

Chapter 4

Force Measurements between Activated Silica Surfaces in Sodium Oleate Solutions

4.1 Introduction

Hydrophobic forces were first observed by Israelachvili and Pashley (1) between two curved mica surfaces in $2.5 \times 10^{-5} \text{M}$ cetyltrimethylammonium bromide (CTAB) solution. Since mica surfaces are negatively charged, cationic surfactants adsorbed on the surface and hydrophobized it. The measured hydrophobic force (F_h) were fit using the following expression:

$$F_h/R = -C_o \exp(-H/D_o) \quad (1)$$

where R is the radius of the mica cylinders used for the measurement, $D_o = 1.2 \text{ nm}$ and $C_o = 30 \text{ mN/m}$. Hydrophobic forces have also been observed with other cationic surfactants, such as dodecylammoniumhydrochloride (DAHCl), octadecyl-ammoniumhydrochloride (OAHCl) and dihexadecyl-dimethyl ammonium acetate (DHDAA), in equilibrium with mica and silica surfaces (2,3). The C_o and D_o values obtained with these surfactants are comparable to those obtained with CTAB.

Pugh and Rutland (4) conducted force measurements with mica surfaces in equilibrium with an anionic surfactant, sodium laurate, at pH 10.3. They added Ca^{2+} ions to the solution to reverse the surface charge of the mica, which is referred to as 'activation' in flotation literature. However, the authors failed to detect any hydrophobic force, indicating that the anionic surfactant did not adsorb on the activated mica surface.

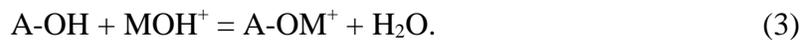
Hydrolyzable cations are commonly used as activators for the adsorption of anionic

surfactants. For example, Ca^{2+} ions are added to activate anatase (TiO_2) during the process of separating this mineral from kaolin clay by flotation using fatty acids. Fuerstenau *et al.* (5) measured the ζ -potentials of silica and alumina in the presence of various hydrolyzable cations. They found that with a given activator, ζ -potentials are reversed at a pH where the concentration of the singly charged hydroxo-complexes, $\text{M}(\text{OH})_{(z-1)}^+$, becomes appreciable in solution. They also conducted microflotation tests with silica in the presence of alkyl sulfonate (6) using various hydrolyzable cations (Fe^{3+} , Pb^{2+} and Mn^{2+}), and found that the flotation is most favored at the pH of maximum concentration of the singly charged hydroxo-complexes. Fuerstenau and Cummins (7) reported similar results for the flotation of quartz with sodium oleate in the presence of CaCl_2 .

To explain the effectiveness of the singly-charged hydroxo-complexes as activator, Palmer *et al.* (8) suggested two possible mechanisms: i) hydroxo complexes are hydrogen bonded to the surface hydroxyl groups:



where A is the metal ion in the substrate and M is the hydrolyzable cation, and ii) hydroxo-complexes chemisorb on the surface as follows:



Both of these mechanisms suggest that the adsorption should give rise to large negative free energies of adsorption ($\Delta G_{\text{ads}}^{\circ}$), which may provide an explanation for the experimental observation that singly charged hydroxo-complexes have higher adsorbability than the unhydrolyzed cations (M^{Z+}).

It should be noted here that while the mechanisms proposed by Palmer *et al.* (8) explains why the adsorption of the singly charged hydroxo-complexes should be favored, they do not explain why the unhydrolyzed cations do not adsorb on negatively charged surfaces. For the latter, James and Healy (9) subdivided $\Delta G_{\text{ads}}^{\circ}$ into three components:

$$\Delta G_{ads}^o = \Delta G_{elec}^o + \Delta G_{solv}^o + \Delta G_{chem}^o \quad (4)$$

where ΔG_{elec}^o represents the component from coulombic attraction, ΔG_{solv}^o represents the unfavorable free energy changes associated with the removal of the water molecules from secondary hydration sheath during the adsorption process, and ΔG_{chem}^o represents the component due to specific chemical interactions such as Reaction (3) above. They suggested that the ΔG_{solv}^o of the unhydrolyzed cations is much larger than that of the singly charged hydroxo-complexes. As a result, ΔG_{ads}^o of the former is more unfavorable than that of the latter. Indeed, it is well known that ΔG_{solv}^o decreases quadratically with charge (9). By calculating all three of the free energy component in Eq.(4), James and Healy were able to explain why unhydrolyzed cations cannot adsorb on silica. However, they assumed the same values of ΔG_{chem}^o for both the unhydrolyzed cations and the singly charged hydroxo-complexes. Also, the values of ΔG_{elec}^o were calculated using the values of surface potentials calculated from the Nernst equation, which are nearly an order of magnitude higher than those obtained from surface force measurements.

Schindler *et al.* (10) developed a completely different model, in which unhydrolyzed metal ions chemisorb on silica via an exchange mechanism involving the protons of the silanol groups. They simply assigned different equilibrium constants for the adsorption isotherms of various cations (i.e., Fe^{3+} , Pb^{2+} , Al^{3+} and Cu^{2+}). Their model does not provide a mechanistic explanation for the fact that the adsorption is a maximum at the pH where the singly charged hydroxo-complexes is maximum.

In the present communication, an atomic force microscope (AFM) was used to measure the surface forces between two silica surfaces in $MgCl_2$ and $CaCl_2$ solutions. The adsorption of these cations should not only change repulsive double-layer forces, but also hydration forces. It was shown previously that silica surfaces exhibit hydration forces in pure water (11-19), that diminishes in the presence of cations (14-18) and short-chain alcohols (13). The surface force measurements were also conducted between activated silica surfaces in the presence of sodium oleate. It is of particular interest to see if hydrophobic forces are observed under conditions of

maximum adsorption of oleate, which in turn may be related to the optimum conditions for activation.

4.2 Experimental

4.2a Sample Preparation

Optically smooth fused silica plates were obtained from Heraeus Amersil, Inc. They were cleaned by boiling in a nitric acid solution for 12 to 15 hours, and then equilibrated in conductivity water from a Barnstead Nanopure II water treatment unit. The clean silica plates thus obtained exhibited equilibrium contact angles (θ_{eq}) of less than 5° , as measured using a Rame-Hart Model 100 goniometer. Glass spheres of radius 10 to 30 μm from Duke Scientific were mounted on cantilever tips by means of an adhesive (Epon R Resin 1004, Shell Chemical Company) (19,20), placed in NaOH solutions (pH~10) overnight, and equilibrated in Nanopure water before each experiment. Reagent grade CaCl_2 and MgCl_2 (both 99.3% purity) were obtained from Fischer Scientific. Purified sodium oleate from Fischer Scientific was used as received. A 2×10^{-3} M stock solution of oleate was prepared in 50% ethanol solution. Since the concentration of oleate used in the experiments was 1×10^{-4} M, the ethanol content never exceeded 2.5%.

4.2b Methods and Apparatus

A Digital Instruments Nanoscope III AFM was used to measure the forces between a silica plate and a glass sphere mounted at the tip of a cantilever. Both standard triangular silicon nitride and TESP were obtained from Digital Instruments Co. The stiffness of the cantilevers was measured by the technique of Senden and Ducker (19). This consists of placing a heavy tungsten sphere on the cantilever tip, turning the cantilever upside down and measuring its deflection from the AFM signal. The glass spheres were mounted on the cantilever using a micromanipulator. The mounted spheres, equilibrated in NaOH solutions and Nanopure water as described in the foregoing last section, were then transferred to a fluid cell for force measurements. The force measurements were conducted in Nanopure water and in solutions containing known

concentrations of CaCl_2 (or MgCl_2) and oleate. The surfaces were equilibrated for 15 minutes in each solution prior to the force measurements. The pH was adjusted by NaOH and measured by means of a glass electrode. The radius of the glass spheres was measured using a Kontron SEM-IPS image processing system.

4.3 Results

4.3a Force Measurements with MgCl_2 and Oleate

Figure 4.1 shows the effect of pH on the equilibrium contact angles (θ_{eq}) of silica plates in the presence of 2×10^{-4} M MgCl_2 and 10^{-4} M sodium oleate. The dotted lines represent the concentrations of the various Mg-species present in solution at 2×10^{-4} M MgCl_2 . The thermodynamic data used for the calculations were taken from reference 21. As shown in the figure, θ_{eq} exhibits a maximum value of 70° in the pH range of 11.0-11.3. At $\text{pH} < 11.0$ and $\text{pH} > 12.0$, the silica surface remains hydrophilic with $\theta_{\text{eq}} \sim 5^\circ$. One can see that the changes in θ_{eq} with pH correspond to the same in the concentration of the $\text{Mg}(\text{OH})^+$ ion, suggesting that silica is activated by the singly charged hydroxo-complexes, rather than the Mg^+ ions, and that oleate adsorption is closely related to the activation mechanism. These findings are in agreement with the results of Fuerstenau *et al.* (6), who showed that the flotation of quartz is maximum at pHs where the concentrations of singly-charged hydroxo complexes are maximum.

Figure 4.2 shows the results of the direct surface force measurements conducted between glass sphere and silica plate using an AFM. The measurements were conducted i) in Nanopure water, ii) in solutions containing 2×10^{-4} M MgCl_2 , and iii) in solutions containing the same amount of MgCl_2 and 10^{-4} M oleate. All the measurements were conducted at pH 5.7. The measured forces, F , normalized by the radius of the sphere, R , are plotted as a function of the separation distance, H . The dashed lines represent the constant charge model of the DLVO theory (19,20) fitted to the experimental data. The DLVO curves were obtained using the algorithm of Chan *et al.* (21). Force data obtained in Nanopure water can be fitted to the constant charge model of the DLVO theory using surface potential (ψ_0) of -57 mV, Debye length (κ^{-1}) of 96 nm, and the

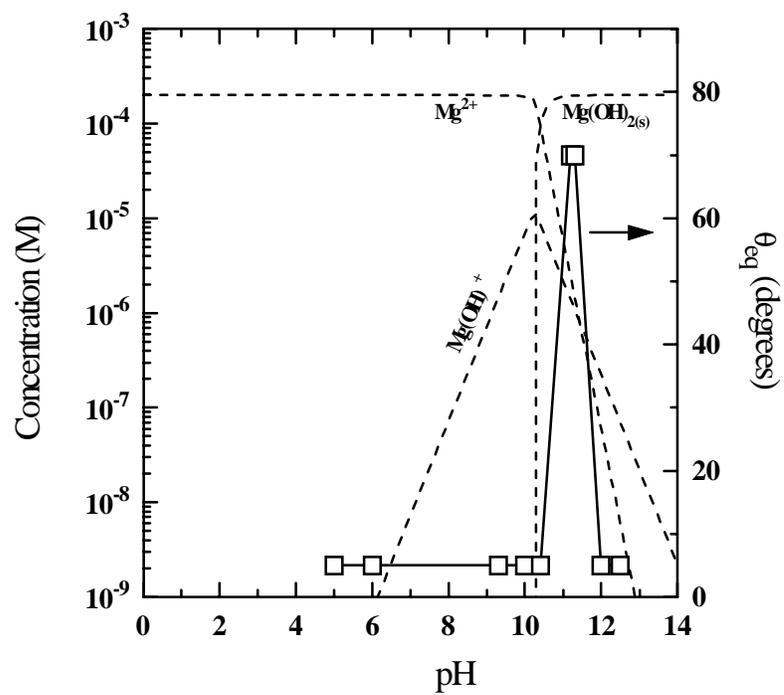


Figure 4.1. θ_{eq} vs pH (\square) curve for a silica plate immersed in 2×10^{-4} M $MgCl_2$ and 10^{-4} M Na-Oleate. The dotted lines represent the concentrations of the various species present in solution.

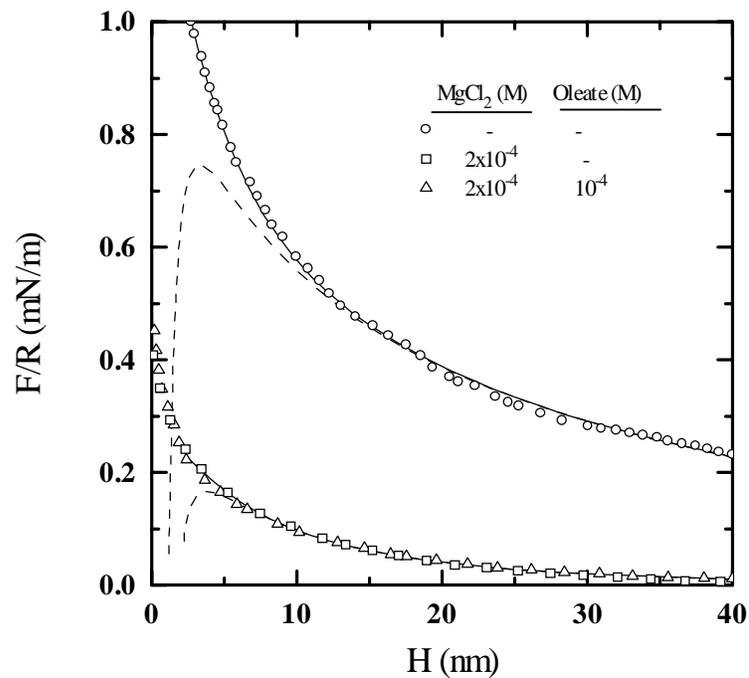


Figure 4.2. F/R vs H curves obtained between silica plate and glass sphere in i) Nanopure water (○), ii) 2×10^{-4} M MgCl_2 (□), and iii) 2×10^{-4} M MgCl_2 and 10^{-4} M sodium oleate (△) at pH 5.7. The dashed lines represent the constant charge model of the DLVO theory fitted to the data using the parameters listed in Table I. The solid lines represent the extended DLVO theory (Eq. [5]), incorporating a hydration force term.

Hamaker constant (A_{131}) of 6.3×10^{-21} J. These values are close to those reported previously (19).

At $H < 12$ nm, the force curve obtained in Nanopure water shows repulsive forces larger than those predicted by the DLVO theory. This extraneous force, which has also been observed previously (22-24), is generally referred to as the hydration force. Therefore, the force data are fitted to the extended DLVO theory:

$$F_t = F_e + F_d + F_{\text{hyd}} \quad (5)$$

which includes contributions from the hydration force (F_{hyd}) in addition to those from the electrostatic (F_e) and dispersion (F_d) forces considered in the DLVO theory. The hydration force is represented by the following expression:

$$F/R = C_1 \exp(-H/D_1) + C_2 \exp(-H/D_2), \quad (6)$$

where $C_1 = 7.7$ mJ/m², $D_1 = 0.50$ nm, $C_2 = 0.50$ mJ/m², and $D_2 = 2.70$ nm. These values are close to those estimated by Yoon and Yotsumoto (15) based on turbidity measurements.

Note in Figure 4.2 that the force data obtained in 2×10^{-4} M MgCl₂ can be fitted with $\psi_0 = -14$ mV, which is much lower than $\psi_0 = -57$ mV obtained in Nanopure water (Table 4.1). At 2×10^{-4} M MgCl₂, the theoretically calculated double layer thickness (κ^{-1}) is 11.0 nm. This value is comparable to that obtained by fitting the force data, which is 13.4 nm. The changes in ψ_0 from -57 mV to -14 can be attributed to the double-layer compression rather than the neutralization of the negative charge sites by the adsorption of Mg²⁺ ions. This can be seen from the following approximate expression:

$$\psi_0 = \sigma \kappa^{-1} / \epsilon \epsilon_0 \quad (7)$$

where σ is the surface charge, ϵ is the permittivity of free space and ϵ_0 the relative dielectric constant for water. It can be seen from Eq.[7] that if κ^{-1} decreases, ψ_0 should also reduce.

Table 4.1

The Contact Angle and Surface Force Data obtained with silica in the presence of MgCl_2 and Na-Oleate at different pHs

pH	MgCl_2 (M)	Oleate (M)	θ_{eq}	ψ_0 (mV)	κ^{-1} (nm)	C_1 (mJ/m ²)	D_1 (nm)	C_2 (mJ/m ²)	D_2 (nm)
5.7	-	-	$<5^\circ$	-57	96.0	7.7	0.50	0.50	2.7
5.7	2×10^{-4}	-	$<5^\circ$	-14	13.4	8.0	0.46	0.40	1.4
5.7	2×10^{-4}	10^{-4}	$<5^\circ$	-14	13.4	8.0	0.46	0.40	1.4
11.2	-	-	$<5^\circ$	-30	11.0	7.0	0.41	0.40	0.8
11.2	2×10^{-4}	-	$<5^\circ$	0	-	-	-	-	-
11.2	2×10^{-4}	10^{-4}	70°	+42	11.0	-30.0	2.10	-	-
12.0	-	-	$<5^\circ$	-12	3.5	8.0	0.42	-	-
12.0	2×10^{-4}	-	$<5^\circ$	-12	3.5	8.0	0.42	-	-
12.0	2×10^{-4}	10^{-4}	$<5^\circ$	-12	3.5	8.0	0.42	-	-

Therefore, decrease in ψ_0 at $2 \times 10^{-4} \text{ M MgCl}_2$ should be attributed to double-layer compression.

It is interesting that in the presence of $2 \times 10^{-4} \text{ M MgCl}_2$, the hydration force is still visible at $H < 5 \text{ nm}$, which further suggests that Mg^{2+} ions do not adsorb on silica at pH 5.7. Eq. [5] can be used to fit the hydration force using the parameters shown in Table I. As shown, D_2 is reduced by $\sim 54\%$ in the presence of $2 \times 10^{-4} \text{ M MgCl}_2$. It is well documented that the primary hydration force observed between silica surfaces decreases in the presence of cations (14-18).

Figure 4.2 shows that the force data obtained in solutions containing $2 \times 10^{-4} \text{ M MgCl}_2$ and 10^{-4} M oleate fall on the same curve as those obtained at $2 \times 10^{-4} \text{ M MgCl}_2$ in the absence of oleate. As seen in Table I, there are no changes in the values of ψ_0 and D_2 that were used for the fit the data. One can conclude, therefore, that oleate does not adsorb on silica at this pH 5.7. This finding is in agreement with the fact that, at these conditions, θ_{eq} remains in the neighborhood of $\sim 5^\circ$.

Figure 4.3 shows the results of the AFM force measurements conducted between the glass sphere and silica plate at pH 11.2. In the absence of MgCl_2 and oleate, the experimental data obtained at $H > 5 \text{ nm}$ can be fitted to the DLVO theory with $\psi_0 = -30 \text{ mV}$. On the other hand, the data obtained at $H < 5 \text{ nm}$ can only be fitted to the extended DLVO theory, which includes contributions from the hydration force. The data obtained in the absence of MgCl_2 and oleate have been fitted using the parameters listed in Table I. Note that both $-\psi_0$ and D_2 are lower than those obtained at pH 5.7, which may be attributed to the Na^+ ions introduced as NaOH for pH adjustment. At pH 11.2, the theoretically expected value of κ^{-1} is 7.5 nm , which is comparable to the value of 7.6 nm obtained by fitting the data to the DLVO theory.

Figure 4.3 also shows the results of force measurements conducted in a $2 \times 10^{-4} \text{ M MgCl}_2$ solution at pH 11.2. No double layer repulsion is observed, suggesting that the charges on the silica surfaces have been neutralized. At $H < 10 \text{ nm}$, the forces are net attractive and the two surfaces jump into contact from a distance of 7 nm . The experimental data can be fitted to the van der Waals force alone with $A_{131} = 6 \times 10^{-21} \text{ J}$, as shown by the dashed line. The value of A_{131}

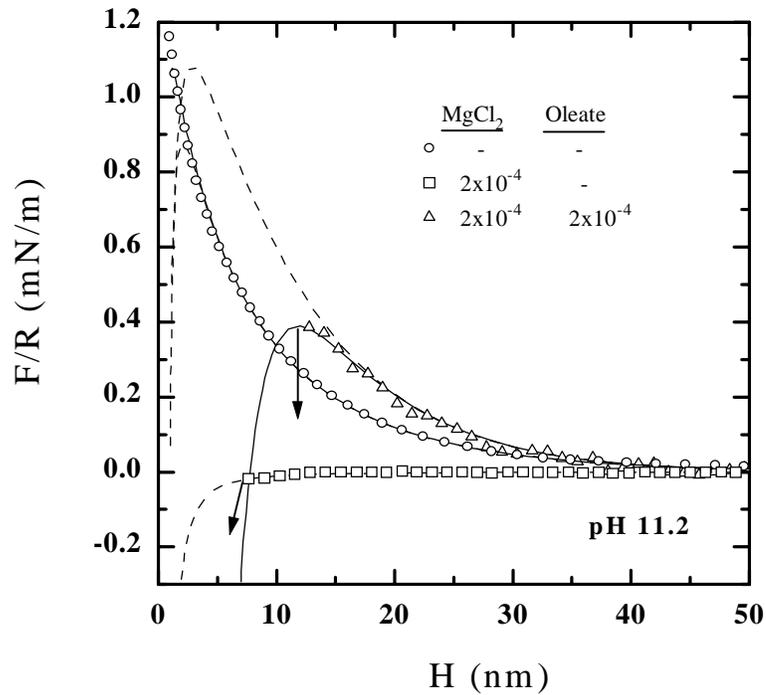


Figure 4.3. F/R vs H curves obtained between silica plate and glass sphere at pH 11.2 i) in the absence of MgCl_2 and sodium oleate (○), ii) in the presence of 2×10^{-4} M MgCl_2 (□), and iii) in the presence of 2×10^{-4} M MgCl_2 and 10^{-4} M sodium oleate (△). The dashed lines represent the constant charge model of the DLVO theory, whereas the solid line represents the extended DLVO theory (Eq. [5]) using the parameters shown in Table I. The arrows indicate the positions at which the two surfaces jump into contact

obtained here agrees well with the reported values in literature (12) for the Hamaker constant of the silica-water-silica system. The fact that $\psi_0=0$ mV under these conditions is not surprising. According to Fuerstenau *et al.* (5), the point of charge reversal (p.c.r.) of silica occurs at $\text{pH}\approx 10.5$ at 10^{-4} M MgCl_2 . Note that the force data obtained at 2×10^{-4} M MgCl_2 show no evidence for the hydration force. This is in sharp contrast to the observation made at $\text{pH } 5.7$, where the hydration force does not disappear in the presence of 2×10^{-4} M MgCl_2 . Both the charge neutralization and the disappearance of the hydration force at $\text{pH } 11.2$ may be attributed to the adsorption of the MgOH^+ ions on the silica surface. As shown by the distribution diagram (dashed lines in Figure 4.3), the concentration of the MgOH^+ ions reaches maximum at $\text{pH } 10.4$.

In the presence of both 2×10^{-4} M MgCl_2 and 10^{-4} M sodium oleate at $\text{pH } 11.2$, the force data can be fitted $\psi_0=-45$ mV. Thus, oleate addition has rendered the surface more negative than at 2×10^{-4} M MgCl_2 alone, which can be attributed to the adsorption of the anionic surfactant. In the presence of oleate, the two surfaces jump into contact from a distance of $H\approx 12$ nm, which is substantially larger than predicted by the DLVO theory. This suggests the presence of a hydrophobic force. The solid line represents an extended DLVO fit to the data, where the hydrophobic force has been represented by the single exponential function (Eq.[1]) with $C_0=-30$ mN/m and $D_0=2.1$ nm. The presence of the hydrophobic force and increase in $-\psi_0$ strongly suggest that oleate adsorbs on silica at $\text{pH } 11.2$. It may be recalled that this is the same pH at which the silica surface is activated by MgOH^+ ions and the hydration force disappears.

Figure 4.4 shows three sets of force data obtained at $\text{pH } 12$, which include those obtained i) in water ii) in 2×10^{-4} M MgCl_2 solution, and iii) in a solution containing both MgCl_2 (2×10^{-4} M) and sodium oleate (10^{-4} M). As can be seen, all the three sets of force data lie on the same curve. The parameters used to fit the data are given in Table I. The hydration force observed at this pH can be adequately represented by a single exponential function with a decay length of 0.42 nm. The short decay length may be attributed to the presence of the Na^+ ions added as NaOH to control pH . The fact that the force curves do not change in the presence of MgCl_2 and/or oleate suggest that neither of these species adsorb on silica at $\text{pH } 12.2$. This finding is in agreement with the fact that the values of θ_{eq} remain $\sim 5^\circ$ in the presence of MgCl_2 and oleate at this pH .

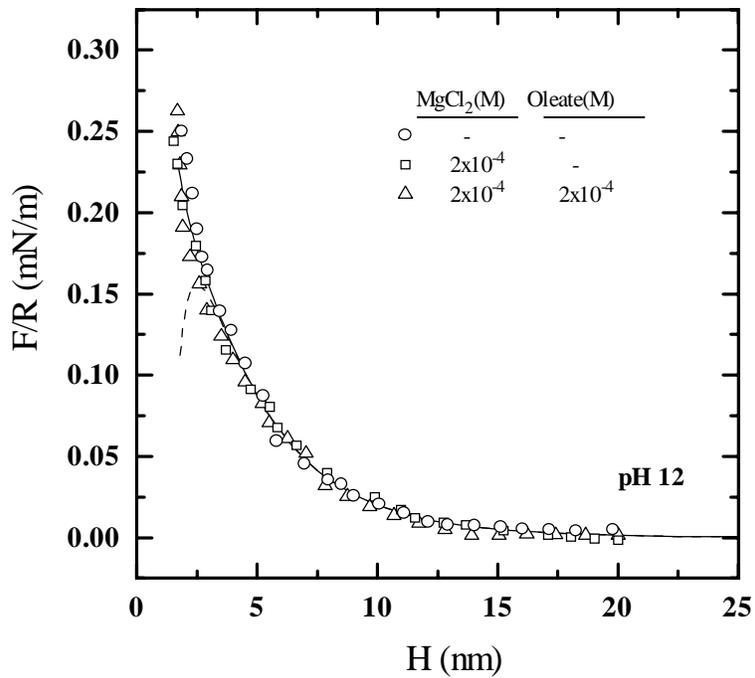


Figure 4.4. F/R vs H obtained between a silica plate and glass sphere at pH 12 in (i) absence of MgCl₂ and oleate (○), (ii) 2x10⁻⁴ M MgCl₂ (□) and (iii) 2x10⁻⁴ M MgCl₂ & 10⁻⁴ M Oleate (△). The dashed line represents the constant charge model of the DLVO theory whereas the solid lines represent the extended DLVO theory.

4.3b Force Measurements with CaCl₂ and Oleate

Figure 4.5 shows the values of θ_{eq} of silica plates in solutions containing both 2×10^{-4} M CaCl₂ and 10^{-4} M sodium oleate at various pHs. Shown by the dashed lines is the species distribution diagram at 2×10^{-4} M CaCl₂. The thermodynamic data used for the calculations were taken from reference 21. It can be seen from the Figure that at $pH < 11$, where the concentration of CaOH⁺ is low, silica remains hydrophilic with $\theta_{eq} = 5^\circ$. At $pH > 12$, where CaOH⁺ is present in significant quantities, the silica plate is rendered hydrophobic and exhibits $\theta_{eq} = 65^\circ$. Contact angle measurements were not conducted at $pH > pH_{ppt}$. It has been reported (25), nevertheless, that silica is not effectively activated by Ca²⁺ ions if the ratio of Na⁺ to Ca²⁺ ions in solution is greater than 10^3 .

Force measurements were conducted for the CaCl₂-oleate-silica system at two different pHs: 5.7 and 12.2. Figure 4.6 shows the effect of CaCl₂ (2×10^{-4} M) and sodium oleate (10^{-4} M) on the surface forces at pH 5.7. The figure also shows the force data obtained in the absence of CaCl₂ and oleate for comparison. It can be seen that CaCl₂ addition reduces both the long-range double layer repulsion and the short-range hydration repulsion. Table 4.2 shows the values of ψ_0 , κ^{-1} , C_1 , D_1 , C_2 and D_2 obtained by fitting the data to the extended DLVO theory. The ψ_0 and D_2 values obtained in the presence of CaCl₂ at pH 5.7 are very similar to those obtained in the presence of MgCl₂ at the same pH and concentration. Figure 4.6 and Table 4.2 show also that the force data obtained in a solution containing both CaCl₂ and oleate lie on the same curve as those obtained in the presence of CaCl₂ alone. This finding suggests that oleate does not adsorb on silica under these conditions.

Figure 4.7 shows the effect of CaCl₂ and sodium oleate at pH 12.2. Also shown for comparison are the force data obtained in the absence of CaCl₂ and oleate. The forces measured at $H > 3$ nm are due to double layer repulsion and can be fitted to the DLVO theory with $\psi_0 = -11$ mV and $\kappa^{-1} = 3.5$ nm. The theoretical value of κ^{-1} at this pH is 2.8 nm, which compares very well with that obtained from fitting the force data. At $H < 3$ nm, repulsive hydration forces are

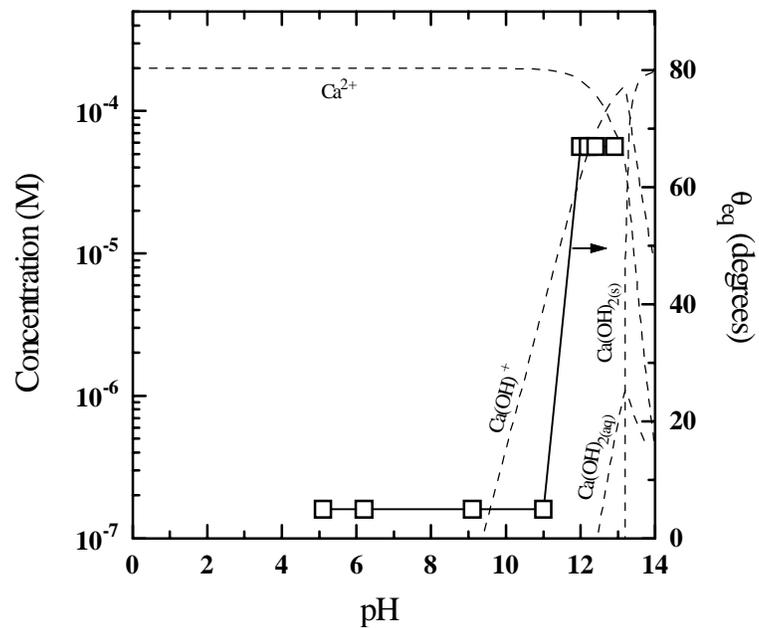


Figure 4.5. θ_{eq} vs pH curve (\square) for a silica plate immersed at 2×10^{-4} M CaCl_2 and 10^{-4} M Na-Oleate. The dashed lines represent the solution concentrations of the various species in solution.

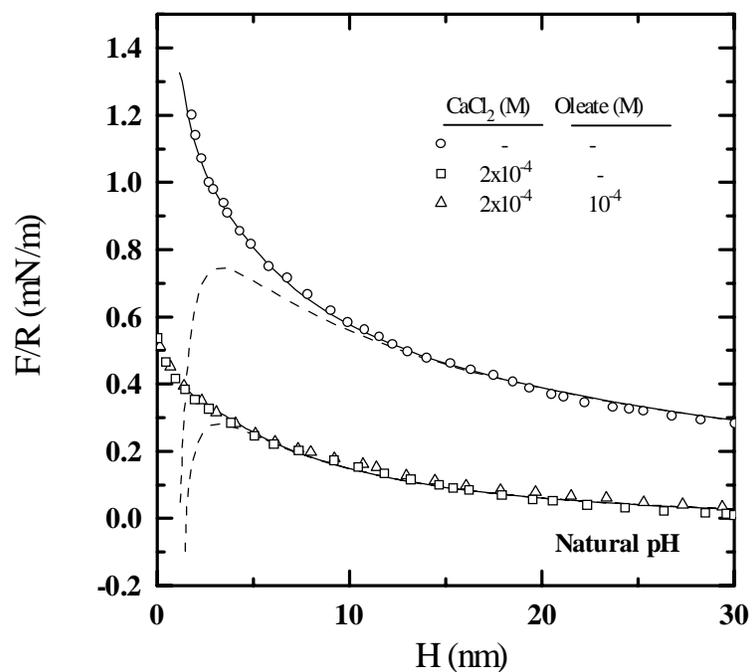


Figure 4.6. F/R vs H curves obtained between silica plate and glass sphere in i) Nanopure water (○), ii) at 2×10^{-4} M CaCl_2 (□), and iii) at 2×10^{-4} M CaCl_2 10^{-4} M sodium olate (△) at pH 5.7. The dashed lines represent the constant charge model of the DLVO theory and the solid lines represent the extended DLVO theory.

Table 4.2

Contact Angles and Parameters obtained from Fitting Force Curves between silica surfaces in solutions of CaCl₂ and Na-Oleate at different pHs

pH	CaCl ₂ (M)	Oleate (M)	θ_{eq} (degrees)	ψ_0 (mV)	κ^{-1} (nm)	C_1 (mJ/m ²)	D_1 (nm)	C_2 (mJ/m ²)	D_2 (nm)
5.7	-	-	<5 ⁰	-57	96.0	7.7	0.50	0.50	2.70
5.7	2x10 ⁻⁴	-	<5 ⁰	-18	15.0	7.6	0.44	0.37	1.25
5.7	2x10 ⁻⁴	10 ⁻⁴	<5 ⁰	-18	15.0	7.6	0.44	0.37	1.25
12.2	-	-	<5 ⁰	-11	3.5	8.0	0.45	-	-
12.2	2x10 ⁻⁴	-	<5 ⁰	+7	3.5	-	-	-	-
12.2	2x10 ⁻⁴	10 ⁻⁴	<65 ⁰	-	-	-30.0	1.05	-	-

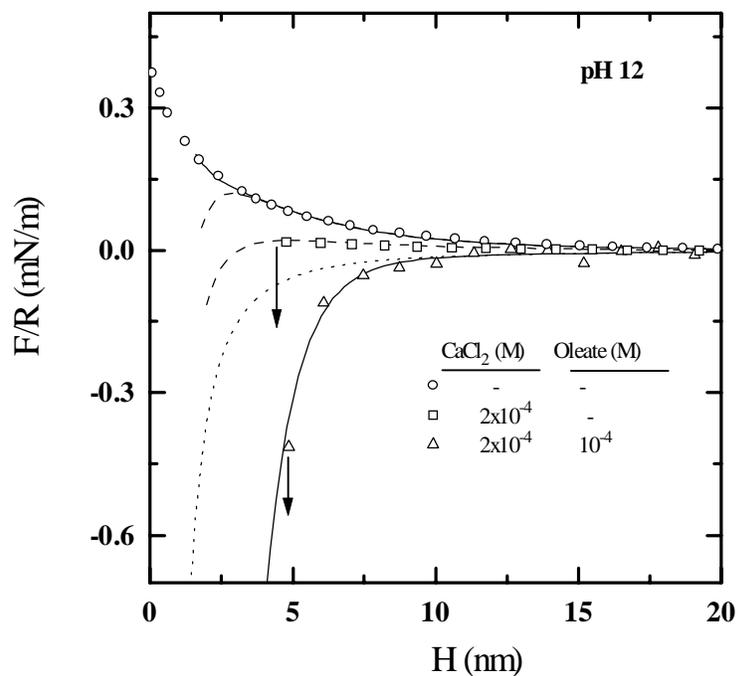


Figure 4.7. F/R vs H curves obtained between silica plate and glass sphere at pH 12.2 i) in Nanopure water (○), ii) at 2×10^{-4} M CaCl_2 (□), and iii) at 2×10^{-4} M CaCl_2 and 10^{-4} M sodium oleate (△). The dashed lines represent the constant charge model of the DLVO theory fitted to the force data. The solid line represents the extended DLVO theory (Eq. [5]) using the parameters shown in Table II. Also shown for comparison is the van der Waals force alone (dotted line).

observed, which can be fitted using the parameters shown in Table 4.2. Note that at 2×10^{-4} M CaCl_2 , the magnitude of ψ_0 is reduced to 7mV. Although the sign of ψ_0 cannot be ascertained from force measurements, it is likely that the silica surface is positively charged under the conditions employed. Many investigators (9,12) have reported charge reversal of silica at pHs close to the pH_{ppt} of the hydrolyzable cation. Note also that the two surfaces jump into contact at $H=4$ nm, and that there is no evidence of hydration repulsion. Since the concentration of CaOH^+ ions are high at pH 12.2, both the charge reversal and the disappearance of hydration force at this pH may be attributed to the specific adsorption of the singly charged hydroxo-complex on silica.

Also shown in figure 4.7 are the forces measured at 2×10^{-4} M CaCl_2 and 10^{-4} M Na-oleate at pH 12.2. No repulsive forces are observed at any separation distance. At $H < 10$ nm, the experimental force data are more attractive than the van der Waals force, which indicates the presence of the hydrophobic force. The solid line represents the extended DLVO theory with the parameters shown in Table II. The D_2 value obtained here is lower than that obtained with MgCl_2 and oleate at pH 11.2. This is in agreement with the fact that the maximum θ_{eq} obtained for the CaCl_2 -oleate-silica system is lower than that obtained for the MgCl_2 -oleate-silica system. The results obtained in the presence of CaCl_2 suggest that oleate adsorbs on silica at pH 12.2 and renders the surface hydrophobic. The fact that hydrophobic force appears at the pH where the concentration of CaOH^+ ions is close to its maximum indicates that silica is activated by the singly charged hydroxo-complex.

4.4 Discussions

In the present work, hydrophobic forces have been observed for the first time between two negatively charged (silica) surfaces immersed in anionic surfactant (sodium oleate) solutions. To achieve this, it was necessary to activate the silica by reversing its surface charge, so that the anionic surfactant can adsorb on the surface and render it hydrophobic. Two hydrolyzable cationic salts, CaCl_2 and MgCl_2 , were used as activators. An important observation made in the present work was that strong hydrophobic forces were measured only at the pH values where the divalent

cations are hydrolyzed to form singly charged hydroxo-complexes, i.e., CaOH^+ and MgOH^+ . This is evident from Figures 4.8 and 4.9, where the changes in the decay lengths (D_0 of Eq. [1]) of the hydrophobic forces are plotted versus pH along with the species distribution diagrams at 2×10^{-4} M MgCl_2 and CaCl_2 . Also shown in these two figures are the decay length (D_0 or D_2 of Eq. [6]) of the hydration forces measured in MgCl_2 and CaCl_2 solutions as a function of pH. Interestingly, the hydration force disappears completely at the pH where the concentrations of CaOH^+ and MgOH^+ reach maximum. These findings suggest, therefore, that the singly charged hydroxo-complexes are the activating species, and that their adsorption on silica causes the hydration force to disappear and, at the same time, causes the anionic surfactant to adsorb and, hence, the hydrophobic force to appear.

The results obtained in the present communication may be compared with those of Pugh and Rutland (4), who conducted force measurements between mica surfaces in solutions of CaCl_2 and sodium laurate at pH 10.3. Unlike the silica surfaces used in the present work, the mica surfaces did not exhibit hydrophobic force, indicating that oleate did not adsorb on the surface. On the other hand, the authors observed disappearance of the 'secondary hydration force', which had been observed in the absence of CaCl_2 . (Secondary hydration forces are observed when force measurements are conducted between mica surfaces in high concentrations of electrolytes (25,26).) The authors attributed the appearance of the hydration force to the presence of the Na^+ ions introduced to the system during pH adjustment, and its disappearance to the adsorption of CaOH^+ ions on mica. Displacement of the strongly hydrated Na^+ ions with the CaOH^+ ions from the mica surface should cause a decrease in hydration force, because the latter may be less strongly hydrated than the former owing to its large ionic size.

An interesting question to be raised here is why laurate/oleate ions adsorb on silica but not on mica despite the fact that CaOH^+ ions adsorb on both. Before embarking on the discussion, it should be noted here that the pH (=10.3) where Pugh and Rutland (4) conducted their force measurements with mica was considerably lower than the pH (=12.8) where the concentration of

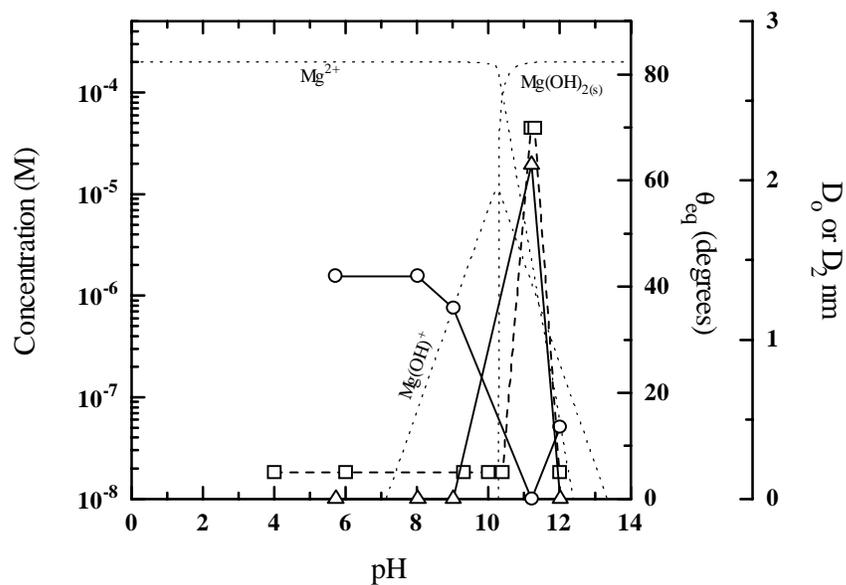


Figure 4.8. Effects of pH on D_o or D_2 of the hydration forces observed at 2×10^{-4} M MgCl_2 (○) and D_o of the hydrophobic forces observed at 2×10^{-4} M MgCl_2 and 10^{-4} M sodium oleate (Δ). The dashed lines represent the concentrations of the species present in solution. Also shown are the θ_{eq} (□) values.

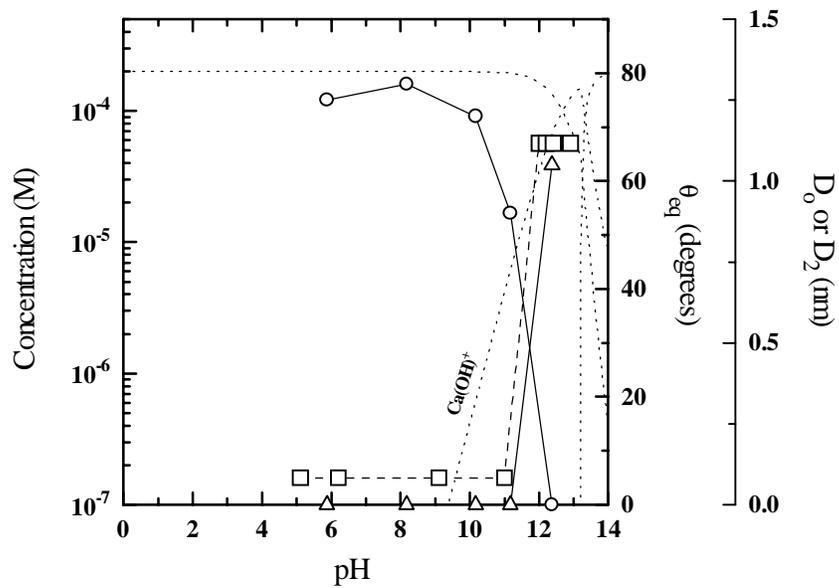


Figure 4.9. Effects of pH on D_o or D_2 of the hydration forces observed at 2×10^{-4} M CaCl_2 (O) and D_o of the hydrophobic force at 2×10^{-4} M CaCl_2 and 10^{-4} M oleate (Δ). The dashed lines represent the concentrations of the species in solution. Also shown are the θ_{eq} (\blacksquare) values.

the CaOH^+ ions reaches maximum (see Figure 4.9.) Thus, it is possible that the low concentration of CaOH^+ ions at pH 10.3 is responsible for oleate not adsorbing on mica at this pH. To test this hypothesis, we measured water contact angles on mica with solutions containing CaCl_2 and oleate solutions in the pH range of 9.0-12.4. The results showed that the equilibrium contact angles (θ_{eq}) were in the range of $34\text{-}36^\circ$ and did not change with pH. These results show that oleate indeed does not adsorb on mica as suggested by Pugh and Rutland.

When sodium oleate is added to a solution containing calcium ions, a basic calcium oleate (CaOHOl) should form in solution at pH where the concentration of CaOH^+ ions is high enough to exceed its solubility limit. The basic calcium oleate cannot adsorb on the mica surface *via* a coulombic attraction as it is a neutral species. Nor is it possible for it to adsorb on the surface *via* chemical interaction or H-bonding mechanism, because the basal surface of mica consists of siloxane bonds and, therefore, is chemically inert. Silica, on the other hand, is capable of forming H-bonds with the CaOHOl species. Fuerstenau (28) suggested that the CaOHOl and MgOHOl species formed in solution adsorb on oxide surfaces *via* H-bonding.

The surface force measurements conducted in the present work show that silica is activated by CaOH^+ and MgOH^+ species rather than by Ca^{2+} and Mg^{2+} ions. Similar observations have also been made with other oxides using different techniques (5,8,29). In general, singly charged hydroxo-complexes adsorb on oxides in preference to unhydrolyzed divalent cations. Figure 4.10 shows an adsorption isotherm for Ca(II) species on silica as obtained by James and Healy (9). Dotted lines represent the concentrations of Ca^{2+} and CaOH^+ ions as a function of pH at 1.4×10^{-4} M Ca(II) , which is the concentration at which the adsorption isotherm was established. As shown, the adsorption density increases with increasing concentration of CaOH^+ ions. No significant adsorption occurs at $\text{pH} < 9.0$, where most of the Ca(II) species exist as Ca^{2+} ions.

The dashed line given in Figure 4.10 represents a theoretical isotherm derived from the James and Healy's model (Eq. [4]). It was obtained by calculating $\Delta G_{\text{ads}}^{\circ}$ from its three components, i.e., $\Delta G_{\text{elec}}^{\circ}$, $\Delta G_{\text{solv}}^{\circ}$, and $\Delta G_{\text{chem}}^{\circ}$. Once the value of $\Delta G_{\text{ads}}^{\circ}$ is known, one can obtain

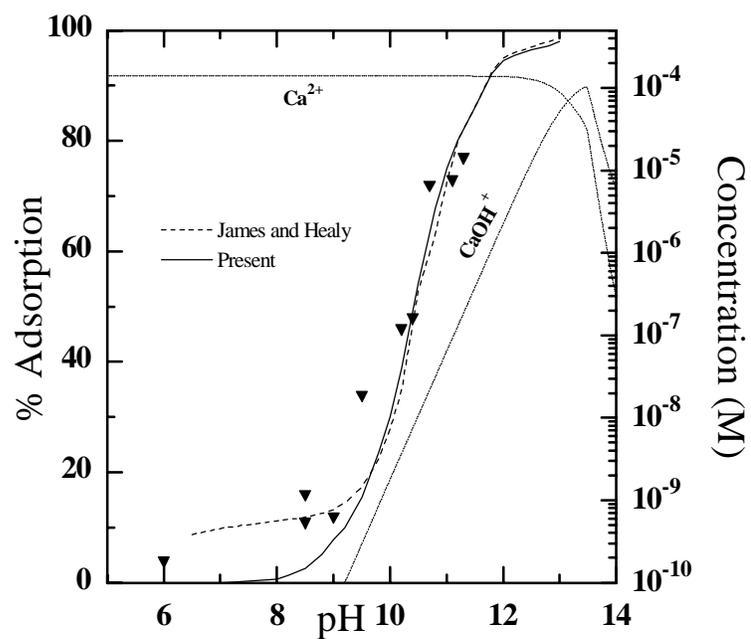


Figure 4.10. Adsorption of CaCl_2 (▼) as a function of pH, with the dashed line representing a theoretical fit using the James and Healy model (ref. 9). The solid line represents the model developed in the present work. The dotted lines represent the concentrations of Ca^{2+} and CaOH^+ ions at $2 \times 10^{-4} \text{ M}$ CaCl_2 .

the value of the equilibrium constant (K_i) of the adsorption equilibria using the following relationship:

$$K_i = \exp(-\Delta G_{\text{ads}}^{\circ}/RT) \quad (8)$$

where R is the universal gas constant and T is the absolute temperature. The adsorption equilibria were given by

$$(1 - \Theta) + M_i^{K_i} = \theta_i \quad (9)$$

where M_i is the concentration of species i , θ_i the fractional surface coverage by species i , Θ is the fraction of the surface sites occupied by all adsorbates, i.e. $\Theta = \sum_i \theta_i$. Eq. [9] yields:

$$\Theta = \frac{\sum_i K_i M_i}{1 + \sum_i K_i M_i} \quad (10)$$

which represents the theoretical adsorption isotherm of James and Healy given in Figure 4.10.

For $\Delta G_{\text{elec}}^{\circ}$, James and Healy used the following relationship:

$$\Delta G_{\text{elec}}^{\circ} = -zF\psi_x \quad (11)$$

where z is the valence of the species under consideration, F the Faraday constant, and ψ_x is the potential at a distance of x away from the surface. The distance x used in this calculation was the sum of the radius of the adsorbing ion and the diameter of water (0.14 nm). Thus, the authors considered that ionic species adsorb on silica without losing the water molecules from the primary hydration sheath. It should be noted here that ψ_x is a function of surface potential (ψ_0), the values of which were obtained using the Nernst equation:

$$\psi_0 = 2.3RT/zF (\text{pH}_{\text{pzc}} - \text{pH}) \quad (12)$$

where pH_{pzc} is the pH of pzc of silica. Finally, James and Healy fitted the theoretical isotherm to the experimental data using $\Delta G^{\circ}_{\text{chem}}$ as a fitting parameter. They obtained the value of -29.4 kJ/mole for both Ca^{2+} and CaOH^+ ions.

A problem associated with the James and Healy's model is that the values of ψ_0 obtained using the Nernst equation were higher than those determined in the present work from direct surface force measurement by nearly an order of magnitude (see Figure 4.11). According to Hunter (30), Nernst equation should not be used in calculating ψ_0 for oxides. It may be useful, therefore, to calculate $\Delta G^{\circ}_{\text{elec}}$ using the more realistic values of ψ_0 . Also, it may be reasonable to consider $\Delta G^{\circ}_{\text{sol}}$ to be part of $\Delta G^{\circ}_{\text{ads}}$. The reason is that dehydration is a prerequisite for chemisorption. Eq. [4] may then be reduced to:

$$\Delta G^{\circ}_{\text{ads}} = \Delta G^{\circ}_{\text{elec}} + \Delta G^{\circ}_{\text{chem}}. \quad (13)$$

In the present work, the values of $\Delta G^{\circ}_{\text{elec}}$ were calculated using Eq. [11]. The values of ψ_0 were obtained from the experimental data given in Figure 4.10. Note that ψ_0 decreases with increasing pH. This may be attributed to both double layer compression (due to the Na^+ ions added for pH adjustment) and neutralization of the SiO^- sites by Na^+ ions to form SiONa (31).

The model based on Eqs. [8], [11] and [13] was then fitted to the James and Healy's experimental data given in Figure 4.10 using $\Delta G^{\circ}_{\text{chem}}$ as an adjustable parameter. In this calculation, the CaOH^+ ions were considered as the activating species. The value of $\Delta G^{\circ}_{\text{chem}}$ used to fit the experimental data was -42 kJ/mole at all pHs. Such a large value of $\Delta G^{\circ}_{\text{chem}}$ cannot be accounted for by the H-bonding mechanism discussed earlier for the adsorption of CaOH on silica. In general, the free energy changes associated with H-bonding are in the neighborhood of -10 kJ/mole (32). Therefore, there may be additional mechanism(s) involved in the activation of silica by CaOH^+ ions. Subtracting the free energy of adsorption due to H-bonding from the value of $\Delta G^{\circ}_{\text{chem}}$ obtained from the model, one obtains a value of -32 kJ/mole, which should be accounted for by additional mechanism(s). One possible mechanism may be the lateral interactions between adsorbed hydroxo- complexes on the substrate. Figure 4.12 depicts how

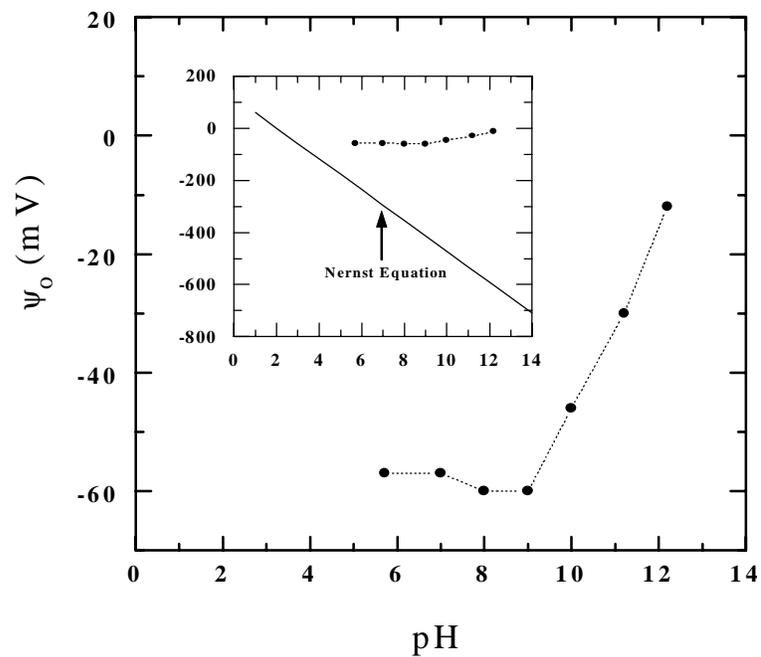


Figure 4.11. The values of ψ_0 (●) obtained from the AFM force measurements. The inset compares the them with those obtained using the Nernst Equation.

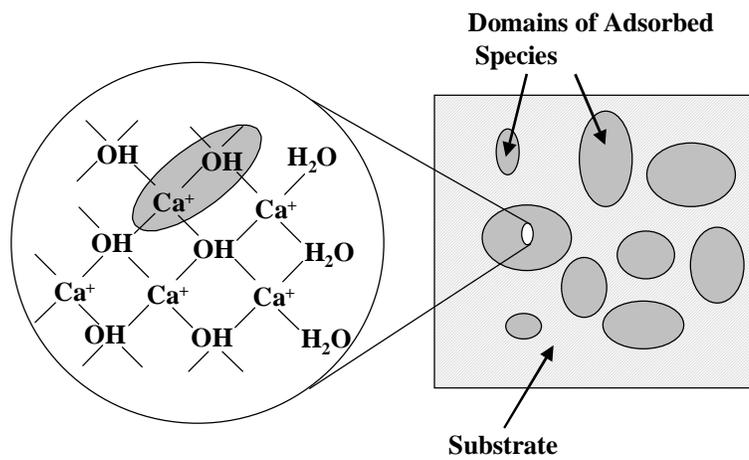


Figure 4.12. A conceptual model for the adsorption of CaOH^+ ions on a substrate. Domains (or clusters) of adsorbed species are formed due to lateral interactions.

CaOH⁺ species may adsorb on silica surface by forming lateral bonds. Neighboring CaOH⁺ species may interact with each other to form two-dimensional ‘crystals’, clusters or domains with free energy changes of -32 kJ/mole.

Thus, $\Delta G_{\text{chem}}^{\circ}$ may be considered to consist of two components, i.e.,

$$\Delta G_{\text{chem}}^{\circ} = \Delta G_{\text{chem}}^{\circ, \text{H}} + \Delta G_{\text{chem}}^{\circ, \text{l}} \quad (14)$$

where $\Delta G_{\text{chem}}^{\circ, \text{H}}$ is due to H-bonding with the silanol groups on the surface and $\Delta G_{\text{chem}}^{\circ, \text{l}}$ is due to the lateral interactions with neighboring adsorbates. $\Delta G_{\text{chem}}^{\circ, \text{l}}$ may be due to H-bonding or covalent bonding, the latter being more likely. For the CaOH⁺ ions adsorbing on silica, $\Delta G_{\text{chem}}^{\circ, \text{l}} = -32$ kJ as discussed in the foregoing paragraph. This value is close to the free energy of the following reaction:



whose ΔG° is -30 kJ/mole (21). It is also close to the value of the following reaction:



whose ΔG° is -34.3 kJ/mole (33). In both reactions [15] and [16], bulk of the free energy changes are due to covalent bonding. Therefore, $\Delta G_{\text{chem}}^{\circ, \text{l}}$ of -32 kJ for CaOH⁺ ions adsorbing on silica is probably due to covalent bonding.

The adsorption model presented here is conceptually similar to the hemimicelle theory of Gaudin and Fuerstenau (34), who proposed that the specific adsorption of ionic surfactants may be due to the van der Waals interaction between adjacent hydrocarbon chains. Similarly, the adsorption of the singly-charged hydroxo-complexes may be favored due to the large value of $-\Delta G_{\text{chem}}^{\circ}$, the major part of which seems to come from the lateral interaction between adsorbed hydroxo-complexes. This may provide an explanation for the observation that CaOH⁺ and

MgOH⁺ ions adsorb on silica but not Ca²⁺ and Mg²⁺ ions. Inability of the latter to adsorb on silica may also be attributed to the unfavorable electrical repulsion between adsorbed species.

An interesting observation made in the present work was that activated silica does not exhibit hydration force. One possible explanation would be that the surface of activated silica is less strongly hydrated than that of unactivated silica. This view may be supported by the fact that the standard free energy of hydration ($\Delta G_{\text{hyd}}^{\circ}$) for Si⁴⁺ ions is -26,220 kJ/mole, while that of Ca²⁺ ion is -2,932 kJ/mole. These values were calculated from the formula (35): $\Delta G_{\text{hyd}}^{\circ} = -136.9z^2/r$, where z is the valence of the cation under consideration and r is its radius. The weakly hydrated surface may fail to support a strong H-bonded network of water molecules in the vicinity of silica, which is probably the origin of the primary hydration force observed with unactivated silica.

4.5 Conclusions

An atomic force microscope (AFM) was used to measure the surface forces between silica surfaces coated with self-assembled monolayers of sodium oleate. Since the surface charge of silica is of the same sign as the charge of oleate ions, it was necessary to reverse the charge of the silica substrate (which is a process referred to as “activation”) by adding CaCl₂ or MgCl₂ to the solution. The results showed the presence of the long-range hydrophobic forces that have not been considered in the DLVO theory. The hydrophobic forces become strongest at the pH where the concentration of CaOH⁺ or MgOH⁺ ions reaches maximum. It is, therefore, suggested that calcium hydroxy-oleate (CaOHOL) and magnesium hydroxy-oleate (MgOHOL) adsorb on the silica surface possibly *via* H-bonding.

The force measurements were also conducted in CaCl₂ and MgCl₂ solutions in the absence of sodium oleate. The results showed that the primary hydration forces, which are observed in Nanopure water, disappear completely at the pH where the concentration of the MgOH⁺ or CaOH⁺ ions reaches maximum. This may be attributed to the fact that the silica surface is covered by the hydroxo-complexes that are not as strongly hydrated as the silanol groups on the surface of

silica. Thus, the AFM force measurements showed that the singly charged hydroxo-complexes are the activators.

A theoretical model was developed to better understand the mechanisms involved in the adsorption of the singly charged hydroxo-complexes on silica. In the present work, the free energy of adsorption ($\Delta G_{\text{ad}}^{\circ}$) was back-calculated from the adsorption isotherm for the hydroxo-complexes on silica. From this, one can subtract the free energy of adsorption ($\Delta G_{\text{elec}}^{\circ}$) due to electrostatic interaction, which can be readily calculated from the surface potentials (ψ_0) of silica determined from the AFM force measurements, to obtain the free energy of adsorption ($\Delta G_{\text{chem}}^{\circ}$) due to chemical interaction. For the adsorption of CaOH^+ ions on silica, one obtains that $\Delta G_{\text{chem}}^{\circ} = -42$ kcal/mole. By subtracting the value of H-bonding of the hydroxo-complex with the silanol groups on the surface, which may be estimated to be -10 kcal/mole, one obtains -32 kcal/mole. This value is close to the free energy of polymerization of $\text{Ca}(\text{OH})_2(\text{aq})$ species, which is -30 kcal/mole. This finding suggests that the lateral interactions between neighboring adsorbates contribute most significantly to the adsorption of singly charged hydroxo-complexes on silica, followed by H-bonding with silanol groups and electrostatic interaction.

4.6 References

1. Israelachvili, J., and Pashley, R. M. *Nature* **300**, 341 (1982).
2. Herder, P. C., *J. Colloid Interface Sci.* **134**, 346 (1990).
3. Pashley, R. M., McGuiggan, P. M., Ninham, B. W. and Evans, D. F. *Science* **229**, 1088 (1985).
4. Pugh, R. J. and Rutland, M. *Proc. XX IMPC*, 583 (1997).
5. Fuerstenau, M. C., Elgillani, D. A. and Miller J. D. *Trans. AIME*, **247**, 11 (1970).
6. Fuerstenau, M. C. *AICHE Symposium Series.* **150**, 16 (1975)
7. Fuerstenau, M. C. and Cummins, W. F. *Trans. AIME*, **238**, 196 (1967).
8. Palmer, B. R., Gutierrez, G. B. and Fuerstenau, M. C. *Trans. AIME* **258**, 257 (1975).
9. James, R. O. and Healy, T. W. *J. Colloid Interface Sci*, **40**, no. 1, 65 (1972).
10. Schindler, P. W., Furst, B., Dick, R. and Wolf, P. U. *J. Colloid Interface Sci.* **55**, 469 (1976).
11. Ducker, W. A., Senden, T. J., and Pashley. R. M. *Nature* **353**, 239 (1993).
12. Grabbe, A. and Horn, R. G. *J. Colloid Interface Sci.* **157**, 375 (1993).
13. Yoon, R-H. and Subramanian, V. *J. Colloid Interface Sci.*, accepted for publication.
14. Peschel, G., Belouschek, P., Muller, M. M., Muller, R. M., and Konig, R. *Colloid Polymer Sci.* **260**, 444 (1982).
15. Yotsumoto, H., and Yoon, R-H. *J. Colloid Interface Sci.* **157**, 434 (1993).
16. Chapel, J. *Langmuir*, **10**, 4237 (1994).
17. Meagher, L. *J. Colloid Interface Sci.* **152**, No.1, 293 (1992)
18. Rabinovich, Ya. I., Derjaguin, B. V. and Churaev, N. V. *Adv. Colloid Interface Sci.* **16**, 63 (1982).
19. Senden, T. J. and Ducker, W. A. *Langmuir* **10**, 1003 (1994).
20. Rabinovich, Ya. I. and Yoon, R.-H. *Colloids and Surfaces*, **93**, 263 (1994).
21. *Standard Potentials in Aqueous Solutions* (A.J. Bard, R. Parsons and J. Jordan, Eds.), IUPAC, Marcel Dekker, Inc., New York (1985).
22. Derjaguin, B. V., and Landau, L. *USSR Acta Physicochim.* **14**, 633 (1941).
23. Verwey, E. J., and Overbeek, J. Th. G. *Theory of the Stability of Lyophobic Colloids* Elsevier, New York, 1947.
24. Chan, D. Y. C., Pashley, R. M., and White, L. R. *J. Colloid Interface Sci.* **77**, 283 (1980).

25. Cooke, S. R. B. and Digre, M. *AIME Trans.* **232**, 24 (1965).
26. Pashley, R. M. *J. Colloid Interface Sci.* **83**, 531 (1981).
27. Yotsumoto, H. and Yoon, R. H. *J. Colloid Interface Sci.* **157**, 426 (1993).
28. Fuerstenau, M.C. and Elgillani, D. A. *Trans. AIME* **235**, 405 (1966).
29. Fuerstenau, M.C., Rice, D. A., Somasundaran, P. and Fuerstenau, D. W. *Bull. Inst. Min. and Met, London.* **701**, 381 (1965).
30. Hunter, R. J. *Foundations of Colloid Science*, Calrendon Press, New York (1987).
31. Allen, L. H. and Matijevic, E. J. *Colloid and Interface Sci.* **31**, 287 (1969).
32. Vinnogradov, S. N. and Linnell, R. H. *Hydrogen Bonding*, Van Nostrand Reinhold Co. (1971).
33. Baes, C. F. and Mesmer, R. E. *The Hydrolysis of Cations*, John Wiley and Sons, Inc., 1976.
34. Gaudin, A. M. and Fuerstenau, D. W. *Trans. AIME* **202**, 958 (1955)
35. Harvey, K. B. and Porter, G. B. *Introduction to Physical Inorganic Chemistry*, Addison-Wesley Publishing Co., Inc., 1963.