TRANSIENT ANALYSIS IN PIPE NETWORKS

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Abstract

Power failure of pumps, sudden valve actions, and the operation of automatic control systems are all capable of generating high pressure waves in domestic water supply systems. These transient conditions resulting in high pressures can cause pipe failures by damaging valves and fittings. In this study, basic equations for solving transient analysis problems are derived using method of characteristics. Two example problems are presented. One, a single pipe system which is solved by developing an excel spreadsheet. Second, a pipe network problem is solved using transient analysis program called TRANSNET.

A transient analysis program is developed in Java. This program can handle suddenly-closing valves, gradually-closing valves, pump power failures and sudden demand changes at junctions. A maximum of four pipes can be present at a junction. A pipe network problem is solved using this java program and the results were found to be similar to that obtained from TRANSNET program. The code can be further extended, for example by developing java applets and graphical user interphase to make it more user friendly.

A two dimensional (2D) numerical model is developed using MATLAB to analyze gaseous cavitation in a single pipe system. The model is based on mathematical formulations proposed by Cannizzaro and Pezzinga (2005) and Pezzinga (2003). The model considers gaseous cavitation due to both thermic exhange between gas bubbles and surrounding liquid and during the process of gas release. The results from the model show that during transients, there is significant increase in fluid temperature along with high pressures. In literature pipe failures and noise problems in premise plumbing are attributed to gaseous cavitation.

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Notation

The following symbols are used in this report:

- *A* Cross sectional area of the pipe
- *C* Courant number
- *Ca* Concentration in gas phase
- *CP* Specific heat of air at constant pressure
- C_v Specific heat of air at constant volume
- C_w Molar concentration (*mol* / *L*) of the dissolved gas or aqueous phase molar concentration
- *c* Wave speed of pure liquid in an elastic pipe
- *E* Bulk modulus of the liquid
- *H* Piezometric head
- *h* Enthalpy per unit mass
- K_H Henry's constant ($atm.L/mol$)
- K_N Experimental positive constant (1/sec)
- *k* Ratio between specific heat at constant pressure and constant volume
- *M* Molecular mass of the gas
- *m* Mass of free gas per unit volume
- *n* Number of moles
- *p* Absolute pressure
- p_g Gas pressure above liquid at the beginning of evolution or solution process
- p_l Partial pressure exerted by the liquid
- p_s Equilibrium or saturation pressure
- *pv* Vapor pressure
- *Q* Discharge in the pipe
- *q* Heat transfer per unit mass and time
- *R* Universal gas constant
- R_0 Pipeline radius
- ℜ Initial Reynolds number
- *r* Distance from the axis
- T_0 Absolute liquid temperature assumed to be constant
- *S* Gas entropy
- *U* Internal energy
- *u* Velocity component in longitudinal direction
- u_* Friction velocity
- *V* Volume of gas
- V_l Volume of liquid
- V_V Volume of gas
- *y* Distance from wall
- *z* Elevation head
- ^α Void fraction

ρ_{g} Density of gas

- ρ_l Density of liquid
- $\sigma_{\rm x}$ Normal stress in longitudinal direction
- σ_r Normal stress in radial direction
- $\sigma_{\scriptscriptstyle{\theta}}$ Normal stress in angular direction
- ^τ Wall shear stress
- τ_w Wall shear stress
- v Kinematic viscosity

Chapter 1: Introduction

Devices such as valves, pumps and surge protection equipment exist in a pipe network. Power failure of pumps, sudden valve actions, and the operation of automatic control systems are all capable of generating high pressure waves in domestic water supply systems. These high pressures can cause pipe failures by damaging valves and fittings. Study of pressure and velocity variations under such circumstances is significant for placement of valves and other protection devices. In this study, the role of each of these devices in triggering transient conditions is studied. Analysis is performed on single and multiple pipe systems.

Transient analysis is also important to draw guidelines for future pipeline design standards. These will use true maximum loads (pressure and velocity) to select the appropriate components, rather than a notional factor of the mean operating pressure. This will lead to safer designs with less over-design, guaranteeing better system control and allowing unconventional solutions such as the omission of expensive protection devices. It will also reveal potential problems in the operation of the system at the design stage, at a much lower cost than during commissioning.

1.1 Literature Review

Most of the problems considering unsteady pure liquid flow in pipes are solved using a set of partial differential equations (Wylie and Streeter, 1993) which are discussed in detail in Chapter 2. These equations are valid only when the pressure is greater than the vapor pressure of liquid, and are solved numerically using the method of characteristics which was introduced by Streeter and Wylie (1967). But in many flow regimes, small amount of free gas is present in a liquid. When local pressure during transient drops below saturation pressure, the liquid releases free gas. If the pressure drops to vapor pressure, cavities are formed (Tullis et al., 1976). The former occurrence is called gaseous cavitation where as the latter occurrence is called vaporous cavitation.

Martin et al., (1976) developed a one-dimensional homogeneous bubbly model using a two step Lax-Wendroff scheme. Pressure wave propagation and interactions are handled well by Lax-Wendroff scheme by introducing a pseudo-viscosity term. The results produced by this model compare favorably with experiment than by using fixed grid method of characteristics. An analytical model was developed

by Wiggert and Sundquist (1979) to investigate gaseous cavitation using the method of characteristics. Gas release is assumed mainly due to difference in local unsteady pressure and saturation pressure. Increase in void fraction due to latent heat flow is not considered. Wylie (1984) investigated both gaseous and vapor cavitation using a discrete free gas model. Free gas is lumped at discrete computing locations and pure liquid is assumed in between these locations. Small void fraction and isothermal behavior of fluid are some of the assumptions made. Gaseous cavitation is simulated and it gave close results when compared with other methods. Pezzinga (1999) developed a 2D model, which computes frictional losses in pipes and pipe networks using instantaneous velocity profiles. The extreme values for pressure heads and pressure wave oscillations were well reproduced by this model.

Pezzinga (2004) adopted "second viscosity" to better explain energy dissipation during transient gaseous cavitation. Constant mass of free gas is assumed at constant temperature. Second viscosity or bulk viscosity coefficient accounts for other forms (other than frictional losses) of energy dissipation such as gas release and heat exchange between gas bubbles and surrounding liquid.

Cannizzaro and Pezzinga (2005) considered the effects of thermic exchange between gas bubbles and surrounding liquid and gas release and solution process separately to study energy dissipation during gaseous cavitation. Separate 2D models were considered. The results of numerical runs shows that 2d model with gas release allows for a good simulation of the experimental data.

1.2 Objective

The objectives of this research are to:

- 1) Study unsteady flow in pipes and pipe networks carrying pure liquid, including evaluation of pressure and velocity heads at nodes and junctions at different time intervals,
- 2) Develop a program using object oriented technology to analyze transients in pipes considering single phase flow, and
- 3) Study the effects of gaseous cavitation on fluid transients using equations developed by Cannizzaro and Pezzinga (2005).

1.3 Organization

This thesis is divided into five chapters. Chapter 1 includes a brief introduction to transients, review of literature, and objectives of the study. Basic equations of transient flow analysis in pipe networks are

discussed in Chapter 2. Two example problems are solved using excel spreadsheet to demonstrate the method of characteristics. Chapter 3 is devoted to use of object oriented technology for analyzing transient problems in a pipe network. Comparison is drawn between procedural language and object oriented approach of analyzing transients in a pipe network. Chapter 4 is about gaseous cavitation in pipes where energy dissipation due to gas release and solution process is studied. Here, thermal exchange between gas bubbles and surrounding liquid is also considered. A comprehensive model to obtain the amount of gas release is developed. Chapter 5 presents the summary of work presented in this thesis, and also discusses its potential application.

Chapter 2: Basic Equations of Transient Flow Analysis in Closed Conduits

2.1 Introduction

Initial studies on water hammer are done assuming single phase flow of fluid (Wylie et al., 1993). The method of characteristics is most widely used for modeling water hammer. First, the fundamental equations involved in water hammer analysis are discussed, following which two example problems are solved to highlight the analytical technique.

2.2 Unsteady Flow Equations

Study of transient flow includes fluid inertia and also elasticity or compressibility of the fluid and the conduit. The analysis of transient flow in either of these cases requires the application of Newton's second law which leads to the Euler equation as discussed below.

2.2.1 The Euler Equation

Consider a small cylindrical control volume of fluid at the pipe centerline as shown in figure 1.

Figure 1. Cylindrical fluid element with all forces shown.

The Euler equation (Wylie, et al.,1993) is derived by applying Newton's second law to this control volume.

$$
\sum F_s = ma_s = m\frac{dv}{dt} \tag{2.1}
$$

Where m is the fluid mass and *dt* $\frac{dv}{dt}$ is the total derivative of the fluid velocity.

Substituting the applied forces into Eq. 2.1., and writing mass in terms of density and volume results in

$$
pdA - \left(p + \frac{\partial p}{\partial s} ds\right) dA - dW \sin \theta - \tau \pi d(ds) = \frac{dW}{g} \frac{dv}{dt}
$$
\n(2.2)

If we divide by *dW* and on simplification yields

$$
-\frac{\partial p}{\partial s}ds dA \frac{1}{dW} - \sin \theta - \tau \pi d(ds) \frac{1}{dW} = \frac{1}{g} \frac{dv}{dt}
$$
(2.3)

We can further substitute *s z* ∂ $\sin \theta = \frac{\partial z}{\partial t}$ and $dW = \gamma ds (dA)$ which yields:

$$
-\frac{1}{\gamma}\frac{\partial p}{\partial s} - \frac{\partial z}{\partial s} - \tau \pi d \frac{1}{\gamma(dA)} = \frac{1}{g}\frac{dv}{dt}
$$
\n(2.4)

on further substituting of $dA = \frac{4}{4}$ $dA = \frac{\pi d^2}{4}$ we get:

$$
-\frac{1}{\gamma}\frac{\partial p}{\partial s} - \frac{\partial z}{\partial s} - \frac{4\tau}{\gamma d} = \frac{1}{g}\frac{dv}{dt}
$$
(2.5)

If we expand the cross section of the element to fill the pipe cross section and introduce average velocity

V, we obtain:

$$
-\frac{1}{\gamma}\frac{\partial p}{\partial s} - \frac{\partial z}{\partial s} - \frac{4\tau_0}{\gamma D} = \frac{1}{g}\frac{dV}{dt}
$$
(2.6)

Where: τ_0 is the shear stress at the wall. Larock et al., (2000) expressed τ_0 as:

$$
\tau_0 = \frac{1}{8} f \rho V |V| \tag{2.7}
$$

Where f is the Darcy-Weisbach friction factor, ρ is the density of the fluid and *V* is its velocity.

After substituting Eq. 2.7 in Eq. 2.6 and under the assumption that the local elevation of the pipe can be described solely as a function of location s, we obtain the Euler equation of motion:

$$
\frac{dV}{dt} + \frac{1}{\rho} \frac{\partial p}{\partial s} + g \frac{dz}{ds} + \frac{f}{2D} V|V| = 0
$$

Or,

$$
\frac{\partial V}{\partial t} + V \frac{\partial V}{\partial s} + \frac{1}{\rho} \frac{\partial p}{\partial s} + g \sin \alpha + \frac{f}{2D} V|V| = 0
$$
\n(2.8)

2.2.2 Conservation of Mass

Applying conservation of mass to a control volume that coincides with the interior of the pipe and is of

length ds: See figure 2.

Figure 2. Control volume coinciding with the interior surface of the pipe

In above figure, ρ is the density of the fluid, A is the cross section area of control volume and V is the velocity of the fluid.

From Wylie, et al., (1993):

$$
\rho AV - \left[\rho AV + \frac{\partial}{\partial s}(\rho AV)ds\right] = \frac{\partial}{\partial t}(\rho Ads)
$$
\n(2.9)

Or
$$
-\frac{\partial}{\partial s}(\rho AV)ds = \frac{\partial}{\partial t}(\rho Ads)
$$
 (2.10)

Expanding the parenthesis of Eq. 2.10 yield:

$$
-\left(\rho A \frac{\partial V}{\partial s} ds + \rho V \frac{\partial A}{\partial s} ds + AV \frac{\partial \rho}{\partial s} ds\right) = \rho A \frac{\partial}{\partial t} (ds) + \rho ds \frac{\partial A}{\partial t} + Ads \frac{\partial \rho}{\partial t}
$$
(2.11)

Dividing both sides by the control volume mass ^ρ*Ads*,

$$
-\left(\frac{\partial V}{\partial s} + \frac{1}{A}V\frac{\partial A}{\partial s} + \frac{1}{\rho}V\frac{\partial \rho}{\partial s}\right) = \frac{1}{ds}\frac{\partial}{\partial t}(ds) + \frac{1}{A}\frac{\partial A}{\partial t} + \frac{1}{\rho}\frac{\partial \rho}{\partial t}
$$
(2.12)

Regrouping above Eq we get:

$$
\frac{1}{\rho} \left(\frac{\partial \rho}{\partial t} + V \frac{\partial \rho}{\partial s} \right) + \frac{1}{A} \left(\frac{\partial A}{\partial t} + V \frac{\partial A}{\partial s} \right) + \frac{1}{ds} \frac{\partial}{\partial t} (ds) + \frac{\partial V}{\partial s} = 0
$$
\n(2.13)

Recognizing that $\frac{\partial \rho}{\partial t} + V \frac{\partial \rho}{\partial s} = \frac{d\rho}{dt}$ *s V* $\frac{\partial \rho}{\partial t} + V \frac{\partial \rho}{\partial s} = \frac{d\rho}{dt}$ ∂ $\frac{\partial \rho}{\partial t} + V \frac{\partial \rho}{\partial t} = \frac{d\rho}{dt}$ and *dt dA s* $V\frac{\partial A}{\partial r}$ $\frac{A}{\partial t} + V \frac{\partial A}{\partial s} =$ ∂ $\frac{\partial A}{\partial x} + V \frac{\partial A}{\partial y} = \frac{dA}{dt}$, Eq. 2.13 becomes

$$
\frac{1}{\rho} \frac{d\rho}{dt} + \frac{1}{A} \frac{dA}{dt} + \frac{1}{ds} \frac{d}{dt} (ds) + \frac{\partial V}{\partial s} = 0
$$
\n(2.14)\n
\nTerms: T(1) T(2) T(3) T(4)

Let's split this Eq. 2.14 in various terms as below, we get:

$$
T(1): \quad \frac{1}{\rho} \frac{d\rho}{dt}
$$

Terms:

Bulk modulus of elasticity for a liquid (K) is expressed as (Larock et al., 2000):

$$
K = -\frac{dp}{d\nu/\nu} = \frac{dp}{d\rho/\rho}, \text{ Now } \mathbf{T(1)} \text{ becomes:}
$$
\n
$$
\frac{1}{\rho} \frac{d\rho}{dt} = \frac{1}{K} \frac{dp}{dt}
$$
\n
$$
\mathbf{T(2)} = \frac{1}{K} \frac{dA}{dt}
$$
\n(2.15)

$$
T(2): \ \frac{1}{A}\frac{dA}{dt}
$$

The above term can be expressed as (Larock et al., 2000):

$$
\frac{1}{A}\frac{dA}{dt} = (1 - \mu^2)\frac{D}{eE}\frac{dp}{dt}
$$
\n(2.16)

Where,

 μ = Poisson's ratio (of the pipe)

e = Pipe wall thickness

 $E = Young's Modulus of the pipe$

$$
\mathbf{T(3)}: \frac{1}{ds}\frac{d}{dt}(ds)
$$

Considering longitudinal expansion of the pipe (Larock et al., 2000):

$$
d(ds) = d\varepsilon_1 ds \tag{2.17}
$$

Where ε_1 is the strain along the pipe axis.

For all the buried pipes, axial movement is restrained. So differential change in strain along pipe axis

$$
(d\varepsilon_1)=0.
$$

Therefore, $\frac{1}{ds} \frac{d}{dt} (ds) = 0$ *ds*

Making all these substitutions in Eq. 2.14 yields,

$$
\frac{1}{K}\frac{dp}{dt} + (1 - \mu^2)\frac{D}{eE}\frac{dp}{dt} + \frac{\partial V}{\partial s} = 0
$$
\n(2.18)

$$
\frac{dp}{dt}\left(\frac{1}{K} + (1 - \mu^2)\frac{D}{eE}\right) + \frac{\partial V}{\partial s} = 0\tag{2.19}
$$

Wave speed (Larock et al., 2000) can be defined as the time taken by the pressure wave generated by instantaneous change in velocity to propogate from one point to another in a closed conduit. Wave $speed(c)$ can be expressed as:

$$
c^2 \rho \left[\frac{1}{K} + \frac{D}{e} \left(\frac{1 - \mu^2}{E} \right) \right] = 1 \tag{2.20}
$$

Hence by using Eq. 2.20 we can write Eq. 2.19 as,

$$
\frac{dp}{dt}\left(\frac{1}{c^2\rho}\right) + \frac{\partial V}{\partial s} = 0\tag{2.21}
$$

$$
\frac{dp}{dt} + c^2 \rho \frac{\partial V}{\partial s} = 0
$$
\n(2.22)

Or,

$$
\frac{\partial p}{\partial t} + V \frac{\partial p}{\partial s} + c^2 \rho \frac{\partial V}{\partial s} = 0
$$
\n(2.23)

2.3 Method of characteristics

The significance of method of characteristics is the successful replacement of a pair of *partial* differential equations by an equivalent set of *ordinary* differential equations. The method of characteristics is developed from assuming that the Eq. 2.9 and Eq. 2.23 can be replaced by a linear combination of themselves (Wylie, et al.,1993)

$$
\lambda \left[\frac{\partial V}{\partial t} + V \frac{\partial V}{\partial s} + \frac{1}{\rho} \frac{\partial p}{\partial s} + g \sin \alpha + \frac{f}{2D} V |V| \right] + \frac{\partial p}{\partial t} + V \frac{\partial p}{\partial s} + c^2 \rho \frac{\partial V}{\partial s} = 0
$$
\n(2.24)

Rearranging the terms in above, we have

$$
\lambda \left[\frac{\partial V}{\partial t} + \left(V + \frac{c^2 \rho}{\lambda} \right) \frac{\partial V}{\partial s} \right] + \left[\frac{\partial p}{\partial t} + \left(V + \frac{\lambda}{\rho} \right) \frac{\partial p}{\partial s} \right] + \lambda \left[g \sin \alpha + \frac{f}{2D} V |V| \right] = 0 \tag{2.25}
$$

Assuming,

$$
\frac{ds}{dt} = V + \frac{c^2 \rho}{\lambda} = V + \frac{\lambda}{\rho}
$$
\n(2.26)

The equality in Eq. 2.26 leads to

$$
\lambda = \pm c\rho \tag{2.27}
$$

Which when substituted back into Eq. 2.26 leads to

$$
\frac{ds}{dt} = V \pm c \tag{2.28}
$$

Use of proper positive and negative signs of lambda permits writing of the following sets of ordinary

differential equations (Wylie, et al.,1993):

$$
C^{+}\begin{cases} \frac{dV}{dt} + \frac{1}{\rho c} \frac{dp}{dt} + g \sin \alpha + \frac{fV|V|}{2D} = 0\\ \frac{ds}{dt} = V + c \end{cases}
$$
(2.29)

and

$$
C^{-}\begin{cases} \frac{dV}{dt} - \frac{1}{\rho c} \frac{dp}{dt} + g \sin \alpha + \frac{fV|V|}{2D} = 0\\ \frac{ds}{dt} = V - c \end{cases}
$$
(2.30)

Finally we replace the pressure in favor of total head using (Wylie, et al.,1993) $p = \gamma(H - z)$ and we assume that the entire pipe network is in the same horizontal plane. i.e., $\sin \alpha = 0$. The new set of characteristic equations can be written as (Wylie, et al.,1993)

$$
C^+ : \frac{dV}{dt} + \frac{g}{c} \frac{dH}{dt} + \frac{fV|V|}{2D} = 0
$$
 Only when $\frac{ds}{dt} = V + c$ (2.31)

$$
C^{-} : \frac{dV}{dt} - \frac{g}{c} \frac{dH}{dt} + \frac{fV|V|}{2D} = 0
$$
 Only when $\frac{ds}{dt} = V - c$ (2.32)

2.3.1 Finite Difference Approximation:

We assume that the characteristic curves can be approximated as straight lines over each single ∆*t* interval. But as a result the slopes of C^+ and C^- characteristic lines will not be the same. This can be seen in Fig3 (Larock et al.,2000) as shown below

Figure 3. Parameters in the interpolation procedure (Reproduced from Larock et al., 2000)

We can see that the characteristics intersecting at P no longer pass through the grid points L_e and R_i. But instead they pass through the points L and R somewhere between L_e and R_i . Hence, the finite difference approximations to Eqs. 2.31 and 2.32 become

$$
\frac{V_P - V_L}{\Delta t} + \frac{g}{c} \frac{H_P - H_L}{\Delta t} + \frac{f}{2D} V_L |V_L| = 0
$$
\n(2.33)

$$
\frac{V_P - V_R}{\Delta t} - \frac{g}{c} \frac{H_P - H_R}{\Delta t} + \frac{f}{2D} V_R |V_R| = 0
$$
\n(2.34)

This results in 4 new unknowns $V_L H_L$, V_R and H_R . But we overcome this problem by choosing Δt so that the point L is near L_e and R is near R_{i} . Now linear interpolation becomes an accurate way to evaluate the values of H and V at points L and R.

Along C^+ characteristic (Larock et al., 2000),

$$
\frac{\Delta x}{\Delta s} = \frac{V_L - V_C}{V_{Le} - V_C} = \frac{H_L - H_C}{H_{Le} - H_C}
$$
\n(2.35)

with

$$
\frac{\Delta x}{\Delta t} = \frac{c + V_L}{1}
$$
 (2.36)

Solving above two equations for V_L and H_L yields,

$$
V_L = (V_{Le} - V_c) \frac{\Delta x}{\Delta s} + V_c \quad \text{and} \quad H_L = (H_{Le} - H_c) \frac{\Delta x}{\Delta s} + H_c \tag{2.37}
$$

Replacing ∆*x* in these equations using the same relation for *t x* ∆ $\frac{\Delta x}{\Delta x}$ now produces (Larock et al., 2000)

$$
V_{L} = \frac{V_{C} + c \frac{\Delta t}{\Delta s} (V_{Le} - V_{c})}{1 - \frac{\Delta t}{\Delta s} (V_{Le} - V_{c})}
$$
(2.38)

and

$$
H_L = H_C + \frac{\Delta t}{\Delta s} (H_{Le} - H_C)(c + V_L)
$$
\n(2.39)

Similarly along the C^- characteristic (Larock et al., 2000),

$$
V_R = \frac{V_c + c\frac{\Delta t}{\Delta s}(V_{Ri} - V_c)}{1 - \frac{\Delta t}{\Delta s}(V_{Ri} - V_c)}
$$
(2.40)

And

$$
H_R = H_C + \frac{\Delta t}{\Delta s} (H_{Ri} - H_C)(c - V_R)
$$
\n(2.41)

Note that $\frac{\Delta t}{\Delta s} (V_{Le} - V_c)$ is on the order of $c + V$ *V* + , which is very small compared to 1, hence the second

terms in the denominator of Eqs. 2.38 and 2.40 can be neglected. This would result in

$$
V_L = V_C + c \frac{\Delta t}{\Delta s} (V_{Le} - V_c) \tag{2.42}
$$

and

$$
V_R = V_C + c \frac{\Delta t}{\Delta s} (V_{Ri} - V_c) \tag{2.43}
$$

As we now have the known values for V_L *H_L*, V_R and *H_R*, we can solve the Eqs. 2.33 and 2.34

simultaneously for velocity and head at point P.

The solutions are (Larock et al., 2000):

$$
V_P = \frac{1}{2} \bigg[(V_L + V_R) + \frac{g}{c} (H_L - H_R) - \frac{f \Delta t}{2D} (V_L |V_L| + V_R |V_R|) \bigg] \tag{2.44}
$$

$$
H_P = \frac{1}{2} \left[\frac{c}{g} (V_L - V_R) + (H_L + H_R) - \frac{c}{g} \frac{f \Delta t}{2D} (V_L |V_L| - V_R |V_R|) \right]
$$
(2.45)

2.4. Summary of Equations

The two basic equations of fluid flow in closed conduits are a pair of quasilinear partial differential equations (Wylie et al., 1993):

$$
\frac{\partial H}{\partial x} + \frac{1}{g} \frac{\partial V}{\partial t} + \frac{\lambda}{2gD} V|V| = 0 \quad \text{(Momentum equation)}
$$
\n
$$
\frac{\partial H}{\partial t} + \frac{c^2}{g} \frac{\partial V}{\partial x} = 0 \quad \text{(Continuity equation)}
$$
\n(2.47)

Using method of characteristics the above pair of partial differential equations can be replaced by an equivalent set (two pairs) of ordinary differential equations which are grouped and identified as C^+ and [−] *C* characteristic (Wylie et al., 1993)

$$
C^{+}\begin{cases} \frac{dV}{dt} + \frac{g}{c} \frac{dH}{dt} + g \sin \alpha + \frac{\lambda V|V|}{2D} = 0\\ \frac{dx}{dt} = c \end{cases}
$$
\n
$$
C^{-}\begin{cases} \frac{dV}{dt} - \frac{g}{c} \frac{dH}{dt} + \frac{\lambda V|V|}{2D} = 0\\ \frac{dx}{dt} = -c \end{cases}
$$
\n(2.49)

Using finite difference scheme, the above equations can be expressed as (Larock et al., 2000):

$$
C^{+}: H_{P} = H_{A} - B(Q_{P} - Q_{A}) - RQ_{P} |Q_{A}|
$$

\n
$$
C^{-}: H_{P} = H_{B} + B(Q_{P} - Q_{B}) + RQ_{P} |Q_{B}|
$$
\n(2.50)

where,

$$
B = \frac{c}{gA} \text{ And } R = \frac{\lambda \Delta x}{2gDA^2}
$$

The above set of characteristic equations can be further simplified as (Larock et al., 2000):

$$
C^*: H_P = C_P - B_P Q_P
$$

\n
$$
C^-: H_P = C_M + B_M Q_P
$$
\n(2.51)

Where,
$$
C_p = H_A + BQ_A
$$
, $B_p = B + RQ_A$ and

(2.52)

$$
C_{_M} = H_{_B} - BQ_{_B}, B_{_M} = B + RQ_{_B}
$$

 C^+ and C^- equations are solved to compute the velocity V_p and the head H_p at each section j. This results in:

 Figure 4. Characterstics shown on a finite difference grid

$$
V_{j,i+1} = \frac{1}{2} \bigg[(V_{j-1,i} + V_{j+1,i}) + \frac{g}{c} (H_{j-1,i} - H_{j+1,i}) - \frac{\lambda \Delta t}{D} (V_{j-1,i} | V_{j-1,i} | + V_{j+1,i} | V_{j+1,i} |) \bigg] \tag{2.53}
$$

And

$$
H_{j,i+1} = \frac{1}{2} \left[H_{j-1,i} + H_{j+1,i} + \frac{c}{g} (V_{j-1,i} - V_{j+1,i}) - \frac{c}{g} \frac{\lambda \Delta t}{D} (V_{j-1,i} | V_{j-1,i} | - V_{j+1,i} | V_{j+1,i} |) \right]
$$
(2.54)

The basic water hammer conditions for a single pipeline provide the fundamental elements that are necessary for the treatment of more complex piping systems. When the system contains more than one pipeline, the interior sections of each pipeline are treated independently of other parts of the system at each instant of time. The end conditions for each pipeline must interface with adjoining pipelines or with other boundary elements. Again each boundary condition is treated independently of other parts of the system.

2.5 Boundary Conditions

Analyzing the following boundary conditions is important part of water hammer study:

2.5.1 Reservoir

Consider a single pipe system with reservoir at the upstream of the pipe. Hence, a C^- characteristic can be drawn from a node closest to the reservoir towards the junction. Therefore from Eq. 2.51:

$$
C^{-}:H_{P}=C_{M}+B_{M}Q_{P}
$$
\n
$$
(2.55)
$$

Here, H_p is the head at the reservoir junction and is assumed constant.

Hence, $H_p = H_R$ and discharge at this junction can be obtained as:

$$
Q_P = \frac{H_R - C_M}{B_M} \tag{2.56}
$$

2.5.2 Valve

Consider a single pipe system with reservoir upstream and valve at downstream. From the node closest to the valve and towards its left, a C^+ characteristic can be drawn towards the valve junction. Therefore from Eq. 2.51

$$
C^*: H_p = C_p - B_p Q_p. \tag{2.57}
$$

During steady state flow through the valve, orifice equation can be applied as (Watters, 1984):

$$
Q_P = (C_d A_v) \sqrt{2gH_P} \tag{2.58}
$$

Solving above two equations, we obtain

$$
Q_P = -gB_P(C_d A_v)^2 + \sqrt{(gB_P(C_d A_v)^2)^2 + 2gC_P(C_d A_v)^2}
$$
\n(2.59)

2.5.3 Pumps and Turbines

Pump/Turbine characteristic curves are needed for finding the water hammer caused by them. Let H_s be the head at suction side of the pump. Let H_D be the head at discharge side of the pump. The difference in head (H_D - H_S) gives the total head (H_p) developed by the pump.

$$
H_p = f(n, Q_p) \tag{2.60}
$$

The above equation (Chaudhry, 1988) is a typical form of pump/turbine characteristic curve. Where, *n* is the pump speed and Q_P is the pump discharge.

Also in this case, the characteristic equations are:

$$
C^*: H_s = C_p - B_p Q_p
$$

\n
$$
C^-: H_p = C_M + B_M Q_p
$$
\n(2.61)

Solving above three basic equations, we obtain discharge and head at the pump during transients.

2.5.4 Junctions

Analyzing network junctions is very important part of transient analysis. Consider a four pipe junction.

Figure 5. Four pipe junction with valve downstream of one pipe

From figure 5, junction J is the boundary condition for the four pipes. But, if we consider each pipe separately, the junction is present downstream to the pipes P1 and P3, and it is upstream to the pipes P2 and P4 (with respect to the initial flow direction). Hence for pipes where junction is downstream, C^+ line is drawn from pipe nodes towards the junction. Similarly, for pipes where junction is upstream a C- line is drawn.

The equations (Larock et al., 2000) for the four characteristic lines are:

Pipe 1, C⁺ $V_{p_1} = C_1 - C_2 H_{p_1}$ (2.62)

$$
\text{Pipe 2, C} \qquad V_{p_2} = C_3 + C_4 H_{p_2} \tag{2.63}
$$

$$
\text{Pipe 3, } C^+ \quad V_{p_3} = C_5 - C_6 H_{p_3} \tag{2.64}
$$

Pipe 4, C'
$$
V_{p_4} = C_7 + C_8 H_{p_4}
$$
 (2.65)

In above equations,

$$
C_2 = \frac{g}{c_1}, C_4 = \frac{g}{c_2}, C_6 = \frac{g}{c_3}, C_8 = \frac{g}{c_4} \text{ and } (2.66)
$$

$$
C_1 = (V_{j-1,i} + \frac{g}{c_1} H_{j-1,i}) - \frac{\lambda_1 \Delta t}{2D} V_{j-1,i} |V_{j-1,i}|
$$
\n(2.67)

$$
C_{3} = (V_{j+1,i} - \frac{g}{c_{2}} H_{j+1,i}) - \frac{\lambda_{2} \Delta t}{2D} V_{j+1,i} |V_{j+1,i}|
$$

\n
$$
C_{5} = (V_{j-1,i} + \frac{g}{c_{3}} H_{j-1,i}) - \frac{\lambda_{3} \Delta t}{2D} V_{j-1,i} |V_{j-1,i}|
$$

\n
$$
C_{7} = (V_{j+1,i} + \frac{g}{c_{4}} H_{j+1,i}) - \frac{\lambda_{4} \Delta t}{2D} V_{j+1,i} |V_{j+1,i}|
$$
\n(2.68), (2.69) and (2.70)

Conservation of Mass
$$
V_{P_1}A_1 + V_{P_3}A_3 = V_{P_4}A_4 + V_{P_2}A_2
$$
 (2.71)

Work Energy neglecting head loss at Junction: $H_{P_1} = H_{P_2} = H_{P_3} = H_{P_4}$ (2.72)

Solving these equations for the head values at the junction (Larock et.al, 2000):

$$
H_{P_1} = H_{P_2} = H_{P_3} = H_{P_4} = \frac{C_1 A_1 + C_5 A_3 - C_3 A_2 - C_7 A_4}{C_2 A_1 + C_4 A_2 + C_6 A_3 + C_8 A_4}
$$
\n(2.73)

Back substitution of these heads into Eqs. 2.62, 2.63, 2.64 and 2.65 yields the velocities in each of the pipes.

Consider a distribution network which is in steady state. Let the valve at the downstream of pipe2 be closed suddenly (as shown in Fig 5) at time $t = 0$ sec. Each pipe connected to the junction is divided into N number of sections. A rectangular grid can be drawn showing N sections each of length ∆*s* on xaxis and ∆*t* increments on y-axis. At each ∆*t* time step increments head and discharge at each node are calculated using method of characteristics (finite difference approximation).

At the node closest to the valve at upstream end of pipe 2, head and discharge are calculated using valve boundary conditions and at the node closest to the junction, boundary conditions of the junction are used to evaluate head and discharge.

Now consider pipe3, which is an inflow pipe to the Junction:

At the node (of pipe 3) closest to the junction, a negative characteristic from the junction and positive characteristic from the left adjacent node are drawn to meet at the first time step. And subsequently head and discharge are evaluated. So, is the case with all other pipes attached to the junction.

2.6 Transient Flow Analysis

2.6.1 Single pipe system

Consider a single pipe system having reservoir at upstream end (see figure 6) and a valve at the downstream end. The valve is subjected to sudden closure (at time $t = 0$ sec) and transients in the system is calculated for 2.5 sec.

Figure 6. Single pipeline with reservoir upstream and valve downstream

In figure 6, 'A' represents the reservoir and 'V' represents the valve.

The data required to carry out transient analysis using method of characteristics is shown below.

Length of the pipe $= 340$ m Diameter of the pipe $= 0.1$ m Darcey friction factor $(f) = 0.025$ Wave speed in the pipe $(a) = 1200$ m/s Valve opening $CdA = 0.00015$ m² (amount of valve opened represented as a factor of pipe area) Reservoir head $= 120$ m

Table 1. Transient analysis of single pipe using method of characteristics

QM, HM represent discharge and the pressure head at the middle of the pipe

QV, HV represent discharge and the pressure head at the upstream of the valve

From above table, the maximum pressure head reached as a result of valve closure = **232m**

Using results given in Table 1, the temporal variation in pressure head during transient condition is shown in figure 7.

 Figure 7. Single pipe with reservoir upstream and valve downstream

From above plots, it is evident that just upstream of the valve (i.e., at the end of the pipe), high pressures are maintained for a longer duration as compared to the middle of the pipe. Hence, pressure surge devices should be placed at pipe joints to avoid failures because they are more susceptible to high pressures.

Analytical solution using Joukowski relation

Joukowski relation (Popescu, 2003) is given by:

$$
\Delta H = \pm \left(\frac{a}{g}\right) \Delta V \tag{2.74}
$$

Where,

 ΔH = change in pressure head due to valve closure.

 $a = Wave$ speed in the pipe.

- $g =$ acceleration due to gravity = 9.8 m/s^2
- ΔV = change in fluid velocity = 0.91 m/s

Using the same example problem as above, we can find the change in pressure head due to sudden valve closure using Joukowski relation as:

$$
\Delta H = \pm \left(\frac{1200}{9.806}\right) 0.91 = 111.4 \text{ m}
$$

Steady state pressure head at the valve $=$ Head in the Reservoir $= 120$ m (neglecting friction losses in the pipe)

Hence, maximum pressure head at the valve $= 111.4 \text{ m} + 120 \text{ m} = 231 \text{ m}$

Results

The maximum pressure head at the valve from method of characteristics $= 232$ m The maximum pressure head at the valve from Joukowski relation $= 231$ m

The values obtained from both these analysis are very close. This can be explained by the simplicity of the pipe system we have chosen (single pipe with reservoir at upstream end and valve at downstream end). Simple expressions, such as the Joukowski relation are only applicable under restricted circumstances. The two most important restrictions are:

- 1) There should be no head loss resulting from friction.
- 2) There should be no wave reflections (i.e., there is no interaction between devices or boundary conditions in the system).

The above restrictions are narrowly met by the single pipe system discussed above. Hence, the results are comparable.

In above section, transient analysis in a single pipe system is solved using method of characteristics and found that the resulting maximum pressure head value and that obtained by Joukowski relation gave close results.

Most of the times, transient analysis performed on single pipe sytem is useful only for experimental purposes. In practice, most of the distribution systems have multiple pipes and transient flow analysis plays a significant role at pipe joints and at locations where pressure control devices such as valves and pumps exist.

2.6.2 Network Distribution

Consider the network given below (Larock et al., 2000). The Hazen-Williams roughness coefficient is 120 for all pipes. This network experiences a transient that is caused by the sudden closure of a valve at the downstream end of pipe 5. Wave speed is 2850 ft/s for all the pipes. Transient analysis is obtained for this network.

Figure 8. Network distribution system

The pipe information is given in Table 2.

Table 2. Pipe data

The node information is given in Table 3.

Node	Elevation(ft)	Demand(gal/min)
	3800	475
2	3830	317
	3370	790
	4050	
5	4000	
	4010	

Table 3. Node data

Steady state analysis:

First, the above network distribution system is solved for steady state analysis.Following this, transient analysis is performed. A program called "NETWK" was adopted for steady state analysis of above network. "NETWK" is a FORTRAN 95 code (Larock et al., 2000) which can perform steady state analysis on complex pipe systems. The results are shown in Table 4.

Pipe Discharge(gpm) Velocity(fps) Head Loss(ft) 1 340.13 0.96 1.32 2 272.99 1.74 15.7 3 138.11 0.88 1.79 4 1109.95 3.15 17.49 5 -458.06 -5.2 66.89 6 | 1699.93 | 3.54 | 9.64

 Table 4. Pipe data-results of steady flow analysis

In Table 4, the negative values for velocity and discharge in pipe 5 show that the actual direction of steady flow is opposite to the initial assumption as shown in figure 8.

Table 5. Node data-results of steady flow analysis

Node	Demand(gpm)	Head(ft)
	475	398.68
2	317	384.38
3	790	826.89
	-340.1	150
5	458.1	130
հ	-1699.9	214.03

In Table 5, the negative values for demand indicate that the flow direction is into the node. Detailed steady state flow output obtained from "NETWK" program is shown in Appendix A-1. This output is used as one of the input files for the transient analysis.

Transient Analysis:

In order to perform transient analysis, "TRANSNET" program is adopted. "TRANSNET" is the FORTRAN 95 code (Larock et al., 2000) which can perform transient analysis caused by sudden valve closures or pump failures. In pipe network shown in figure 8, the valve located on pipe 5 is closed suddenly at time 0.0 s and below is the summary of results for pipe 5 at time 7.73 s when column separation occurs.

Pipe	X/L	Head(ft)	Velocity(fps)
5		822	-1.1
	0.2	693	-1.14
	0.4	493	-0.56
	0.6	306	-0.35
	0.8	120	-0.27
	$4*$	-40	

Table 6. Results of transient flow analysis at pipe 5(t =7.73 s)

* at location of the valve

From Table 6, it can be seen that column separation has occurred at location 1 of pipe 5 where pressure head has gone below 0. Detailed transient analysis output table is shown in Appendix A-2. The minimum head attained during transients is -40 ft is already shown in Table 6 and the maximum head is also in pipe 5 at location 0. Its value is 822 ft at time 7.73 sec.

In Appendix B, the author has compared the results of transient analysis with that of WHAMO model. WHAMO (water hammer and model oscillation) model is a hydraulic transient flow model based on an implicit finite difference method and the results where found comparable.

In this chapter, the author has derived the basic equations for solving transient analysis problems using method of characteristics. Two problems are solved. One is a single pipe system which is solved using spreadsheet, and second is a pipe network which is solved using transient analysis program.

In literature most of the transients programs are in FORTRAN language. In chapter 3, the author used Java to code transient analysis program using an object oriented approach. The results of Java program are compared with that obtained from "TRANSNET" program.

CHAPTER 3: An Object Oriented Approach for Transient Analysis in Water Distribution Systems using JAVA programming

3.1 Introduction

Most of the algorithms in computational hydraulics discipline are written in procedural language (FORTRAN, Pascal and C). Procedural programming was found to be adequate for coding moderately extensive programs until 90's (Madan, 2004). In procedural programming, the strategy is based on dividing the computational task into smaller groups termed as functions, procedures or subroutines which perform well-defined operations on their input arguments and have well defined interfaces to other subprograms in the main program.

However, procedural programming approach can get challenging when the code needs to be extended for enhancing the scope of the ptogram. A detailed knowledge of the program is required to work on a small part of the code and poor equivalence between program variables and physical entities further makes it difficult. Integrity of data is another area of concern in procedural programs because, the emphasis is on functions and data is considered secondary. All the functions of a program have access to data and as a result data is highly susceptible to get corrupted when dealing with complex programs. In addition, there are difficulties related to reusability and maintenance of code as procedural programs are platform and version dependent.

Object oriented programming is developed with the objective of addressing some of the typical difficulties associated with procedural programming approach (Madan, 2004). The concept of Object-oriented programming (OOP) began in 1970s and found its first convivial concretization with the Small Talk (Fenves, 1992) language at the beginning of the eighties. Object oriented applications in scientific computing appeared around 1990 and since then its use has been spreading.

Madan (2004) discussed the modeling and design of structural analysis programs using OOP approach. He developed a C_{++} program for matrix analysis of a space frame. Krishnamoorthy et al., (2002) used object oriented approach to design and develop a genetic algorithm library
for solving optimization problems. They later solved a space truss optimization problem by implementing this library. A window-based finite-element analysis program was developed by Ju and Hosain (1996). Graphical finite element objects developed using C++ language were implemented in the program.

Liu et al., (2003) developed an object oriented framework for structural analysis and design. They implemented this framework in optimal design of energy dissipation device (EDD) configurations for best structural performance under earthquake conditions. Tisdale (1996) followed object oriented approach to organize South Florida hydrologic system information as object, dynamic and functional modules. These modules can be used for hydrologic software development.

A real-time flood forecasting system called RIBS (Real time Interactive Basin Simulator) is developed by Garrote and Becchi (1997) using object-oriented framework. Primary applications of RIBS include rainfall forecasting, estimate runoff generating potential of watersheds, and to develop hydrographs. Solomatine (1996) presented a water distribution modeling system (HIS) as an example of object oriented design. He developed two models namely, LinHIS – to solve for linear system of pipes, and NetHIS – to solve for a distributed network of pipes. In both cases, flow is considered either steady or quasi-steady.

Figure 9 depicts the basic difference between the organization of an object –oriented program and a procedure-oriented program. In procedure-oriented program, there are different functions, each of which perform specific action by accessing the data common to all functions whereas in object-oriented program, each object has its own data and methods.

Figure 9. Different types of programming methods (reproduced from www.agiledata.org)

3.2 Object-oriented programming in Java

A Java program describes a community of objects arranged to interact in well-defined ways for a common purpose. None of the objects is sufficient on its own. Each object provides specific services required by other objects in the community to fulfill the program's promise. When an object requires a specific service, it sends a request (called a message) to another object capable of providing that service. The object that receives the message responds by performing actions that often involve additional messages being sent to other objects. This results in a vibrant cascade of messages among a network of objects.

Next the author discusses the use of Java in developing a program for transient analysis in a pipe network. Each of the pipe network elements is represented by a separate class. A detailed description of these classes and primary methods they use are discussed in this chapter. The program can handle pump power failure, suddenly-closing valves, one gradually-closing valve, and sudden demand changes at junctions. A maximum of four pipes can be present at a junction. The program solves for pressure heads and velocity values at any section of the network system after inception of transients. The method of characteristics is used to obtain new pressure head and velocity values from current time step values. The program implements analytical equations discussed in Chapter 1. Figure 10 shows various steps involved in the transient analysis program.

Figure 10. Flow chart to solve for pressure and velocity heads using method of characteristics

The pipe network problem discussed in Section 2.6.2 in Chapter 2 is solved using Java program and results are compared with TRANSNET program output are discussed next.

3.3 Object oriented design

Unified Modeling Language (UML) is widely used tool for object oriented analysis and design. UML diagrams are used to visualize a system by showing interrelationships between classes. Following UML class diagram is used by the JAVA code to solve above discussed pipe network problem:

Figure 11. UML class diagram for Transient analysis program

This symbol represents '*composition relationship'* between any two classes. In above class diagram, Pressure Analyzer *is composed of* several junctions, and many pipes. We can also say that pipes and junctions are *part of* Pressure Analyzer class.

This symbol represents '*inheritance or dependency relationship'* between any two classes. In above class diagram, Reservoir junction and Standard Junction are child classes which *inherit* properties from parent class Junction. i.e., Reservoir Junction *is a* type of Junction.

- 1..4 implies that there can be one or a maximum of four instances of Pipe at each Junction.
- 0..1 implies that there can be no or one instance of Pump associtated with each Pipe

The program can handle multiple number of pipes and junctions in the network. A maximum of four pipes can meet at a junction. Also, each pipe can have one pump.

The class diagram shown in figure 11 portrays a number of relationships as given in Table 7.

Class1	Class 2	Relationship
Reservoir Junction	Junction	Reservoir Junction "is a" Junction
Standard Junction	Junction	Standard Junction "is a" Junction
Junction	Pressure Analyzer	Junction is " <i>part of</i> " Pressure Analyzer
Pipe	Pressure Analyzer	Pipe is " <i>part of</i> " Pressure Analyzer

Table 7. Relationships between classes

3.4 Classes used in Object oriented program

Following are the classes used in Transient analysis program written in Java. Corresponding methods are listed under each class:

Pipe class

A pipe network may contain more than one pipe. Each of the pipe objects is initialized in the pressure analyzer class. The primary attributes of the pipe class are length, diameter, wavespeed and friction factor. The constructor for the pipe class is shown below.

```
public Pipe(double length, double diameter, double waveSpeed,
                  double frictionFactor) {
            super();
            this.length = length;
            this.diameter = diameter;
            this.waveSpeed = waveSpeed;
            this.frictionFactor = frictionFactor;
      }
```
Some of the methods present in the pipe class are:

Method	Description
computeNextPressureVelocityValues()	Calculates the pressure and velocity heads at
	each node of the pipe for the current
	timestep using the previous timestep values.
calculateTimePerLength()	Calculates time step with respect to each
	pipe which in turn is used to determine time
	increments for the finite difference scheme.
calAverageLevel()	Calculates slope of each pipe based on start
	and ending node elevations.
setInitialVelocity(double)	Updates velocity at each node of the pipe
	with initial steady state velocity value.

Table 8. Methods in Pipe class

Junction Class

This is an abstract class which only defines the general form of its subclases: Reservoir Junction class and Standard Junction class. The constructor for Junction class is shown below: **public** Junction(**double** elevation) {

```
super();
      this.elevation = elevation;
}
```
It also has two abstract methods getHead() and calculateNextHeadVelocity(), which are overridden in the subclasses.

Pressure Analyser Class

This is main class of the program in which all other classes are initialized and transient anaysis is prefromed. The constructor for Pressure Analyser Class is shown below:

```
PressureAnalyser() {
            doReadMethod();
      }
```
The following methods of this class help to solve transient's problem in any pipe network comprehensively.

Table 9. Methods in Pressure Analyser class

Method	Description
doReadMethod()	Reads input data file and updates all the elements of the
	network such as Pipe, Junction, Reservoir and Pumps.
Transient_analysis()	Calculates pressure and velocity heads for each timestep in
	the network
doWriteMethod1()	Writes results to an output text file.
locateMaxMineadsHVal()	Locates maximum and minimum heads for each time step

The main() method of this class where the program will begin executing is shown below:

```
public static void main(String[] args) {
            PressureAnalyser pan = new PressureAnalyser();
            try {
                  pan.doWriteMethod1();
                  pan.transient_analysis();
                  if (pan.transient_analysis() != 1)
                        pan.doWriteMethod2(); //prints extreme values at 
                                                        each pipe node
            }
            catch (Exception ioe) { //Open catch
                  ioe.printStackTrace();
            } // Close catch
            finally { //Open finally
                  pan.closeStreams();
            } // Close finally
      }
```
Pump Class

The java program can handle pumps in pipe network. Typically one pump is associated with each pipe. The following constructor is used in Pressure Analyser class to intialise pump object.

```
public Pump(double speed, double torque, Pipe pipe) {
            super();
            this.speed = speed;
            this.torque = torque;
            this.pipe = pipe;
      }
```
The following methods are present in pump class:

Reservoir Class

Reservoir object is initialized in Pressure Analyser class. Its pressure head is constant

throughout the simulation; get and set methods are used to obtain Reservoir head values.

Reservoir Junction Class

This is subclass of Junction Class. Here is the constructor for Reservoir Junction class:

```
public ReservoirJunction(double elevation, Pipe pipe, Reservoir reservoir) {
            super(elevation);
            this.pipe = pipe;
            this.reservoir = reservoir;
      }
```
Methods in Reservoir Junction class are shown below:

Method	Description	
calculateNextHeadVelocity()	Computes current Head and Velocity at pipe node	
	closest to the reservoir using previous timestep values.	
Rundown()	Computes current Head and Velocity at pipe node if a	
	pump is present in the pipeline.	
isUpStream()	Checks if the reservoir object is at upstream or	
	downstream of a pipe	

Table 11. Methods in Reservoir Junction class

Standard Junction Class:

A junction where two or more pipes meet is termed as Standard Junction. Maximum pipes that can meet at a junction is four. The constructor for the Standard Junction Class is:

```
public StandardJunction(double elevation, double demand, ArrayList 
pipes) {
            super(elevation);
            this.demand = demand;
            this.pipes = pipes;
```
}

Following are the methods in Standard Junction Class:

Method	Description	
calculateNextHeadVelocity()	Computes current Head and Velocity at pipe node	
	closest to the Junction using previous timestep values.	
isUpStream()	Checks if the pipe object is at upstream or downstream	
	of the Junction	

Table 12. Methods in Standard Junction class

3.5 Comparison between TRANSNET and JAVA

TRANSNET is the Fortran code for transient analysis. The following portion of TRANSNET code computes Pressure and Velocity heads for current time step at upstream end (for any of reservoir, demand or pump boundary conditions) using previous time step values.

```
 80 I=IABS(NODEP(III,1))
   VRITE=V(I,1)-RATIO(I)*A(I)*(V(I,1)-V(I,2))HRITE=H(I,1)-RATIO(I)*A(I)*(H(I,1)-H(I,2)) IF(NTYPE(III).EQ.99) GO TO 85
    CC=VRITE-C(I)*HRITE-C(I)*DELT*VRITE*SINE(I)-AK(I)*VRITE*ABS(VRITE)
    IF(NSHUT.EQ.0) GO TO 800 //check for suddenly closed valves //
    II=NODEP(III,1)
    DO 801 N=1,NSHUT
    IF(ISHUT(N).NE.II) GO TO 801
   VNEW(I, 1)=0.
   HNEW(I, 1) = -CC/C(I) HNODE(III)=HNEW(I,1)-ELG(III) //calculates new Head at each node //
    ELHGL(III)=HNEW(I,1)
    GO TO 40
 801 CONTINUE
 800 CONTINUE
     IF(ABS(QNEW(III)).LE..00001) GO TO 802
    VNEW(I,1)=QDEM(III)/AREA(I) //calculates new velocity at each node //
   HNEW(I,1) = (VNEW(I,1) - CC)/C(I)
    ELHGL(III)=HNEW(I,1)
```

```
 GO TO 40
 802 HNEW(I,1)=ELHGL(III)
    VNEW(I, 1) = CC + C(I) * HNEW(I, 1) GO TO 40
```
The corresponding code in Java for Transient analysis program from "*ReservoirJunction"* class is shown below

```
Pump pump = pipe.getPump();
                  pipeJuncVel = pipe.getVelocityAtParts()[0];
                  pipeNextPrevVel = pipe.getVelocityAtParts()[1];
                  pipeJuncPress = pipe.getPressureAtParts()[0];
                  this.head= pipe.getPressureAtParts()[0];
                  pipeNextPrevPress = pipe.getPressureAtParts()[1];
                  VRITE = pipeJuncVel - timeIntPerLen * pipeWaveSpeed
                              * (pipeJuncVel - pipeNextPrevVel);
                  HRITE = pipeJuncPress - timeIntPerLen * pipeWaveSpeed
                              * (pipeJuncPress - pipeNextPrevPress);
             double CC;
                  if(pump == null){
                         CC = VRITE - inverseSpeed * HRITE - inverseSpeed
                              * timeInterval * VRITE * avgLevel -
                              frictionRatio * VRITE* Math.abs(VRITE);
                  if (pipe.getSide() > 0) {
                        pipe.getNewVelocityAtParts()[0] = 0.0;
                        pipe.getNewPressureAtParts()[0] = -CC /inverseSpeed;
                              setHeadAtNode( pipe.getNewPressureAtParts()[0] 
                               - getElevation()); 
                        head = pipe.getNewPressureAtParts()[0];
                        return;
                  }
                  if (Math.abs(getNewDemand()) > .00001) {
                        pipe.getNewVelocityAtParts()[0] = demand / 
                              pipe.getArea();
                        pipe.getNewPressureAtParts()[0] = 
                        (pipe.getVelocityAtParts()[0] - CC)/inverseSpeed;
                        head = pipe.getNewPressureAtParts()[0];
                        return;
                  }
                  pipe.getNewPressureAtParts()[0] = head;
                  pipe.getNewVelocityAtParts()[0] = CC + inverseSpeed
                              * pipe.getNewPressureAtParts()[0];
                  return;
                  }
```
Table 13 compares various features of TRANSNET and Transient analysis program.

TRANSNET	Transient analysis program
1) The variables program uses does not	1) Direct correlation exists between the
relate to the pipe network attributes. Hence,	variables and the pipe network elements.
secondary reference is necessary to a	e.g., getNewVelocityAtParts,
understand the code. e.g., VNEW, HNEW.	getNewPressureAtParts.
2) "go to" is the control statement used very	2) "go to" is very rarely used in Java. "if" is
often to cause the flow of program execution	the control statement used. All the code which
to advance and to make changes to the state	needs to be executed is specified in a single
of the program. This makes the code hard to	block.
comprehend as the reader needs to juggle	
between different parts of the program	
3) Expansion of code is very difficult as	3) This is very flexible. Adding new methods
developer needs to make changes at various	and new classes is possible without modifying
portions of code.	the entire program.

Table 13. Comparative advantages of the object- oriented programming

3.6 Test problem and verification of results

The pipe network problem described in Section 2.6.2 in Chapter 2 is solved using the Java program. The input and output data for this program are in text file format, which are shown in Appendices C-1 and C-2 respectively.

The results obtained from Transient analysis program are compared with TRANSNET program, which are shown in Table 14.

			o	. .
Pipe	Node Location*	Result from Transient analysis	Results from TRANSNET	Percent error
Pipe 1		359.8	359.7	0.028
Pipe 2	0.583	538.2	538.2	0.000
Pipe 3	0.4	696.9	697.2	0.043
Pipe 4	0.286	635.7	635.8	0.016
Pipe 5	0.8	761.7	762.4	0.092
Pipe 6	0.75	431.8	431.6	0.046

 Table 14. Maximum pressure head (in feet) values during transients in pipe network

*represents distance in terms of pipe length from upstream end of pipe

The results from Transient analysis program and TRANSNET program are quite comparable. However, the slight difference in values between these two programs is due to the rounding error. All the values calculated in Java are third decimal accurate where as in TRANSNET computations are done until first decimal only.

Chapter 4: Modeling Transient Gaseous Cavitation in Pipes

4.1 Introduction

During transient flow of liquid in pipelines, pressures sufficiently less than the saturation pressure of dissolved gas can be reached. As a result, gas bubbles are formed due to diffusion of cavitation nuclei. In this process of gas release, free gas volume increases. Consequently, the mixture celerity is reduced due to added compressibility of the gas, which in turn may give rise to significant pressure wave dispersion.

The decision to account for gas release (Wiggert and Sundquist, 1979) during a pressure transient in a pipe depends upon the system dimensions, type of fluid mixture being transported, extent of saturation of the gas, and low pressure residence times. In long pipelines where big elevation difference exists at different sections of pipe, transient pressures below gas saturation pressure is possible. Also, in highly soluble solutions such as water and carbon dioxide or hydraulic oil and air, significant gas release can take place and should be considered to correctly simulate the transient. Examples of gaseous cavitation are found in large-scale cooling water units, aviation fuel lines, and hydraulic control systems.

In present work, a two dimensional (2D) numerical model is developed to analyze gaseous cavitation in a single pipe system. The model is based on mathematical formulations proposed by Cannizzaro and Pezzinga (2005) and Pezzinga (2003). The model considers energy dissipation due to both thermic exhange between gas bubbles and surrounding liquid and during the process of gas release where as Cannizzaro and Pezzinga (2005) and Pezzinga (2003) present separate models for thermic exchange and gas release.

The model considers liquid flow with a small amount of free gas. Gas release expression is derived from Henry's law. The energy expression in complete form is from Anderson (1995) and Saurel and Le Metayer (2001). The thermic exchange is expressed from Newton's law of cooling (Ewing, 1980). These are the basic equations model uses which undergo several transformations. Finally, four constitutive equations – rate of mass increase of free gas, the mixture continuity equation, mixture energy equation, and the mixture momentum – yield a set of differential equations that are solved by the MacCormack (1969) explicit finite-difference method.

Later, an example is presented to demonstrate this model to analyse transients for the closure of a downstream valve in a pipe containing an air-water mixture with reservoir upstream.

4.2 Free Gas in Liquids

The model assumes free gas to be distributed throughout the pipe as a homogeneous bubblyfluid mixture with gas bubbles and liquid moving at same velocity. A void fraction α , is used to describe the ratio of volume of free gas, V_g , to the mixture volume, *V*, and for a given mass of free gas, it is pressure dependant (Wylie, 1984).

$$
\alpha = \frac{V_s}{V} \tag{4.1}
$$

Dalton's law states that the total pressure exerted by a mixture of gases is equal to sum of partial pressures of various components. If air were the free gas distributed in water in a pipeline the bubbles would contain a mixture of air and water vapor. The total volume occupied by air would be the same volume as occupied by the water vapor. Dalton's law is expressed (Wylie, 1984):

$$
P^* = P_g^* + P_v^* \tag{4.2}
$$

in which P^* is the total absolute pressure, P_g^* is the absolute partial pressure of air and P_v^* is the absolute vapor pressure.

For small void fractions isothermal behavior of free gas is a valid assumption. With a given mass of free gas, M_g , and use of the perfect gas law (Wylie, 1984):

$$
P_s^* V_g = M_g R_g T \tag{4.3}
$$

$$
\text{and } P_v^* V_v = M_v R_v T \tag{4.4}
$$

In which V_g is the gas volume, R_g is the gas constant, T is the absolute temperature, V_g is the vapor volume and is equal to the free gas volume, M_{v} is mass of vapor and R_{v} is the vapor gas constant.

We define mass of free gas per unit volume of the mixture as:

$$
m = \frac{M_s}{V} \tag{4.5}
$$

But from above, $M_g = P_g^* V_g / R_g T$. Therefore, *m* can be rewritten as:

$$
m = \frac{P_g^* V_g}{R_g T V} \tag{4.6}
$$

Hence, from above:

$$
\alpha = \frac{V_s}{V} = \frac{mR_s T}{P_s^*} \tag{4.7}
$$

From (4.5) and (4.7)

$$
\frac{M_s}{V_s} = \rho_s = \frac{m}{\alpha} \tag{4.8}
$$

4.3 Rate of Gas Release

During low pressure periods, gas release takes place and should be considered to correctly simulate the transient. The present model accommodates for gas release using Henry's law.

Henry's law: The amount of gas dissolved in a liquid is proportional to the partial pressure of gas above the liquid(www).

$$
P = K_H C_W \tag{4.9}
$$

Where K_H is the Henry's law constant, P is the partial pressure of gas and C_w is concentration of dissolved gas in *mol* / *L*.

The dimensionless Henry's law constant, K_H is defined as

$$
K_H = \frac{C_a}{C_w} \tag{4.10}
$$

C_a is the concentration in gas phase. Using $PV = nRT$ and $\frac{n}{V} = C_a$ we have

$$
C_a = \frac{P}{RT} \tag{4.11}
$$

From Eqs. 4.9, 4.10 and 4.11 we have

$$
K_H = \frac{K_H}{RT} \tag{4.12}
$$

Mass of dissolved gas in aqueous phase, $m_{aq} = C_w V_l M$

$$
m_{aq} = C_w V_l M = \frac{P}{K_H} V_l M = \frac{P}{K_H R T} V_l M
$$
\n(4.13)

The increase in mass of free gas = $m_{aq,s}$ – m_{aq}

$$
m_{aq,s} - m_{aq} = \frac{V_l M}{K_H R T} [p_s - p]
$$
\n(4.14)

Where $m_{aq,s}$ is the initial mass of free gas present at equilibrium, and p_s is the saturation or equilibrium pressure. Hence, the rate of increase in mass of free gas is given by:

$$
\frac{\partial m}{\partial t} = \frac{m_{aq,s} - m_{aq}}{\Delta t} = \frac{V_l M}{\Delta t} \frac{1}{K_H RT} [p_s - p]
$$
\n(4.15)

4.4 Energy Equation of gas as function of Temperature and Pressure

The energy equation for a gas of unit mass is written as (Anderson, 1995):

$$
dq - pdV = dU \tag{4.16}
$$

Where dq is the heat added in Joules, $pdV = dW$ is the work done by the gas in Joules, dU is the change in internal energy in Joules.

$$
\frac{dq}{dt} = \frac{dU}{dt} + p\frac{dV}{dt} \tag{4.17}
$$

From the definition of enthalpy (Anderson, 1995)

$$
h = U + pV \tag{4.18}
$$

Where *h* is the enthalpy per unit mass.

$$
\frac{dh}{dt} = \frac{dU}{dt} + p\frac{dV}{dt} + V\frac{dp}{dt}
$$
\n(4.19)

As we are considering unit mass of gas, $V = 1/\rho_g$ where, ρ_g is the density of gas.

Hence, above equation can be written as:

$$
\frac{dh}{dt} = \frac{dU}{dt} + p\frac{dV}{dt} + \frac{1}{\rho_g}\frac{dp}{dt}
$$
\n(4.20)

From (4.17) and (4.20)

$$
\frac{dh}{dt} = \frac{dq}{dt} + \frac{1}{\rho_g} \frac{dp}{dt} \tag{4.21}
$$

The above is another form of expressing energy equation of a gas.

From the definition of specific heat at constant pressure, heat Q added to raise the temperature of a mass of m from T_1 to T is (Saurel and Le Metayer, 2001):

$$
Q = mc_p (T - T_1) \tag{4.22}
$$

For unit mass:
$$
q = c_p (T - T_1)
$$
 (4.23)

$$
\frac{dq}{dt} = c_p \frac{dT}{dt} \tag{4.24}
$$

4.4.1 Relation between c_p and c_v

The energy equation for a reversible process of a closed system, ideal gas is (Anderson, 1995):

$$
dq = dU + pdV \tag{4.25}
$$

Where q is the heat transfer per unit mass, U is the internal energy per unit mass.

Ideal gas equation for unit mass of gas can be written as (Anderson, 1995):

$$
pV = RT \tag{4.26}
$$

At constant pressure, $pdV = RdT$ (4.27)

From above:

$$
dq = dU + RdT \tag{4.28}
$$

$$
\frac{dq}{dT} = \frac{dU}{dT} + R\tag{4.29}
$$

Specific heat(www) of a gas is defined as the amount of heat required to change unit mass of gas by one degree in temperature.

Hence, specific heat of gas at constant volume c_v is $\frac{dq}{dT}$ $\frac{dq}{dr}$ when $dV = 0$. From energy equation

$$
dq = dU \tag{4.30}
$$

$$
\frac{dq}{dT} = \frac{dU}{dT} = c_v \tag{4.31}
$$

$$
d(pV) = pdV + Vdp \tag{4.32}
$$

$$
pdV = d(pV) - Vdp \tag{4.33}
$$

Hence, the above energy equation can be rewritten as:

$$
dq = dU + d(pV) - Vdp \tag{4.34}
$$

$$
dq = d(U + pV) - Vdp \tag{4.35}
$$

From the definition of enthalpy, *H* (Anderson, 1995)

$$
H = U + pV \tag{4.36}
$$

Therefore at constant pressure, the energy equation is:

$$
dq = dH \tag{4.37}
$$

Specific heat of gas at constant pressure c_p :

$$
\frac{dq}{dT} = \frac{dh}{dT} = c_p \tag{4.38}
$$

Hence from Eq.s 4.29, 4.31 and 4.38

$$
c_p = c_v + R \tag{4.39}
$$

Let *k* be the ratio of specific heats,

$$
k = \frac{c_p}{c_v} \tag{4.40}
$$

From Eq.s 4.39 and 4.40

$$
c_p(1 - \frac{1}{k}) = R \tag{4.41}
$$

$$
c_p = \frac{k}{(k-1)}R\tag{4.42}
$$

4.4.2 Newton's law of cooling

Newton's Law of Cooling states that the rate of change of temperature of an object is proportional to the difference between its own temperature and the ambient temperature (Ewing, 1980).

$$
\frac{dT}{dt} = -K_N(T - T_0) \tag{4.43}
$$

Where T is the absolute temperature of the body (gas), T_0 is the absolute temperature of the surrounding (liquid), K_N is the experimental positive constant (1/sec). Hence, from above:

$$
\frac{dq}{dt} = -c_p K_N (T - T_0) \tag{4.44}
$$

From Eq. 4.38,
$$
\frac{dh}{dt} = c_p \frac{dT}{dt}
$$
 (4.45)

Hence, energy Eq. 4.21 can be written as:

$$
c_p \frac{dT}{dt} = -c_p K_N (T - T_0) + \frac{1}{\rho_g} \frac{dp}{dt}
$$
 (4.46)

Using equation (4.42) in above along with the gas law for ρ_{ϱ} , we have

$$
\frac{dT}{dt} - \frac{(k-1)T}{k} \frac{dp}{dt} + K_N(T - T_0) = 0
$$
\n(4.47)

4.5 One-Dimensional Two-Phase Flow

Present model considers the fluid as a homogeneous two-phase air-water mixture. The properties of the fluid will be average of both components. Following are the assumptions used in the model:

- 1) Gas bubbles are distributed throughout the pipe and they are very small compared to pipe diameter.
- 2) Difference in pressure due to surface tension across a bubble surface can be neglected; and
- 3) Momentum exchange between gas bubbles and surrounding liquid is negligible, so that gas bubbles and liquid have the same velocity.

4.5.1 Mixture Density

The density of a liquid gas mixture, ρ can be expressed as (Wiggert and Sundquist, 1979)

$$
\rho = \alpha \rho_g + (1 - \alpha) \rho_l \tag{4.48}
$$

Where ρ_l is the liquid density. Using Eq. 4.8, above can be written as:

$$
\rho = m + (1 - \alpha)\rho_l \tag{4.49}
$$

$$
\frac{\partial \rho}{\partial t} = \frac{\partial m}{\partial t} + \frac{\partial \rho_l (1 - \alpha)}{\partial t} - \rho_l \frac{\partial \alpha}{\partial t}
$$
(4.50)

Bulk modulus of elasticity of a liquid is given by (Larock, 2000):

$$
E = -\frac{dp}{d\nu/\nu} = \frac{dp}{d\rho_l/\rho_l} \tag{4.51}
$$

Also $E / \rho_l = c^2$ (4.52)

Hence, we can write:

$$
\frac{\partial \rho_l}{\partial t} = \frac{\rho_l}{E} \frac{\partial p}{\partial t} = \frac{1}{c^2} \frac{\partial p}{\partial t}
$$
(4.53)

Using Eq. 4.53 in Eq. 4.50 we have

$$
\frac{\partial \rho}{\partial t} = \frac{1}{c^2} \frac{\partial p}{\partial t} (1 - \alpha) - \rho_l \left[-\frac{mRT}{p^2} \frac{\partial p}{\partial t} + \frac{mR}{p} \frac{\partial T}{\partial t} + \frac{RT}{p} \frac{\partial m}{\partial t} \right] + \frac{\partial m}{\partial t}
$$
(4.54)

$$
\frac{1}{\rho} \frac{\partial \rho}{\partial t} = \frac{\partial p}{\partial t} \left[\frac{1}{\rho c^2} (1 - \alpha) + \frac{\rho_l}{\rho} \frac{mRT}{p^2} \right] - \frac{\partial T}{\partial t} \left(\frac{mR\rho_l}{p\rho} \right) + \frac{\partial m}{\partial t} \left[\frac{1}{\rho} - \frac{RT}{p} \frac{\rho_l}{\rho} \right]
$$
(4.55)

For $\alpha \ll 1$ and $\rho_l \approx \rho$ we have

$$
\frac{1}{\rho} \frac{\partial \rho}{\partial t} = \frac{\partial p}{\partial t} \left[\frac{1}{\rho c^2} (1 - \alpha) + \frac{\rho_l}{\rho} \frac{mRT}{p^2} \right] - \frac{\partial T}{\partial t} \left(\frac{mR\rho_l}{p\rho} \right) + \frac{\partial m}{\partial t} \left[\frac{1}{\rho} - \frac{RT}{p} \frac{\rho_l}{\rho} \right]
$$
(4.56)

4.5.2 Continuity Equation

From one dimensional continuity equation given in Wylie et al.,(1993)

$$
\frac{1}{\rho}\frac{d\rho}{dt} + \frac{1}{A}\frac{dA}{dt} + \frac{\partial u}{\partial x} = 0
$$
\n(4.57)

The second term in above equation deals with elasticity of the pipe and its rate of deformation with pressure. Assuming pipe is rigid $\frac{dA}{dt}$ can be scaled to zero. As a consequence, we have constant area and

$$
\frac{\partial u}{\partial x} = \frac{1}{A} \frac{\partial Q}{\partial x}.
$$
 Where $Q = AV$ (4.58)

Using above expression in continuity equation Eq. 4.57 we have:

$$
\frac{1}{\rho} \left[\frac{\partial \rho}{\partial t} + u \frac{\partial \rho}{\partial x} \right] + \frac{1}{A} \frac{\partial Q}{\partial x} = 0
$$
\n(4.59)

Also neglecting the spatial derivative of density in above continuity equation we have,

$$
\frac{1}{\rho} \frac{\partial \rho}{\partial t} + \frac{1}{A} \frac{\partial Q}{\partial x} = 0
$$
\n(4.60)

From Eqs. 4.56 and 4.60 continuity equation is finally expressed as:

$$
\frac{\partial p}{\partial t} \left[1 + \frac{\rho c^2 mRT}{p^2} \right] + \frac{\partial m}{\partial t} \left[c^2 - \frac{\rho c^2 RT}{p} \right] - \frac{\partial T}{\partial t} \left[\frac{\rho c^2 mR}{p} \right] + \frac{\rho c^2}{A} \frac{\partial Q}{\partial x} = 0 \tag{4.61}
$$

4.6 Conservation form of Mixture Continuity and Energy Equations

$$
Define \phi = \frac{p}{\rho g} - \frac{c^2 mRT}{\rho g} + \frac{mc^2}{\rho g}
$$
\n
$$
\frac{\partial \phi}{\partial t} = \frac{1}{\rho g} \frac{\partial p}{\partial t} - \frac{1}{\rho^2} \frac{p}{g} \frac{\partial \rho}{\partial t} + \frac{c^2}{p^2} (\frac{mRT}{g}) \frac{\partial p}{\partial t} - \frac{c^2 RT}{pg} \frac{\partial m}{\partial t} - \frac{c^2 mR}{pg} \frac{\partial T}{\partial t} - \frac{mc^2}{\rho^2 g} \frac{\partial \rho}{\partial t} + \frac{c^2}{\rho g} \frac{\partial m}{\partial t} (4.63)
$$
\n
$$
Assuming \frac{\partial \rho}{\partial t} \approx 0
$$
\n
$$
\frac{\partial \phi}{\partial t} = \frac{\partial p}{\partial t} \left[\frac{1}{\rho g} + \frac{c^2}{g} \frac{mRT}{p^2} \right] + \frac{\partial m}{\partial t} \left[\frac{c^2}{\rho g} - \frac{c^2}{g} \frac{RT}{p} \right] - \frac{\partial T}{\partial t} \left[\frac{c^2}{g} \frac{mR}{p} \right]
$$
\n
$$
(4.64)
$$

Using above, mixture continuity equation, Eq. 4.61 is expressed as:

$$
\rho g \frac{\partial \phi}{\partial t} + \frac{\rho c^2}{A} \frac{\partial Q}{\partial x} = 0
$$
\n(4.65)

Finally, the following conservation form of mixture continuity equation is obtained (Cannizzaro and Pezzinga, 2005):

$$
\frac{\partial \phi}{\partial t} + \frac{c^2}{gA} \frac{\partial Q}{\partial x} = 0
$$
\n(4.66)

Define gas entropy $s = c_p \left[\ln T - \frac{k-1}{k} \ln p \right]$ (4.67)

$$
\frac{\partial s}{\partial t} = c_p \left[\frac{1}{T} \frac{\partial T}{\partial t} - \frac{k - 1}{k} \frac{1}{p} \frac{\partial p}{\partial t} \right]
$$
(4.68)

From energy equation, Eq. 4.47, which is:

$$
\frac{dT}{dt} - \frac{(k-1)T}{k} \frac{dp}{dt} + K_N(T - T_0) = 0
$$
\n
$$
T \left[\frac{1}{T} \frac{dT}{dt} - \frac{(k-1)T}{k} \frac{dp}{dt} \right] + K_N(T - T_0) = 0
$$
\n(4.69)

From Eq.s 4.68 and 4.69 we have

$$
T\left[\frac{1}{c_p}\frac{\partial s}{\partial t}\right] + K_N(T - T_0) = 0\tag{4.70}
$$

$$
\frac{\partial s}{\partial t} + c_p K_N (1 - \frac{T_0}{T}) = 0 \tag{4.71}
$$

The conservation form of Energy equation is (Cannizzaro and Pezzinga, 2005):

$$
\frac{\partial s}{\partial t} = c_p K_N \left(\frac{T_0}{T} - 1 \right) \tag{4.72}
$$

4.7 Mixture Momentum Equation

The following momentum equations are written for an elastic pipe with circular cross section, and assume a 2D flow field with axial symmetry (Pezzinga, 1999).

Along longitudinal direction, x-axis momentum equation:

$$
\rho \left[\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial r} \right] = -\frac{\partial p}{\partial x} - \frac{\partial \sigma_x}{\partial x} - \frac{1}{r} \frac{\partial (r\tau)}{\partial r}
$$
(4.73)

Along radial direction, r-component momentum equation:

$$
\rho \left[\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial r} \right] = -\frac{\partial p}{\partial r} - \frac{\partial \tau}{\partial x} - \frac{1}{r} \frac{\partial (r \sigma_r)}{\partial r} + \frac{\sigma_\theta}{r}
$$
(4.74)

For laminar and turbulent unsteady flows, radial velocity component has values of order10 − 20µ*m* /*s*. Hence, in above momentum equations radial velocity *v* and its derivatives are neglected. Also, in the momentum equation, in the longitudinal direction the normal stress is assumed to be equal to the pressure and residual convection term is neglected as usually done also in 1D models.

The x-axis momentum equation can be rewritten as:

$$
\rho \frac{\partial u}{\partial t} + \frac{\partial p}{\partial x} + \frac{1}{r} \frac{\partial (r\tau)}{\partial r} = 0
$$
\n(4.75)

Neglecting the stress terms in the radial momentum equation can be written as:

$$
\frac{\partial p}{\partial r} = 0\tag{4.76}
$$

4.7.1 Stress Model

Selection of appropriate stress model is important to estimate the initial velocity profile. In the present gaseous cavitation model for turbulent flow, a two-zone model is considered. In the viscous sub layer Newton's law is used and in turbulent zone, Prandtl mixing length hypothesis (Pezzinga, 1999), namely:

$$
\tau = -\rho v \frac{\partial u}{\partial r} - \rho l^2 \left| \frac{\partial u}{\partial r} \right| \frac{\partial u}{\partial r}
$$
\n(4.77)

is used. In above, $v =$ kinematic viscosity and $l =$ mixing length. For the mixing length the following expression was adopted (Pezzinga, 1999):

$$
\frac{l}{R_0} = \kappa \frac{y}{R_0} e^{-(y/R_0)} \tag{4.78}
$$

Where *y* = distance from wall, R_0 = pipe radius; and κ = constant. The following expression (Pezzinga, 1999) is used to obtain κ .

$$
\kappa = 0.374 + 0.0132 \ln(1 + \frac{83,100}{9})
$$
\n(4.79)

The value of κ was evaluated with initial Reynolds number and maintained constant throughout the simulation.

4.7.2 Evaluating Thickness of Viscous Sub layer

Thickness of viscous sub layer δ is distance from wall to the intersection between the velocity profiles in the viscous sub layer and the turbulent zone. Linear velocity profile is assumed in the viscous sub layer and in the turbulent zone the profile is locally logarithmic.

Figure 12. Shear stress distribution and velocity profile for turbulent flow (reproduced from Wylie et al., 1993)

If $u_* = \sqrt{\tau_w/\rho}$ is the friction velocity and τ_w is the wall shear stress then the equation for velocity profile in the viscous sub layer is (Pezzinga, 1999):

$$
\frac{u}{u_*} = \frac{u_* y}{v} \tag{4.80}
$$

The logarithmic velocity profile in turbulent zone is (Pezzinga G., 1999)

$$
\frac{u}{u_*} = 2.5 \ln \frac{y}{\varepsilon} + 5.0 \tag{4.81}
$$

Viscous sub layer thickness is obtained from the following expression (Streeter and Wylie 1985)

$$
\delta = 11.63 \nu / u_* \tag{4.82}
$$

4.7.3 Boundary Conditions

With regard to the wall condition, if first grid mesh close to the wall is in viscous sub layer then Eq. 4.80 is used to calculate the wall shear stress τ_w otherwise, Eq. 4.81 is used.

On the pipe axis, symmetry condition is adopted. The initial velocity profile was determined by an iterative procedure, using the unsteady-flow momentum equation and constant discharge. The initial values of discharge in the pipe and heads at the nodes were obtained from steadystate conditions.

4.8 Finite Difference Scheme

The pipe is divided into cylindrical grid elements of constant length ∆*x* in longitudinal direction and constant area ∆*A* in radial direction. Velocities are defined at center of each radial mesh, and shear stresses on internal and external sides. Other parameters such as pressure head H , mass m , Temperature T, ϕ and s vary along longitudinal direction only and are defined at each grid point.

Figure 13. Cylindrical grid element (reproduced from Cannizzaro and Pezzinga, 2005)

The time step Δt is defined as: $\Delta t = C \Delta x/c$. Where *C* is the Courant number and *c* the wave speed of the fluid.

4.8.1 Mac Cormack Method

MacCormack finite difference method (MacCormack, 1969) is used in this model. This technique is second order accurate in both space and time. In the *predictor* step, forward finite differences are used (Cannizzaro and Pezzinga, 2005):

 Figure 14. 2-D Finite difference grid (reproduced from Anderson, 1995)

Predictor Step:

$$
\frac{m_i^p - m_i^{n-1}}{\Delta t} = \frac{V_i M}{\Delta t} \frac{1}{K_H^i RT} (p_{si}^{n-1} - p_i^{n-1})
$$
\n(4.83)

$$
\frac{\phi_i^p - \phi_i^{n-1}}{\Delta t} + \frac{c^2}{gA} \frac{Q_{i+1}^{n-1} - Q_i^{n-1}}{\Delta x} = 0
$$
\n(4.84)

$$
\frac{s_i^p - s_i^{n-1}}{\Delta t} - c_p K_N \left(\frac{T_0}{T_i^{n-1}} - 1 \right) = 0 \tag{4.85}
$$

$$
\rho \frac{u_{i,j}^p - u_{i,j}^{n-1}}{\Delta t} + \frac{p_{i+1}^{n-1} - p_i^{n-1}}{\Delta x} + 2\Pi \frac{(r_{j+1} \tau_{i,j+1}^* - r_j \tau_{i,j}^*)}{\Delta A} = 0
$$
\n(4.86)

Corrector Step:

Backward finite differences are used in the *corrector* step (Cannizzaro and Pezzinga, 2005):

$$
\frac{m_i^c - m_i^{n-1}}{\Delta t} = \frac{V_i M}{\Delta t} \frac{1}{K_H^i RT} (p_{si}^p - p_i^p)
$$
\n(4.87)

$$
\frac{\phi_i^c - \phi_i^{n-1}}{\Delta t} + \frac{c^2}{gA} \frac{Q_i^p - Q_{i-1}^p}{\Delta x} = 0
$$
\n(4.88)

$$
\frac{s_i^c - s_i^{n-1}}{\Delta t} - c_p K_N \left(\frac{T_0}{T_i^p} - 1\right) = 0\tag{4.89}
$$

$$
\rho \frac{u_{i,j}^c - u_{i,j}^{n-1}}{\Delta t} + \frac{p_i^p - p_{i-1}^p}{\Delta x} + 2\Pi \frac{(r_{j+1} \tau_{i,j+1}^* - r_j \tau_{i,j}^*)}{\Delta A} = 0
$$
\n(4.90)

In above equations, indices i, j, and n refer, respectively, to directions *x* , and *r* , and time *t* . Indices p and c refer to predictor and corrector steps. $\tau_{i,j}^*$ is average between values of stress in previous and in current time steps. The value at time *n* is computed as the average of predictor and corrector values.

The flowchart below explains the gaseous cavitation problem modeled in MATLAB:

Figure 15. Flowchart to show steps for solving gaseous cavitation

The above model is applied to the single pipe system setup as shown in Figure 6 in Chapter 2. Reservoir is at upstream of the pipe and the valve present downstream is closed suddenly. An initial amount of free gas (air) is assumed in the pipe along with water. Below is the set of values used in the problem:

Steady state Discharge	$0.227*10^{\circ} - 3$ cu.m/s	
Pipe diameter	$53.9*10^{\circ} - 3$ m	
Kinematic Viscosity	$0.837 * 10^{-6}$ m/s2	
of water		
Pipe roughness	$0.1*10^{(-)}$	
Initial mass of free gas	$4.5*10^{-6}$ kg/cu.m	
Initial temperature of	28 deg.C	
water		
Wave speed	1324 m/s	
Length of the pipe	36m	
Reservoir Head	4.53 m	
Time Interval	0.002 sec	

 Table 15. Data used by gaseous cavitation model

A detailed descitption of all the parameters used in the model can be found in Appendix D.

The following matlab results are obtained:

Figure 16. Head vs Time plot near the valve

Figure 17. Temperature vs Time plot near the valve

Figure 18. Mass of released gas vs time plot near the valve

As shown above, the results representing head and temperature near the valve, and total mass of gas release are plotted for 2500 timesteps (i.e., 5 s). As shown in figure 16, a maximum pressure head of 52 m is reached near the valve during the first 5 s. Also from figure17, a peak temperature of 140 (deg c) is attained during the first 5 s resulting in vaporization of liquid near the valve. As shown in figure 18, the total amount of free gas is reduced during this time. These results can be verified by conducting an experiment with similar conditions.

The peak value for head near the valve is attained only once, and further oscillation in pressure head cannot be plotted. This may be due to stability and concurrency issues which are part of numerical schemes. By using a smaller time step, one can increase the stability of the numerical computations. But, as MacCormack method is an explicit finite difference method, time steps cannot be smaller than a certain value decided by the stability criteria or else, the resulting solution will not be accurate.

It can be observed that high values of temperature can be reached during cavitation process. The combination of high values of pressure and temperature in a pipeline may give rise to disastrous consequences. Accidents arisen from operationg errors in pumping plants of combustible liquids are mentioned in literature.

Chapter 5: Summary and Conclusions

5.1 Summary

In this thesis, first one dimensional flow in closed conduits was considered to analyse the changes in pressure and velocity in water distribution system during transient conditions. Equations related to method of characteristics and finite difference approximations are discussed. Various boundary conditions such as reservoir, valve, pumps/turbines, and junctions are considered in the pipe system.

In order to demonstrate the use of method of characterisitics for transient analysis two problems are solved.

1) A single pipe with reservoir upstream and open valve downstream is considered. In order to simulate transients, the valve is closed suddenly. An excel spreadsheet is developed to calculate pressure and velocity heads at each time step. The maximum pressure heads are then compared with Joukowski's equation and the results were found comparable. The excel spreadsheet can be used by any researcher to estimate the critical pressures reached when a valve is closed suddenly and to measure the total time it takes for pressure waves to die down and for water to become steady.

2) In order to perform transient analysis in pipe network, a distribution system with six pipes, three reservoirs, a pump located upstream of one of the reservoirs and a valve at downstream of one of the pipes is considered. In this network there is a constant demand at three junctions. Valve is closed suddenly to simulate transient analysis and variations in pressure and velocity heads are studied in each of the connecting pipes. TRANSNET program is used for transient analysis and results are compared with WHAMO which gave comparable results for pressure and velocity head values in pipes.

In Chapter 3, a transient analysis program is developed in Java. This program can analyze transients caused by suddenly-closing valves, gradually-closing valves, pump power failures

and sudden demand changes at junctions. A maximum of four pipes can be present at a junction. This program is used to solve the pipe network problem discussed above; the results were found to be similar to that of TRANSNET program.

In Chapter 4 a two – dimensional numerical model is developed using MATLAB to analyse gaseous cavitation in pipes. Two phase flow is assumed, with free gas being distributed throughout the pipe as a homogeneous fluid mixture with gas bubbles and liquid moving at same velocity. Equations used in the model such as Henry's law, energy equation of a gas, mixture continuity and momentum equations are derived and adopted from literature.

The results from the model show that during transients, there is significant increase in fluid temperature along with high pressures. As a result some of the fluid gets vaporized and pockets of air are formed in distribution systems. In literature pipe failures and noise problems in premise plumbing are atributed to gaseous cavitation.

5.2 Conclusion

In this study, an excel spreadsheet is developed to critically analyse transients that can occur in closed conduits. The Java program developed as part of this work is an attempt to introduce object oriented technology for analyzing problems in hydraulic engineering field. The code can be further extended, for example by developing java applets and graphical user interphase to make it more user friendly. A MATLAB program was developed to analyse gaseous cavitation using standard equations proposed by Cannizzaro and Pezzinga (2005). Though the gaseous cavitation program seemed stable, there are issues with accuracy and concurrency of the code. Based on author's experience in this work, it is recommended to use separate models for gas release and for thermic exchange so that parameters can be better estimated.

Appendix A-1: Steady state analysis results

STEADY STATE ANALYSIS OF THE NETWORK ------------------------------------ Example Problem 2 ALL DEMAND FLOWS ARE MULTIPLIED BY 1.0000 PIPES 6
NODES 6 NODES 6
SOURCE PUMPS 1 SOURCE PUMPS BOOSTER PUMPS 0
RESERVOIRS 2 RESERVOIRS 2
MINOR LOSSES 0 MINOR LOSSES 0
PRVS 0 PRVS 0
NOZZLES 0 NOZZLES 0
CHECK VALVE 0 CHECK VALVE BACK PRES. V. 0 DIF. HEAD DEV 0
SPECIFIED 0 0 SPECIFIED O SPECIFIED PRES 0 NODES AT S. P. & RES. WHICH HAVE BEEN ELIMINATED 4 5 6 PIPE 2ND ORDER COEF LINEAR COEF SHUT-OFF HEAD SUMP ELEV
6 - 504 1.122 97.00 4130.00 $-$.504 RES.(NOZZLE) PIPES & THEIR ELEV. ARE 1 4200.0 -5 4130.0 $N9 = 6 N8 = 3$ JUNCTION EXT. FLOW PIPES AT JUNCTION $\begin{array}{ccccccccc}\n1 & 1 & 1.058 & -1 & -2 & 3 \\
2 & 2 & .706 & 2 & 4 & -6\n\end{array}$ 2 2 .706 2 4 -6 3 3 1.760 -3 -4 -5

FLOW FROM PUMPS AND RESERVOIRS EQUALS 1582.000

THE FLOWRATE 3.787 IS QUITE DIFFERENT THAN THE NORMAL CAPACITY 6.684 AS GIVEN BY THE PUMP CHARACTERISTICS FOR PUMP 1 IN PIPE 6 HAVE YOU ERRED.

PUMPS:

RESULTS OF SOLUTION

UNITS OF SOLUTION ARE DIAMETERS - inch LENGTH - feet HEADS - feet ELEVATIONS - feet PRESSURES - (psi) FLOWRATES - (gpm) HAZEN-WILLIAMS FORMULA USED FOR COMPUTING HEAD LOSS PIPE DATA

NODE DATA:

Appendix A-2: Transient Analysis results

 * NETWORK TRANSIENT ANALYSIS * ******************************

> DEMONSTRATION OF PROGRAM TRANSNET - INPUT DATA FILE "EPB12_5.DAT" NETWORK OF EXAMPLE 2 - SUDDENLY-CLOSED VALVE AT THE DS END OF PIPE 5

 $IOUT = 100$ NPARTS = 4 $NPIPES = 6$ HATM = 30.0 FT TMAX = 20.00 SEC $DELT = .227$ SEC

TRANSIENT CONDITIONS IMPOSED

SUDDENLY CLOSED VALVE AT DOWNSTREAM END OF PIPE 5

PIPE INPUT DATA

PIPE DIAM-IN LENGTH-FT WAVE SPD-FPS PIPEZ-FT C-VALUE VEL-FPS

NODE INPUT DATA

** PUMP INFORMATION **

 PRESSURE HEADS, H-VALUES AND VELOCITIES AS FUNCTIONS OF TIME --

PIPE 6 PUMP SPEED = 1180.0 RPM PUMP DISCHARGE = 1700.6 GPM EACH PUMP HEAD = 94.0 FT

COLUMN SEPARATION HAS OCCURRED AT 7.73 SEC IN PIPE 5 AT LOCATION 1.000

1.000 379. 4209. 3.10

PIPE 6 PUMP SPEED = 1180.0 RPM PUMP DISCHARGE = 1471.4 GPM EACH PUMP HEAD = 98.9 FT

* TABLE OF EXTREME VALUES *

* HEAD AND VELOCITY VS TIME FOR SELECTED POINTS * ***

PIPE 5 NODE 5

PIPE 5 NODE 3

PIPE 1 NODE 1

PIPE 6 NODE 2

Appendix B: Comparison between WHAMO and TRANSNET results

Figure A.1 Network distribution system

(1) Comparing the results at steady state condition:

Time = 0.0 sec Total head (ft)

Discharge(cfs)

(2) At time $t = 7.73$ sec, (just at the time of column separation noted by transient program)

Time =7.73 sec

Total head(ft)

Discharge(cfs)

(3) Comparing maximum and minimum heads (ft):

Max heads:

Min heads:

Appendix C-1: Java Program input

4 100 30 120 0 false true true false false false true 0 0 $1 - 5$ 0 6 1 6 1 4198.68 3800 2 4214.38 3830 3 4196.89 3370 4 4200.00 4050 5 4130.00 4000 6 4224.03 4010 1 12 3300 120 0.96 4 1 2850 2 8 8200 120 1.74 2 1 2850 3 8 3300 120 0.88 1 3 2850 4 12 4900 120 3.15 2 3 2850 5 6 3300 120 5.20 3 5 2850 6 14 2600 120 3.54 6 2 2850 -1 -2 3 0 $24 - 60$ $-3 -4 5 0$ 1 0 0 0 -5 0 0 0 6 0 0 0 1 6 1 1 1180 4130 50 0.00 118.0 57.0 2000.76 92.0 68.0 3001.14 82.0 77.0 4001.52 67.0 80.0 4501.70 52.0 76.0 5302.01 0.0 60.0 3 1 3 1 2 5 3 4 1 5

Appendix C-2: Java Program Output

NETWORK TRANSIENT ANALYSIS TITLE 1 TITLE 2 IOUT=100 NPARTS=4 NPIPES=6 HATM=30.0 TMAX=120.0 DELT=0.227 TRANSIENT CONDITIONS IMPOSED ------------------------------ SUDDENLY CLOSED VALVE AT DOWNSTREAM END OF PIPE 5 PIPE INPUT DATA --------------- PIPE DIAM-IN LENGTH-FT WAVE SPD-FPS PIPEZ-FT F-VALUE VEL-FPS ---- ------- --------- ------------ -------- ------- ------- 1 12.0 3300.0 2850.0 4050.0 0.028 0.96 2 8.0 8200.0 2850.0 3830.0 0.027 1.74 3 8.0 3300.0 2850.0 3800.0 0.03 0.88 4 12.0 4900.0 2850.0 3830.0 0.023 3.15 5 6.0 3300.0 2850.0 3370.0 0.024 5.2 6 14.0 2600.0 2850.0 4010.0 0.022 3.54 PIPE DELT-SEC PARTS SINE L/A-SEC INTERPOLATION ---- -------- ----- ---- ------- ------------- 1 0.289 5 -0.076 1.158 0.019 2 0.716 12 -0.0040 2.877 0.052 3 0.289 5 -0.13 1.158 0.019 4 0.428 7 -0.094 1.719 0.075 5 0.282 5 0.191 1.158 0.019 6 0.227 4 -0.069 0.912 0.0030 NODE INPUT DATA --------------- NODE ELHGL-FT GRNDEL-FT 1 2 3 4 DEMAND-CFS ---- -------- --------- -- -- -- -- ---------- 1 4198.68 3800.0 -1 -2 3 0 1.054 2 4214.38 3830.0 2 4 -6 0 0.703 3 4196.89 3370.0 -3 -4 5 0 1.76 4 4200.0 4050.0 1 0 0 0 -0.754 5 4130.0 4000.0 -5 0 0 0 1.021 6 4224.03 4010.0 6 0 0 0 -3.784 PUMP INFORMATION ---------------- Q-GPM HEAD/STAGE-FT HP/STAGE-HP ----- ------------- ----------- $\begin{array}{rcl} \texttt{LINE} & = & 6 \\ \texttt{PUMPS} & = & 1 \end{array}$ 0.0 8.474576271186441E-5 1.8220609699376222E-4
 $\begin{array}{rcl} \texttt{PUMPS} & = & 1 \end{array}$ 0.0037763013853761656 6.607296753806377E-5 0.0037763013853761656 6.607296753806377E-5 2.1736867711536547E-4

PRESSURE HEADS, H-VALUES AND VELOCITIES AS FUNCTIONS OF TIME ---

PIPE2

 0.0 384.38 4214.38 1.74 0.083 385.561 4213.061 1.74 0.167 386.741 4211.741 1.74 0.25 387.922 4210.422 1.74 0.333 389.103 4209.103 1.74 0.417 390.284 4207.784 1.74 0.5 391.464 4206.464 1.74 0.583 392.645 4205.145 1.74 0.667 393.826 4203.826 1.74 0.75 395.006 4202.506 1.74 0.833 396.187 4201.187 1.74 0.917 397.368 4199.868 1.74 1.0 398.549 4198.549 1.74

PIPE3

PIPE4

 0.0 384.38 4214.38 3.15 0.143 447.573 4211.859 3.15 0.286 510.766 4209.337 3.15 0.429 573.959 4206.816 3.15 0.571 637.152 4204.295 3.15 0.714 700.345 4201.773 3.15 0.857 763.538 4199.252 3.15 1.0 826.73 4196.73 3.15

PIPE5

0.0 826.89 4196.89 5.2

 0.0 462.825 4.546 378.335 5.455 4262.825 4178.335 0.2 607.189 2.5 439.735 5.682 4321.189 4153.735 0.4 696.87 2.727 513.464 5.909 4324.87 4141.464 0.6 784.781 2.955 597.316 6.137 4326.781 4139.316 0.8 874.243 3.182 687.47 5.909 4330.243 4143.47

PIPE4

0.0 476.223 4.546 337.37 6.818 4306.223 4167.37

PIPE5

 0.0 962.444 3.409 780.201 6.591 4332.444 4150.201 0.2 1118.828 1.364 577.869 3.636 4614.828 4073.869 0.4 1000.988 1.591 440.55 3.864 4622.988 4062.55 0.6 881.546 1.818 295.009 7.727 4629.546 4043.009 0.8 761.691 2.046 119.601 7.727 4635.691 3993.601

PIPE6

 0.0 227.294 5.909 214.03 0.0 4237.294 4224.03 0.25 330.086 3.864 225.045 6.137 4295.086 4190.045 0.5 386.569 4.091 256.424 6.364 4306.569 4176.424 0.75 431.84 4.318 296.573 6.591 4306.84 4171.573

MAXIMUM HEAD=1118.8284208516707 FT IN PIPE 1 AT X=0.2 AT TIME =1.3636676485626478 SEC

MINIMUM HEAD=-39.03264828702777 FT IN PIPE1 AT X=1.0 AT TIME=7.727450008521666 SEC

Appendix D: Matlab program

%% Initial parameters

% All units are in SI system

 $Q0 = 0.227*10^2-3$; % Discharge in cu.m/s

Dia = $53.9*10^{\circ}$ -3; % Diameter of the pipe in m

Ki_Visc0= $0.837*10^{\circ}$ -6; % Kinematic Viscosity in m/s^{γ}2

 $e = 0.1*10^{\circ} - 3$; % Pipe roughness in m

DeltaT=0.002; % Time Interval in sec

m0=4.5*10^-6; % Initial mass of gas in Kg/m^3

Theta_T=0.06; % Calibrated value of relaxation time for Temperature variation

Theta_m=400; % Calibrated value of relaxation time for Variable Mass

Gas_Const=287.05;%Gas Constant of Air in J/Kg.K

T0_c=28; % Temperature in Deg.C

T0_k=301.15; % Absolute Temperature in Deg.K

p_atm=1.013*10^5; % Atmospheric pressure in Pa

c=1324; % Wave speed in m/s

g=9.81; % gravity in m/s^2

 $L=36.0;$ %Pipe length in m

N_x=5;% Number of Nodes in longitudinal direction

N_y=6; % Number of Radial Nodes

 $cp=1000$;% Specific heat of air at constant pressure in J/Kg.K

sp_ratio=1.4; % specific heat ratio

Delta $x=L/(N x-1);$ % Length of each longitudinal segment in m

Delta_y= $Di\frac{a}{2*(N_y-1)}$; % Length of each radial mesh in m

H0=4.53; % Reservoir Head in m

Area=pi*Dia^2/4;% Cross section area in m^2

%%% %%%%%%%%%%%%%

% PHYSICAL PROPERTIES OF AIR AND WATER AS FUNCTION OF TEMPERATURE

% Saturation vapor pressure in Pa

 $p_s_f=inline('610.78*exp((t/(t+238.3))*17.2694))$;

% Water Density in Kg/m^3

row_w_f=inline('1000*(1-((t+288.9414)*(t-3.9863)^2/(508929.2*(t+68.12963))))');

% Air Density in Kg/m^3

row_g_f=inline('(288.16*1.2255/(273.15+T0_c))');

% Colebrook-White Formula to calculate friction factor

%% $1/\sqrt{(f)} = -2 \cdot \log(((e/D)/3.7) + (2.51/(R \cdot \sqrt{sqrt(f)})))$;

f=0.0371; % This value is obtained from above formula

%%% %%%%%%%%%%%%%%

% INITIAL PRESSURE HEAD AND MIXTURE DENSITY CALCULATIONS

for $i=1:1:N$ x

H(i,1)= H0-((8*f*(i-1)*Delta_x*Q0^2)/(pi^2*g*Dia^5))-Q0^2/((pi*Dia^2/4)^2*2*g);;

end

% Initialising the boundary grid point near resevoir

$$
p(1,1)=1.418*10^5;
$$

\n
$$
row_w(1,1)=row_w_f(T0_c);
$$

\n
$$
row_g(1,1)=row_g_f(T0_c);
$$

\n
$$
row(1,1)=996.26;
$$

\n
$$
p_s(1,1)=p_s_f(T0_c);
$$

\n
$$
Ki_visc(1,1)=.837*10^5-6;
$$

for i=2:1:N_x-1

row_w(i,1)=row_w_f(T0_c); row_g(i,1)=row_g_f(T0_c); $p_s(i,1)=p_s_f(T0_c);$ $m(i,1)=(m0/(N_x-1));$

% Solve Quadratic Expression for Absolute Pressure

a=1; $b = -(g * H(i,1) * row_w(i,1) + p_atm-p_s(i,1));$ $c0 = g * m(i,1) * Gas_Const * (273.15 + T0_c) * H(i,1) * (row_w(i,1) - row_g(i,1));$

p(i,1)=(-b+(b^2-4*a*c0)^.5)/(2*a);

% Mixture Density

row(i,1)=(p(i,1)+p_s(i,1)-p_atm)/(H(i,1)*g);

end

%%% %%%%%%%%%%%%%%

% INITIAL VELOCITY PROFILE CALCULATIONS

% Grid points in Radial Direction

for $i=1:1:N_y$

 $r_N(i,1)=Di a/2-(Delta_y*(N_y-i));$ if $i < N_y$ $r_D(i,1)=r_N(i,1)+Delta_y/2;$ else $r_D(i,1)=r_N(i,1);$

end

end

% Reynold's Number R=round((4*Q0)/(pi*Dia*Ki_Visc0));

% Constant K throughout simulation k=0.374+(.0132*log(1+(83100/R)));

% Delta_A Delta_A= (1/(N_y-1))*(pi/4)*Dia^2;

for i=2:1:N_x-1 % Friction Velocity ustar(i,1)= $((4*Q0)/(pi*Dia^2))*sqrt(f/8);$

 % Sublayer thickness Delta $(i,1)=11.63*Ki$ _Visc0/ustar $(i,1);$

% Special Parameters

Phi(i,1)=($p(i,1)/(row(i,1)*g)$)- $(c^2 * m(i,1)*Gas_Const*(273.15+TO_c)/(g*p(i,1))) + c^2 * m(i,1)/(row(i,1)*g);$ $s(i,1)=cp*(log(273.15+TO_c)-(sp_ratio-1)*log(p(i,1))/sp_ratio);$ end

% Validating if a Grid point is in viscous sublayer or in Transient region % and then obtain the initial velocity profile

for $j=N_y:-1:1$ $u(1,N_y)=0;$

end

for i=2:1:N_x-1

for $j=N_y:-1:1$

if $Dia/2-r_D(i,1) \le Delta(i,1)$ $u(i,j)$ =ustar(i,1)^2*(Dia/2-r_D(j,1))/Ki_Visc0;

else

 $u(i,j) = u\frac{\sin(i,1) * (2.5 * \log((Di a/2-r_D(j,1))/e) + 5.0)}{$;

end

```
l(j,1)=k*(Dia/2-r_N(j,1))*exp(-(Dia/2-r_N(j,1))/(Dia/2));
```
end

end

 $t=1$; status=1; $TOL=10^{\circ}-4;$ error=1; while status <= 4

```
for i=2:1:N x-1
   count=0;
   flag=0;
  Ki Visc(i,1)=.837*10^-6;
  for j=N_y-1:-1:1if Dia/2-r N(i,1) \le Delta(i,1)
```
if j== N_y-1

```
D(1,i)=(1/DeltaT)+(pi/row(i,t))*(1/Delta_A)*Ki\_Visc(i,t)*row_w(i,t)*(((r_N(N_y,1)/(r_D(N_y,1))))),1)-r_D(N_y-1,1)))+(r_N(N_y-1,1)/(r_D(N_y-1,1)-r_D(N_y-2-1,1))));
```

```
A(1,i)=pi^*r_N(N_y-1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(N_y-1,1))1,1)-r_D(N_y-2,1)));
```

```
C(1,i)=-((-u(i,N_y-1)/DeltaT)+(g*(H(i+1,1)-
```

```
H(i,1)/\text{Delta}_X +(pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*((r_N(N_y,1)*u(i,N_y-
```

```
1)/(r_D(N_y,1)-r_D(N_y-1,1)))-(r_N(N_y-1,1)*u(i,N_y-2)/(r_D(N_y-1,1)-r_D(N_y-
```

```
(2,1)) + (r_N(N_y-1,1)*u(i,N_y-1)/(r_D(N_y-1,1)-r_D(N_y-2,1))));
```
flag= $flag+1$;

else

 $B(2+count,i)=-$

```
pi*Ki_Visc(i,t)*row_w(i,t)*r_N(j+1,1)/(row(i,t)*Delta_A*(r_D(j+1,1)-r_D(j,1)));
```

```
D(2+\text{count},i)=(1/DeltaT)+(\text{pi}/\text{row}(i,t))^*(1/Delta_A)*Ki\_Visc(i,t)*row_w(i,t)*((r_N(i+1,1)/(r_L));D(j+1,1)-r\_D(j,1))+(r\_N(j,1)/(r\_D(j,1)-r\_D(j-1,1))));
```

```
A(2+\text{count},i)=-pi*r_N(j,1)*Ki\_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-r_D*(r_I));r D(i-1,1));
            C(2+count,i)=((-u(i,j)/DeltaT)+(g*(H(i+1,1)-
```

```
H(i,1)/\text{Delta}_x + (pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*((r_N(j+1,1)*(u(i,j)-
```

```
u(i,j+1)/(r_D(j+1,1)-r_D(j,1))(r_N(j,1)*(u(i,j-1)-u(i,j))/(r_D(j,1)-r_D(j-1,1))));
```

```
 count=count+1;
flag = flag + 1;
```
end

```
 end 
   if Dia/2-r N(j,1)>Delta(i,1)
```
if $j == N_y-1$

```
D(1,i)=(1/DeltaT)+(pi*Ki-Visc(i,t)*row(w(i,t)*r_N(i+1,1))/(row(i,t)*DeltaA* (r_D(i+1,1)-r_D(i+1,1))r_D(i,1))))+((pi*r_N(i,1)*row_w(i,t))/(row(i,t)*Delta_A*(r_D(j,1)-r_D(j-
(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
           A(1,i)=-(pi*r_N(i,1)*row_N(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-r_D(i-))(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
           C(1,i)=-((-u(i,j)/DeltaT)+(g*(H(i+1,1)-H(i,1)/\text{Delta}_x + (pi/row(i,t))* (Ki_Visc(i,t)*row_w(i,t)/Delta_A)*(r_N(j+1,1)/r_D(j+1,1)-
r_D(i,1))^*u(i,j)+(pi^*r_N(i,1)*row_w(i,t)*(u(i,j)-u(i,j-1))/(row(i,t)*DeltaA*(r_D(i,1)-r_D(i-1)))(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1)))));
           flag=flag+1;
```

```
elseif Dia/2-r N(i+1,1) Delta(i,1)
```
 $B(flag+1,i)=$

 $pi^*r_N(j+1,1)*Ki_Nisc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(j+1,1)-r_D(j,1)));$

 $D(flag+1,i)=$

```
(1/DeltaT) + ((pi/row(i,t))^*(1/Delta_A)*Ki_Visc(i,t)*row_w(i,t)*( (r_N(j+1,1)/(r_D(j+1,1)-1))
```

```
r_D(i,1))+(r_N(i,1)/(r_D(i,1)-r_D(i-1))
```

```
1,1)))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
```

```
(1,1)))*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
```

```
A(flag+1,i) = (-pi^*r_N(i,1)*Ki_Nsc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_N(i,1)-i))r_D(j-1,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
(1,1))<sup>*</sup>abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
```

```
C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H(i+1,1)-
```

```
H(i,1)/\text{Delta}_x + ((pi/row(i,t))^*(Ki\_Visc(i,t)*row_w(i,t)/\text{Delta}_A)*(r_N(j+1,1)*(u(i,j)-\text{angle}_x(i,t))/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{Delta}_y(i,t)/\text{
```

```
u(i,j+1)/(r_D(j+1,1)-r_D(j,1)) + (r_N(j,1)*(u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1))
```

```
1,1)))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-
r\_D(i-1,1))<sup>*</sup>((u(i,j)-u(i,j-1))/(r\_D(i,1)-r\_D(i-1,1))));
```
 $flag = flag + 1;$

elseif r_N(j,1) $\sim= 0$

 $B(flag+1,i)=(-pi^*r_N(i+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i+1,1)-1))$

r_D(j,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1)-

 $r_D(i,1))$ ^{*}abs((u(i,j+1)-u(i,j))/(r_{_}D(j+1,1)-r_D(j,1))));

 $D(flag+1,i)=(1/DeltaT)+((pi/row(i,t))^*(1/Delta_A)*Ki_Visc(i,t)*row_w(i,t)*(r_N(i+1,1)/(r_D))$ $(j+1,1)-r_D(j,1))+(r_N(j,1)/(r_D(j,1)-r_D(j-1))$

1,1)))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*((l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1)-r_D(j-

1,1)))*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1))))+(l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-

 $(1,1))$ ^{*}abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))))));

 $A(flag+1,i)=(-pi*r_N(i,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-r_D(i-1,1))))$ $((pi/row(i,t))^*(1/Delta_A)*row_w(i,t)^*[i,j]^2*r_N(i,1)^*(1/(r_D(i,1)-r_D(i-1,1)))^*abs((u(i,j)-r_D(i+1),i^*)$ $u(i,j-1)/(r_D(j,1)-r_D(j-1,1))))$;

 $C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H(i+1,1)-H(i,1))/Delta_x)$ -

```
((pi/row(i,t))*(r_N(i+1,1)/\Delta^*)*(Ki_Ni,t)*row_w(i,t)*((u(i,j+1)-u(i,j))/(r_N(i+1,1)-i))
```

```
r_D(i,1)) + row_w(i,t)*l(j+1,1)^2*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1)))*((u(i,j+1)-
```
 $u(i,j)/(r_D(j+1,1)-$

```
r_D(i,1)))) + ((pi/row(i,t))^*(r_N(i,1)/Delta_A)*(Ki_iv_isc(i,t)*row_w(i,t)*(u(i,j)-u(i,j-1))))
```
1))/(r_D(j,1)-r_D(j-1,1)))+row_w(i,t)*l(j,1)^2*abs((u(i,j)-u(i,j))/(r_D(j,1)-r_D(j-1,1)))*((u(i,j) $u(i,j-1)/(r_D(j,1)-r_D(j-1,1))))$;

 $flag = flag + 1$;

elseif $r_N(i,1) == 0$

```
D(flag+1,i)=(1/DeltaT)+(pi^*r_N(i+1,1)*Ki_Nisc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_N(i+1,1))))-r_D(j,1))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1)-
r_D(j,1)))*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1))));
B(flag+1,i)=(-pi^*r_N(j+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(j+1,1)-1))r_D(i,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1)-
r_D(i,1))<sup>*</sup>abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1))));
C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H(i+1,1)-H(i,1))/DeltaX)-
((pi/row(i,t))*(r_N(i+1,1)/Delta_A)*(Ki_Nisc(i,t)*row_w(i,t)*(u(i,j+1)-u(i,j))/(r_N(i+1,1)-u(i,j))
```
 $r_D(j,1))$ + row $_w(i,t)*l(j+1,1)^2*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1)))*(u(i,j+1)-r_D(j,1)))$ $u(i,j)/(r_D(j+1,1)-r_D(j,1))))$);

end

end

end

end

% Thomas Algorithm

 $B(1,1)=0;$ for i=2:1:N_x-1 $D1(1,i)=D(1,i);$ $C1(1,i)=C(1,i);$ $B(1,i)=0;$

for j=2:1:N_y-1

 $D1(j,i)=D(j,i)-(B(j,i)*A(j-1,i)/D1(j-1,i));$ $C1(j,i)=C(j,i)-(C1(j-1,i)*B(j,i)/D1(j-1,i));$

end

 $u(i,1)=C1(N_y-1,i)/D1(N_y-1,i);$

for $j=2:1:N_y-1$ $u(i,j)=(C1(N_y-j,i)-A(N_y-j,i)*u(i,j-1))/D1(N_y-j,i);$

end

end

```
status=status+1;
```
end

% OBTAIN DISCHARGE AT A SECTION FROM VELOCITY PROFILE:

```
for i=2:1:N_x-1
temp(i,1)=u(i,1)*pi*r_N(2,1)^2;
tempo=0; 
for j=2:1:N_y-1
tempo=tempo+u(i,j)*pi*(r_N(j+1,1)^2-r_N(j,1)^2);
end
Q(i,1)=temp(i,1)+tempo;end
```
 $Q(1,1)=2*Q(2,1)-Q(3,1);$ $Q(N_x,1)=0;$ $m(N_x,1)=m0/(N_x-1);$

% BEGIN OF MACCORMACK SCHEME

%Setting up the initial guess values

for $i=2:1:N_x-1$ $p_p(i,1)=p(i,1);$ $p_c(i,1)=p(i,1);$

end

for $t=1:1:40$

% Predictor Step

t

```
 for i=2:1:N_x-1
  T(i,1)=T0_k;Phi_p(i,1)=Phi(i,t)+(c^2/(g*Area))*(Q(i,t)-Q(i+1,t))*(DeltaT/Delta_x);
  s_p(i,1)=s(i,t)+(cp/Theta_T)*((T0_k/T(i,t))-1)*DeltaT;
```

```
 count=0;
flag=0;
Ki Visc(i,t)=.837*10^-6;
```
for $j=N_y-1:-1:1$ if Dia/2-r $N(i,1) \leq Delta(i,1)$

if $j == N_y-1$

```
D(1,i)=(1/DeltaT)+(pi/row(i,t))^*(1/Delta_A)*Ki\_Visc(i,t)*row_w(i,t)^*((r_N(N_y,1)/(r_D(N_y,1))^2)(t_N,1)(r_N,1)^*,1)-r_D(N_y-1,1)))+(r_N(N_y-1,1)/(r_D(N_y-1,1)-r_D(N_y-2-1,1))));
            A(1,i)=pi^*r_N(N_y-1,1)*K_i/Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(N_y-1,1))1,1)-r_D(N_y-2,1)));
            C(1,i)=-((-u(i,N_v-1)/DeltaT)+(g*(H(i+1,1)-H(i,1)/\text{Delta}_x + (pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*((r_N(N_y,1)*u(i,N_y-
1)/(r_D(N_y,1)-r_D(N_y-1,1)))-(r_N(N_y-1,1)*u(i,N_y-2)/(r_D(N_y-1,1)-r_D(N_y-
2,1)))+(r_N(N_y-1,1)*u(i,N_y-1)/(r_D(N_y-1,1)-r_D(N_y-2,1)))));
            flag=flag+1;
           else
            B(2+count,i)=-pi*Ki\_Visc(i,t)*row_w(i,t)*r_N(j+1,1)/(row(i,t)*Delta_A*(r_D(j+1,1)-r_D(j,1)));D(2+count,i)=(1/DeltaT)+(pi/row(i,t))*(1/Delta A)*Ki_Visc(i,t)*row_w(i,t)*((r_N(j+1,1)/(r_
D(j+1,1)-r\_D(j,1))+(r\_N(j,1)/(r\_D(j,1)-r\_D(j-1,1))));
            A(2+\text{count},i)=-pi^*r_N(i,1)*Ki_Nisc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-r_D*(r_I));r\_D(j-1,1));
            C(2+count,i)=((-u(i,j)/DeltaT)+(g*(H(i+1,1)-
```

```
H(i,1)/\text{Delta}_x + (pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*((r_N(j+1,1)*(u(i,j)-
u(i,j+1)/(r_D(j+1,1)-r_D(j,1)) (r_N(j,1)*(u(i,j-1)-u(i,j))/(r_D(j,1)-r_D(j-1,1))));
```

```
 count=count+1;
     flag = flag + 1;
   end
 end 
if Dia/2-r N(i,1)>Delta(i,1)
```

```
if j == N_y-1
```

```
D(1,i)=((1/DeltaT)+(pi*Ki_Visc(i,t)*row_w(i,t)*r_N(i+1,1))/(row(i,t)*Delta_A*(r_D(i+1,1)-r_N(i+1,1))r_D(i,1))))+((pi*r_N(j,1)*row_w(i,t))/(row(i,t)*Delta_A*(r_D(j,1)-r_D(j-
(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
             A(1,i)=-(pi*r_N(i,1)*row_N(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-r_D(i-))(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
             C(1,i)=-((-u(i,j)/\text{DeltaT})+(g^*(H(i+1,1)-H(i,1)/\text{Delta}_x + (pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*(r_N(j+1,1)/r_D(j+1,1)-
r_D(i,1))^*u(i,j)+(pi^*r_N(i,1)*row_w(i,t)*(u(i,j)-u(i,j-1))/(row(i,t)*DeltaA*(r_D(i,1)-r_D(i-1)))1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1)))));
             flag = flag + 1;
```
elseif $Dia/2-r_N(i+1,1) < Delta(i,1)$

 $B(flag+1,i)=$

```
pi^*r_N(i+1,1)*Ki_Nisc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i+1,1)-r_D(i,1)));
```

```
D(flag+1,i)=
```

```
(1/DeltaT) + ((pi/row(i,t)) * (1/Delta_A) * Ki_Visc(i,t) * row_w(i,t) * ((r_N(j+1,1)/(r_D(j+1,1)-1))
```

```
r_D(j,1))+(r_N(j,1)/(r_D(j,1)-r_D(j-1))
```

```
1,1)))))+((pi/row(i,t))*(1/Delta A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
```
 $(1,1))$ ^{*}abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));

```
A(flag+1,i)= (-pi^*r_N(i,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_N(i,1)-
```

```
r_D(j-1,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
```

```
(1,1)))*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
```

```
C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H(i+1,1)-
```
 $H(i,1)/\text{Delta}_x$ + $((pi/row(i,t))^*(Ki_Visc(i,t)*row_w(i,t)/\text{Delta}_A)*(r_N(i+1,1)*(u(i,j)-\text{angle}_x(i+1))$

```
u(i,j+1)/(r_D(j+1,1)-r_D(j,1))+(r_N(j,1)*(u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1))1,1)))))+((pi/row(i,t))*(1/Delta A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-
r\_D(i-1,1))<sup>*</sup>((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1)))));
             flag=flag+1;
```
elseif r_N(j,1) $\sim= 0$

 $B(flag+1,i)=(-$

pi*r_N(j+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(j+1,1)-r_D(j,1))))- $((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(i+1,1)^2*r_N(i+1,1)*(1/(r_D(i+1,1)-1))$ $r_D(i,1))$ ^{*}abs((u(i,j+1)-u(i,j))/(r_{_}D(j+1,1)-r_D(j,1))));

```
D(flag+1,i)=(1/DeltaT)+((pi/row(i,t))*(1/Delta_A)*Ki\_Visc(i,t)*row_w(i,t)*( (r_N(i+1,1)/(r_D))(j+1,1)-r_D(j,1))+(r_N(j,1)/(r_D(j,1)-r_D(j-1))1,1)))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*((l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1)-r_D(j-
1,1)))*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1))))+(l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
(1,1))<sup>*</sup>abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))))));
```

```
A(flag+1,i)=(-pi^*r_N(i,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-n))r_D(j-1,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
(1,1))<sup>*</sup>abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
```

```
C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H(i+1,1)-H(i,1))/Delta_x)((pi/row(i,t))^*(r_N(i+1,1)/Delta_A)*(Ki_Nisc(i,t)*row_w(i,t)*(u(i,j+1)-u(i,j))/(r_N(i+1,1)-u(i+1))r_D(j,1)))+row_w(i,t)*l(j+1,1)^2*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1)))*((u(i,j+1)-
u(i,j)/(r_D(j+1,1)-
```

```
r_D(i,1)))))+((pi/row(i,t))*(r_N(j,1)/Delta_A)*(Ki_Visc(i,t)*row_w(i,t)*((u(i,j)-u(i,j-
```

```
1))/(r_D(j,1)-r_D(j-1,1)))+row_w(i,t)*l(j,1)^2*abs((u(i,j)-u(i,j))/(r_D(j,1)-r_D(j-1,1)))*((u(i,j)-
u(i,j-1)/(r\_D(j,1)-r\_D(j-1,1))));
```

```
flag=flag+1;
elseif r_N(j,1) == 0
```
D(flag+1,i)=(1/DeltaT)+(pi*r_N(j+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(j+1,1))-r_D(j,1))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1) $r_D(i,1))$ ^{*}abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1))));

```
B(flag+1,i)=(-pi^*r_N(i+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i+1,1)-r_D(i,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(i+1,1)^2*r_N(i+1,1)*(1/(r_D(i+1,1)-1))r\_D(j,1))<sup>*</sup>abs((u(i,j+1)-u(i,j))/(r\_D(j+1,1)-r\_D(j,1))));
             C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H(i+1,1)-H(i,1))/Delta_x)((pi/row(i,t))*(r_N(j+1,1)/Delta_A)*(Ki_Visc(i,t)*row_w(i,t)*(u(i,j+1)-u(i,j))/(r_N(j+1,1)-u(i,j))r_D(i,1)) + row_w(i,t)*l(j+1,1)^2*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1)))*((u(i,j+1)-
u(i,j)/(r\_D(j+1,1)-r\_D(j,1))));
```
 end end end

end

% Thomas Algorithm

 $B(1,1)=0$; for $i=2:1:N_x-1$ $D1(1,i)=D(1,i);$ $C1(1,i)=C(1,i);$ $B(1,i)=0;$

for $j=2:1:N_y-1$

 $D1(j,i)=D(j,i)-(B(j,i)*A(j-1,i)/D1(j-1,i));$ $C1(i,i)=C(i,i)-(C1(i-1,i)*B(i,i)/D1(i-1,i));$

end

 $u_p(i,1)=C1(N_y-1,i)/D1(N_y-1,i);$

for $j=2:1:N_y-1$ $u_p(i,j)=(C1(N_y-j,i)-A(N_y-j,i)*u(i,j-1))/D1(N_y-j,i);$ end

end

```
 % Finding predictor value of Discharge Q
 for i=2:1:N_x-1
  temp(i,1)=u_p(i,1)*pi*r_N(2,1)^2;
   tempo=0; 
  for j=2:1:N_y-1tempo=tempo+u_p(i,j)*pi*(r_N(j+1,1)^2-r_N(j,1)^2);
   end
  Q_p(i,1)=temp(i,1)+tempo; end
```
% Finding predictor value of Pressure p and Temperature T

for $i=2:1:N$ x-1

cff1= $(1/(row(i,t)*g))$; cff2=- $(c^2*m_p(i,1)*Gas_Const/g);$ cff3=((c^2*m_p(i,1)/(row(i,t)*g))-Phi_p(i,1)); cff4= $(s_p(i,1)/cp)$; cff5=((sp_ratio-1)/sp_ratio);

```
f01=inline('cf1*x^2+cf2*exp(cf5*log(x)+cf4)+cf3*x','x','cf1','cf2','cf3','cf4','cf5');
[p_p(i,1), fval, exitFlag] = fzero(f01, p_p(i,1), [], cff1, cff2, cff3, cff4, cff5);
```

```
 if(exitFlag<0)
 % fprintf('You faced a NaN value') 
p_p(i,1)=fzero(f01,100,[],cff1,cff2,cff3,cff4,cff5);
 end
T_p(i,1)=exp(cff5*log(p_p(i,1))+cff4);
```
if $T_p(i,1) < 273.15$ $p_s_p(i,1)=0;$ else p_s_p(i,1)=p_s_f(T_p(i,1)-273.15); end

 $H_p(i,1)=(p_p(i,1)+p_s_p(i,1)-p_atm)/(row(i,t)*g);$ end

 %Interpolating to obtain the values at Boundary points $H_p(1,1)=H0;$ $Q_p(1,1)=2*Q_p(2,1)-Q_p(3,1);$

```
 % Corrector Step
```
for $i=2:1:N_x-1$

Phi_c(i,1)=Phi(i,t)+(c^2/(g*Area))*(Q_p(i-1,1)-Q_p(i,1))*(DeltaT/Delta_x); s_c(i,1)=s(i,t)+(cp/Theta_T)*((T0_k/T_p(i,1))-1)*DeltaT; $m_c(i,1)=m(i,t)+0.02*(Delta T/(Theta_m*Gas_Const*T(i,t)))*(p_s_p(i,1)-p_p(i,1));$

 count=0; flag=0; Ki_Visc(i,t)=.837*10^-6;

for $j=N_y-1:-1:1$

if $Dia/2-r_N(j,1) \leq Delta(i,1)$

if $j == N_y-1$
```
D(1,i)=(1/DeltaT)+(pi/row(i,t))^*(1/Delta_A)^*Ki Visc(i,t)*row_w(i,t)*((r_N(N_y,1)/(r_D(N_y
,1)-r_D(N_y-1,1)))+(r_N(N_y-1,1)/(r_D(N_y-1,1)-r_D(N_y-2-1,1))));
```

```
A(1,i)=pi^*r_N(N_y-1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(N_y-1,1))1,1)-r_D(N_y-2,1)));
```

```
C(1,i)=-((-u(i,N_y-1)/DeltaT)+(g*(H_p(i,1)-H_p(i-1)))
```

```
1,1))/Delta_x)+(pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*((r_N(N_y,1)*u(i,N_y-
```

```
1)/(r_D(N_y,1)-r_D(N_y-1,1)))-(r_N(N_y-1,1)*u(i,N_y-2)/(r_D(N_y-1,1)-r_D(N_y-
```

```
2,1)))+(r_N(N_y-1,1)*u(i,N_y-1)/(r_D(N_y-1,1)-r_D(N_y-2,1)))));
```
flag= $flag+1$;

else

 $B(2+count,i)=-$

```
pi*Ki_Visc(i,t)*row_w(i,t)*r_N(j+1,1)/(row(i,t)*Delta_A*(r_D(j+1,1)-r_D(j,1)));
```

```
D(2+\text{count},i)=(1/DeltaT)+(\text{pi}/\text{row}(i,t))^*(1/Delta_A)*Ki\_Visc(i,t)*row_w(i,t)*((r_N(i+1,1)/(r_N(i+1,1))^2))D(j+1,1)-r\_D(j,1))+(r\_N(j,1)/(r\_D(j,1)-r\_D(j-1,1))));
             A(2+\text{count},i)=-pi^*r_N(i,1)*Ki_Nisc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-r_D));r\_D(i-1,1));
             C(2+count,i)=((-u(i,j)/DeltaT)+(g*(H-p(i,1)-H-p(i-1)))1,1))/Delta_x)+(pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*((r_N(j+1,1)*(u(i,j)-
u(i,j+1)/(r_D(j+1,1)-r_D(j,1))(r_N(j,1)*(u(i,j-1)-u(i,j))/(r_D(j,1)-r_D(j-1,1))));
              count=count+1;
             flag = flag + 1;
           end
         end 
        if Dia/2-r_N(i,1) > Delta(i,1)
```

```
if j == N y-1
```

```
D(1,i)=((1/DeltaT)+(pi*Ki_Visc(i,t)*row_w(i,t)*r_N(j+1,1))/(row(i,t)*Delta_A*(r_D(j+1,1)-r_N(i,t)*Delta_A*(r_D(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(r_N(i,t)*Deta_A*(rr_D(i,1))))+((pi*r_N(i,1)*row_w(i,t))/(row(i,t)*Delta_A*(r_D(j,1)-r_D(j-
(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
```

```
A(1,i)=-(pi^*r_N(i,1)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-r_D(i-))(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
             C(1,i)=-((-u(i,j)/\text{DeltaT})+(g*(H-p(i,1)-H-p(i-1,1))/Delta_x)+(pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*(r_N(j+1,1)/r_D(j+1,1)-
r_D(i,1))^*u(i,j)+(pi*r_N(i,1)*row_w(i,t)*(u(i,j)-u(i,j-1))/(row(i,t)*Delta_A*(r_D(i,1)-r_D(i-1)))(1,1))))*(Ki_Visc(i,t)+l(j,1)^2*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1)))));
```
flag= $flag+1$;

elseif $Dia/2-r_N(i+1,1) < Delta(i,1)$

 $B(flag+1,i)=$

```
pi^*r_N(i+1,1)*Ki_Nisc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i+1,1)-r_D(i,1)));
```
 $D(flag+1,i)=$

```
(1/DeltaT) + ((pi/row(i,t)) * (1/Delta_A) * Ki_Visc(i,t) * row_w(i,t) * ((r_N(j+1,1)/(r_D(j+1,1)-1))
```
 $r_D(j,1))+(r_N(j,1)/(r_D(j,1)-r_D(j-1))$

```
1,1)))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
```
 $(1,1))$ ^{*}abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));

```
A(flag+1,i)= (-pi^*r_N(j,1)*Ki_Visc(i,t)*row_N(i,t)/(row(i,t)*Delta_A*(r_N(i,1)-r_D(j-1,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
(1,1))<sup>*</sup>abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
```
 $C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H_p(i,1)-H_p(i-1)))$

```
1,1))/Delta_x)+((pi/row(i,t))*(Ki_Visc(i,t)*row_w(i,t)/Delta_A)*((r_N(j+1,1)*(u(i,j)-
```

```
u(i,j+1)/(r_D(j+1,1)-r_D(j,1)) + (r_N(j,1)*(u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1))
```

```
1,1)))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-
```

```
r\_D(i-1,1))<sup>*</sup>((u(i,j)-u(i,j-1))/(r\_D(j,1)-r\_D(i-1,1))));
```
flag= $flag+1$;

elseif r_N(j,1) $\sim= 0$

 $B(flag+1,i)=(-$

 $pi^*r_N(j+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(j+1,1)-r_D(j,1))))-$

 $((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(i+1,1)-1))$

 $r_D(i,1))$ ^{*}abs((u(i,j+1)-u(i,j))/($r_D(i+1,1)$ -r $_D(i,1))$));

```
D(flag+1,i)=(1/DeltaT)+(p_i/row(i,t))^*(1/Delta_A)*KiVisc(i,t)*row(w(i,t))^*(rN(i+1,1)/(rD))(j+1,1)-r_D(j,1))+(r_N(j,1)/(r_D(j,1)-r_D(j-1))1,1)))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*((l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1)-r_D(j-
1,1)))*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1))))+(l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
(1,1))<sup>*</sup>abs((u(i,j)-u(i,j-1))/(r_D(i,1)-r_D(i-1,1))));
             A(flag+1,i)=(-pi^*r_N(i,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i,1)-n))r_D(i-1,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j,1)^2*r_N(j,1)*(1/(r_D(j,1)-r_D(j-
(1,1)))*abs((u(i,j)-u(i,j-1))/(r_D(j,1)-r_D(j-1,1))));
             C(flag+1,i)=-((-u(i,j)/DeltaT)+(g*(H_p(i,1)-H_p(i-1,1))/Delta_x)((pi/row(i,t))*(r_N(i+1,1)/\Delta^*)*(Ki_Nisc(i,t)*row_w(i,t)*((u(i,j+1)-u(i,j))/(r_N(i+1,1)-u(i+1,1))r_D(j,1)) + row_w(i,t)*l(j+1,1)^2*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1)))*((u(i,j+1)-
u(i,j)/(r\_D(j+1,1)-
r\_D(i,1))))+(p i/row(i,t))^*(r\_N(j,1)/Delta_A)*(Ki\_Visc(i,t)*row\_w(i,t)*( (u(i,j)-u(i,j-s))1))/(r_D(j,1)-r_D(j-1,1)))+row_w(i,t)*l(j,1)^2*abs((u(i,j)-u(i,j))/(r_D(j,1)-r_D(j-1,1)))*((u(i,j)-
u(i,j-1)/(r\_D(j,1)-r\_D(j-1,1))));
```
flag= $flag+1$; elseif r_N(j,1) = = 0

```
D(flag+1,i)=(1/DeltaT)+(pi*r_N(i+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(i+1,1)))-r_D(j,1))))+((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(j+1,1)^2*r_N(j+1,1)*(1/(r_D(j+1,1)-
r_D(j,1)))*abs((u(i,j+1)-u(i,j))/(r_D(j+1,1)-r_D(j,1))));
             B(flag+1,i)=(-pi^*r_N(j+1,1)*Ki_Visc(i,t)*row_w(i,t)/(row(i,t)*Delta_A*(r_D(j+1,1)-r_D(i,1))))-((pi/row(i,t))*(1/Delta_A)*row_w(i,t)*l(i+1,1)^2*r_N(i+1,1)*(1/(r_D(i+1,1)-1))r\_D(i,1))<sup>*</sup>abs((u(i,j+1)-u(i,j))/(r\_D(i+1,1)-r\_D(i,1))));
             C(flag+1,i)=((-u(i,j)/DeltaT)+(g*(H-p(i,1)-H-p(i-1,1))/DeltaX)-
((pi/row(i,t))^*(r_N(j+1,1)/Delta_A)*(Ki_Visc(i,t)*row_w(i,t)*(u(i,j+1)-u(i,j))/(r_N(j+1,1)-u(i,j))r\_D(i,1)) + row (w(i,t)*1(i+1,1)^2* abs((u(i,j+1)-u(i,j))/r\_D(i+1,1)-r\_D(i,1)))*((u(i,j+1)-u(i,j)/r\_D(i+1,1)-r\_D(i,1))u(i,j)/(r\_D(j+1,1)-r\_D(j,1))));
```
 end end

end

end

% Thomas Algorithm

 $B(1,1)=0;$ for i=2:1:N_x-1 $D1(1,i)=D(1,i);$ $C1(1,i)=C(1,i);$ $B(1,i)=0;$

for j=2:1:N_y-1

 $D1(j,i)=D(j,i)-(B(j,i)*A(j-1,i)/D1(j-1,i));$ $C1(j,i)=C(j,i)-(C1(j-1,i)*B(j,i)/D1(j-1,i));$ end

 $u_c(i,1)=C1(N_y-1,i)/D1(N_y-1,i);$

for j=2:1:N_y-1

 $u_c(i,j)=(C1(N_y-j,i)-A(N_y-j,i)*u(i,j-1))/D1(N_y-j,i);$

end

end

```
 % Finding predictor value of Discharge Q
```
 for i=2:1:N_x-1 temp(i,1)=u_c(i,1)*pi*r_N(2,1)^2; tempo=0; for $j=2:1:N_y-1$

```
tempo=tempo+u_c(i,j)*pi*(r_N(j+1,1)^2-r_N(j,1)^2);
```
end

```
Q_c(i,1)=temp(i,1)+tempo;
```
end

% Finding corrector value of Pressure p and Temperature T

```
for i=2:1:N_x-1
```

```
cff1=(1/(row(i,t)*g));cff2=-(c^2*m_c(i,1)*Gas\_Const/g);cff3=((c^2*m_c(i,1)/(row(i,t)*g))-Phi_c(i,1));
\text{cff4}=(s_c(i,1)/cp); cff5=((sp_ratio-1)/sp_ratio);
```

```
f01=inline('cf1*x^2+cf2*exp(cf5*log(x)+cf4)+cf3*x','x','cf1','cf2','cf3','cf4','cf5');
[p_c(i,1), fval, exitFlag] = fzero(f01, p_c(i,1),[], cff1, cff2, cff3, cff4, cff5);
```

```
 if(exitFlag<0)
```

```
 % fprintf('You faced a NaN value') 
p_c(i,1)=fzero(f01,100,[],eff1,eff2,eff3,eff4,eff5); end
```

```
T_c(i,1)=exp(cff5*log(p_c(i,1))+cff4);
```

```
if T_c(i,1) < 273.15
```

```
p_s_c(i,1)=0;
```
else

```
p_s_c(i,1)=p_s_f(T_c(i,1)-273.15); end
```

```
H_c(i,1)=(p_c(i,1)+p_s_c(i,1)-p_atm)/(row(i,t)*g);
```

```
 end
```
 %Interpolating to obtain the values at Boundary points $H_c(1,1)=H0;$ $Q_c(1,1)=2^*Q_c(2,1)-Q_c(3,1);$

% OBTAINING THE PARAMETERS

```
for i=2:1:N_x-1
```

```
Phi(i,t+1)=0.5*(Phi(p_i,1)+Phi(c_i,1));s(i,t+1)=0.5*(s_p(i,1)+s_c(i,1));m(i,t+1)=0.5*(m_p(i,1)+m_c(i,1));Q(i,t+1)=0.5*(Q_p(i,1)+Q_c(i,1));p(i,t+1)=0.5*(p_p(i,1)+p_c(i,1));T(i,t+1)=0.5*(T_p(i,1)+T_c(i,1));p_s(i,t+1)=0.5*(p_s_p(i,1)+p_s_c(i,1));H(i,t+1)=0.5*(H_p(i,1)+H_c(i,1));for j=2:1:N_y-1u(i,j)=0.5*(u_p(i,j)+u_c(i,j));
```
end

```
row_w(i,t+1)=row_w_f(T(i,t+1)-273.15);
  row_g(i,t+1)=row_g_f(T(i,t+1)-273.15);
 end
```

```
Q(1,t+1)=2*Q(2,t+1)-Q(3,t+1);H(1,t+1)=4.53;
H(N_x,t+1)=2*H(N_x-1,t+1)-H(N_x-2,t+1);T(N_x,t+1)=2*T(N_x-1,t+1)-T(N_x-2,t+1);m(N_x,t+1)=2*m(N_x-1,t+1)-m(N_x-2,t+1);
```


```
 %Find Mass of gas released
 hold=0;
for i=2:1:N_x hold=hold+m(i,t); 
 end 
mass(t,1)=hold;
```

```
end
```
%% THE GRAPHS AND RESULTS figure(1) t=[1:1:40]';

plot(t,H(5,[1:1:40])); grid on; xlabel('t'); ylabel('H(m)'); title('Head Near the Valve');

figure(2) plot(t,T(5,[1:1:40])-273.15); grid on; xlabel('t'); ylabel('T(deg c)'); title('Temperature Near the Valve');

figure(3) plot(t,mass([1:1:40],1)*10^6); grid on; xlabel('t'); ylabel('m(mg/m3)'); title('Mass of gas release');

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