PX(I)-PARM
IF(PY(I).NE.0.000) PY(I)=PY(I)-PARM
IF(PZ(I).NE.0.000) PZ(I)=PZ(I)-PARM
4050 CONTINUE
IF(PX(I).NE.0.000) PX(I)=PX(I)-(PARM/FLCAD)
RETURN
END

```

```

TRN 0650
TRN 0660
TRN 0670
TRN 0680
TRN 0690
TRN 0700
TRN 0710
TRN 0720
TRN 0730
TRN 0740
TRN 0750
TRN 0760
TRN 0770
TRN 0780
TRN 0790
TRN 0800

```

```

SUBROUTINE FIRST(X,Q,N)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION X(N,N),Q(N)
DIMENSION DLU(42),DLV(42),DLW(42)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
DO 1000 I=1,N
DO 1000 J=1,N
X(I,J)=0.0D0
1000 CONTINUE
DO 6000 M=1,NM
J=INP(M)
K=INQ(M)
IF(J.GT.NF.AND.K.GT.NF) GO TO 6000
IF(J.GT.NF) GO TO 4000
IF(K.GT.NF) GO TO 5000
M1=3#J-2
M2=3#J-1
M3=3#J
M4=3#K-2
M5=3#K-1
M6=3#K
DU(M)=Q(M4)-Q(M1)
DV(M)=Q(M5)-Q(M2)
DW(M)=Q(M6)-Q(M3)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
1ST 0010
1ST 0020
1ST 0030
1ST 0040
1ST 0050
1ST 0060
1ST 0070
1ST 0080
1ST 0090
1ST 0100
1ST 0110
1ST 0120
1ST 0130
1ST 0140
1ST 0150
1ST 0160
1ST 0170
1ST 0180
1ST 0190
1ST 0200
1ST 0210
1ST 0220
1ST 0230
1ST 0240
1ST 0250
1ST 0260
1ST 0270
1ST 0280
1ST 0290
1ST 0300
1ST 0310
1ST 0320

```

```

DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
X(M1,M1)=X(M1,M1)+(DL(M)*(1.-DLU(M)*DLU(M))/XLP(M)
$ + DLU(M)*DLU(M))/XL(M)
X(M1,M2)=X(M1,M2) + (DLV(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M2,M1)=X(M2,M1) + (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M1,M3)=X(M1,M3) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M1)=X(M3,M1) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M2,M2)=X(M2,M2) + (DL(M)*(1.-DLV(M)*DLV(M))/XLP(M)
$ + DLV(M)*DLV(M))/XL(M)
X(M2,M3)=X(M2,M3) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M2)=X(M3,M2) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M3)=X(M3,M3)+(DL(M)*(1.-DLW(M)*DLW(M))/XLP(M)
$ + DLW(M)*DLW(M))/XL(M)
X(M4,M4)=X(M4,M4) + (DL(M)*(1.-DLU(M)*DLU(M))/XLP(M)
$ + DLU(M)*DLU(M))/XL(M)
X(M4,M5)=X(M4,M5) + (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M4,M6)=X(M4,M6) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M5,M4)=X(M5,M4) + (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M5,M5)=X(M5,M5) + (DL(M)*(1.-DLV(M)*DLV(M))/XLP(M)
$ + DLV(M)*DLV(M))/XL(M)
X(M5,M6)=X(M5,M6) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M4)=X(M6,M4) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M5)=X(M6,M5) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M6)=X(M6,M6) + (DL(M)*(1.-DLW(M)*DLW(M))/XLP(M)
$ + DLW(M)*DLW(M))/XL(M)
X(M1,M4)=X(M1,M4) - (DL(M)*(1.-DLU(M)*DLU(M))/XLP(M)
$ + DLU(M)*DLU(M))/XL(M)
X(M1,M5)=X(M1,M5) - (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M1,M6)=X(M1,M6) - (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)

```

```

1ST 0330
1ST 0340
1ST 0350
1ST 0360
1ST 0370
1ST 0380
1ST 0390
1ST 0400
1ST 0410
1ST 0420
1ST 0430
1ST 0440
1ST 0450
1ST 0460
1ST 0470
1ST 0480
1ST 0490
1ST 0500
1ST 0510
1ST 0520
1ST 0530
1ST 0540
1ST 0550
1ST 0560
1ST 0570
1ST 0580
1ST 0590
1ST 0600
1ST 0610
1ST 0620
1ST 0630
1ST 0640

```

```

X(M2,M4)=X(M2,M4) - (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M2,M5)=X(M2,M5) -(DL(M)*(1.-DLV(M)*DLV(M))/XLP(M)
$ + DLV(M)*DLV(M))/XL(M)
X(M2,M6)=X(M2,M6) - (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M4)=X(M3,M4) - (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M5)=X(M3,M5) - (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M6)=X(M3,M6) - (DL(M)*(1.-DLW(M)*DL(M))/XLP(M)
$ + DLW(M)*DLW(M))/XL(M)
X(M4,M1)=X(M4,M1) - (DL(M)*(1.-DLU(M)*DLU(M))/XLP(M)
$ + DLU(M)*DLU(M))/XL(M)
X(M5,M1)=X(M5,M1) - (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M1)=X(M6,M1) - (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M4,M2)=X(M4,M2) - (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M5,M2)=X(M5,M2) -(DL(M)*(1.-DLV(M)*DLV(M))/XLP(M)
$ + DLV(M)*DLV(M))/XL(M)
X(M6,M2)=X(M6,M2) - (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M4,M3)=X(M4,M3) - (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M5,M3)=X(M5,M3) - (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M3)=X(M6,M3) - (DL(M)*(1.-DLW(M)*DL(M))/XLP(M)
$ + DLW(M)*DLW(M))/XL(M)
GO TO 5500
4000 M4=3*K-2
M5=3*K-1
M6=3*K
DU(M)=Q(M4)
DV(M)=Q(M5)
DW(M)=Q(M6)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
1ST 0650
1ST 0660
1ST 0670
1ST 0680
1ST 0690
1ST 0700
1ST 0710
1ST 0720
1ST 0730
1ST 0740
1ST 0750
1ST 0760
1ST 0770
1ST 0780
1ST 0790
1ST 0800
1ST 0810
1ST 0820
1ST 0830
1ST 0840
1ST 0850
1ST 0860
1ST 0870
1ST 0880
1ST 0890
1ST 0900
1ST 0910
1ST 0920
1ST 0930
1ST 0940
1ST 0950
1ST 0960

```



```

DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
X(M4,M4)=X(M4,M4) + (DL(M)*(1.-DLU(M)*DLU(M))/XLP(M)
$ + DLU(M)*DLU(M))/XL(M)
X(M4,M5)=X(M4,M5) + (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M5,M4)=X(M5,M4) + (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M4,M6)=X(M4,M6) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M4)=X(M6,M4) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M5,M5)=X(M5,M5) + (DL(M)*(1.-DLV(M)*DLV(M))/XLP(M)
$ + DLV(M)*DLV(M))/XL(M)
X(M5,M6)=X(M5,M6) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M5)=X(M6,M5) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M6,M6)=X(M6,M6) + (DL(M)*(1.-DLW(M)*DLW(M))/XLP(M)
$ + DLW(M)*DLW(M))/XL(M)
GO TO 5500
5000 M1=3#J-2
M2=3#J-1
M3=3#J
DU(M)=-Q(M1)
DV(M)=-Q(M2)
DW(M)=-Q(M3)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
X(M1,M1)=X(M1,M1)+(DL(M)*(1.-DLU(M)*DLU(M))/XLP(M)
$ + DLU(M)*DLU(M))/XL(M)
1ST 0970
1ST 0980
1ST 0990
1ST 1000
1ST 1010
1ST 1020
1ST 1030
1ST 1040
1ST 1050
1ST 1060
1ST 1070
1ST 1080
1ST 1090
1ST 1100
1ST 1110
1ST 1120
1ST 1130
1ST 1140
1ST 1150
1ST 1160
1ST 1170
1ST 1180
1ST 1190
1ST 1200
1ST 1210
1ST 1220
1ST 1230
1ST 1240
1ST 1250
1ST 1260
1ST 1270
1ST 1280

```

```

X(M1,M2)=X(M1,M2) + (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M2,M1)=X(M2,M1) + (DLU(M)*DLV(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M1,M3)=X(M1,M3) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M1)=X(M3,M1) + (DLU(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M2,M2)=X(M2,M2) + (DL(M)*(1.-DLV(M)*DLV(M))/XLP(M)
$ + DLV(M)*DLV(M))/XL(M)
X(M2,M3)=X(M2,M3) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M2)=X(M3,M2) + (DLV(M)*DLW(M)*(1.-DL(M)/XLP(M)))/XL(M)
X(M3,M3)=X(M3,M3) + (DL(M)*(1.-DLW(M)*DLW(M))/XLP(M)
$ + DLW(M)*DLW(M))/XL(M)
5500 CONTINUE
6000 CONTINUE
      RETURN
      END
1ST 1290
1ST 1300
1ST 1310
1ST 1320
1ST 1330
1ST 1340
1ST 1350
1ST 1360
1ST 1370
1ST 1380
1ST 1390
1ST 1400
1ST 1410
1ST 1420

```

```

SUBROUTINE SECOND(S,Q,N,IORDER)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION A(3,3),B(3,3),C(3,3)
DIMENSION RAI(3),RBI(3),RA2(3),RB2(3),RA4(3),RB4(3),RA3(3),RB3(3)
DIMENSION S(N,IORDER),Q(N)
DIMENSION DLU(42),DLV(42),DLW(42)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
DO 2000 I=1,N
DO 2000 J=1,IORDER
S(I,J)=0.000
2000 CONTINUE
DO 2080 M=1,NM
J1=INP(M)
K1=INQ(M)
IF(J1.GT.NF.AND.K1.GT.NF) GO TO 2080
IF(J1.GT.NF) GO TO 2060
IF(K1.GT.NF) GO TO 2070
M1=3*J1-2
M2=3*J1-1
M3=3*J1
M4=3*K1-2
M5=3*K1-1
M6=3*K1
RA1(1)=Q1(M1)
RA1(2)=Q1(M2)
RA1(3)=Q1(M3)
RBI(1)=Q1(M4)
RBI(2)=Q1(M5)
2ND 0010
2ND 0020
2ND 0030
2ND 0040
2ND 0050
2ND 0060
2ND 0070
2ND 0080
2ND 0090
2ND 0100
2ND 0110
2ND 0120
2ND 0130
2ND 0140
2ND 0150
2ND 0160
2ND 0170
2ND 0180
2ND 0190
2ND 0200
2ND 0210
2ND 0220
2ND 0230
2ND 0240
2ND 0250
2ND 0260
2ND 0270
2ND 0280
2ND 0290
2ND 0300
2ND 0310
2ND 0320

```

```

RB1(3)=Q1(M6)
RA2(1)=Q2(M1)
RA2(2)=Q2(M2)
RA2(3)=Q2(M3)
RB2(1)=Q2(M4)
RB2(2)=Q2(M5)
RB2(3)=Q2(M6)
RA3(1)=Q3(M1)
RA3(2)=Q3(M2)
RA3(3)=Q3(M3)
RB3(1)=Q3(M4)
RB3(2)=Q3(M5)
RB3(3)=Q3(M6)
RA4(1)=Q4(M1)
RA4(2)=Q4(M2)
RA4(3)=Q4(M3)
RB4(1)=Q4(M4)
RB4(2)=Q4(M5)
RB4(3)=Q4(M6)
DU(M)=Q(M4)-Q(M1)
DV(M)=Q(M5)-Q(M2)
DW(M)=Q(M6)-Q(M3)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
D11=(1.-DLU(M)**2)/XLP(M)
D22=(1.-DLV(M)**2)/XLP(M)
2ND 0330
2ND 0340
2ND 0350
2ND 0360
2ND 0370
2ND 0380
2ND 0390
2ND 0400
2ND 0410
2ND 0420
2ND 0430
2ND 0440
2ND 0450
2ND 0460
2ND 0470
2ND 0480
2ND 0490
2ND 0500
2ND 0510
2ND 0520
2ND 0530
2ND 0540
2ND 0550
2ND 0560
92ND 0570
2ND 0580
2ND 0590
2ND 0600
2ND 0610
2ND 0620
2ND 0630
2ND 0640

```

D33=(1.-DLW(M)**2)/XLP(M) 2ND 0650
 D12=-DLU(M)*DLV(M)/XLP(M) 2ND 0660
 D13=-DLU(M)*DLW(M)/XLP(M) 2ND 0670
 D23=-DLV(M)*DLW(M)/XLP(M) 2ND 0680
 D111=(-3.*DLU(M)*(1.-DLU(M)**2))/XLP(M)**2 2ND 0690
 D222=(-3.*DLV(M)*(1.-DLV(M)**2))/XLP(M)**2 2ND 0700
 D333=(-3.*DLW(M)*(1.-DLW(M)**2))/XLP(M)**2 2ND 0710
 D112=(-DLV(M)*(1.-3.*DLU(M)**2))/XLP(M)**2 2ND 0720
 D113=(-DLW(M)*(1.-3.*DLU(M)**2))/XLP(M)**2 2ND 0730
 D223=(-DLW(M)*(1.-3.*DLV(M)**2))/XLP(M)**2 2ND 0740
 D122=(-DLU(M)*(1.-3.*DLV(M)**2))/XLP(M)**2 2ND 0750
 D133=(-DLU(M)*(1.-3.*DLW(M)**2))/XLP(M)**2 2ND 0760
 D233=(-DLV(M)*(1.-3.*DLW(M)**2))/XLP(M)**2 2ND 0770
 D123=(3.*DLU(M)*DLV(M)*DLW(M))/XLP(M)**2 2ND 0780
 A(1,1)=(DL(M)*D111+3.*DLU(M)*D11)/XL(M) 2ND 0790
 A(1,2)=(DL(M)*D112+DLV(M)*D11+2.*DLU(M)*D12)/XL(M) 2ND 0800
 A(1,3)=(DL(M)*D113+DLW(M)*D11+2.*DLU(M)*D13)/XL(M) 2ND 0810
 A(2,2)=(DL(M)*D122+DLU(M)*D22+2.*DLV(M)*D12)/XL(M) 2ND 0820
 A(2,3)=(DL(M)*D123+DLW(M)*D12+DLU(M)*D23+DLV(M)*D13)/XL(M) 2ND 0830
 A(3,3)=(DL(M)*D133+2.*DLW(M)*D13+DLU(M)*D33)/XL(M) 2ND 0840
 A(2,1)=A(1,2) 2ND 0850
 A(3,1)=A(1,3) 2ND 0860
 A(3,2)=A(2,3) 2ND 0870
 B(1,1)=A(1,2) 2ND 0880
 B(1,2)=A(2,2) 2ND 0890
 B(1,3)=A(3,2) 2ND 0900
 B(2,2)=(DL(M)*D222+3.*DLV(M)*D22)/XL(M) 2ND 0910
 B(2,3)=(DL(M)*D223+DLW(M)*D22+2.*DLV(M)*D23)/XL(M) 2ND 0920
 B(3,3)=(DL(M)*D233+2.*DLW(M)*D23+DLV(M)*D33)/XL(M) 2ND 0930
 B(2,1)=B(1,2) 2ND 0940
 B(3,1)=B(1,3) 2ND 0950
 B(3,2)=B(2,3) 2ND 0960

```

C(1,1)=A(1,3)
C(1,2)=A(2,3)
C(1,3)=A(3,3)
C(2,1)=C(1,2)
C(2,2)=B(2,3)
C(2,3)=B(3,3)
C(3,3)=(DL(M)*D333+3.*DLW(M)*D33)/XL(M)
C(3,1)=C(1,3)
C(3,2)=C(2,3)

DO 2050 J=1,3
DO 2050 K=1,3
S(M1,1)=S(M1,1) + (RAL(J)*A(J,K)*RAL(K)
$ -2.*RAL(J)*A(J,K)*RB1(K) + RB1(J)*A(J,K)*RB1(K))
S(M2,1)=S(M2,1) + (RAL(J)*B(J,K)*RAL(K)
$ -2.*RAL(J)*B(J,K)*RB1(K) + RB1(J)*B(J,K)*RB1(K))
S(M3,1)=S(M3,1) + (RAL(J)*C(J,K)*RAL(K)
$ -2.*RAL(J)*C(J,K)*RB1(K) + RB1(J)*C(J,K)*RB1(K))
S(M4,1)=S(M4,1) - (RAL(J)*A(J,K)*RAL(K)
$ -2.*RAL(J)*A(J,K)*RB1(K) + RB1(J)*A(J,K)*RB1(K))
S(M5,1)=S(M5,1) - (RAL(J)*B(J,K)*RAL(K)
$ -2.*RAL(J)*B(J,K)*RB1(K) + RB1(J)*B(J,K)*RB1(K))
S(M6,1)=S(M6,1) - (RAL(J)*C(J,K)*RAL(K)
$ -2.*RAL(J)*C(J,K)*RB1(K) + RB1(J)*C(J,K)*RB1(K))
S(M1,2)=S(M1,2) + 3.*(RAL(J)*A(J,K)*RA2(K) - RAL(J)*A(J,K)*RB2(K)
$ - RB1(J)*A(J,K)*RA2(K) + RB1(J)*A(J,K)*RB2(K))
S(M2,2)=S(M2,2) + 3.*(RAL(J)*B(J,K)*RA2(K) - RAL(J)*B(J,K)*RB2(K)
$ - RB1(J)*B(J,K)*RA2(K) + RB1(J)*B(J,K)*RB2(K))
S(M3,2)=S(M3,2) + 3.*(RAL(J)*C(J,K)*RA2(K) - RAL(J)*C(J,K)*RB2(K)
$ - RB1(J)*C(J,K)*RA2(K) + RB1(J)*C(J,K)*RB2(K))
S(M4,2)=S(M4,2) - 3.*(RAL(J)*A(J,K)*RA2(K) - RAL(J)*A(J,K)*RB2(K)
$ - RB1(J)*A(J,K)*RA2(K) + RB1(J)*A(J,K)*RB2(K))

```

C

```

2ND 0970
2ND 0980
2ND 0990
2ND 1000
2ND 1010
2ND 1020
2ND 1030
2ND 1040
2ND 1050
2ND 1060
2ND 1070
2ND 1080
2ND 1090
2ND 1100
2ND 1110
2ND 1120
2ND 1130
2ND 1140
2ND 1150
2ND 1160
2ND 1170
2ND 1180
2ND 1190
2ND 1200
2ND 1210
2ND 1220
2ND 1230
2ND 1240
2ND 1250
2ND 1260
2ND 1270
2ND 1280

```



```

$ + 10.*(RA2(J)*B(J,K)*RA3(K)-RA2(J)*B(J,K)*RB3(K) 2ND 1610
$ - RB2(J)*B(J,K)*RA3(K) + RB2(J)*B(J,K)*RB3(K) 2ND 1620
S(M3,4)=S(M3,4) + 5.*(RA1(J)*C(J,K)*RA4(K) 2ND 1630
$ - RA1(J)*C(J,K)*RB4(K)-RB1(J)*C(J,K)*RA4(K)+RB1(J)*C(J,K)*RB4(K) 2ND 1640
$ + 10.*(RA2(J)*C(J,K)*RA3(K)-RA2(J)*C(J,K)*RB3(K) 2ND 1650
$ - RB2(J)*C(J,K)*RA3(K) + RB2(J)*C(J,K)*RB3(K) 2ND 1660
S(M4,4)=S(M4,4)-5.*(RA1(J)*A(J,K)*RA4(K)-RA1(J)*A(J,K)*RB4(K) 2ND 1670
$ -RB1(J)*A(J,K)*RA4(K)+RB1(J)*A(J,K)*RB4(K)-10.*(RA2(J)*A(J,K) 2ND 1680
$ *RA3(K)-RA2(J)*A(J,K)*RB3(K)-RB2(J)*A(J,K)*RA3(K)+RB2(J)*A(J,K) 2ND 1690
$ *RB3(K)) 2ND 1700
S(M5,4)=S(M5,4) - 5.*(RA1(J)*B(J,K)*RA4(K) 2ND 1710
$ - RA1(J)*B(J,K)*RB4(K)-RB1(J)*B(J,K)*RA4(K)+RB1(J)*B(J,K)*RB4(K) 2ND 1720
$ - 10.*(RA2(J)*B(J,K)*RA3(K)-RA2(J)*B(J,K)*RB3(K) 2ND 1730
$ - RB2(J)*B(J,K)*RA3(K) + RB2(J)*B(J,K)*RB3(K) 2ND 1740
S(M6,4)=S(M6,4) - 5.*(RA1(J)*C(J,K)*RA4(K) 2ND 1750
$ - RA1(J)*C(J,K)*RB4(K)-RB1(J)*C(J,K)*RA4(K)+RB1(J)*C(J,K)*RB4(K) 2ND 1760
$ - 10.*(RA2(J)*C(J,K)*RA3(K)-RA2(J)*C(J,K)*RB3(K) 2ND 1770
$ - RB2(J)*C(J,K)*RA3(K) + RB2(J)*C(J,K)*RB3(K)) 2ND 1780
2050 CONTINUE 2ND 1790
GO TO 2080 2ND 1800
2060 M4=3*K1-2 2ND 1810
M5=3*K1-1 2ND 1820
M6=3*K1 2ND 1830
RB1(1)=Q1(M4) 2ND 1840
RB1(2)=Q1(M5) 2ND 1850
RB1(3)=Q1(M6) 2ND 1860
RB2(1)=Q2(M4) 2ND 1870
RB2(2)=Q2(M5) 2ND 1880
RB2(3)=Q2(M6) 2ND 1890
RB3(1)=Q3(M4) 2ND 1900
RB3(2)=Q3(M5) 2ND 1910
RB3(3)=Q3(M6) 2ND 1920

```



```

RB4(1)=Q4(M4)
RB4(2)=Q4(M5)
RB4(3)=Q4(M6)
DU(M)=Q(M4)
DV(M)=Q(M5)
DW(M)=Q(M6)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
D11=(1.-DLU(M)**2)/XLP(M)
D22=(1.-DLV(M)**2)/XLP(M)
D33=(1.-DLW(M)**2)/XLP(M)
D12=-DLU(M)*DLV(M)/XLP(M)
D13=-DLU(M)*DLW(M)/XLP(M)
D23=-DLV(M)*DLW(M)/XLP(M)
D111=(-3.*DLU(M)*(1.-DLU(M)**2))/XLP(M)**2
D222=(-3.*DLV(M)*(1.-DLV(M)**2))/XLP(M)**2
D333=(-3.*DLW(M)*(1.-DLW(M)**2))/XLP(M)**2
D112=(-DLV(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
D113=(-DLW(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
D223=(-DLW(M)*(1.-3.*DLV(M)**2))/XLP(M)**2
D122=(-DLU(M)*(1.-3.*DLV(M)**2))/XLP(M)**2
D133=(-DLU(M)*(1.-3.*DLW(M)**2))/XLP(M)**2
D233=(-DLV(M)*(1.-3.*DLW(M)**2))/XLP(M)**2
D123=(3.*DLU(M)*DLV(M)*DLW(M))/XLP(M)**2
A(1,1)=(DL(M)*D111+3.*DLU(M)*D11)/XL(M)
A(1,2)=(DL(M)*D112+DLV(M)*D11+2.*DLU(M)*D12)/XL(M)
2ND 1930
2ND 1940
2ND 1950
2ND 1960
2ND 1970
2ND 1980
2ND 1990
2ND 2000
2ND 2010
2ND 2020
2ND 2030
2ND 2040
2ND 2050
2ND 2060
2ND 2070
2ND 2080
2ND 2090
2ND 2100
2ND 2110
2ND 2120
2ND 2130
2ND 2140
2ND 2150
2ND 2160
2ND 2170
2ND 2180
2ND 2190
2ND 2200
2ND 2210
2ND 2220
2ND 2230
2ND 2240

```

```

A(1,3)=(DL(M)*D113+DLW(M)*D11+2*DLU(M)*D13)/XL(M)
A(2,2)=(DL(M)*D122+DLU(M)*D22+2.*DLV(M)*D12)/XL(M)
A(2,3)=(DL(M)*D123+DLW(M)*D12+DLU(M)*D23+DLV(M)*D13)/XL(M)
A(3,3)=(DL(M)*D133+2.*DLW(M)*D13+DLU(M)*D33)/XL(M)
A(2,1)=A(1,2)
A(3,1)=A(1,3)
A(3,2)=A(2,3)
B(1,1)=A(1,2)
B(1,2)=A(2,2)
B(1,3)=A(3,2)
B(2,2)=(DL(M)*D222+3.*DLV(M)*D22)/XL(M)
B(2,3)=(DL(M)*D223+DLW(M)*D22+2.*DLV(M)*D23)/XL(M)
B(3,3)=(DL(M)*D233+2.*DLW(M)*D23+DLV(M)*D33)/XL(M)
R(2,1)=B(1,2)
B(3,1)=B(1,3)
B(3,2)=B(2,3)
C(1,1)=A(1,3)
C(1,2)=A(2,3)
C(1,3)=A(3,3)
C(2,1)=C(1,2)
C(2,2)=B(2,3)
C(2,3)=B(3,3)
C(3,3)=(DL(M)*D333+3.*DLW(M)*D33)/XL(M)
C(3,1)=C(1,3)
C(3,2)=C(2,3)
DO 2065 J=1,3
DO 2065 K=1,3
S(M4,1)=S(M4,1) - RB1(J)*A(J,K)*RB1(K)
S(M5,1)=S(M5,1) - RB1(J)*B(J,K)*RB1(K)
S(M6,1)=S(M6,1) - RB1(J)*C(J,K)*RB1(K)
S(M4,2)=S(M4,2) - 3.*RB1(J)*A(J,K)*RB2(K)
S(M5,2)=S(M5,2) - 3.*RB1(J)*B(J,K)*RB2(K)
2ND 2250
2ND 2260
2ND 2270
2ND 2280
2ND 2290
2ND 2300
2ND 2310
2ND 2320
2ND 2330
2ND 2340
2ND 2350
2ND 2360
2ND 2370
2ND 2380
2ND 2390
2ND 2400
2ND 2410
2ND 2420
2ND 2430
2ND 2440
2ND 2450
2ND 2460
2ND 2470
2ND 2480
2ND 2490
2ND 2500
2ND 2510
2ND 2520
2ND 2530
2ND 2540
2ND 2550
2ND 2560

```

```

S(M6,2)=S(M6,2) - 3.*RB1(J)*C(J,K)*RB2(K)
S(M4,3)=S(M4,3)-3.*RB2(J)*A(J,K)*RB2(J)-RB1(J)*A(J,K)*RB3(K)*4.
S(M5,3)=S(M5,3)-3.*RB2(J)*B(J,K)*RB2(J)-RB1(J)*B(J,K)*RB3(K)*4.
S(M6,3)=S(M6,3)-3.*RB2(J)*C(J,K)*RB2(J)-RB1(J)*C(J,K)*RB3(K)*4.
S(M4,4)=S(M4,4)-5.*RB1(J)*A(J,K)*RE4(K)-10.*RB2(J)*A(J,K)*RB3(K)
S(M5,4)=S(M5,4)-5.*RB1(J)*B(J,K)*RE4(K)-10.*RB2(J)*B(J,K)*RB3(K)
S(M6,4)=S(M6,4)-5.*RB1(J)*C(J,K)*RE4(K)-10.*RB2(J)*C(J,K)*RB3(K)
2065 CONTINUE
GO TO 2080
C
2070 M1=3*J1-2
M2=3*J1-1
M3=3*J1
RA1(1)=Q1(M1)
RA1(2)=Q1(M2)
RA1(3)=Q1(M3)
RA2(1)=Q2(M1)
RA2(2)=Q2(M2)
RA2(3)=Q2(M3)
RA3(1)=Q3(M1)
RA3(2)=Q3(M2)
RA3(3)=Q3(M3)
RA4(1)=Q4(M1)
RA4(2)=Q4(M2)
RA4(3)=Q4(M3)
DU(M)=-Q(M1)
DV(M)=-Q(M2)
DW(M)=-Q(M3)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
2ND 2570
2ND 2580
2ND 2590
2ND 2600
2ND 2610
2ND 2620
2ND 2630
2ND 2640
2ND 2650
2ND 2660
2ND 2670
2ND 2680
2ND 2690
2ND 2700
2ND 2710
2ND 2720
2ND 2730
2ND 2740
2ND 2750
2ND 2760
2ND 2770
2ND 2780
2ND 2790
2ND 2800
2ND 2810
2ND 2820
2ND 2830
2ND 2840
2ND 2850
2ND 2860
2ND 2870
2ND 2880
7

```

```

DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
D11=(1.-DLU(M)**2)/XLP(M)
D22=(1.-DLV(M)**2)/XLP(M)
D33=(1.-DLW(M)**2)/XLP(M)
D12=-DLU(M)*DLV(M)/XLP(M)
D13=-DLU(M)*DLW(M)/XLP(M)
D23=-DLV(M)*DLW(M)/XLP(M)
D111=(-3.*DLU(M)*(1.-DLU(M)**2))/XLP(M)**2
D222=(-3.*DLV(M)*(1.-DLV(M)**2))/XLP(M)**2
D333=(-3.*DLW(M)*(1.-DLW(M)**2))/XLP(M)**2
D112=(-DLV(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
D113=(-DLW(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
D223=(-DLW(M)*(1.-3.*DLV(M)**2))/XLP(M)**2
D122=(-DLU(M)*(1.-3.*DLV(M)**2))/XLP(M)**2
D133=(-DLU(M)*(1.-3.*DLW(M)**2))/XLP(M)**2
D233=(-DLV(M)*(1.-3.*DLW(M)**2))/XLP(M)**2
D123=(3.*DLU(M)*DLV(M)*DLW(M))/XLP(M)**2
A(1,1)=(DL(M)*D111+3.*DLU(M)*D11)/XL(M)
A(1,2)=(DL(M)*D112+DLV(M)*D11+2.*DLU(M)*D12)/XL(M)
A(1,3)=(DL(M)*D113+DLW(M)*D11+2.*DLU(M)*D13)/XL(M)
A(2,2)=(DL(M)*D122+DLU(M)*D22+2.*DLV(M)*D12)/XL(M)
A(2,3)=(DL(M)*D123+DLW(M)*D12+DLU(M)*D23+DLV(M)*D13)/XL(M)
A(3,3)=(DL(M)*D133+2.*DLW(M)*D13+DLU(M)*D33)/XL(M)
A(2,1)=A(1,2)
A(3,1)=A(1,3)
A(3,2)=A(2,3)
B(1,1)=A(1,2)
B(1,2)=A(2,2)
B(1,3)=A(3,2)
2ND 2890
2ND 2900
2ND 2910
2ND 2920
2ND 2930
2ND 2940
2ND 2950
2ND 2960
2ND 2970
2ND 2980
2ND 2990
2ND 3000
2ND 3010
2ND 3020
2ND 3030
2ND 3040
2ND 3050
2ND 3060
2ND 3070
2ND 3080
2ND 3090
2ND 3100
2ND 3110
2ND 3120
2ND 3130
2ND 3140
2ND 3150
2ND 3160
2ND 3170
2ND 3180
2ND 3190
2ND 3200

```

```

B(2,2)=(DL(M)*D222+3.*DLV(M)*D22)/XL(M)
B(2,3)=(DL(M)*D223+DLW(M)*D22+2.*DIV(M)*D23)/XL(M)
B(3,3)=(DL(M)*D233+2.*DLW(M)*D23+DLV(M)*D33)/XL(M)
B(2,1)=8(1,2)
B(3,1)=8(1,3)
B(3,2)=8(2,3)
C(1,1)=A(1,3)
C(1,2)=A(2,3)
C(1,3)=A(3,3)
C(2,1)=C(1,2)
C(2,2)=B(2,3)
C(2,3)=8(3,3)
C(3,3)=(DL(M)*D333+3.*DLW(M)*D33)/XL(M)
C(3,1)=C(1,3)
C(3,2)=C(2,3)
DO 2075 J=1,3
DO 2075 K=1,3
S(M1,1)=S(M1,1) + RAL(J)*A(J,K)*RAL(K)
S(M2,1)=S(M2,1) + RAL(J)*B(J,K)*RAL(K)
S(M3,1)=S(M3,1) + RAL(J)*C(J,K)*RAL(K)
S(M1,2)=S(M1,2) + 3.*RAL(J)*A(J,K)*RA2(K)
S(M2,2)=S(M2,2) + 3.*RAL(J)*B(J,K)*RA2(K)
S(M3,2)=S(M3,2) + 3.*RAL(J)*C(J,K)*RA2(K)
S(M1,3)=S(M1,3)+3.*RA2(J)*A(J,K)*RA2(K)+4.*RAL(J)*A(J,K)*RA3(K)
S(M2,3)=S(M2,3)+3.*RA2(J)*B(J,K)*RA2(K)+4.*RAL(J)*B(J,K)*RA3(K)
S(M3,3)=S(M3,3)+3.*RA2(J)*C(J,K)*RA2(K)+4.*RAL(J)*C(J,K)*RA3(K)
S(M1,4)=S(M1,4)+5.*RAL(J)*A(J,K)*RA4(K)+10.*RA2(J)*A(J,K)*RA3(K)
S(M2,4)=S(M2,4)+5.*RAL(J)*B(J,K)*RA4(K)+10.*RA2(J)*B(J,K)*RA3(K)
S(M3,4)=S(M3,4)+5.*RAL(J)*C(J,K)*RA4(K)+10.*RA2(J)*C(J,K)*RA3(K)
2075 CONTINUE
2080 CONTINUE
RETURN
2ND 3210
2ND 3220
2ND 3230
2ND 3240
2ND 3250
2ND 3260
2ND 3270
2ND 3280
2ND 3290
2ND 3300
2ND 3310
2ND 3320
2ND 3330
2ND 3340
2ND 3350
2ND 3360
2ND 3370
2ND 3380
2ND 3390
2ND 3400
2ND 3410
2ND 3420
2ND 3430
2ND 3440
2ND 3450
2ND 3460
2ND 3470
2ND 3480
2ND 3490
2ND 3500
2ND 3510
2ND 3520

```

END

2ND 3530

```

SUBROUTINE THIRD(T,Q,N)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION A11(3,3),A12(3,3),A13(3,3),A22(3,3),A23(3,3),A33(3,3)
DIMENSION RA1(3),RB1(3),RA2(3),RB2(3),RA3(3),RB3(3)
DIMENSION T(N,3),Q(N)
DIMENSION DLU(42),DLV(42),DLW(42)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KCOUNT
DO 3000 I=1,N
DO 3000 J=1,3
T(I,J)=0.000
3000 CONTINUE
DO 3300 M=1,NM
J1=INP(M)
K1=INQ(M)
IF(J1.GT.NF.AND.K1.GT.NF) GO TO 3300
IF(J1.GT.NF) GO TO 3100
IF(K1.GT.NF) GO TO 3200
M1=3*J1-2
M2=3*J1-1
M3=3*J1
M4=3*K1-2
M5=3*K1-1
M6=3*K1
RA1(1)=Q1(M1)
RA1(2)=Q1(M2)
RA1(3)=Q1(M3)
RB1(1)=Q1(M4)
RB1(2)=Q1(M5)
3RD 0010
3RD 0020
3RD 0030
3RD 0040
3RD 0050
3RD 0060
3RD 0070
3RD 0080
3RD 0090
3RD 0100
3RD 0110
3RD 0120
3RD 0130
3RD 0140
3RD 0150
3RD 0160
3RD 0170
3RD 0180
3RD 0190
3RD 0200
3RD 0210
3RD 0220
3RD 0230
3RD 0240
3RD 0250
3RD 0260
3RD 0270
3RD 0280
3RD 0290
3RD 0300
3RD 0310
3RD 0320

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```

RB1(3)=Q1(M6)
RA2(1)=Q2(M1)
RA2(2)=Q2(M2)
RA2(3)=Q2(M3)
RB2(1)=Q2(M4)
RB2(2)=Q2(M5)
RB2(3)=Q2(M6)
RA3(1)=Q3(M1)
RA3(2)=Q3(M2)
RA3(3)=Q3(M3)
RB3(1)=Q3(M4)
RB3(2)=Q3(M5)
RB3(3)=Q3(M6)
DU(M)=Q(M4)-Q(M1)
DV(M)=Q(M5)-Q(M2)
DW(M)=Q(M6)-Q(M3)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
D11=(1.-DLU(M)**2)/XLP(M)
D22=(1.-DLV(M)**2)/XLP(M)
D33=(1.-DLW(M)**2)/XLP(M)
D12=-DLU(M)*DLV(M)/XLP(M)
D13=-DLU(M)*DLW(M)/XLP(M)
D23=-DLV(M)*DLW(M)/XLP(M)
D111=(-3.*DLU(M)*(1.-DLU(M)**2))/XLP(M)**2
D222=(-3.*DLV(M)*(1.-DLV(M)**2))/XLP(M)**2
3RD 0330
3RD 0340
3RD 0350
3RD 0360
3RD 0370
3RD 0380
3RD 0390
3RD 0400
3RD 0410
3RD 0420
3RD 0430
3RD 0440
3RD 0450
3RD 0460
3RD 0470
3RD 0480
3RD 0490
3RD 0500
3RD 0510
3RD 0520
3RD 0530
3RD 0540
3RD 0550
3RD 0560
3RD 0570
3RD 0580
3RD 0590
3RD 0600
3RD 0610
3RD 0620
3RD 0630
3RD 0640

```



```

D333=(-3.*DLW(M)*(1.-DLW(M)**2)/XLP(M)**2      3RD 0650
D112=(-DLV(M)*(1.-3.*DLU(M)**2)/XLP(M)**2      3RD 0660
D113=(-DLW(M)*(1.-3.*DLU(M)**2)/XLP(M)**2      3RD 0670
D223=(-DLW(M)*(1.-3.*DLV(M)**2)/XLP(M)**2      3RD 0680
D122=(-DLU(M)*(1.-3.*DLV(M)**2)/XLP(M)**2      3RD 0690
D133=(-DLU(M)*(1.-3.*DLW(M)**2)/XLP(M)**2      3RD 0700
D233=(-DLV(M)*(1.-3.*DLW(M)**2)/XLP(M)**2      3RD 0710
D123=(3.*DLU(M)*DLV(M)*DLW(M))/XLP(M)**2      3RD 0720
D111=(-3.+18.*DLU(M)**2-15.*DLU(M)**4)/XLP(M)**3  3RD 0730
D222=(-3.+18.*DLV(M)**2-15.*DLV(M)**4)/XLP(M)**3  3RD 0740
D333=(-3.+18.*DLW(M)**2-15.*DLW(M)**4)/XLP(M)**3  3RD 0750
D1112=(DLU(M)*DLV(M)*(9.-15.*DLU(M)**2))/XLP(M)**3  3RD 0760
D1113=(DLU(M)*DLW(M)*(9.-15.*DLU(M)**2))/XLP(M)**3  3RD 0770
D1222=(DLU(M)*DLV(M)*(9.-15.*DLV(M)**2))/XLP(M)**3  3RD 0780
D2223=(DLW(M)*DLV(M)*(9.-15.*DLV(M)**2))/XLP(M)**3  3RD 0790
D1333=(DLW(M)*DLU(M)*(9.-15.*DLW(M)**2))/XLP(M)**3  3RD 0800
D2333=(DLW(M)*DLV(M)*(9.-15.*DLW(M)**2))/XLP(M)**3  3RD 0810
D1122=(-1.+3.*(DLU(M)**2+DLV(M)**2)          3RD 0820
$ -15.*DLU(M)**2*DLV(M)**2)/XLP(M)**3      3RD 0830
D1133=(-1.+3.*(DLU(M)**2+DLW(M)**2)          3RD 0840
$ -15.*DLU(M)**2*DLW(M)**2)/XLP(M)**3      3RD 0850
D2233=(-1.+3.*(DLV(M)**2+DLW(M)**2)          3RD 0860
$ -15.*DLV(M)**2*DLW(M)**2)/XLP(M)**3      3RD 0870
D1123=DLV(M)*DLW(M)*(3.-15.*DLU(M)**2)/XLP(M)**3  3RD 0880
D1223=DLU(M)*DLW(M)*(3.-15.*DLV(M)**2)/XLP(M)**3  3RD 0890
D1233=DLU(M)*DLV(M)*(3.-15.*DLW(M)**2)/XLP(M)**3  3RD 0900
A11(1,1)=(DL(M)*D1111+4.*DLU(M)*D111+3.*D11**2)/XL(M)  3RD 0910
A11(1,2)=(DL(M)*D1112+DLV(M)*D111          3RD 0920
$ +3.*DLU(M)*D112+3.*D11*D12)/XL(M)      3RD 0930
A11(1,3)=(DL(M)*D1113+DLW(M)*D111          3RD 0940
$ +3.*DLU(M)*D113+3.*D11*D13)/XL(M)      3RD 0950
A11(2,2)=(DL(M)*D1122+2.*DLV(M)*D112+D11*D22      3RD 0960

```

```

$ +2.*DLU(M)*D122+2.*D12**2)/XL(M) 3RD 0970
A11(2,3)= (DL(M)*D1123+DLW(M)*D112+DLV(M)*D113 3RD 0980
$ +D23*D11+2.*DLU(M)*D123+2.*D13*D12)/XL(M) 3RD 0990
A11(3,3)= (DL(M)*D1133+2.*DLW(M)*D113+D11*D33 3RD 1000
$ +2.*DLU(M)*D113+2.*D13**2)/XL(M) 3RD 1010
A11(2,1)=A11(1,2) 3RD 1020
A11(3,1)=A11(1,3) 3RD 1030
A11(3,2)=A11(2,3) 3RD 1040
A12(1,1)=A11(1,2) 3RD 1050
A12(1,2)=A11(2,2) 3RD 1060
A12(1,3)=A11(3,2) 3RD 1070
A12(2,1)=A12(1,2) 3RD 1080
A12(2,2)=(DL(M)*D1222+DLU(M)*D222+3.*DLV(M)*D122+3.*D22*D12)/XL(M) 3RD 1090
A12(2,3)= (DL(M)*D1223+DLW(M)*D122+DLU(M)*D223 3RD 1100
$ +D22*D13+2.*D23*D12+2.*DLV(M)*D123)/XL(M) 3RD 1110
A12(3,2)=A12(2,3) 3RD 1120
A12(3,1)=A12(1,3) 3RD 1130
A12(3,3)= (DL(M)*D1233+DLV(M)*D133+DLU(M)*D233+D33*D12 3RD 1140
$ +2.*D23*D13+2.*DLW(M)*D123)/XL(M) 3RD 1150
A13(1,1)=A11(1,3) 3RD 1160
A13(1,2)=A11(2,3) 3RD 1170
A13(1,3)=A11(3,3) 3RD 1180
A13(2,1)=A13(1,2) 3RD 1190
A13(2,2)=A12(2,3) 3RD 1200
A13(2,3)=A12(3,3) 3RD 1210
A13(3,1)=A13(1,3) 3RD 1220
A13(3,2)=A13(2,3) 3RD 1230
A13(3,3)=(DL(M)*D1333+DLU(M)*D333+3.*DLW(M)*D133+3.*D22*D13)/XL(M) 3RD 1240
A22(1,1)=A12(1,2) 3RD 1250
A22(1,2)=A12(2,2) 3RD 1260
A22(1,3)=A12(3,2) 3RD 1270
A22(2,1)=A22(1,2) 3RD 1280

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A22(2,2) = (DL(M)*D2222+4.*DLV(M)*D222+3.*D22**2)/XL(M)
A22(2,3) = (DL(M)*D2223+DLW(M)*D222+3.*DLV(M)*D223
$   +3.*D22*D23)/XL(M)
A22(3,1) = A22(1,3)
A22(3,2) = A22(2,3)
A22(3,3) = (DL(M)*D2233+2.*DLV(M)*D233+D22*D33
$   +2.*DLW(M)*D223+2.*D23**2)/XL(M)
A23(1,1) = A12(1,3)
A23(1,2) = A12(2,3)
A23(1,3) = A12(3,3)
A23(2,1) = A23(1,2)
A23(2,2) = A22(2,3)
A23(2,3) = A22(3,3)
A23(3,1) = A23(1,3)
A23(3,2) = A23(2,3)
A23(3,3) = (DL(M)*D2333+DLV(M)*D333+3.*DLW(M)*D233+3.*D33*D23)/XL(M)
A33(1,1) = A13(1,3)
A33(1,2) = A13(2,3)
A33(1,3) = A13(3,3)
A33(2,1) = A33(1,2)
A33(2,2) = A23(2,3)
A33(2,3) = A23(3,3)
A33(3,1) = A33(1,3)
A33(3,2) = A33(2,3)
A33(3,3) = (DL(M)*D3333+4.*DLW(M)*D333+3.*D33**2)/XL(M)
DO 3050 I=1,3
DO 3050 J=1,3
T(M1,1) = T(M1,1) + ((Q1(M1)-Q1(M4))*A11(I,J)
$   + (Q1(M2)-Q1(M5))*A12(I,J) + (Q1(M3)-Q1(M6))*A13(I,J))
$   *(RA1(I)*RA1(J)-2.*RA1(I)*RA1(J)+RB1(I)*RB1(J))
T(M2,1) = T(M2,1) + ((Q1(M1)-Q1(M4))*A12(I,J)
$   + (Q1(M2)-Q1(M5))*A22(I,J) + (Q1(M3)-Q1(M6))*A23(I,J))

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3RD 1290

3RD 1300

3RD 1310

3RD 1320

3RD 1330

3RD 1340

3RD 1350

3RD 1360

3RD 1370

3RD 1380

3RD 1390

3RD 1400

3RD 1410

3RD 1420

3RD 1430

3RD 1440

3RD 1450

3RD 1460

3RD 1470

3RD 1480

3RD 1490

3RD 1500

3RD 1510

3RD 1520

3RD 1530

3RD 1540

3RD 1550

3RD 1560

3RD 1570

3RD 1580

3RD 1590

3RD 1600

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$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M3,1)=T(M3,1) + ((Q1(M1)-Q1(M4))*A13(I,J)
$ +(Q1(M2)-Q1(M5))*A23(I,J) + (Q1(M3)-Q1(M6))*A33(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M4,1)=T(M4,1) - ((Q1(M1)-Q1(M4))*A11(I,J)
$ +(Q1(M2)-Q1(M5))*A12(I,J) + (Q1(M3)-Q1(M6))*A13(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M5,1)=T(M5,1) - ((Q1(M1)-Q1(M4))*A12(I,J)
$ +(Q1(M2)-Q1(M5))*A22(I,J) + (Q1(M3)-Q1(M6))*A23(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M6,1)=T(M6,1) - ((Q1(M1)-Q1(M4))*A13(I,J)
$ +(Q1(M2)-Q1(M5))*A23(I,J) + (Q1(M3)-Q1(M6))*A33(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M1,2)=T(M1,2) + 4.*((Q1(M1)-Q1(M4))*A11(I,J)
$ +(Q1(M2)-Q1(M5))*A12(I,J) + (Q1(M3)-Q1(M6))*A13(I,J))
$ *(RA1(I)*RA2(J)-RA1(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J))
$ + 2.*((Q2(M1)-Q2(M4))*A11(I,J)+(Q2(M2)-Q2(M5))*A12(I,J)
$ +(Q2(M3)-Q2(M6))*A13(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M2,2)=T(M2,2) + 4.*((Q1(M1)-Q1(M4))*A12(I,J)
$ +(Q1(M2)-Q1(M5))*A22(I,J) + (Q1(M3)-Q1(M6))*A23(I,J))
$ *(RA1(I)*RA2(J)-RA1(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J))
$ + 2.*((Q2(M1)-Q2(M4))*A12(I,J)+(Q2(M2)-Q2(M5))*A22(I,J)
$ +(Q2(M3)-Q2(M6))*A23(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M3,2)=T(M3,2) + 4.*((Q1(M1)-Q1(M4))*A13(I,J)
$ +(Q1(M2)-Q1(M5))*A23(I,J) + (Q1(M3)-Q1(M6))*A33(I,J))
$ *(RA1(I)*RA2(J)-RA1(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J))
$ + 2.*((Q2(M1)-Q2(M4))*A13(I,J)+(Q2(M2)-Q2(M5))*A23(I,J)
$ +(Q2(M3)-Q2(M6))*A33(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M4,2)=T(M4,2) - 4.*((Q1(M1)-Q1(M4))*A11(I,J)

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3RD 1610
3RD 1620
3RD 1630
3RD 1640
3RD 1650
3RD 1660
3RD 1670
3RD 1680
3RD 1690
3RD 1700
3RD 1710
3RD 1720
3RD 1730
3RD 1740
3RD 1750
3RD 1760
3RD 1770
3RD 1780
3RD 1790
3RD 1800
3RD 1810
3RD 1820
3RD 1830
3RD 1840
3RD 1850
3RD 1860
3RD 1870
3RD 1880
3RD 1890
3RD 1900
3RD 1910
3RD 1920

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$ + (Q1(M2)-Q1(M5))*A12(I,J) + (Q1(M3)-Q1(M6))*A13(I,J)
$ *(RA1(I)*RA2(J)-RA1(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J))
$ - 2.*((Q2(M1)-Q2(M4))*A11(I,J)+(Q2(M2)-Q2(M5))*A12(I,J)
$ + (Q2(M3)-Q2(M6))*A13(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M5,2)=T(M5,2) - 4.*((Q1(M1)-Q1(M4))*A12(I,J)
$ + (Q1(M2)-Q1(M5))*A22(I,J) + (Q1(M3)-Q1(M6))*A23(I,J))
$ *(RA1(I)*RA2(J)-RA1(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J))
$ - 2.*((Q2(M1)-Q2(M4))*A12(I,J)+(Q2(M2)-Q2(M5))*A22(I,J)
$ + (Q2(M3)-Q2(M6))*A23(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
T(M6,2)=T(M6,2) - 4.*((Q1(M1)-Q1(M4))*A13(I,J)
$ + (Q1(M2)-Q1(M5))*A23(I,J) + (Q1(M3)-Q1(M6))*A33(I,J))
$ *(RA1(I)*RA2(J)-RA1(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J))
$ - 2.*((Q2(M1)-Q2(M4))*A13(I,J)+(Q2(M2)-Q2(M5))*A23(I,J)
$ + (Q2(M3)-Q2(M6))*A33(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
E1=
20./3.*((Q1(M1)-Q1(M4))*A11(I,J)
$ +(Q1(M2)-Q1(M5))*A12(I,J)+(Q1(M3)-Q1(M6))*A13(I,J))*
$ (RA1(I)*RA3(J)-RA1(I)*RB3(J)-RB1(I)*RA3(J)+RB1(I)*RB3(J))
E2= 10./3.*((Q3(M1)-Q3(M4))*A11(I,J)+(Q3(M2)-Q3(M5))*A12(I,J)
$ +(Q3(M3)-Q3(M6))*A13(I,J))
$ *(RA1(I)*RA1(J)-2.*RA1(I)*RB1(J)+RB1(I)*RB1(J))
E3= 10.*((Q2(M1)-Q2(M4))*A11(I,J)
$ +(Q2(M2)-Q2(M5))*A12(I,J)+(Q2(M3)-Q2(M6))*A13(I,J))*
$ (RA1(I)*RA2(J)-RA1(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J))
E4= 5.*((Q1(M1)-Q1(M4))*A11(I,J)
$ +(Q1(M2)-Q1(M5))*A12(I,J)+(Q1(M3)-Q1(M6))*A13(I,J))
$ *(RA2(I)*RA2(J)-2.*RA2(I)*RB2(J)+RB2(I)*RB2(J))
T(M1,3)=T(M1,3) + E1 + E2 + E3 + E4
F1=
20./3.*((Q1(M1)-Q1(M4))*A12(I,J)
$ +(Q1(M2)-Q1(M5))*A22(I,J)+(Q1(M3)-Q1(M6))*A23(I,J))*

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3RD 1930

3RD 1940

3RD 1950

3RD 1960

3RD 1970

3RD 1980

3RD 1990

3RD 2000

3RD 2010

3RD 2020

3RD 2030

3RD 2040

3RD 2050

3RD 2060

3RD 2070

3RD 2080

3RD 2090

3RD 2100

3RD 2110

3RD 2120

3RD 2130

3RD 2140

3RD 2150

3RD 2160

3RD 2170

3RD 2180

3RD 2190

3RD 2200

3RD 2210

3RD 2220

3RD 2230

3RD 2240

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$ (RAL(I)*RA3(J)-RAL(I)*RB3(J)-RB1(I)*RA3(J)+RB1(I)*RB3(J)) 3RD 2250
F2= 10./3.*((Q3(M1)-Q3(M4))*A12(I,J)+(Q3(M2)-Q3(M5))*A22(I,J) 3RD 2260
$ +(Q3(M3)-Q3(M6))*A23(I,J)) 3RD 2270
$ *(RAL(I)*RA1(J)-2.*RAL(I)*RB1(J)+RB1(I)*RB1(J)) 3RD 2280
F3= 10.*((Q2(M1)-Q2(M4))*A12(I,J) 3RD 2290
$ +(Q2(M2)-Q2(M5))*A22(I,J)+(Q2(M3)-Q2(M6))*A23(I,J))* 3RD 2300
$ (RAL(I)*RA2(J)-RAL(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J)) 3RD 2310
F4= 5.*((Q1(M1)-Q1(M4))*A12(I,J) 3RD 2320
$ +(Q1(M2)-Q1(M5))*A22(I,J)+(Q1(M3)-Q1(M6))*A23(I,J)) 3RD 2330
$ *(RA2(I)*RA2(J)-2.*RA2(I)*RB2(J)+RB2(I)*RB2(J)) 3RD 2340
T(M2,3)=T(M2,3) + F1 + F2 +F3 + F4 3RD 2350
G1= 20./3.*((Q1(M1)-Q1(M4))*A13(I,J) 3RD 2360
$ +(Q1(M2)-Q1(M5))*A13(I,J)+(Q1(M3)-Q1(M6))*A33(I,J))* 3RD 2370
$ (RAL(I)*RA3(J)-RAL(I)*RB3(J)-RB1(I)*RA3(J)+RB1(I)*RB3(J)) 3RD 2380
G2= 10./3.*((Q3(M1)-Q3(M4))*A13(I,J)+(Q3(M2)-Q3(M5))*A13(I,J) 3RD 2390
$ +(Q3(M3)-Q3(M6))*A33(I,J)) 3RD 2400
$ *(RAL(I)*RA1(J)-2.*RAL(I)*RB1(J)+RB1(I)*RB1(J)) 3RD 2410
G3= 10.*((Q2(M1)-Q2(M4))*A13(I,J) 3RD 2420
$ +(Q2(M2)-Q2(M5))*A13(I,J)+(Q2(M3)-Q2(M6))*A33(I,J))* 3RD 2430
$ (RAL(I)*RA2(J)-RAL(I)*RB2(J)-RB1(I)*RA2(J)+RB1(I)*RB2(J)) 3RD 2440
G4= 5.*((Q1(M1)-Q1(M4))*A13(I,J) 3RD 2450
$ +(Q1(M2)-Q1(M5))*A13(I,J)+(Q1(M3)-Q1(M6))*A33(I,J)) 3RD 2460
$ *(RA2(I)*RA2(J)-2.*RA2(I)*RB2(J)+RB2(I)*RB2(J)) 3RD 2470
T(M3,3)=T(M3,3) + G1 + G2 + G3 + G4 3RD 2480
T(M4,3)=T(M4,3) - E1 - E2 -E3 -E4 3RD 2490
T(M5,3)=T(M5,3) - F1 - F2 -F3 - F4 3RD 2500
T(M6,3)=T(M6,3) - G1 - G2 - G3 - G4 3RD 2510
3050 CONTINUE 3RD 2520
GO TO 3300 3RD 2530
3100 M4=3*K1-2 3RD 2540
M5=3*K1-1 3RD 2550
M6=3*K1 3RD 2560

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R81(1)=Q1(M4)
R81(2)=Q1(M5)
R81(3)=Q1(M6)
R82(1)=Q2(M4)
R82(2)=Q2(M5)
R82(3)=Q2(M6)
R83(1)=Q3(M4)
R83(2)=Q3(M5)
R83(3)=Q3(M6)
DU(M)=Q(M4)
DV(M)=Q(M5)
DW(M)=Q(M6)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
D11=(1.-DLU(M)**2)/XLP(M)
D22=(1.-DLV(M)**2)/XLP(M)
D33=(1.-DLW(M)**2)/XLP(M)
D12=-DLU(M)*DLV(M)/XLP(M)
D13=-DLU(M)*DLW(M)/XLP(M)
D23=-DLV(M)*DLW(M)/XLP(M)
D111=(-3.*DLU(M)*(1.-DLU(M)**2))/XLP(M)**2
D222=(-3.*DLV(M)*(1.-DLV(M)**2))/XLP(M)**2
D333=(-3.*DLW(M)*(1.-DLW(M)**2))/XLP(M)**2
D112=(-DLV(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
D113=(-DLW(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
D223=(-DLW(M)*(1.-3.*DLV(M)**2))/XLP(M)**2
3RD 2570
3RD 2580
3RD 2590
3RD 2600
3RD 2610
3RD 2620
3RD 2630
3RD 2640
3RD 2650
3RD 2660
3RD 2670
3RD 2680
3RD 2690
3RD 2700
3RD 2710
3RD 2720
3RD 2730
3RD 2740
3RD 2750
3RD 2760
3RD 2770
3RD 2780
3RD 2790
3RD 2800
3RD 2810
3RD 2820
3RD 2830
3RD 2840
3RD 2850
3RD 2860
3RD 2870
3RD 2880

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D122=(-DLU(M)*(1,-3.*DLV(M)**2))/XLP(M)**2
 D133=(-DLU(M)*(1,-3.*DLW(M)**2))/XLP(M)**2
 D233=(-DLV(M)**2)/XLP(M)**2
 D123=(3.*DLU(M)*DLV(M)*DLW(M))/XLP(M)**2
 D1111=(-3.+18.*DLU(M)**2-15.*DLU(M)**4)/XLP(M)**3
 D2222=(-3.+18.*DLV(M)**2-15.*DLV(M)**4)/XLP(M)**3
 D3333=(-3.+18.*DLW(M)**2-15.*DLW(M)**4)/XLP(M)**3
 D1112=(DLU(M)*DLV(M)*(9,-15.*DLU(M)**2))/XLP(M)**3
 D1113=(DLU(M)*DLW(M)*(9,-15.*DLU(M)**2))/XLP(M)**3
 D2223=(DLW(M)*DLV(M)*(9,-15.*DLV(M)**2))/XLP(M)**3
 D1333=(DLW(M)*DLU(M)*(9,-15.*DLW(M)**2))/XLP(M)**3
 D2333=(DLW(M)*DLV(M)*(9,-15.*DLW(M)**2))/XLP(M)**3
 D1122=(-1.+3.*(DLU(M)**2+DLV(M)**2)
 \$ -15.*DLU(M)**2*DLV(M)**2)/XLP(M)**3
 D1133=(-1.+3.*(DLU(M)**2+DLW(M)**2)
 \$ -15.*DLU(M)**2*DLW(M)**2)/XLP(M)**3
 D2233=(-1.+3.*(DLV(M)**2+DLW(M)**2)
 \$ -15.*DLV(M)**2*DLW(M)**2)/XLP(M)**3
 D1123=DLV(M)*DLW(M)*(3,-15.*DLU(M)**2)/XLP(M)**3
 D1223=DLU(M)*DLW(M)*(3,-15.*DLV(M)**2)/XLP(M)**3
 D1233=DLU(M)*DLV(M)*(3,-15.*DLW(M)**2)/XLP(M)**3
 A11(1,1)=(DL(M)*D1111+4.*DLU(M)*D111+3.*D11**2)/XL(M)
 A11(1,2)=(DL(M)*D1112+DLV(M)*D111
 \$ +3.*DLU(M)*D112+3.*D11*D12)/XL(M)
 A11(1,3)=(DL(M)*D1113+DLW(M)*D111
 \$ +3.*DLU(M)*D113+3.*D11*D13)/XL(M)
 A11(2,2)=(DL(M)*D1122+2.*DLV(M)*D112+D11*D22
 \$ +2.*DLU(M)*D122+2.*D12**2)/XL(M)
 A11(2,3)=(DL(M)*D1123+DLW(M)*D112+DLV(M)*D113
 \$ +D23*D11+2.*DLU(M)*D123+2.*D13*D12)/XL(M)
 A11(3,3)=(DL(M)*D1133+2.*DLW(M)*D113+D11*D33
 \$ +2.*DLU(M)*D113+2.*D13**2)/XL(M)

3RD 2890

3RD 2900

3RD 2910

3RD 2920

3RD 2930

3RD 2940

3RD 2950

3RD 2960

3RD 2970

3RD 2980

3RD 2990

3RD 3000

3RD 3010

3RD 3020

3RD 3030

3RD 3040

3RD 3050

3RD 3060

3RD 3070

3RD 3080

3RD 3090

3RD 3100

3RD 3110

3RD 3120

3RD 3130

3RD 3140

3RD 3150

3RD 3160

3RD 3170

3RD 3180

3RD 3190

3RD 3200

A11(2,1)=A11(1,2) 3RD 3210
 A11(3,1)=A11(1,3) 3RD 3220
 A11(3,2)=A11(2,3) 3RD 3230
 A12(1,1)=A11(1,2) 3RD 3240
 A12(1,2)=A11(2,2) 3RD 3250
 A12(1,3)=A11(3,2) 3RD 3260
 A12(2,1)=A12(1,2) 3RD 3270
 A12(2,2)=(DL(M)*D1222+DLU(M)*D222+3.*DLV(M)*D122+3.*D22*D12)/XL(M) 3RD 3280
 A12(2,3)=(DL(M)*D1223+DLW(M)*D122+DLU(M)*D223 3RD 3290
 \$ +D22*D13+2.*D23*D12+2.*DLV(M)*D123)/XL(M) 3RD 3300
 A12(3,2)=A12(2,3) 3RD 3310
 A12(3,1)=A12(1,3) 3RD 3320
 A12(3,3)=(DL(M)*D1233+DLV(M)*D133+DLU(M)*D233+D33*D12 3RD 3330
 \$ +2.*D23*D13+2.*DLW(M)*D123)/XL(M) 3RD 3340
 A13(1,1)=A11(1,3) 3RD 3350
 A13(1,2)=A11(2,3) 3RD 3360
 A13(1,3)=A11(3,3) 3RD 3370
 A13(2,1)=A13(1,2) 3RD 3380
 A13(2,2)=A12(2,3) 3RD 3390
 A13(2,3)=A12(3,3) 3RD 3400
 A13(3,1)=A13(1,3) 3RD 3410
 A13(3,2)=A13(2,3) 3RD 3420
 A13(3,3)=(DL(M)*D1333+DLU(M)*D333+3.*DLW(M)*D133+3.*D22*D13)/XL(M) 3RD 3430
 A22(1,1)=A12(1,2) 3RD 3440
 A22(1,2)=A12(2,2) 3RD 3450
 A22(1,3)=A12(3,2) 3RD 3460
 A22(2,1)=A22(1,2) 3RD 3470
 A22(2,2)=(DL(M)*D222+4.*DLV(M)*D222+3.*D22**2)/XL(M) 3RD 3480
 A22(2,3)=(DL(M)*D2223+DLW(M)*D222+3.*DLV(M)*D223 3RD 3490
 \$ +3.*D22*D23)/XL(M) 3RD 3500
 A22(3,1)=A22(1,3) 3RD 3510
 A22(3,2)=A22(2,3) 3RD 3520

```

A22(3,3) = (DL(M)*D2233+2.*DLV(M)*D233+D22*D33
$ +2.*DLW(M)*D223+2.*D23**2)/XL(M)
A23(1,1)=A12(1,3)
A23(1,2)=A12(2,3)
A23(1,3)=A12(3,3)
A23(2,1)=A23(1,2)
A23(2,2)=A22(2,3)
A23(2,3)=A22(3,3)
A23(3,1)=A23(1,3)
A23(3,2)=A23(2,3)
A23(3,3)=(DL(M)*D2333+DLV(M)*D333+3.*DLW(M)*D233+3.*D33*D23)/XL(M)
A33(1,1)=A13(1,3)
A33(1,2)=A13(2,3)
A33(1,3)=A13(3,3)
A33(2,1)=A33(1,2)
A33(2,2)=A23(2,3)
A33(2,3)=A23(3,3)
A33(3,1)=A33(1,3)
A33(3,2)=A33(2,3)
A33(3,3)=(DL(M)*D3333+4.*DLW(M)*D333+3.*D22**2)/XL(M)
DO 3150 I=1,3
DO 3150 J=1,3
T(M4,1)=T(M4,1) + (Q1(M4)*A11(I,J)+Q1(M5)*A12(I,J)
$ + Q1(M6)*A13(I,J))*RB1(I)*RB1(J)
T(M5,1)=T(M5,1) + (Q1(M4)*A12(I,J)+Q1(M5)*A22(I,J)
$ + Q1(M6)*A13(I,J))*RB1(I)*RB1(J)
T(M6,1)=T(M6,1) + (Q1(M4)*A13(I,J)+Q1(M5)*A23(I,J)
$ + Q1(M6)*A13(I,J))*RB1(I)*RB1(J)
T(M4,2)=T(M4,2)+4.*(Q1(M4)*A11(I,J)+Q1(M5)*A12(I,J)
$ + Q1(M6)*A13(I,J))*RB1(I)*RB2(J)
$ +2.*(Q2(M4)*A11(I,J)+Q2(M5)*A12(I,J)
$ + Q2(M6)*A13(I,J))*RB1(I)*RB1(J)

```

3RD 3530

3RD 3540

3RD 3550

3RD 3560

3RD 3570

3RD 3580

3RD 3590

3RD 3600

3RD 3610

3RD 3620

3RD 3630

3RD 3640

3RD 3650

3RD 3660

3RD 3670

3RD 3680

3RD 3690

3RD 3700

3RD 3710

3RD 3720

3RD 3730

3RD 3740

3RD 3750

3RD 3760

3RD 3770

3RD 3780

3RD 3790

3RD 3800

3RD 3810

3RD 3820

3RD 3830

3RD 3840

```

T(M5,2)=T(M5,2)+4.*(Q1(M4)*A12(I,J)+Q1(M5)*A22(I,J)
$ + Q1(M6)*A23(I,J))*RB1(I)*RB2(J)
$ +2.*(Q2(M4)*A12(I,J)+Q2(M5)*A22(I,J)
$ + Q2(M6)*A23(I,J))*RB1(I)*RB1(J)
T(M6,2)=T(M6,2)+4.*(Q1(M4)*A13(I,J)+Q1(M5)*A23(I,J)
$ + Q1(M6)*A33(I,J))*RB1(I)*RB2(J)
$ +2.*(Q2(M4)*A13(I,J)+Q2(M5)*A23(I,J)
$ + Q2(M6)*A33(I,J))*RB1(I)*RB1(J)
T(M4,3)=T(M4,3) + 20./3.*(Q1(M4)*A11(I,J)+Q1(M5)*A12(I,J)
$ + Q1(M6)*A13(I,J))*RB1(I)*RB3(J)
$ +10./3.*(Q3(M4)*A11(I,J)+Q3(M5)*A12(I,J)
$ + Q3(M6)*A13(I,J))*RB1(I)*RB1(J)
$ +10.*(Q2(M4)*A11(I,J)+Q2(M5)*A12(I,J)
$ +Q2(M6)*A13(I,J))*RB1(I)*RB2(J)
$+5.*(Q1(M4)*A11(I,J)+Q1(M5)*A12(I,J)+Q1(M6)*A13(I,J))
$ *RB2(I)*RB2(J)
T(M5,3)=T(M5,3) + 20./3.*(Q1(M4)*A12(I,J)+Q1(M5)*A22(I,J)
$ + Q1(M6)*A23(I,J))*RB1(I)*RB3(J)
$ +10./3.*(Q3(M4)*A12(I,J)+Q3(M5)*A22(I,J)
$ + Q3(M6)*A23(I,J))*RB1(I)*RB1(J)
$ +10.*(Q2(M4)*A12(I,J)+Q2(M5)*A22(I,J)
$ +Q2(M6)*A23(I,J))*RB1(I)*RB2(J)
$+5.*(Q1(M4)*A12(I,J)+Q1(M5)*A22(I,J)+Q1(M6)*A23(I,J))
$ *RB2(I)*RB2(J)
T(M6,3)=T(M6,3) + 20./3.*(Q1(M4)*A13(I,J)+Q1(M5)*A23(I,J)
$ + Q1(M6)*A33(I,J))*RB1(I)*RB3(J)
$ +10./3.*(Q3(M4)*A13(I,J)+Q3(M5)*A23(I,J)
$ + Q3(M6)*A33(I,J))*RB1(I)*RB1(J)
$ +10.*(Q2(M4)*A13(I,J)+Q2(M5)*A23(I,J)
$ +Q2(M6)*A33(I,J))*RB1(I)*RB2(J)
$+5.*(Q1(M4)*A13(I,J)+Q1(M5)*A23(I,J)+Q1(M6)*A33(I,J))
$ *RB2(I)*RB2(J)

```

```

3RD 3850
3RD 3860
3RD 3870
3RD 3880
3RD 3890
3RD 3900
3RD 3910
3RD 3920
3RD 3930
3RD 3940
3RD 3950
3RD 3960
3RD 3970
3RD 3980
3RD 3990
3RD 4000
3RD 4010
3RD 4020
3RD 4030
3RD 4040
3RD 4050
3RD 4060
3RD 4070
3RD 4080
3RD 4090
3RD 4100
3RD 4110
3RD 4120
3RD 4130
3RD 4140
3RD 4150
3RD 4160

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```

3150 CONTINUE
GO TO 3300
3200 M1=3*J1-2
M2=3*J1-1
M3=3*J1
RA1(1)=Q1(M1)
RA1(2)=Q1(M2)
RA1(3)=Q1(M3)
RA2(1)=Q2(M1)
RA2(2)=Q2(M2)
RA2(3)=Q2(M3)
RA3(1)=Q3(M1)
RA3(2)=Q3(M2)
RA3(3)=Q3(M3)
DU(M)=-Q(M1)
DV(M)=-Q(M2)
DW(M)=-Q(M3)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
DLU(M)=-SUMX(M)/XLP(M)
DLV(M)=-SUMY(M)/XLP(M)
DLW(M)=-SUMZ(M)/XLP(M)
D11=(1.-DLU(M)**2)/XLP(M)
D22=(1.-DLV(M)**2)/XLP(M)
D33=(1.-DLW(M)**2)/XLP(M)
D12=-DLU(M)*DLV(M)/XLP(M)
D13=-DLU(M)*DLW(M)/XLP(M)
D23=-DLV(M)*DLW(M)/XLP(M)
D111=(-3.*DLU(M)*(1.-DLU(M)**2))/XLP(M)**2
3RD 4170
3RD 4180
3RD 4190
3RD 4200
3RD 4210
3RD 4220
3RD 4230
3RD 4240
3RD 4250
3RD 4260
3RD 4270
3RD 4280
3RD 4290
3RD 4300
3RD 4310
3RD 4320
3RD 4330
3RD 4340
3RD 4350
3RD 4360
3RD 4370
3RD 4380
3RD 4390
3RD 4400
3RD 4410
3RD 4420
3RD 4430
3RD 4440
3RD 4450
3RD 4460
3RD 4470
3RD 4480

```

D222=(-3.*DLV(M)*(1.-DLV(M)**2))/XLP(M)**2
 D333=(-3.*DLW(M)*(1.-DLW(M)**2))/XLP(M)**2
 D112=(-DLV(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
 D113=(-DLW(M)*(1.-3.*DLU(M)**2))/XLP(M)**2
 D223=(-DLW(M)*(1.-3.*DLV(M)**2))/XLP(M)**2
 D122=(-DLU(M)*(1.-3.*DLV(M)**2))/XLP(M)**2
 D133=(-DLU(M)*(1.-3.*DLW(M)**2))/XLP(M)**2
 D233=(-DLV(M)*(1.-3.*DLW(M)**2))/XLP(M)**2
 D123=(3.*DLU(M)*DLV(M)*DLW(M))/XLP(M)**2
 D111=(-3.+18.*DLU(M)**2-15.*DLU(M)**4)/XLP(M)**3
 D222=(-3.+18.*DLV(M)**2-15.*DLV(M)**4)/XLP(M)**3
 D333=(-3.+18.*DLW(M)**2-15.*DLW(M)**4)/XLP(M)**3
 D112=(DLU(M)*DLV(M)*(9.-15.*DLU(M)**2))/XLP(M)**3
 D113=(DLU(M)*DLW(M)*(9.-15.*DLU(M)**2))/XLP(M)**3
 D223=(DLW(M)*DLV(M)*(9.-15.*DLV(M)**2))/XLP(M)**3
 D133=(DLW(M)*DLU(M)*(9.-15.*DLW(M)**2))/XLP(M)**3
 D233=(DLW(M)*DLV(M)*(9.-15.*DLW(M)**2))/XLP(M)**3
 D1122=(-1.+3.*(DLU(M)**2+DLV(M)**2)
 \$ -15.*DLU(M)**2*DLV(M)**2)/XLP(M)**3
 D1133=(-1.+3.*(DLU(M)**2+DLW(M)**2)
 \$ -15.*DLU(M)**2*DLW(M)**2)/XLP(M)**3
 D2233=(-1.+3.*(DLV(M)**2+DLW(M)**2)
 \$ -15.*DLV(M)**2*DLW(M)**2)/XLP(M)**3
 D1123=DLV(M)*DLW(M)*(3.-15.*DLU(M)**2)/XLP(M)**3
 D1223=DLU(M)*DLW(M)*(3.-15.*DLV(M)**2)/XLP(M)**3
 D1233=DLU(M)*DLV(M)*(3.-15.*DLW(M)**2)/XLP(M)**3
 A11(1,1)=(DL(M)*D1111+4.*DLU(M)*D111+3.*D11**2)/XL(M)
 A11(1,2)=(DL(M)*D1112+DLV(M)*D111
 \$ +3.*DLU(M)*D112+3.*D11*012)/XL(M)
 A11(1,3)=(DL(M)*D1113+DLW(M)*D111
 \$ +3.*DLU(M)*D113+3.*D11*013)/XL(M)
 A11(2,2)=(DL(M)*D1122+2.*DLV(M)*D112+D11*D22

3RD 4490
 3RD 4500
 3RD 4510
 3RD 4520
 3RD 4530
 3RD 4540
 3RD 4550
 3RD 4560
 3RD 4570
 3RD 4580
 3RD 4590
 3RD 4600
 3RD 4610
 3RD 4620
 3RD 4630
 3RD 4640
 3RD 4650
 3RD 4660
 3RD 4670
 3RD 4680
 3RD 4690
 3RD 4700
 3RD 4710
 3RD 4720
 3RD 4730
 3RD 4740
 3RD 4750
 3RD 4760
 3RD 4770
 3RD 4780
 3RD 4790
 3RD 4800

```

$ +2.*DLU(M)*D122+2.*D12**2)/XL(M) 3RD 4810
A11(2,3)= (DL(M)*D1123+DLW(M)*D112+DLV(M)*D113 3RD 4820
$ +D23*D11+2.*DLU(M)*D123+2.*D13*D12)/XL(M) 3RD 4830
A11(3,3)= (DL(M)*D1133+2.*DLW(M)*D113+D11*D33 3RD 4840
$ +2.*DLU(M)*D113+2.*D13**2)/XL(M) 3RD 4850
A11(2,1)=A11(1,2) 3RD 4860
A11(3,1)=A11(1,3) 3RD 4870
A11(3,2)=A11(2,3) 3RD 4880
A12(1,1)=A11(1,2) 3RD 4890
A12(1,2)=A11(2,2) 3RD 4900
A12(1,3)=A11(3,2) 3RD 4910
A12(2,1)=A12(1,2) 3RD 4920
A12(2,2)=(DL(M)*D1222+DLU(M)*D222+3.*DLV(M)*D122+3.*D22*D12)/XL(M) 3RD 4930
A12(2,3)= (DL(M)*D1223+DLW(M)*D122+DLU(M)*D223 3RD 4940
$ +D22*D13+2.*D23*D12+2.*DLV(M)*D123)/XL(M) 3RD 4950
A12(3,2)=A12(2,3) 3RD 4960
A12(3,1)=A12(1,3) 3RD 4970
A12(3,3)= (DL(M)*D1233+DLV(M)*D133+DLU(M)*D233+D33*D12 3RD 4990
$ +2.*D23*D13+2.*DLW(M)*D123)/XL(M) 3RD 4990
A13(1,1)=A11(1,3) 3RD 5000
A13(1,2)=A11(2,3) 3RD 5010
A13(1,3)=A11(3,3) 3RD 5020
A13(2,1)=A13(1,2) 3RD 5030
A13(2,2)=A12(2,3) 3RD 5040
A13(2,3)=A12(3,3) 3RD 5050
A13(3,1)=A13(1,3) 3RD 5060
A13(3,2)=A13(2,3) 3RD 5070
A13(3,3)=(DL(M)*D1333+DLU(M)*D333+3.*DLW(M)*D133+3.*D22*D13)/XL(M) 3RD 5080
A22(1,1)=A12(1,2) 3RD 5090
A22(1,2)=A12(2,2) 3RD 5100
A22(1,3)=A12(3,2) 3RD 5110
A22(2,1)=A22(1,2) 3RD 5120

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A22(2,2) = (DL(M)*D2222+4.*DLV(M)*D222+3.*D22**2)/XL(M)
A22(2,3) = (DL(M)*D2223+DLW(M)*D222+3.*DLV(M)*D223
$ +3.*D22*D23)/XL(M)
A22(3,1)=A22(1,3)
A22(3,2)=A22(2,3)
A22(3,3) = (DL(M)*D2233+2.*DLV(M)*D233+D22*D033
$ +2.*DLW(M)*D223+2.*D23**2)/XL(M)
A23(1,1)=A12(1,3)
A23(1,2)=A12(2,3)
A23(1,3)=A12(3,3)
A23(2,1)=A23(1,2)
A23(2,2)=A22(2,3)
A23(2,3)=A22(3,3)
A23(3,1)=A23(1,3)
A23(3,2)=A23(2,3)
A23(3,3)=(DL(M)*D2333+DLV(M)*D333+3.*DLW(M)*D233+3.*D33*D23)/XL(M)
A33(1,1)=A13(1,3)
A33(1,2)=A13(2,3)
A33(1,3)=A13(3,3)
A33(2,1)=A33(1,2)
A33(2,2)=A23(2,3)
A33(2,3)=A23(3,3)
A33(3,1)=A33(1,3)
A33(3,2)=A33(2,3)
A33(3,3)=(DL(M)*D3333+4.*DLW(M)*D333+3.*D22**2)/XL(M)
D0 3250 I=1,3
D0 3250 J=1,3
T(M1,1)=T(M1,1) + (Q1(M1)*A11(I,J)+Q1(M2)*A12(I,J)
$ + Q1(M3)*A13(I,J))*RA1(I)*RA1(J)
T(M2,1)=T(M2,1) + (Q1(M1)*A12(I,J)+Q1(M2)*A22(I,J)
$ + Q1(M3)*A23(I,J))*RA1(I)*RA1(J)
T(M3,1)=T(M3,1) + (Q1(M1)*A13(I,J)+Q1(M2)*A23(I,J)
3RD 5130
3RD 5140
3RD 5150
3RD 5160
3RD 5170
3RD 5180
3RD 5190
3RD 5200
3RD 5210
3RD 5220
3RD 5230
3RD 5240
3RD 5250
3RD 5260
3RD 5270
3RD 5280
3RD 5290
3RD 5300
3RD 5310
3RD 5320
3RD 5330
3RD 5340
3RD 5350
3RD 5360
3RD 5370
3RD 5380
3RD 5390
3RD 5400
3RD 5410
3RD 5420
3RD 5430
3RD 5440

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```

$ + Q1(M3)*A33(I,J)*RA1(I)*RA1(J) 3RD 5450
T(M1,2)=T(M1,2) + 4.*(Q1(M1)*A11(I,J)+Q1(M2)*A12(I,J) 3RD 5460
$ + Q1(M3)*A13(I,J)*RA1(I)*RA2(J) 3RD 5470
$+2.*(Q2(M1)*A11(I,J)+Q2(M2)*A12(I,J)+Q2(M3)*A13(I,J)) 3RD 5480
$ *RA1(I)*RA1(J) 3RD 5490
T(M2,2)=T(M2,2) + 4.*(Q1(M1)*A12(I,J)+Q1(M2)*A22(I,J) 3RD 5500
$ + Q1(M3)*A23(I,J)*RA1(I)*RA2(J) 3RD 5510
$+2.*(Q2(M1)*A12(I,J)+Q2(M2)*A22(I,J)+Q2(M3)*A23(I,J)) 3RD 5520
$ *RA1(I)*RA1(J) 3RD 5530
T(M3,2)=T(M3,2) + 4.*(Q1(M1)*A13(I,J)+Q1(M2)*A23(I,J) 3RD 5540
$ + Q1(M3)*A33(I,J)*RA1(I)*RA2(J) 3RD 5550
$+2.*(Q2(M1)*A13(I,J)+Q2(M2)*A23(I,J)+Q2(M3)*A33(I,J)) 3RD 5560
$ *RA1(I)*RA1(J) 3RD 5570
T(M1,3)=T(M1,3) + 20./3.*(Q1(M1)*A11(I,J)+Q1(M2)*A12(I,J) 3RD 5580
$ + Q1(M3)*A13(I,J)*RA1(I)*RA3(J) 3RD 5590
$ +10./3.*(Q3(M1)*A11(I,J)+Q3(M2)*A12(I,J)+Q3(M3)*A13(I,J)) 3RD 5600
$ *RA1(I)*RA1(J) 3RD 5610
$ +10.*(Q2(M1)*A11(I,J)+Q2(M2)*A12(I,J)+Q2(M3)*A13(I,J)) 3RD 5620
$ *RA1(I)*RA2(J) 3RD 5630
$ +5.*(Q1(M1)*A11(I,J)+Q1(M2)*A12(I,J)+Q1(M3)*A13(I,J)) 3RD 5640
$ *RA2(I)*RA2(J) 3RD 5650
T(M2,3)=T(M2,3) + 20./3.*(Q1(M1)*A12(I,J)+Q1(M2)*A22(I,J) 3RD 5660
$ + Q1(M3)*A23(I,J)*RA1(I)*RA3(J) 3RD 5670
$ +10./3.*(Q3(M1)*A12(I,J)+Q3(M2)*A22(I,J)+Q3(M3)*A23(I,J)) 3RD 5680
$ *RA1(I)*RA1(J) 3RD 5690
$ +10.*(Q2(M1)*A12(I,J)+Q2(M2)*A22(I,J)+Q2(M3)*A23(I,J)) 3RD 5700
$ *RA1(I)*RA2(J) 3RD 5710
$ +5.*(Q1(M1)*A12(I,J)+Q1(M2)*A22(I,J)+Q1(M3)*A23(I,J)) 3RD 5720
$ *RA2(I)*RA2(J) 3RD 5730
T(M3,3)=T(M3,3) + 20./3.*(Q1(M1)*A13(I,J)+Q1(M2)*A23(I,J) 3RD 5740
$ + Q1(M3)*A33(I,J)*RA1(I)*RA3(J) 3RD 5750
$ +10./3.*(Q3(M1)*A13(I,J)+Q3(M2)*A23(I,J)+Q3(M3)*A33(I,J)) 3RD 5760

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$ *RA1(I)*RA1(J)
$ +10.*(Q2(M1)*A13(I,J)+Q2(M2)*A23(I,J)+Q2(M3)*A23(I,J))
$ *RA1(I)*RA2(J)
$ +5.*(Q1(M1)*A13(I,J)+Q1(M2)*A23(I,J)+Q1(M3)*A33(I,J))
$ *RA2(I)*RA2(J)
3250 CONTINUE
3300 CONTINUE
RETURN
END
3RD 5770
3RD 5780
3RD 5790
3RD 5800
3RD 5810
3RD 5820
3RD 5830
3RD 5840
3RD 5850

```

```

SUBROUTINE TAYLOR(A,B,C,N,INS,IORDER,IW)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION A(N),B(N),GRAD(39),X(39)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DU(42),DV(42),DW(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
IF(IW.EQ.0) RETURN
IF(IORDER.EQ.1) WRITE(6,10)
IF(IORDER.EQ.2) WRITE(6,15)
10 FORMAT('1',///5X,'FIRST ORDER INCREMENTAL SOLUTION FROM PREVIOUS
$EQUILIBRIUM',/5X,'
$
$-----'//)
15 FORMAT('1',///5X,'SECOND ORDER INCREMENTAL SOLUTION FROM PREVIOUS
$EQUILIBRIUM',/5X,'
$
$-----'//)
IF(IORDER.EQ.3) WRITE(6,25)
25 FORMAT('1',///5X,'THIRD ORDER INCREMENTAL SOLUTION FROM PREVIOUS
$EQUILIBRIUM',/5X,'
$
$-----'//)
IF(INS.EQ.1) WRITE(6,20)
IF(INS.NE.1) WRITE(6,30)
20 FORMAT(17X,'BY DISPLACEMENT-INCREMENT APPROACH'/
$17X,'
$22X,'
$-----'//)
30 FORMAT(22X,'BY LOAD-INCREMENT APPROACH'/
$22X,'
$-----'//)
WRITE(6,40)
40 FORMAT(1X,'JOINT NO',5X,'X-DISPL',10X,'Y-DISPL',10X,'Z-DISPL',
$10X,'LOAD-X',11X,'LOAD-Y',11X,'LOAD-Z'//)
DO 90 I=1,NJ
I1=3*I-2
TLR 0010
TLR 0020
TLR 0030
TLR 0040
TLR 0050
TLR 0060
TLR 0070
TLR 0080
TLR 0090
TLR 0100
TLR 0110
TLR 0120
TLR 0130
TLR 0140
TLR 0150
TLR 0160
TLR 0170
TLR 0180
TLR 0190
TLR 0200
TLR 0210
TLR 0220
TLR 0230
TLR 0240
TLR 0250
TLR 0260
TLR 0270
TLR 0280
TLR 0290
TLR 0300
TLR 0310
TLR 0320

```

```

12=3*I-1
13=3*I
IF(I.GT.NF) GO TO 70
IF(I.EQ.1) WRITE(6,50) I,A(I1),A(I2),A(I3),C
IF(I.NE.1) WRITE(6,60) I,A(I1),A(I2),A(I3)
50 FORMAT(I5,3(6X,D12.5),35X,D12.5)
60 FORMAT(I5,3(6X,D12.5))
70 IF(I.LE.NF) GO TO 90
   WRITE(6,80) I
80 FORMAT(I5)
90 CONTINUE
C
   WRITE(6,85) SDET
85 FORMAT(//10X,'STABILITY DETERMINANT =',D12.5)
C
IF(IORDER.NE.1) RETURN
DO 100 I=2,NF
IF(PX(I).NE.0.000) PX(I)=PX(I)+C
IF(PY(I).NE.0.000) PY(I)=PY(I)+C
IF(PZ(I).NE.0.000) PZ(I)=PZ(I)+C
100 CONTINUE
IF(PX(1).NE.0.000) PX(1)=PX(1)+(C/FLOAD)
DO 200 I=1,N
X(I)=A(I)+B(I)
200 CONTINUE
CALL V(N,X,FUNT,GRAD)
CALL OUTP(X,GRAD,FUNT,N,IORDER,INS,IW)
DO 300 I=2,NF
IF(PX(I).NE.0.000) PX(I)=PX(I)-C
IF(PY(I).NE.0.000) PY(I)=PY(I)-C
IF(PZ(I).NE.0.000) PZ(I)=PZ(I)-C
300 CONTINUE
TLR 0330
TLR 0340
TLR 0350
TLR 0360
TLR 0370
TLR 0380
TLR 0390
TLR 0400
TLR 0410
TLR 0420
TLR 0430
TLR 0440
TLR 0450
TLR 0460
TLR 0470
TLR 0480
TLR 0490
TLR 0500
TLR 0510
TLR 0520
TLR 0530
TLR 0540
TLR 0550
TLR 0560
TLR 0570
TLR 0580
TLR 0590
TLR 0600
TLR 0610
TLR 0620
TLR 0630
TLR 0640

```

```
IF(PX(1) .NE. 0.000) PX(1)=PX(1)-(C/FLOAD)  
RETURN  
END
```

```
TLR 0650  
TLR 0660  
TLR 0670
```

```

SUBROUTINE OUTP(Q,GR,FUN,N,INS,INT,IW)
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION Q(N),GR(N)
COMMON PX(13),PY(13),PZ(13),DX(42),DY(42),DZ(42)
COMMON XL(42),XLP(42),DUP(42),DV(42),DH(42),DL(42)
COMMON SUMX(42),SUMY(42),SUMZ(42)
COMMON Q1(39),Q2(39),Q3(39),Q4(39),FLOAD,SDET
COMMON INP(42),INQ(42),NM,NJ,NF,KOUNT
IF(IW.EQ.0) RETURN
IF(INS.EQ.0.AND.INT.EQ.0) WRITE(6,400)
IF(INS.EQ.1.AND.INT.EQ.1) WRITE(6,500)
IF(INS.EQ.1.AND.INT.EQ.2) WRITE(6,600)
IF(INS.EQ.2.AND.INT.EQ.1) WRITE(6,700)
IF(INS.EQ.2.AND.INT.EQ.2) WRITE(6,800)
IF(INS.EQ.3.AND.INT.EQ.1) WRITE(6,900)
IF(INS.EQ.3.AND.INT.EQ.2) WRITE(6,1000)
400 FORMAT('1'///10X,'SOLUTION OBTAINED FROM MINIMIZATION VIA CONJUGATE')
$E GRADIENT METHOD//10X,'-----'//OUT 0180
$-----'//)
500 FORMAT('1'///10X,'FIRST ORDER PERTURBATION SOLUTION VIA DISPLACEMENT')
$NT INCREMENT//10X,'-----'//OUT 0200
$-----'//)
600 FORMAT('1'///10X,'FIRST ORDER PERTURBATION SOLUTION VIA LOAD INCREMENT')
$MENT//10X,'-----'//OUT 0220
$)
700 FORMAT('1'///10X,'SECOND ORDER PERTURBATION SOLUTION VIA DISPLACEMENT')
$ENT INCREMENT//10X,'-----'//OUT 0250
$-----'//)
800 FORMAT('1'///10X,'SECOND ORDER PERTURBATION SOLUTION VIA LOAD INCREMENT')
$EMENT//10X,'-----'//OUT 0280
$)
900 FORMAT('1'///10X,'THIRD ORDER PERTURBATION SOLUTION VIA DISPLACEMENT')
$)

```

```

$NT INCREMENT',/10X,'-----OUT 0330
$-----'///)OUT 0340
1000 FORMAT('1'///10X,'THIRD ORDER PERTURBATION SOLUTION VIA LOAD INCREDEUT 0350
$MENT',/10X,'-----'///)OUT 0360
$)OUT 0370
WRITE(6,10)OUT 0380
10 FORMAT(/1X,'JOINT NO',5X,'X-DISPL',10X,'Y-DISPL',10X,'Z-DISPL',OUT 0390
$10X,'LOAD-X',11X,'LOAD-Y',11X,'LOAD-Z'///)OUT 0400
DO 100 J=1,NJOUT 0410
J1=3*J-2OUT 0420
J2=3*J-1OUT 0430
J3=3*JOUT 0440
IF(J.LE.NF) WRITE(6,20) J,Q(J1),Q(J2),Q(J3),PX(J),PY(J),PZ(J)OUT 0450
IF(J.GT.NF) WRITE(6,30) JOUT 0460
20 FORMAT(15,6(6X,D12.5))OUT 0470
30 FORMAT(15,7X,'0.0',5(15X,'0.0'))OUT 0480
100 CONTINUEOUT 0490
WRITE(6,40)OUT 0500
40 FORMAT(/1X,'JOINT NO',5X,'GRAD-X',12X,'GRAD-Y',12X,'GRAD-Z'///)OUT 0510
DO 300 J=1,NJOUT 0520
J1=3*J-2OUT 0530
J2=3*J-1OUT 0540
J3=3*JOUT 0550
IF(J.LE.NF) WRITE(6,50) J,GR(J1),GR(J2),GR(J3)OUT 0560
IF(J.GT.NF) WRITE(6,60) JOUT 0570
50 FORMAT(15,3(6X,D12.5))OUT 0580
60 FORMAT(15,7X,'0.0',2(15X,'0.0'))OUT 0590
300 CONTINUEOUT 0600
WRITE(6,350) FUNOUT 0610
350 FORMAT(/1X,'FUNCTIONAL VALUE = ',D12.5/10X,'-----'///)OUT 0620
$)OUT 0630
RETURNOUT 0640

```

END

OUT 0650

```

C===== DATA FOR UNIFORM LOADING =====
C
42 19
1 0.0 0.0 0.0
2 2.0 0.0 25.0
3 2.0 21.65 12.5
4 2.0 21.65 -12.5
5 2.0 0.0 -25.0
6 2.0 -21.65 -12.5
7 2.0 -21.65 12.5
8 8.1609 25.0 43.3
9 8.1609 50.0 0.0
10 8.1609 25.0 -43.3
11 8.1609 -25.0 -43.3
12 8.1609 -50.0 0.0
13 8.1609 -25.0 43.3
14 10.87 0.0 50.0
15 10.87 50.0 25.0
16 10.87 50.0 -25.0
17 10.87 0.0 -50.0
18 10.87 -50.0 -25.0
19 10.87 -50.0 25.0
1 1 2 1
2 1 3 1
3 1 4 1
4 1 5 1
5 1 6 1
6 1 7 1
7 7 6 7
8 2 7 2
9 3 2 3
INP 0010
INP 0020
INP 0030
INP 0040
INP 0050
INP 0060
INP 0070
INP 0080
INP 0090
INP 0100
INP 0110
INP 0120
INP 0130
INP 0140
INP 0150
INP 0160
INP 0170
INP 0180
INP 0190
INP 0200
INP 0210
INP 0220
INP 0230
INP 0240
INP 0250
INP 0260
INP 0270
INP 0280
INP 0290
INP 0300
INP 0310
INP 0320

```


INP 0330
 INP 0340
 INP 0350
 INP 0360
 INP 0370
 INP 0380
 INP 0390
 INP 0400
 INP 0410
 INP 0420
 INP 0430
 INP 0440
 INP 0450
 INP 0460
 INP 0470
 INP 0480
 INP 0490
 INP 0500
 INP 0510
 INP 0520
 INP 0530
 INP 0540
 INP 0550
 INP 0560
 INP 0570
 INP 0580
 INP 0590
 INP 0600
 INP 0610
 INP 0620
 INP 0630
 INP 0640

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 37
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 39
 40
 41

42	10	17
1	0.00001	
2	0.00001	
3	0.00001	
4	0.00001	
5	0.00001	
6	0.00001	
7	0.00001	
1.0		
DISP		LOAD

INP 0650
INP 0660
INP 0670
INP 0680
INP 0690
INP 0700
INP 0710
INP 0720
INP 0730
INP 0740


```

C
C
C
      INPUT DATA FOR MEMBER NUMBERS, NM AND NUMBER OF JOINTS, NJ
      READ(5,20) NM,NJ
      20 FORMAT(2I5)
      WRITE(6,30) NM,NJ
      30 FORMAT(//10X,'NUMBER OF ELEMENTS, NM =',I3//10X,'NUMBER OF JOINTS,
      $ NJ =',I3)
C
C
      JOINT COORDINATES AND CONSTRAINT OF THE JOINT
      ICON(J):      MEANS CONSTRAINT AT JOINT 'J'
      ITS VALUES ARE:      ICON(J)=1 -COMPLETE CONSTRAINED OR FIXED
                           ICON(J)=0 -COMPLETELY FREE
C
C
      NF= NUMBER OF UNCONSTRAINED JOINTS
      NF=NJ
      WRITE(6,40)
      40 FORMAT(/////20X,'JOINT COORDINATES AND CONSTRAINTS'//////11X,'JOINT
      $ NO',5X,'X-COORD',10X,'Y-COORD',10X,'Z-COORD',8X,'CONSTRAINT'//)
      DO 70 J=1,NJ
      READ(5,50) K,X(K),Y(K),Z(K),ICON(K)
      50 FORMAT(I5,3F15.6,4X,I1)
      WRITE(6,60) K,X(K),Y(K),Z(K),ICON(K)
      60 FORMAT(10X,I5,6X,D12.5,6X,D12.5,5X,D12.5,10X,I1)
      NF=NF-ICON(K)
      70 CONTINUE
C
C
      MEMBER INCIDENCE TABLE
      WRITE(6,80)

```

JOY 0330

JOY 0340

JOY 0350

JOY 0360

JOY 0370

JOY 0380

JOY 0390

JOY 0400

JOY 0410

JOY 0420

JOY 0430

JOY 0440

JOY 0450

JOY 0460

JOY 0470

JOY 0480

JOY 0490

JOY 0500

JOY 0510

JOY 0520

JOY 0530

JOY 0540

JOY 0550

JOY 0560

JOY 0570

JOY 0580

JOY 0590

JOY 0600

JOY 0610

JOY 0620

JOY 0630

JOY 0640

```

80 FORMAT('1',///10X, 'MEMBER INCIDENCE TABLE AND ITS LENGTH'/////5X, JOY 0650
$, 'MEMBER NO', 10X, 'MEMBER INCIDENCE', 13X, 'MEMBER LENGTH'//)
DO 110 M=1, NM
READ(5, 90) J, INP(J), INQ(J)
90 FORMAT(3I5)
I=INP(J)
K=INQ(J)
DX(M)=X(K)-X(I)
DY(M)=Y(K)-Y(I)
DZ(M)=Z(K)-Z(I)
XL(M)=DSQRT(DX(M)**2+DY(M)**2+DZ(M)**2)
WRITE(6, 100) J, INP(J), INQ(J), XL(J)
100 FORMAT(4X, I6, 15X, 2I5, 15X, D15.5)
110 CONTINUE
C
C NUMBER OF DEGREE OF FREEDOM IS EQUAL TO THREE
C TIMES NUMBER OF UNCONSTRAINED JOINTS, N=3*NF
C
N=3*NF
WRITE(6, 120) NF, N
120 FORMAT(///5X, 'NUMBER OF FREE JOINTS, NF =', I3//5X, 'NUMBER OF DEJOY 0850
$GREE OF FREEDOM, N =', I3)
C
C SOLUTION FOR THE EQUILIBRIUM PATH VIA MINIMIZATION TECHNIQUE
C USING POWELL-FLECHER METHOD, THIS SOLUTION WILL BE USED AS
C A STARTING POINT FOR THE PERTURBATION TECHNIQUE TO DESCRIBE
C THE LOAD VS DISPLACEMENT FOR THE ENTIRE PATH
C
WRITE(6, 130)
130 FORMAT('1',///10X, 'MINIMIZATION VIA CONJUGATE GRADIENT METHOD'////)
C
C INITIAL GUESS OF VARIABLES
JOY 0660
JOY 0670
JOY 0680
JOY 0690
JOY 0700
JOY 0710
JOY 0720
JOY 0730
JOY 0740
JOY 0750
JOY 0760
JOY 0770
JOY 0780
JOY 0790
JOY 0800
JOY 0810
JOY 0820
JOY 0830
JOY 0840
JOY 0860
JOY 0870
JOY 0880
JOY 0890
JOY 0900
JOY 0910
JOY 0920
JOY 0930
JOY 0940
JOY 0950
JOY 0960

```

```

C
WRITE(6,140)
140 FORMAT('//20X, 'INITIAL GUESS OF VARIABLES'///11X, 'JOINT', 8X, 'X-DISP' JOY 0970
5., '10X, 'Y-DISP.', 10X, 'Z-DISP.', '///) JOY 0980
DO 160 I=1, NF JOY 0990
I1=3*I-2 JOY 1000
I2=3*I-1 JOY 1010
I3=3*I JOY 1020
U(I1)=0.000 JOY 1030
U(I2)=0.000 JOY 1040
U(I3)=0.000 JOY 1050
WRITE(6,150) I, U(I1), U(I2), U(I3) JOY 1060
150 FORMAT(10X, I5, 9X, D12.5, 6X, D12.5, 6X, D12.5) JOY 1070
160 CONTINUE JOY 1080
C
C INITIALIZED ALL JOINT LOADS TO BE ZERO JOY 1090
C
DO 170 I=1, NF JOY 1100
PX(I)=0.000 JOY 1110
PY(I)=0.000 JOY 1120
PZ(I)=0.000 JOY 1130
170 CONTINUE JOY 1140
C
C THE ERROR INDICATOR, IER = 9 IS USED AS THE INITIAL INPUT JOY 1150
C THE RETURN VALUE ARE SUCH THAT : JOY 1160
C IER = 0..... MEANS CONVERGENCE IS ACHIEVED WITHIN SPECIFIED ACCURACY JOY 1220
C IER = 1..... MEANS NO CONVERGENCE IN THE LIMIT ITERATION JOY 1230
C IER =-1..... MEANS ERRORS IN THE GRADIENT CALCULATION JOY 1240
C IER = 2..... MEANS IT IS LIKELY THAT THERE EXISTS NO MINIMUM JOY 1250
C IER=9 JOY 1260
ACC=1.0D-10 JOY 1270
JOY 1280

```

```

LIMIT=100
FS=0.0
KOUNT=0
C
C INPUT LOADING AT THE JOINT
C
DO 305 KQ=1,NF
READ(5,180) K,PX(K)
180 FORMAT(I5,E15.5)
305 CONTINUE
CALL DFMCQ(V,N,U,F,G,FS,ACC,LIMIT,IER,H)
WRITE(6,186) LIMIT,KOUNT
186 FORMAT(///5X,'NUMBER OF ITERATION =',I4//
$ 5X,I5,5X,'FUNCTION SUBROUTINE CALLED'/////))
IW=1
CALL OUTP(U,G,F,N,0,0,IW)
KOUNT=0
CALL DYN
STOP
END
JOY 1290
JOY 1300
JOY 1310
JOY 1320
JOY 1330
JOY 1340
JOY 1350
JOY 1360
JOY 1370
JOY 1380
JOY 1390
JOY 1400
JOY 1410
JOY 1420
JOY 1430
JOY 1440
JOY 1450
JOY 1460
JOY 1470
JOY 1480

```

```

C
C
SUBROUTINE V(N,A,VAL,GRAD)
VEE 0010
VEE 0020
VEE 0030
VEE 0040
VEE 0050
VEE 0060
VEE 0070
VEE 0080
VEE 0090
VEE 0100
VEE 0110
VEE 0120
VEE 0130
VEE 0140
VEE 0150
VEE 0160
VEE 0170
VEE 0180
VEE 0190
VEE 0200
VEE 0210
VEE 0220
VEE 0230
VEE 0240
VEE 0250
VEE 0260
VEE 0270
VEE 0280
VEE 0290
VEE 0300
VEE 0310
VEE 0320

IMPLICIT REAL*8(A-H,O-Z)
DIMENSION A(21),GRAD(21),DE(21)
COMMON U(21),XM(21),C(21)
COMMON PX(7),PY(7),PZ(7),DX(24),DY(24),DZ(24)
COMMON XL(24),XLP(24),DUP(24),DV(24),DW(24),DL(24)
COMMON SUMX(24),SUMY(24),SUMZ(24)
COMMON INP(24),INQ(24),NM,NJ,NF,KOUNT
KOUNT=KOUNT+1
WORK=0.000
DO 6000 J=1,NF
  J1=3*J-2
  J2=3*J-1
  J3=3*J
  WORK=WORK+PX(J)*A(J1)+PY(J)*A(J2)+PZ(J)*A(J3)
DE(J1)=0.000
DE(J2)=0.000
DE(J3)=0.000
6000 CONTINUE
STENG=0.000
DO 6001 M=1,NM
  J=INP(M)
  K=INQ(M)
  IF(J.GT.NF.AND.K.GT.NF) GO TO 6001
  M1=3*J-2
  M2=3*J-1
  M3=3*J
  M4=3*K-2
  M5=3*K-1
  M6=3*K

```



```

IF(J.GT.NF) GO TO 700
IF(K.GT.NF) GO TO 750
GO TO 760
700 DU(M)=A(M4)
DV(M)=A(M5)
DW(M)=A(M6)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
STENG=STENG+(0.5*DL(M)**2)/XL(M)
DE(M4)=DE(M4)+(DL(M)*SUMX(M))/(XL(M)*XLP(M))
DE(M5)=DE(M5)+(DL(M)*SUMY(M))/(XL(M)*XLP(M))
DE(M6)=DE(M6)+(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
GO TO 6001
750 DU(M)=-A(M1)
DV(M)=-A(M2)
DW(M)=-A(M3)
SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
STENG=STENG+(0.5*DL(M)**2)/XL(M)
DE(M1)=DE(M1)-(DL(M)*SUMX(M))/(XL(M)*XLP(M))
DE(M2)=DE(M2)-(DL(M)*SUMY(M))/(XL(M)*XLP(M))
DE(M3)=DE(M3)-(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
GO TO 6001
760 DU(M)=A(M4)-A(M1)
DV(M)=A(M5)-A(M2)
DW(M)=A(M6)-A(M3)
VEE 0330
VEE 0340
VEE 0350
VEE 0360
VEE 0370
VEE 0380
VEE 0390
VEE 0400
VEE 0410
VEE 0420
VEE 0430
VEE 0440
VEE 0450
VEE 0460
VEE 0470
VEE 0480
VEE 0490
VEE 0500
VEE 0510
VEE 0520
VEE 0530
VEE 0540
VEE 0550
VEE 0560
VEE 0570
VEE 0580
VEE 0590
VEE 0600
VEE 0610
VEE 0620
VEE 0630
VEE 0640

```

```

SUMX(M)=DX(M)+DU(M)
SUMY(M)=DY(M)+DV(M)
SUMZ(M)=DZ(M)+DW(M)
XLP(M)=DSQRT(SUMX(M)**2+SUMY(M)**2+SUMZ(M)**2)
DL(M)=XLP(M)-XL(M)
STENG=STENG+(0.5*DL(M)**2)/XL(M)
DE(M1)=DE(M1)-(DL(M)*SUMX(M))/(XL(M)*XLP(M))
DE(M2)=DE(M2)-(DL(M)*SUMY(M))/(XL(M)*XLP(M))
DE(M3)=DE(M3)-(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
DE(M4)=DE(M4)+(DL(M)*SUMX(M))/(XL(M)*XLP(M))
DE(M5)=DE(M5)+(DL(M)*SUMY(M))/(XL(M)*XLP(M))
DE(M6)=DE(M6)+(DL(M)*SUMZ(M))/(XL(M)*XLP(M))
6001 CONTINUE
VAL=STENG-WORK
DO 6002 J=1,NF
J1=3*J-2
J2=3*J-1
J3=3*J
GRAD(J1)=DE(J1)-PX(J)
GRAD(J2)=DE(J2)-PY(J)
GRAD(J3)=DE(J3)-PZ(J)
6002 CONTINUE
RETURN
END
VEE 0650
VEE 0660
VEE 0670
VEE 0680
VEE 0690
VEE 0700
VEE 0710
VEE 0720
VEE 0730
VEE 0740
VEE 0750
VEE 0760
VEE 0770
VEE 0780
VEE 0790
VEE 0800
VEE 0810
VEE 0820
VEE 0830
VEE 0840
VEE 0850
VEE 0860
VEE 0870
VEE 0880

```

```

SUBROUTINE OUTP(Q,GR,FUN,N,INS,INT,IW)
C
C
C
C
      THIS SUBROUTINE GIVES THE OUT-PUT
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION Q(N),GR(N)
      COMMON U(21),XM(21),C(21)
      COMMON PX(7),PY(7),PZ(7),DX(24),DY(24),DZ(24)
      COMMON XL(24),XLP(24),DU(24),DV(24),DW(24),DL(24)
      COMMON SUMX(24),SUMY(24),SUMZ(24)
      COMMON INP(24),INQ(24),NM,NJ,NF,KOUNT
      IF(IW.EQ.0) RETURN
      IF(INS.EQ.0.AND.INT.EQ.0) WRITE(6,400)
      IF(INS.EQ.1.AND.INT.EQ.1) WRITE(6,500)
      IF(INS.EQ.1.AND.INT.EQ.2) WRITE(6,600)
      IF(INS.EQ.2.AND.INT.EQ.1) WRITE(6,700)
      IF(INS.EQ.2.AND.INT.EQ.2) WRITE(6,800)
      IF(INS.EQ.3.AND.INT.EQ.1) WRITE(6,900)
      IF(INS.EQ.3.AND.INT.EQ.2) WRITE(6,1000)
      IF(INS.EQ.4.AND.INT.EQ.1) WRITE(6,1100)
      IF(INS.EQ.4.AND.INT.EQ.2) WRITE(6,1200)
      IF(INS.EQ.5.AND.INT.EQ.1) WRITE(6,1300)
      IF(INS.EQ.5.AND.INT.EQ.2) WRITE(6,1400)
      IF(INS.EQ.6.AND.INT.EQ.1) WRITE(6,1500)
      IF(INS.EQ.6.AND.INT.EQ.2) WRITE(6,1600)
      +00 FORMAT('1'//10X,'SOLUTION OBTAINED FROM MINIMIZATION VIA CONJUGATE')
      $E GRADIENT METHOD'//10X,'-----'
      $-----'//)
      500 FORMAT('1'//10X,'FIRST ORDER PERTURBATION SOLUTION VIA DISPLACEMENT')
      $NT INCREMENT'//10X,'-----'
      $-----'//)
      OUTP0010
      OUTP0020
      OUTP0030
      OUTP0040
      OUTP0050
      OUTP0060
      OUTP0070
      OUTP0080
      OUTP0090
      OUTP0100
      OUTP0110
      OUTP0120
      OUTP0130
      OUTP0140
      OUTP0150
      OUTP0160
      OUTP0170
      OUTP0180
      OUTP0190
      OUTP0200
      OUTP0210
      OUTP0220
      OUTP0230
      OUTP0240
      OUTP0250
      OUTP0260
      OUTP0270
      OUTP0280
      OUTP0290
      OUTP0300
      OUTP0310
      OUTP0320

```

```

600 FORMAT('1'///10X,'FIRST ORDER PERTURBATION SOLUTION VIA LOAD INCREOUTP0330
$MENT'//10X,'-----'//OUTP0340
$)
      OUTP0350
700 FORMAT('1'///10X,'SECOND ORDER PERTURBATION SOLUTION VIA DISPLACEMENTOUTP0360
$ENT INCREMENT'//10X,'-----'OUTP0370
$-----'///)
      OUTP0380
800 FORMAT('1'///10X,'SECOND ORDER PERTURBATION SOLUTION VIA LOAD INCREOUTP0390
$EMENT'//10X,'-----'OUTP0400
$//)
      OUTP0410
900 FORMAT('1'///10X,'THIRD ORDER PERTURBATION SOLUTION VIA DISPLACEMENTOUTP0420
$NT INCREMENT'//10X,'-----'OUTP0430
$-----'///)
      OUTP0440
1000 FORMAT('1'///10X,'THIRD ORDER PERTURBATION SOLUTION VIA LOAD INCREOUTP0450
$MENT'//10X,'-----'//OUTP0460
$)
      OUTP0470
1100 FORMAT('1'///10X,'FOURTH ORDER PERTURBATION SOLUTION VIA DISPLACEMENTOUTP0480
$ENT INCREMENT'//10X,'-----'OUTP0490
$-----'///)
      OUTP0500
1200 FORMAT('1'///10X,'FOURTH ORDER PERTURBATION SOLUTION VIA LOAD INCREOUTP0510
EMENT'//10X,'-----'//OUTP0520
$//)
      OUTP0530
1300 FORMAT('1'///10X,'FIFTH ORDER PERTURBATION SOLUTION VIA DISPLACEMENTOUTP0540
$NT INCREMENT'//10X,'-----'OUTP0550
$-----'///)
      OUTP0560
1400 FORMAT('1'///10X,'FIFTH ORDER PERTURBATION SOLUTION VIA LOAD INCREOUTP0570
$MENT'//10X,'-----'//OUTP0580
$)
      OUTP0590
1500 FORMAT('1'///10X,'SIXTH ORDER PERTURBATION SOLUTION VIA DISPLACEMENTOUTP0600
$NT INCREMENT'//10X,'-----'OUTP0610
$-----'///)
      OUTP0620
1600 FORMAT('1'///10X,'SIXTH ORDER PERTURBATION SOLUTION VIA LOAD INCREOUTP0630
$MENT'//10X,'-----'//OUTP0640

```

```

$)
WRITE(6,10)
10 FORMAT(//1X,'JOINT NO',5X,'X-DISPL',10X,'Y-DISPL',10X,'Z-DISPL',
$10X,'LOAD-X',11X,'LOAD-Y',11X,'LOAD-Z'//)
DO 100 J=1,NJ
J1=3#J-2
J2=3#J-1
J3=3#J
IF(J.LE.NF) WRITE(6,20) J,Q(J1),Q(J2),Q(J3),PX(J),PY(J),PZ(J)
IF(J.GT.NF) WRITE(6,30) J
20 FORMAT(15,6(6X,D12.5))
30 FORMAT(15,7X,'0.0',5(15X,'0.0'))
100 CONTINUE
WRITE(6,40)
40 FORMAT(//1X,'JOINT NO',5X,'GRAD-X',12X,'GRAD-Y',12X,'GRAD-Z'//)
DO 300 J=1,NJ
J1=3#J-2
J2=3#J-1
J3=3#J
IF(J.LE.NF) WRITE(6,50) J,GR(J1),GR(J2),GR(J3)
IF(J.GT.NF) WRITE(6,60) J
50 FORMAT(15,3(6X,D12.5))
60 FORMAT(15,7X,'0.0',2(15X,'0.0'))
300 CONTINUE
WRITE(6,350) FUN
350 FORMAT(//10X,'FUNCTIONAL VALUE = ',D12.5/10X,'-----'//)
$)
RETURN
END

```

OUTP0650

OUTP0660

OUTP0670

OUTP0680

OUTP0690

OUTP0700

OUTP0710

OUTP0720

OUTP0730

OUTP0740

OUTP0750

OUTP0760

OUTP0770

OUTP0780

OUTP0790

OUTP0800

OUTP0810

OUTP0820

OUTP0830

OUTP0840

OUTP0850

OUTP0860

OUTP0870

OUTP0880

OUTP0890

OUTP0900

OUTP0910

OUTP0920

OUTP0930

```

C
C
SURROUTINE DYN
DYN 0010
DYN 0020
DYN 0030
DYN 0040
DYN 0050
DYN 0060
DYN 0070
DYN 0080
DYN 0090
DYN 0100
DYN 0110
DYN 0120
DYN 0130
DYN 0140
DYN 0150
DYN 0160
DYN 0170
DYN 0180
DYN 0190
DYN 0200
DYN 0210
DYN 0220
DYN 0230
DYN 0240
DYN 0250
DYN 0260
DYN 0270
DYN 0280
DYN 0290
DYN 0300
DYN 0310
DYN 0320

IMPLICIT REAL*8(A-H,O-Z)
EXTERNAL FCT,OUT
DIMENSION PRMT(5),X(42),DERY(42),AUX(16,42),XW(13)
COMMON U(21),XM(21),C(21)
COMMON PX(7),PY(7),PZ(7),DX(24),DY(24),DZ(24)
COMMON XL(24),XLP(24),DU(24),DV(24),DW(24),DL(24)
COMMON SUMX(24),SUMY(24),SUMZ(24)
COMMON INP(24),INQ(24),NM,NJ,NF,KOUNT
DO 100 J=1,NJ
  XW(J)=0.0D0
100 CONTINUE
DO 200 M=1,NM
  J=INP(M)
  K=INQ(M)
  XW(J)=XW(J)+(0.5*XL(M)*490.)/(32.2*12.**4)
  XW(K)=XW(K)+(0.5*XL(M)*490.)/(32.2*12.**4)
200 CONTINUE
DO 400 J=1,NF
  J1=3*J-2
  J2=3*J-1
  J3=3*J
  XM(J1)=XW(J)
  XM(J2)=XW(J)
  XM(J3)=XW(J)
  READ(5,300) C(J1),C(J2),C(J3)
300 FORMAT(3E20.5)
400 CONTINUE
  WRITE(6,500)
500 FORMAT('1'///30X,'INITIAL CONDITIONS FOR DYNAMIC ANALYSIS'/////

```

```

$ 30X, 'DISPLACEMENT', 45X, 'VELOCITY'//
$ 30X, '-----', 45X, '-----'//
$ 2X, 'JOINTNO', 8X, 'X-DIRCT', 10X, 'Y-DIRCT', 10X, 'Z-DIRCT',
$ 15X, 'X-DIRCT', 10X, 'Y-DIRCT', 10X, 'Z-DIRCT'///
DO 700 J=1,NF
  J1=6*J-5
  J2=6*J-4
  J3=6*J-3
  J4=6*J-2
  J5=6*J-1
  J6=6*J
  READ(5,600) TQ,X(J1),X(J3),X(J5),X(J2),X(J4),X(J6)
  600 FORMAT(F8.5,6D12.5)
  WRITE(6,650) J,X(J1),X(J3),X(J5),X(J2),X(J4),X(J6)
  650 FORMAT(I6,7X,6D15.5)
  700 CONTINUE
  READ(5,800) (PRMT(I),I=1,5)
  WRITE(6,850) (PRMT(I),I=1,5)
  800 FORMAT(5E15.5)
  850 FORMAT(/////5X,5D15.5)
  NDIM=6*NF
  XN=NDIM
  DO 900 I=1,NDIM
    DERY(I)=1./XN
  900 CONTINUE
  IH=5
  WRITE(6,950)
  950 FORMAT('I')
  CALL DHPGCG(PRMT,X,DERY,NDIM,IH,FCT,OUT,AUX)
  RETURN
  END
DYN 0330
DYN 0340
DYN 0350
DYN 0360
DYN 0370
DYN 0380
DYN 0390
DYN 0400
DYN 0410
DYN 0420
DYN 0430
DYN 0440
DYN 0450
DYN 0460
DYN 0470
DYN 0480
DYN 0490
DYN 0500
DYN 0510
DYN 0520
DYN 0530
DYN 0540
DYN 0550
DYN 0560
DYN 0570
DYN 0580
DYN 0590
DYN 0600
DYN 0610
DYN 0620
DYN 0630

```

```

SUBROUTINE FCT(T,X,DERY)
C
C
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION X(42),DERY(42),G(21),O(21)
      COMMON U(21),XM(21),C(21)
      COMMON PX(7),PY(7),PZ(7),DX(24),DY(24),DZ(24)
      COMMON XL(24),XLP(24),DU(24),DV(24),DW(24),DL(24)
      COMMON SUMX(24),SUMY(24),SUMZ(24)
      COMMON INP(24),INQ(24),NM,NJ,NF,KOUNT
      N=3*NF
      DO 100 J=1,NF
      J1=3*J-2
      J2=3*J-1
      J3=3*J
      J4=6*J-5
      J5=6*J-3
      J6=6*J-1
      Q(J1)=U(J1)+X(J4)
      Q(J2)=U(J2)+X(J5)
      Q(J3)=U(J3)+X(J6)
100 CONTINUE
      CALL V(N,Q,VAL,G)
      DO 200 J=1,N
      G(J)=-G(J)*3.0D07
200 CONTINUE
      NDIM=2*N
      DO 300 J=1,N
      I1=2*J-1
      I2=2*J
      DERY(I1)=X(I2)
      IF(J.EQ.1) DERY(I2)=(G(J)-C(J))*X(I2)/XM(J)
FCT 0010
FCT 0020
FCT 0030
FCT 0040
FCT 0050
FCT 0060
FCT 0070
FCT 0080
FCT 0090
FCT 0100
FCT 0110
FCT 0120
FCT 0130
FCT 0140
FCT 0150
FCT 0160
FCT 0170
FCT 0180
FCT 0190
FCT 0200
FCT 0210
FCT 0220
FCT 0230
FCT 0240
FCT 0250
FCT 0260
FCT 0270
FCT 0280
FCT 0290
FCT 0300
FCT 0310
FCT 0320

```



```
IF(J.NE.1) DERY(I2)=(G(J)-C(J)*X(I2))/XM(J)
300 CONTINUE
RETURN
END
```

```
FCT 0330
FCT 0340
FCT 0350
FCT 0360
```

```

C
C
SUBROUTINE OUT(T,X,DERY,IH,N,PRMT)
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION Q(21),G(21)
      DIMENSION X(42),DERY(42),PRMT(5)
      COMMON U(21),XM(21),C(21)
      COMMON PX(7),PY(7),PZ(7),DX(24),DY(24),DZ(24)
      COMMON XL(24),XLP(24),DU(24),DV(24),DW(24),DL(24)
      COMMON SUMX(24),SUMY(24),SUMZ(24)
      COMMON INP(24),INQ(24),NM,NJ,NF,KOUNT
      IF(T.EQ.0.000) GO TO 40
      JN=INT/100
      JM=JN*100
      IF(JM.EQ.INT) GO TO 60
      GO TO 50
40 INT=0
60 CONTINUE
      EK=0.000
      DO 15 J=1,NF
      J1=3*J-2
      J2=3*J-1
      J3=3*J
      J4=6*J-5
      J5=6*J-3
      J6=6*J-1
      K1=6*J-4
      K2=6*J-2
      K3=6*J
      Q(J1)=U(J1)+X(J4)
      Q(J2)=U(J2)+X(J5)
      Q(J3)=U(J3)+X(J6)
      OUT 0010
      OUT 0020
      OUT 0030
      OUT 0040
      OUT 0050
      OUT 0060
      OUT 0070
      OUT 0080
      OUT 0090
      OUT 0100
      OUT 0110
      OUT 0120
      OUT 0130
      OUT 0140
      OUT 0150
      OUT 0160
      OUT 0170
      OUT 0180
      OUT 0190
      OUT 0200
      OUT 0210
      OUT 0220
      OUT 0230
      OUT 0240
      OUT 0250
      OUT 0260
      OUT 0270
      OUT 0280
      OUT 0290
      OUT 0300
      OUT 0310
      OUT 0320

```

```

EK=EK+0.5D0*(XM(J1)*X(K1)**2+XM(J2)*X(K2)**2+XM(J3)*X(K3)**2)
15 CONTINUE
CALL V(21,Q,VAL,G)
EEK=EK/3.0D07
TEG=VAL+EEK
WRITE(6,10) T,VAL,EK,EEK,TEG
10 FORMAT(/2X,5D20.5//)
DO 30 J=1,NF
J1=6#J-5
J2=6#J-4
J3=6#J-3
J4=6#J-2
J5=6#J-1
J6=6#J
WRITE(6,20) J,X(J1),X(J3),X(J5),X(J2),X(J4),X(J6)
20 FORMAT(16X,I5,6(6X,D12.5))
WRITE(7,26) T,X(J1),X(J3),X(J5),X(J2),X(J4),X(J6)
26 FORMAT(F8.5,6D12.5)
30 CONTINUE
50 INT=INT+1
RETURN
END
OUT 0330
OUT 0340
OUT 0350
OUT 0360
OUT 0370
OUT 0380
OUT 0390
OUT 0400
OUT 0410
OUT 0420
OUT 0430
OUT 0440
OUT 0450
OUT 0460
OUT 0470
OUT 0480
OUT 0490
OUT 0500
OUT 0510
OUT 0520
OUT 0530
OUT 0540

```

VITA

Ayodele Olushola Abatan was born on September 16, 1949, in Abeokuta, Nigeria. He attended Baptist Boys' High School, Abeokuta, from January 1962 to December 1966. After a preliminary engineering course at Ahmadu Bello University, he began studying for a degree in Civil Engineering in September 1968 at the same University. In June, 1971, he graduated with a B.S. in Civil Engineering, with first class honors, specializing in Structures.

As a recipient of the African Graduate (AFGRAD) Fellowship program, under the auspices of the African-American Institute in New York, the author began his graduate studies at the University of Illinois at Urbana in September 1971. He received the M.S. degree in Civil Engineering in February 1973 specializing in Structural Mechanics. He then started studying towards a Ph.D. degree in Civil Engineering at the same University. In 1974, he worked as a structural engineer with Westenhoff and Novick Consulting Engineering Company in Chicago. His major assignments were on the design phases of the tunnel sections, plans and profiles for the extension to the Chicago subway transit system.

In January 1975, he entered Virginia Polytechnic Institute and State University to continue his doctoral studies while participating in a research program.

The author is a member of the American Society of Civil Engineers and the Institution of Civil Engineers (London). He is currently employed as Lecturer in the Civil Engineering Department at Ahmadu Bello University, Zaria, Nigeria.


A handwritten signature in cursive script, reading "Abatan", with a horizontal line underneath.