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Magnetoelectric effect in sputtered composites

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The magnetoelectric effect in millimeter size PMN-PT/Terfenol-D composites is known. In an effort towards miniaturization, we report on the magnetoelectric effect in micrometer-size sputtered composites. Multilayers of TbFe/FeCo with a thickness of 4 μm were sputter deposited on both sides of PMN-PT piezoelectric single crystals. The magnetoelectric voltage of samples was measured and reached values of 13 mV/(Oe cm) at dc bias fields of 2 mT, a linear dependence of magnetoelectric voltage on ac amplitude was detected in the range from 1 mT to 1 nT.

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I. INTRODUCTION

There are natural and composite magnetoelectric (ME) solids. Representatives of the former are BiFeO₃ (Refs. 1–4) and YMnO₃.⁵ Studies on ME bulk composites originated in The Netherlands in the 1970s and were based on the idea of mechanically coupling highly strictive ferromagnetic and ferroelectric solids. The coefficients relating the magnetic field H and either the electrical polarization P or the electrical field E through the ME_H charge coefficient α^P or ME_H^E voltage coefficient α^E would then be given by⁶

$$\alpha^P = \frac{\partial \sigma}{\partial H} \frac{\partial P}{\partial \sigma} \quad (1a)$$

and

$$\alpha^E = \frac{\partial \varepsilon}{\partial H} \frac{\partial E}{\partial \varepsilon}. \quad (1b)$$

It was also shown that α^E can be as large as 0.13 V/(cm Oe) in the eutectic composites BaTiO₃/CoFe₂O₄, BaTiO₃/(CoFe₂O₄/Ti₂O₄), and BaTiO₃/Ni_{0.97}Co_{0.03}Mn_{0.10}Fe_{1.90}O₄.^{7,8} Higher coefficients can be achieved by combining materials with larger strictive coefficients such as PMN and Terfenol.^{9–12} The value of α^E depends on the respective orientations of the field vectors¹³ and can be increased in resonant composites.¹⁴

The ferromagnetic components in the above composites possess a significant magnetocrystalline anisotropy. Therefore, a large magnetic bias field is required to develop a sizable ME voltage. This disadvantage can be circumvented by selecting amorphous ferromagnetic components such as amorphous FeCoSiB (Ref. 15) whose magnetostriction, however, is not large.¹⁶ TbFe/FeCo multilayers are both mag-

netically soft and highly strictive.¹⁷ They have thus been deposited onto highly strictive single-crystal PMN-PT substrates to produce candidates for miniaturized magnetic-field sensors. The preparation and ME properties of micron-sized single-crystal PMN-PT||TbFe//FeCo multilayers ME composites are described in this publication.

II. FABRICATION

The investigated ME composites consist of a piezoelectric substrate on both sides of which a magnetostrictive multilayer was sputter deposited. During each deposition the edges of the substrate were covered to inhibit short circuiting. The Tb₄₀Fe₆₀/Fe₅₀Co₅₀ multilayers were magnetron sputtered with a Von Ardenne CS 730S tool using cast targets (TbFe: diameter of 10 cm and FeCo: 20 cm). The base and Ar sputtering pressures were kept at 7×10^{-8} and 4×10^{-3} mbar, respectively. During the room-temperature deposition of TbFe and FeCo layers the rf power was adjusted to 125 and 350 W, resulting in respective deposition rates of 0.55 and 0.22 nm/s. Both materials were deposited with a magnetic bias field of 10 mT to induce an easy axis in the plane of the multilayer. The resulting multilayer had the composition [TbFe(7 nm)/FeCo(10 nm)]₂₆₂ with a total thickness of 4 μm .

The piezoelectric substrate was a 100- μm -thick (001)-oriented 0.68Pb(Mg_{1/3}Nb_{2/3})O₃-0.32PbTiO₃ (PMN-PT) piezoelectric single crystal with lateral dimensions of $5 \times 5 \text{ mm}^2$. It was poled perpendicular to the substrate plane.

III. MEASUREMENT METHOD

The ME composite was electroded and mounted on a carrier. In order to connect the bottom electrode as unconstrained as possible one side of the ME composite was glued at its center to the tip of a small pyramid formed of conduc-

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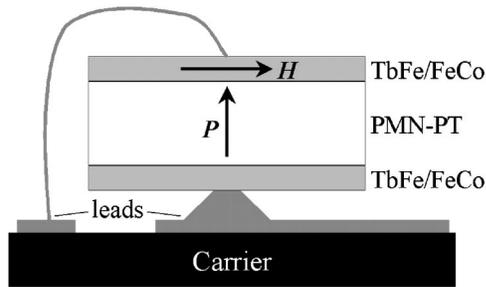


FIG. 1. ME composite layout and measurement setup for longitudinal-transversal (LT) mode alignment.

tive silver sitting on the carrier. A 25- μm -diameter aluminum wire lead to the top electrode was fabricated with the aid of a wire bonder. The carrier was mounted inside an electromagnet where a dc magnetic field up to 1 T could be applied. The Helmholtz coils positioned on the pole shoes of the solenoid allowed superposition of the dc field, H , and an ac field, H_{ac} , of up to 10 Oe (peak). The voltage drop across the sample was recorded via a lock-in amplifier.

During the measurement the ME samples were aligned for the longitudinal-transversal (LT) mode, as shown in Fig. 1.

The voltage that develops at the sample leads consists of a part stemming from the ME effect, U_{ME} , and a part induced by the measuring field H_{ac} into the loop formed by the sample leads. These contributions differ by a phase angle of $\pi/2$,

$$U_{\text{meas}} e^{i(\omega t + \Phi)} = \alpha^E H_{\text{ac}} e^{i\omega t} + \gamma H_{\text{ac}} e^{i\omega t + \pi/2}. \quad (2)$$

The quantity γ indicates the magnitude of the inductively coupled voltage. A phase-sensitive voltage meter was used to separate the ME voltage by adjusting the phase Φ such that $e^{i\Phi} \rightarrow \pm 1$ as $H \rightarrow \pm \infty$ so that $U_{\text{ME}} = U_{\text{meas}}(0)$.

For magnetization measurements, a Lakeshore vibrating-sample magnetometer (VSM) tool was employed.

IV. RESULTS

Figures 2 and 3 show the magnetization and ME voltages of a 4- μm ($\text{Tb}_{40}\text{Fe}_{60}/\text{Fe}_{50}\text{Co}_{50}$)₂₆₂//100- μm (PMN-PT)//4- μm ($\text{Tb}_{40}\text{Fe}_{60}/\text{Fe}_{50}\text{Co}_{50}$)₂₆₂ composite. The graphs identified by the open and closed circles represent measurements in which the applied magnetic bias field was directed parallel to the hard and easy axes, respectively. While the difference in the two resulting M - H characteristics is small

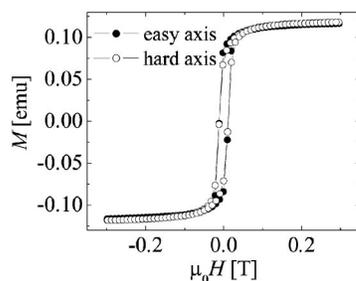


FIG. 2. Room-temperature magnetization of ME sample (4- μm multilayer electrodes) along easy and hard axes measured by vibrating-sample magnetometry (VSM).

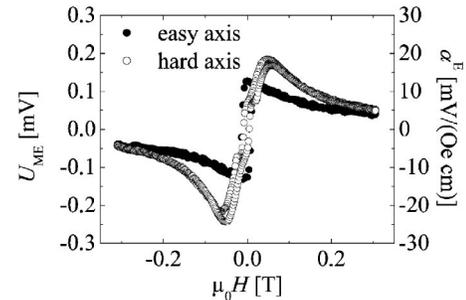


FIG. 3. Dependence of ME voltage on dc magnetic bias field in LT mode, measured with ac field of $f=1$ kHz and amplitude of 1 Oe.

(see Fig. 2); the two ME voltage-bias field characteristics differ significantly (see Fig. 3). Since in the LT mode the latter represents in essence the derivative $\partial \lambda / \partial H \propto \partial M / \partial H$ [see Eq. (3)] the enhancement of the difference is readily understood. Both characteristics in Fig. 3 exhibit the small hysteresis characteristic of a soft magnetic material.

In hard axis alignment, the maximum α^E reaches a value of 19 mV/(Oe cm) at a bias field of 40 mT. Although there is no asymmetry in peak position, the highest value in the opposite sweep direction at -40 mT lies at -24 mV/(Oe cm). The easy axis measurement shows no asymmetry and the maximum ME coefficient is 13 mV/(Oe cm) at a bias field of 2 mT. This asymmetry can be attributed to anisotropies: either the deposition area is asymmetric due to the covering of the edges or both magnetic electrodes are slightly misaligned regarding the sputtering bias field.

Figure 4 shows the linear dependence of the ME voltage drop on the applied ac field amplitude.

V. DISCUSSION

The ME voltage characteristic $U_{\text{ME}} = f(\mu_0 H)$ shown in Fig. 3 demonstrates that the magnitude of the ME voltage depends on the orientation of the easy axis of the TbFe/FeCo multilayer and that its maximum occurs at a small bias field of $H=2$ and 40 mT, respectively. The influence of the orientation of the easy axis with respect to the probing field H_{ac} can be understood by observing that magnetic switching between easy axes involves fewer 90° domain boundaries than switching between hard axes. The small magnitude of the bias field maximizing the ME voltage, H_{bias} , is a consequence of the small anisotropy field of the TbFe/FeCo multilayers, as can be seen from expressing Eq. (1b) as

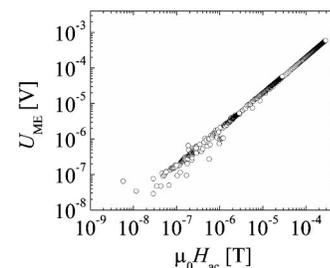


FIG. 4. Dependence of ME voltage on ac magnetic field (peak value). Hard axis alignment measured with ac frequency of 1 kHz at dc field of 60 mT.

$$\alpha^E = \left. \frac{\partial \lambda}{\partial H} \right|_{H_{\text{bias}}} \frac{\partial E}{\partial \lambda}, \quad (3)$$

where the quantity λ represents the magnetostriction of the multilayer.

VI. CONCLUSION

In a step towards miniaturized high-performance ME sensors a composite consisting of 100- μm ferroelectric PMN-PT with double-sided 4- μm magnetostrictive magnetically soft ($\text{Tb}_{40}\text{Fe}_{60}/\text{Fe}_{50}\text{Co}_{50}$)₂₆₂ “electrodes” was fabricated. Due to the excessive sputtering times required to deposit electrodes with a thickness comparable to that of the PMN substrate, only elastically mismatched composites with small ME coefficients were manufactured. It was observed, though, that the ME voltage is proportional to the thickness of the electrodes, as expected, if their thickness is small in comparison to that of the PMN-PT single-crystal substrate. If the observed ME voltage is converted to what would be observed in elastically matched composites the present values are comparable to those observed in bulk composites.

ACKNOWLEDGMENTS

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