

GEOMETRICALLY NONLINEAR FINITE ELEMENT ANALYSIS  
OF SPACE FRAMES

by

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(ABSTRACT)

The displacement method of the finite element is adopted. Both the updated Lagrangian formulation and total Lagrangian formulation of a three-dimensional beam element is employed for large displacement and large rotation, but small strain analysis.

A beam-column element or finite element can be used to model geometrically nonlinear behavior of space frames. The two element models are compared on the basis of their efficiency, accuracy, economy and limitations.

An iterative approach, either Newton-Raphson iteration or modified Riks/Wempner iteration, is employed to trace the nonlinear equilibrium path. The latter can be used to perform postbuckling analysis.

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Chapter I  
INTRODUCTION

1.1 PURPOSE AND SCOPE

The primary purpose of this dissertation is to implement and compare three formulations in geometrically nonlinear finite element analysis of space frames for static analysis

1. The updated Lagrangian (U.L.) formulation: all variables are referred to the current deformed configuration at time  $t$ .
2. The total Lagrangian (T.L.) formulation: all variables are referred to the initial undeformed configuration at time  $0$ .
3. The convected coordinate formulation: all variables are referred to the new incremented configuration at time  $t+\Delta t$ . This formulation utilizes a set of moving rigid convected coordinates that rotate and translate with the element, but do not deform with the element.

To derive the element deformation displacements, four types of coordinate systems are defined (section 3.2). The element may undergo large translations and large rotations, but is restricted to small strains which means the cross sectional area does not change.

The second purpose is to compare two models of geometrically nonlinear space frames. To predict the structural response accurately, it is necessary to select the proper mathematical models, either finite element model or a beam-column model can be used. The finite element model is formulated by the principle of virtual work with Lagrangian and Hermitian interpolation functions used for discretization. The beam-column model, developed by Oran [1], [2], [3], is based on the conventional beam-column theory.

The solution procedure is iterative as well as incremental. The Newton-Raphson method and the modified Riks/Wempner method [27] are employed. The Newton-Raphson method cannot be used to trace equilibrium path beyond the limit point; for this reason the modified Riks/Wempner method was developed.

The final purpose is to develop a computer program for geometrically nonlinear static analysis of space frames. The U.L. formulation was implemented in this program. Several test examples were investigated using the computer program. The results compared well with available data in the literature.

## 1.2 SURVEY OF LITERATURE

The analysis of geometrically nonlinear framed structures has attracted considerable attention during the last two decades. An early paper by Connor [16] presented a nonlinear formulation for a rigid-jointed space frame with small rotations subjected to loads applied only at the joints. Oran [1], [2], [3] derived a tangent stiffness matrix for elastic frame structures based on the conventional beam-column theory [23] in which small relative deformations of the members were assumed, however the rotations and translations of the joints were considered to be arbitrarily large. Belytschko and coworkers [5],[17] employed a convected coordinate system in which the deformation displacements were separated from the rigid body motion, and node orientations were described by unit vectors that only three components of two unit vectors were stored. Mikkola [4] combined Oran's and Belytschko's formulations in which joint displacements and element deformations were described, and derived the tangent stiffness matrix in the different form.

A large displacement problem in structural analysis can be analyzed in three types of formulations: U.L., T.L. and convected coordinates formulations.

References [25],[26],[39],[41] adopted the U.L. formulation: Murray and Wilson [25],[41] investigated the response of thin elastic plates and employed in-plane displacement functions and plate bending displacement functions which maintained boundary compatibility for the in-plane and bending problems, respectively, but violated boundary compatibility when superimposed in the large deflection problem. Yang [26] applied a linearized midpoint tangent incremental approach to predict the nonlinear equilibrium path. Chu and coworkers [39] developed the constant load method to determine buckling loads of space frames based on the large deflection theory in which the iteration may start at any load level and the stiffness matrix developed for small deflection theory could be used directly.

References [7],[12],[19],[20],[32] adopted the T.L. formulation: Rajasekaran and Murray [7] showed that the equilibrium equation and the linear incremental equilibrium equation did not necessarily follow from the total potential energy in the form introduced by Mallett and Marcal [8], and derived the particular forms of the incremental stiffness matrices. Hibbitt and coworkers [12] developed a large displacement, large strain formulation by introducing an additional initial load stiffness matrix into the large dis-

placement, small strain formulation, and took the material to be elastic-plastic. Remseth [19],[32] presented nonlinear static and dynamic analysis of space frames in which the node rotations were limited to 12 to 15 degree, and higher order axial interpolation polynomials were included in order to obtain an appropriate coupling between axial forces and bending. Wood and Zienkiewicz [20] presented the geometrically nonlinear analysis of the two dimensional inplane structures, e.g. beams, frames and arches, which omitted mid-side nodes in the "thickness" direction, thereby reducing the number of degrees of freedom, and employed a parilinear isoparametric element.

Cook [33] gave an basic introduction to the geometric nonlinear problem. Bathe [9],[10],[38] presented an U.L. and a T.L. formulation, derived from the continuum mechanics, of space frame element for large displacement analysis. Tang [18] described these three formulations and explained that in the convected coordinate formulation an incremental concept was impracticable, thus the tangent stiffness matrix was difficult to be established.

Wood and Schrefler [6] gave a correlation between the so-called N-notation and the B-notation of the T.L. formula-

tion of geometrically nonlinear problems. Mallet [8] and Rajasekaran [7] adopted the N-notation, Zienkiewicz and co-workers [31], [20] adopted the B-notation.

Katzenberger [48] derived the secant stiffness matrices for the plane frame element from which element forces can be obtained. Butler [29] compared two models for geometrically nonlinear finite element analysis of plane frames.

To trace nonlinear equilibrium paths into the postbuckling range, Holzer [27] employed the modified Riks/Wempner method. Bathe and Cimento [49] described the practical procedures for the incremental solution of nonlinear finite element equations and proposed specific ways to measure convergence.

Papadrakakis [21] employed the vector iteration methods to study the post-buckling behavior of spatial structures in which there is no need to compute or formulate the tangent stiffness matrix.

## Chapter II

### UPDATED AND TOTAL LAGRANGIAN FORMULATIONS IN GEOMETRICALLY NONLINEAR FINITE ELEMENT ANALYSIS

#### 2.1 INTRODUCTION

Large displacement analysis may be formulated in three types of coordinate systems (Fig. 2.1):

1. The total Lagrangian (T.L.) formulation which refers to the initial undeformed equilibrium configuration  ${}^0C$  at time 0.
2. The updated Lagrangian (U.L.) formulation which refers to the current deformed equilibrium configuration  ${}^tC$  at time  $t$ .
3. The convected coordinate (Eulerian) formulation which refers to the new incremented configuration  ${}^{t+\Delta t}C$  at time  $t+\Delta t$ .

In Fig. 2.1  ${}^0X_i$ ,  ${}^tX_i$ ,  ${}^{t+\Delta t}X_i$  are the global coordinate systems in the configuration at time 0,  $t$ ,  $t+\Delta t$  respectively,  $i=1, 2, 3$ ;  ${}^0x_i$ ,  ${}^tx_i$ ,  ${}^{t+\Delta t}x_i$  are the local coordinate systems at time 0,  $t$ ,  $t+\Delta t$  respectively;  ${}^tu_i$ ,  ${}^{t+\Delta t}u_i$  are the displacement components from initial position at time 0 to configuration at time  $t$ ,  $t+\Delta t$  respectively.



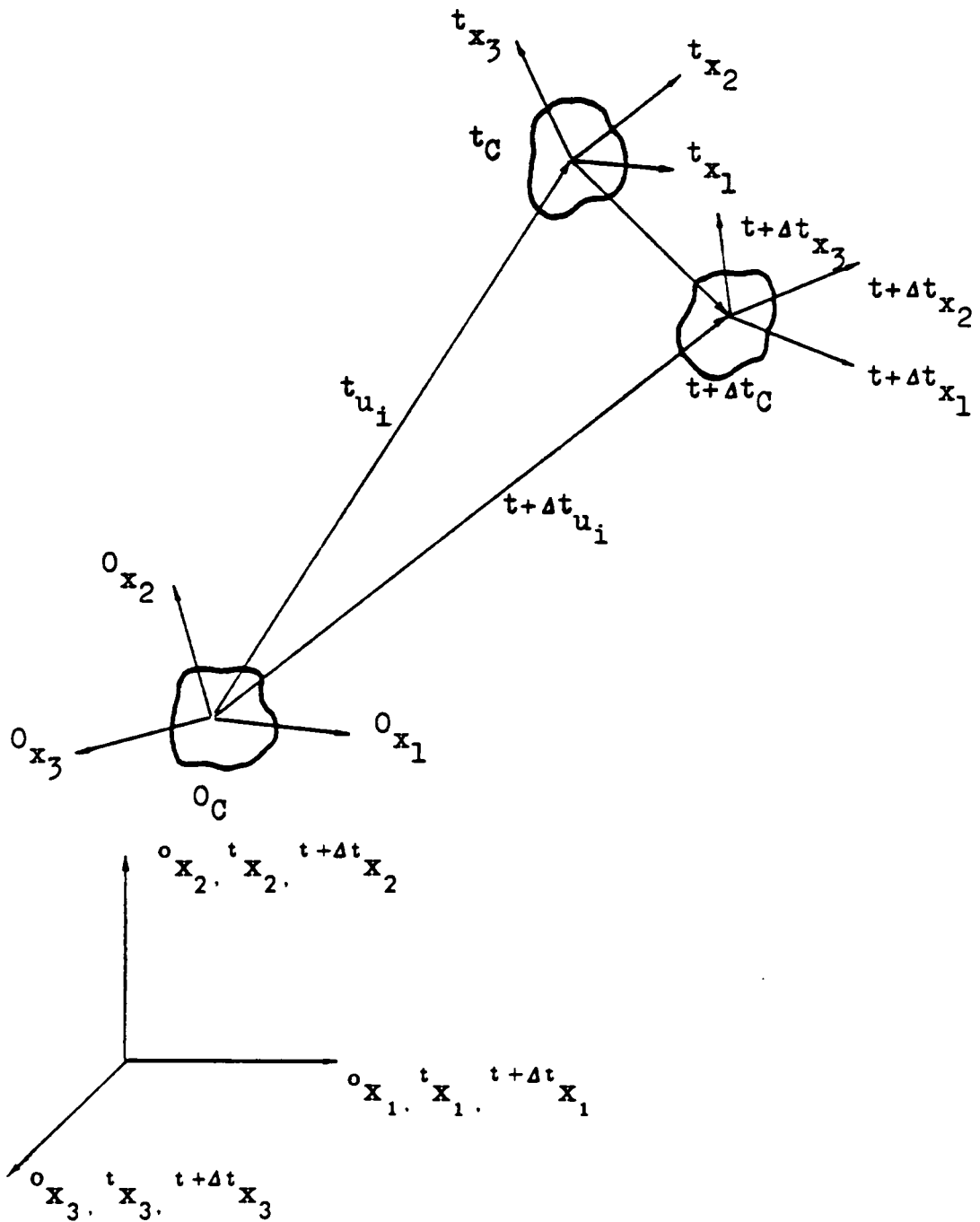


Fig. 2.1 Motion of an Element in Cartesian Coordinate System

Tensor notation is adopted in incremental U.L. and T.L. continuum mechanics formulation. The Green-Lagrange strain tensor used with the 2nd Piola-Kirchhoff stress tensor is defined as

$${}^t_0\varepsilon_{ij} = 1/2({}^t_0u_{i,j} + {}^t_0u_{j,i} + {}^t_0u_{k,i} {}^t_0u_{k,j}) \quad (2.0)$$

in which  ${}^t_0u_{i,j} = \partial {}^t u_i / \partial {}^0 x_j$ . Using the principle of virtual work, we can express the equilibrium equation in the configuration at time  $t+\Delta t$  as [9], [10], [38]

$$\int_{{}^t_0V} {}^{t+\Delta t} {}^t S_{ij} \delta {}^{t+\Delta t} {}^t \varepsilon_{ij} {}^t dV = {}^{t+\Delta t} R \quad \text{in U.L. formulation} \quad (2.1)$$

$$\int_{{}^0_0V} {}^{t+\Delta t} {}^0 S_{ij} \delta {}^{t+\Delta t} {}^0 \varepsilon_{ij} {}^0 dV = {}^{t+\Delta t} R \quad \text{in T.L. formulation} \quad (2.2)$$

where a left superscript indicates in which configuration the quantity occurs; a left subscript indicates the configuration to which the quantity is referred; a right subscript indicates the components of a tensor.

${}^{t+\Delta t} {}^t S_{ij}$ ,  ${}^{t+\Delta t} {}^0 S_{ij}$  = component of 2nd Piola-Kirchhoff stress tensor

$\delta$  = 'variation in '

${}^{t+\Delta t} {}^t \varepsilon_{ij}$ ,  ${}^{t+\Delta t} {}^0 \varepsilon_{ij}$  = component of Green-Lagrange strain tensor

${}^{t+\Delta t} R$  = external virtual work corresponding to

configuration at time  $t+\Delta t$

Since Eqs. (2.1) and (2.2) are nonlinear in the displacement increments  $u_i$ , i.e.  $u_i = {}^{t+\Delta t}u_i - {}^t u_i$ , they cannot be solved directly. Approximate solutions can be obtained by referring all variables to a previously calculated known equilibrium configuration, and linearizing the resulting equation; then the solution can be improved by iteration. Assuming

$${}^t \varepsilon_{ij} = {}^t e_{ij} \quad \text{in U.L. formulation} \quad (2.3)$$

$${}^0 \varepsilon_{ij} = {}^0 e_{ij} \quad \text{in T.L. formulation} \quad (2.4)$$

where

${}^t \varepsilon_{ij}, {}^0 \varepsilon_{ij}$  = component of strain increment tensor (Green-Lagrange) referred to configuration at time  $t$  and  $0$  respectively

${}^t e_{ij}, {}^0 e_{ij}$  = linear part of strain increment  ${}^t \varepsilon_{ij}, {}^0 \varepsilon_{ij}$  respectively

In addition, since in Eqs. (2.1) and (2.2) the 2nd Piola-Kirchhoff stresses and Green-Lagrange strains are unknown, they must be decomposed into

In U.L. formulation:

$${}^{t+\Delta t} S_{ij} = {}^t \tau_{ij} + {}^t S_{ij} \quad (\text{note: } {}^t S_{ij} = {}^t \tau_{ij}) \quad (2.5)$$

$${}^{t+\Delta t}{}^t\varepsilon_{ij} = {}^t\varepsilon_{ij} \quad (2.6)$$

In T.L. formulation:

$${}^{t+\Delta t}{}_0S_{ij} = {}^t{}_0S_{ij} + {}_0S_{ij} \quad (2.7)$$

$${}^{t+\Delta t}{}_0\varepsilon_{ij} = {}^t{}_0\varepsilon_{ij} + {}_0\varepsilon_{ij} \quad (2.8)$$

where

${}^t\tau_{ij} = {}^t{}_tS_{ij}$  = component of Cauchy stress tensor

${}^tS_{ij}$ ,  ${}_0S_{ij}$  = component of 2nd Piola-Kirchhoff stress  
increment at time  $t$

${}^t\varepsilon_{ij}$ ,  ${}_0\varepsilon_{ij}$  = component of strain increment tensor  
(Green-Lagrange)

## 2.2 INCREMENTAL EQUILIBRIUM EQUATION IN U.L. FORMULATION

### 2.2.1 Incremental U.L. Continuum Mechanics Formulation

In the U.L. formulation all variables, i.e. displacement, area, volume, stress, strain, differentiations and integrations, etc., refer to the current deformed configuration  ${}^tC$  at time  $t$  (Fig. 2.1).

Since the Green-Lagrange strain tensor  ${}^{t+\Delta t}{}^t\varepsilon_{ij}$  uses the displacements from the configuration at time  $t$  to the configuration at time  $t+\Delta t$ , from Eq. (2.6) we have

$$\delta_{t^{\varepsilon_{ij}}}^{t+\Delta t} = \delta_{t^{\varepsilon_{ij}}} \quad (2.9)$$

Substituting Eqs. (2.5), (2.9) into Eq. (2.1) yields

$$\int_{t_V} t^{\tau_{ij}} \delta_{t^{\varepsilon_{ij}}} t_{dV} + \int_{t_V} t^S_{ij} \delta_{t^{\varepsilon_{ij}}} t_{dV} = t^{+\Delta t}_R \quad (2.10)$$

The strain increment components can be separated into linear and nonlinear parts

$$t^{\varepsilon_{ij}} = t^e_{ij} + t^n_{ij} \quad (2.11)$$

where

$$t^e_{ij} = 1/2 (t^{u_{i,j}} + t^{u_{j,i}}) \quad (2.12)$$

$$t^n_{ij} = 1/2 t^{u_{k,i}} t^{u_{k,j}} \quad (2.13)$$

$$t^{u_{i,j}} = \partial u_i / \partial x_j \quad (2.14)$$

$$u_i = t^{+\Delta t} u_i - t u_i ; \quad i=1,2,3 \quad (2.15)$$

$t^n_{ij}$  is the nonlinear part of strain increment  $t^{\varepsilon_{ij}}$ ;  $t^{u_{i,j}}$  is the derivative of displacement increment with respect to local coordinate  $t x_j$ ;  $u_i$  are the increments in the displacements from time  $t$  to time  $t+\Delta t$ ;  $t u_i$ ,  $t^{+\Delta t} u_i$  are the displacement components in the local coordinate system from the initial configuration at time 0 to the deformed configuration at time  $t$  and time  $t+\Delta t$  respectively.

$$t_{u_i} = t_{x_i} - 0_{x_i} \quad (2.15a)$$

$$t+\Delta t_{u_i} = t+\Delta t_{x_i} - 0_{x_i} \quad (2.15b)$$

where  $t_{u_i}$  and  $t+\Delta t_{u_i}$  are the functions of position coordinates  $(x_1, x_2, x_3)$  respectively.

The constitutive law is [10]

$$t^{S_{ij}} = t^{C_{ijrs}} t^{\epsilon_{rs}} \quad (2.16)$$

where  $t^{C_{ijrs}}$  is the component of incremental material property tensor at time  $t$  referred to the configuration at time  $t$ .

Using Eqs. (2.11), (2.16), Eq. (2.10) can be transformed to

$$\begin{aligned} & \int_{t_V} t^{C_{ijrs}} t^{\epsilon_{rs}} \delta t^{\epsilon_{ij}} t_{dV} + \int_{t_V} t^{\tau_{ij}} \delta t^{\eta_{ij}} t_{dV} \\ & = t+\Delta t_R - \int_{t_V} t^{\tau_{ij}} \delta t^{\epsilon_{ij}} t_{dV} \end{aligned} \quad (2.17)$$

Eq. (2.17) is nonlinear in the incremental displacements  $u_i$ , and it can be linearized by using the approximations

$$\delta t^{\epsilon_{ij}} = \delta t^{\epsilon_{ij}} \quad (2.18)$$

$$t^{S_{ij}} = t^{C_{ijrs}} t^{\epsilon_{rs}} \quad (2.19)$$

Therefore Eq. (2.17) becomes

$$\begin{aligned} & \int_{t_V} t^{C_{ijrs}} t^{\epsilon_{rs}} \delta t^{\epsilon_{ij}} t_{dV} + \int_{t_V} t^{\tau_{ij}} \delta t^{\eta_{ij}} t_{dV} \\ & = t+\Delta t_R - \int_{t_V} t^{\tau_{ij}} \delta t^{\epsilon_{ij}} t_{dV} \end{aligned} \quad (2.20)$$

Eq. (2.20) is the incremental equilibrium equation of a deformed element, linear in the incremental displacements  $u_i$ , corresponding to the local coordinate system.

In three dimensional beam element for our problem, small deformation and uniaxial state of strain (i.e.  $t^{\epsilon_{11}}$  only) are assumed, and torsion is treated independently from bending and axial force so that it can be obtained from linear theory; in this situation Eq. (2.20) can be specialized as

$$\begin{aligned} E \int_{t_V} t^{\epsilon_{11}} \delta_t \epsilon_{11} t^{\epsilon_{11}} t^{\epsilon_{11}} dV + \int_{t_V} t^{\sigma} \delta_t \eta_{11} t^{\sigma} dV \\ = t^{\Delta t_R} - \int_{t_V} t^{\sigma} \delta_t \epsilon_{11} t^{\sigma} dV \end{aligned} \quad (2.21)$$

where

$$E = t^C_{1111} \quad (2.22)$$

$$t^{\sigma} = t^{\tau}_{11} \quad (2.23)$$

$E$  is the Young's modulus and  $t^{\sigma}$  is the axial Cauchy stress.

Eq. (2.21) can be expressed in matrix form. By using interpolation functions for incremental displacements to evaluate the derivatives of displacements, we will obtain the

linear and nonlinear strain-displacement transformation matrices. For a beam element it is more effective to first evaluate the finite element matrices in the local coordinate axes  $x_i$ , and then transform them to the global coordinate axes  $X_i$  prior to the element assemblage process.

### 2.2.2 Incremental Strain

For our uniaxial strain problem, the  $u_i$  in Eq. (2.15) means the displacement increments along the centroid axis of elements and is only the function of  $x_1$ -axis. By using interpolation function,  $u_i$  can be expressed in the nodal displacement increments.

$$\begin{bmatrix} {}^t u_1 \\ {}^t u_2 \\ {}^t u_3 \end{bmatrix} = \begin{bmatrix} {}^t h_1 & 0 & 0 & 0 & 0 & 0 & | & {}^t h_7 & 0 & 0 & 0 & 0 & 0 \\ 0 & {}^t h_2 & 0 & 0 & 0 & {}^t h_6 & | & 0 & {}^t h_8 & 0 & 0 & 0 & {}^t h_{12} \\ 0 & 0 & {}^t h_3 & 0 & {}^t h_5 & 0 & | & 0 & 0 & {}^t h_9 & 0 & {}^t h_{11} & 0 \end{bmatrix} \begin{bmatrix} \Delta^t d_1 \\ \Delta^t d_2 \\ \cdot \\ \cdot \\ \cdot \\ \Delta^t d_{11} \\ \Delta^t d_{12} \end{bmatrix}$$

(2.24)

where

${}^t u_i$  = increment in displacement component of element from



$t$  to  $t+\Delta t$  measured in the local axes  ${}^t x_i$ ;  $i = 1, 2, 3$   
 ${}^t h_k$  = finite element interpolation function corresponding  
to  $\Delta^t d_k$ ,  $k = 1$  to  $12$   
 $\Delta^t d_k$  = increment in nodal displacement component of  
element from  $t$  to  $t+\Delta t$  measured in the local  
axes  ${}^t x_i$

In Eq. (2.24) we use the Hermitian interpolation functions to describe bending deformations and linear interpolation functions to describe axial and torsional displacements; however, the torsional end displacements do not effect the local element displacements  $u_1, u_2, u_3$  of the centroid axis.

Assuming the cross section of element do not change during deformation for small strain analysis, thus the distances from the centroid axis in the local  $x_2, x_3$ -axes directions respectively, say  $y, z$ , are constant. The term  ${}^t u_{1,1}$  in Eq. (2.13) is always small compared to unity, and the square of  ${}^t u_{1,1}$  is negligible in comparison with  ${}^t u_{1,1}$ . Therefore, from Eq. (2.11) the incremental uniaxial strain along the beam element for small deformation is

$${}^t \varepsilon_{11} = {}^t \varepsilon_{11}^e + {}^t \eta_{11} \quad (2.25)$$

where

$$t^{e11} = \underbrace{t^{u_{1,1}}}_{\text{due to axial force}} - \underbrace{y t^{u_{2,11}} - z t^{u_{3,11}}}_{\text{due to bending}} \quad (2.26)$$

$$t^{n11} = \frac{1}{2} (t^{u_{2,1}})^2 + \frac{1}{2} (t^{u_{3,1}})^2 \quad (2.27)$$

$$t^{u_{i,jj}} = \partial^2 u_i / \partial x_j^2 \quad (2.28)$$

### 2.2.3 Incremental Equilibrium Equation

Eqs. (2.25), (2.26), (2.27) can be expressed in matrix form introduced by Wood and Schrefler [6]

$$t^\varepsilon = t^e + t^n \quad (2.29)$$

where

$$t^e = t^{L^T} t^\theta \quad (2.30)$$

$$t^n = 1/2 t^{\theta^T} H t^\theta \quad (2.31)$$

$t^{L^T}$  is the row vector defining linear strains  $t^e$  from displacement gradients given by

$$t^{L^T} = [ 1 \quad 0 \quad 0 \quad -y \quad -z ] \quad (2.32)$$

$t^\theta$  is the column vector of displacement gradient contributing to the strain  $t^\varepsilon$  given by

$$t^\theta = [ t^{u_{1,1}} \quad t^{u_{2,1}} \quad t^{u_{3,1}} \quad t^{u_{2,11}} \quad t^{u_{3,11}} ]^T \quad (2.33)$$

H is the symmetric matrix containing arrangements of unity and zero given by

$$H = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix} \quad (2.34)$$

Employing the finite element discretization of Eq. (2.24) into Eq. (2.33),  $\theta$  can be expressed in terms of the nodal incremental displacements as

$$t^\theta = t_{NL}^B \Delta^t d \quad (2.35)$$

where

$$t_{NL}^B = \begin{pmatrix} t^{h_{1,1}} & 0 & 0 & 0 & 0 & 0 & t^{h_{7,1}} & 0 & 0 & 0 & 0 & 0 \\ 0 & t^{h_{2,1}} & 0 & 0 & 0 & t^{h_{6,1}} & 0 & t^{h_{8,1}} & 0 & 0 & 0 & t^{h_{12,1}} \\ 0 & 0 & t^{h_{3,1}} & 0 & t^{h_{5,1}} & 0 & 0 & 0 & t^{h_{9,1}} & 0 & t^{h_{11,1}} & 0 \\ 0 & t^{h_{2,11}} & 0 & 0 & 0 & t^{h_{6,11}} & 0 & t^{h_{8,11}} & 0 & 0 & 0 & t^{h_{12,11}} \\ 0 & 0 & t^{h_{3,11}} & 0 & t^{h_{5,11}} & 0 & 0 & 0 & t^{h_{9,11}} & 0 & t^{h_{11,11}} & 0 \end{pmatrix} \quad (2.36)$$

in which

$${}^t h_{i,j} = \partial {}^t h_i / \partial {}^t x_j \quad (2.37)$$

$${}^t h_{i,jj} = \partial^2 {}^t h_i / \partial {}^t x_j^2 \quad (2.38)$$

Eq. (2.21) in matrix form is

$$\begin{aligned} E \int_{t_V} \delta_t e^T {}^t e {}^t dV + \int_{t_V} \delta_t \eta^T {}^t \sigma {}^t dV \\ = {}^{t+\Delta t} R - \int_{t_V} \delta_t e^T {}^t \sigma {}^t dV \end{aligned} \quad (2.39)$$

Taking the variation of Eq. (2.30) and from Eq. (2.35) we have

$$\delta_t e = {}^t B_L \delta \Delta {}^t d \quad \text{where } {}^t B_L = {}^t L^T {}^t B_{NL} \quad (2.40)$$

Taking the variation of Eq. (2.31) and from Eq. (2.35) we have

$$\begin{aligned} \delta_t \eta &= 1/2 ( \delta_t \theta^T H {}^t \theta + {}^t \theta^T H \delta_t \theta ) \\ &= {}^t \theta^T H \delta_t \theta \\ &= {}^t \theta^T H {}^t B_{NL} \delta \Delta {}^t d \end{aligned} \quad (2.41)$$

Substituting Eqs. (2.30), (2.35), (2.40), (2.41) into Eq. (2.39) and eliminating  $\delta \Delta {}^t d^T$  on both sides of equation, we have

$$\begin{aligned}
& \int_{t_V} t_{B_L}^T E t_{B_L}^T \Delta^t_d t_{dV} + \int_{t_V} t_{B_{NL}}^T t_{\sigma} H t_{B_{NL}}^T \Delta^t_d t_{dV} \\
& = t^{+\Delta t}_r - \int_{t_V} t_{B_L}^T t_{\sigma} t_{dV}
\end{aligned} \tag{2.42}$$

Where  $t^{+\Delta t}_r$  is the vector of externally applied element nodal loads at time  $t+\Delta t$  in the local coordinate system.

Let

$$t_{\tau} = t_{\sigma} H = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & t_{\sigma} & 0 & 0 & 0 \\ 0 & 0 & t_{\sigma} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix} \tag{2.43}$$

which is a symmetric Cauchy stress matrix in configuration at time  $t$ .

Substituting Eq. (2.43) into Eq. (2.42) we obtain the incremental equilibrium equation

$$\begin{aligned}
& (t_{B_L}^k + t_{B_{NL}}^k) \Delta^t_d{}^k \\
& = t^{+\Delta t}_r - t^{+\Delta t}_f{}^{k-1}
\end{aligned} \tag{2.44}$$

where

$$t_{B_L}^k = \int_{t_V} t_{B_L}^T E t_{B_L}^T t_{dV} \tag{2.45}$$

$$t_{B_{NL}}^k = \int_{t_V} t_{B_{NL}}^T t_{\tau} t_{B_{NL}}^T t_{dV} \tag{2.46}$$

$$\Delta^t_d{}^k = t^{+\Delta t}_d{}^k - t^t_d{}^{k-1} \tag{2.47}$$

$$\frac{t+\Delta t}{t+\Delta t} f^{k-1} = \int_{t_V} t_{B_L}^T t_{\sigma} t_{dV} \quad (2.48)$$

$$\begin{aligned} t_{B_L} &= t^L{}^T t_{B_{NL}} \\ &= \begin{bmatrix} t^{h_{1,1}} & -y_t^{h_{2,11}} & -z_t^{h_{3,11}} & 0 & \\ & -z_t^{h_{5,11}} & -y_t^{h_{6,11}} & t^{h_{7,1}} & -y_t^{h_{8,11}} \\ & -z_t^{h_{9,11}} & 0 & -z_t^{h_{11,11}} & -y_t^{h_{12,11}} \end{bmatrix} \end{aligned} \quad (2.49)$$

$t_{B_L}^k$  is the linear strain incremental stiffness matrix,  $t_{B_{NL}}^k$  is the nonlinear strain incremental (geometric or initial stress) stiffness matrix,  $\Delta^t d^k$  is the vector of incremental nodal displacements in iteration  $k$ ,  $t^{+\Delta t} r$  is the vector of externally applied element nodal loads (given new load level) at time  $t+\Delta t$ ,  $\frac{t+\Delta t}{t+\Delta t} f^{k-1}$  is the vector of nodal equivalent element forces corresponding to  $\frac{t+\Delta t}{t+\Delta t} d^{k-1}$ ,  $t_{B_L}$  is the linear strain-displacement transformation matrix,  $t_{B_{NL}}$  is the nonlinear strain-displacement transformation matrix defined by Eq. (2.36),  $t_{\sigma}$  is the Cauchy stress matrix in configuration at time  $t$  defined by Eq. (2.43).

For simplicity of presentation, Fig. 2.2 shows the local element incremental displacements in U. L. formulation for a

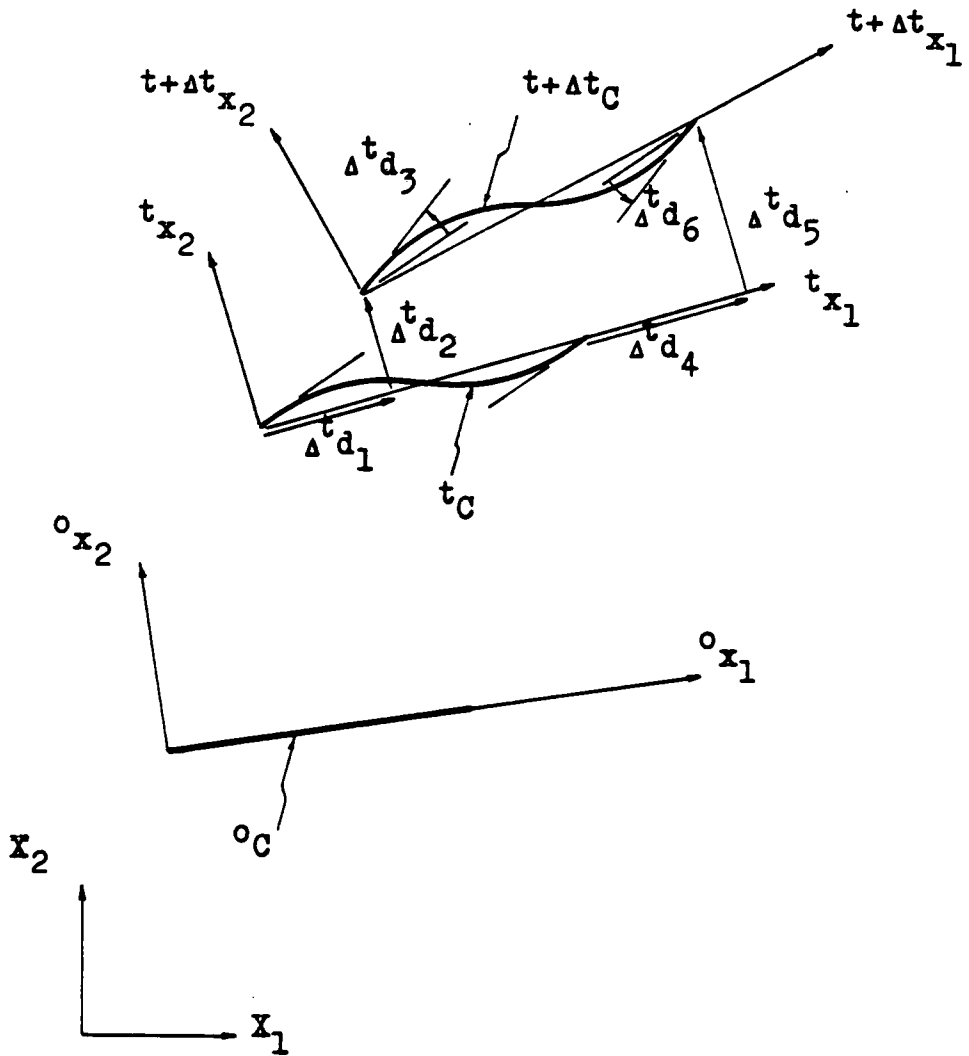


Fig. 2.2 Local Element Displacements in U.L.  
Formulation for Plane Frame

plane frame, although the theoretical development is carried out for a space frame.

The finite element matrices in Eq. (2.44) are transformed to the global coordinate system for element  $i$

$$\begin{matrix} t_{K^i} \\ t^{\Lambda} \end{matrix} = \begin{matrix} t^{\Lambda T} & t_k & t_{\Lambda} \\ t^{\Lambda} & t_k & t_{\Lambda} \end{matrix} \quad (2.50)$$

$$\begin{matrix} t_{F^i} \\ t^{\Lambda} \end{matrix} = \begin{matrix} t^{\Lambda T} & t_f \\ t^{\Lambda} & t_f \end{matrix} \quad (2.51)$$

$$\Delta^t_D = \begin{matrix} t^{\Lambda T} \\ t^{\Lambda} \end{matrix} \Delta^t_d \quad (\text{see sec. 2.2.4}) \quad (2.52)$$

in which

$$\begin{matrix} t_k \\ t^{\Lambda} \end{matrix} = \begin{matrix} t_{k_L} \\ t_{k_{NL}} \end{matrix} + \begin{matrix} t_{k_{NL}} \\ t_{k_{NL}} \end{matrix} \quad (2.53)$$

$$\begin{matrix} t_{\Lambda} \\ t^{\Lambda} \end{matrix} = \left[ \begin{array}{cc|cc} t_{\lambda} & 0 & 0 & 0 \\ 0 & t_{\lambda} & 0 & 0 \\ \hline 0 & 0 & t_{\lambda} & 0 \\ 0 & 0 & 0 & t_{\lambda} \end{array} \right] \quad (2.54)$$

and  $\begin{matrix} t_{\lambda} \\ t^{\lambda} \end{matrix}$  is the element orthogonal orientation matrix in configuration  ${}^t_C$ , same as in Eq. (3.17).



By employing the member code technique [28] (i.e. direct stiffness procedure), the incremental equilibrium equation in U.L. formulation of the whole structure is

$${}^t_K \Delta q^k = {}^{t+\Delta t}_Q - \frac{{}^{t+\Delta t}_F^{k-1}}{{}^{t+\Delta t}_F} \quad (2.55)$$

where

${}^t_K$  = structural strain incremental stiffness matrix  
corresponding to  $q^{k-1}$

$\Delta q^k$  = vector of incremental nodal displacements at  $k^{\text{th}}$   
iteration in configuration at time  $t$

${}^{t+\Delta t}_Q$  = vector of given new load level in configuration  
configuration at time  $t+\Delta t$

$\frac{{}^{t+\Delta t}_F^{k-1}}{{}^{t+\Delta t}_F}$  = vector of nodal equivalent element  
forces corresponding to  $q^{k-1}$

$q^{k-1}$  = vector of nodal displacements at  $k-1$  iteration

Here the structural tangent stiffness matrix  $K$ , is a function of displacements  $q$ , since the problem is nonlinear. In Eq. (2.55) the response of a nonlinear structure may be approximated for incremental nodal displacements by a linear relationship. Because the final configuration is based on equilibrium balance between the nodal element forces and the applied nodal load, the stiffness used to solve Eq. (2.55) need not be exact [41].

#### 2.2.4 Transformation Matrix

The global element displacements are transformed to the local element displacements by the transformation (Fig. 2.2 to Fig. 2.4)

$$\begin{matrix} t_d \\ t \end{matrix} = \begin{matrix} t \\ t^\Lambda \end{matrix} \begin{matrix} t_D \\ t \end{matrix} \quad (2.56)$$

Similarly, the local element forces are transformed to the global element forces by equation

$$\begin{matrix} t_{F^i} \\ t \end{matrix} = \begin{matrix} t \\ t^\Lambda \end{matrix} \begin{matrix} t_f \\ t \end{matrix} \quad (2.57)$$

where  $t^\Lambda$  is transformation matrix defined by Eq. (2.54).

Because all variables refer to the current deformed configuration  $t_C$  in the U.L. formulation, the transformation matrix  $t^\Lambda$  in Eq. (2.54) remains constant from configuration  $t_C$  to  $t+\Delta t_C$ . Eqs. (2.56) and (2.57) yield the increments of local element displacements and global element forces

$$\Delta \begin{matrix} t_d \\ t \end{matrix} = \begin{matrix} t \\ t^\Lambda \end{matrix} \Delta \begin{matrix} t_D \\ t \end{matrix} \quad (2.58)$$

$$\Delta \begin{matrix} t_{F^i} \\ t \end{matrix} = \begin{matrix} t \\ t^\Lambda \end{matrix} \Delta \begin{matrix} t_f \\ t \end{matrix} \quad (2.59)$$

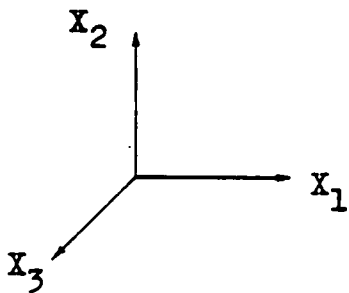
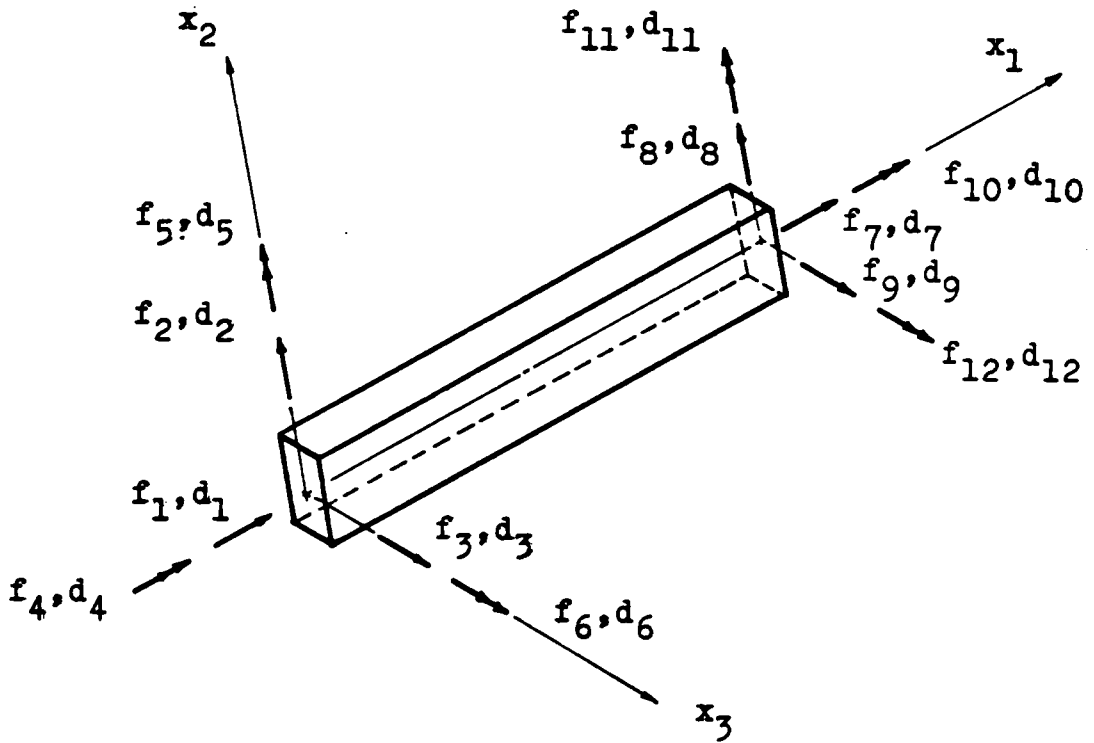


Fig. 2.3 Local Element Forces and Displacements in Space Frame

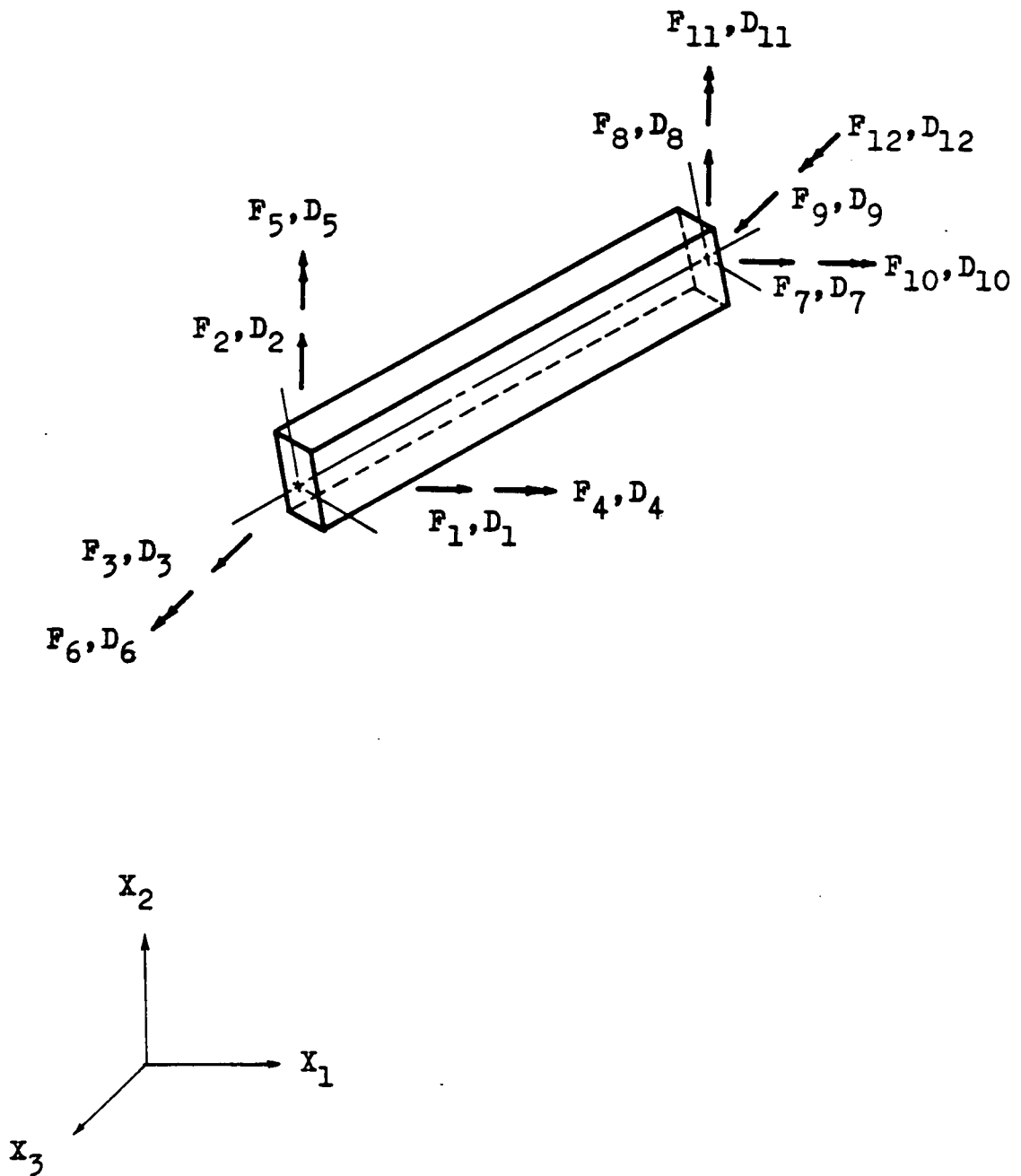


Fig. 2.4 Global Element Forces and Displacements  
in Space Frame

where  $\Delta^t f$  and  $\Delta^{tF^i}$  are the vectors of the local and global incremental element forces from time  $t$  to  $t+\Delta t$  respectively, referred to configuration at time  $t$ .

## 2.3 INCREMENTAL EQUILIBRIUM EQUATION IN T.L. FORMULATION

### 2.3.1 Incremental T.L. Continuum Mechanics Formulation

The T.L. formulation is based on the same procedures that are used in the U.L. formulation, but all variables refer to the initial undeformed configuration  ${}^0C$  at time 0 (Fig. 2.1).

Taking the variation of Eq. (2.8), we have

$$\delta {}^{t+\Delta t} {}_0 \varepsilon_{ij} = \delta {}_0 \varepsilon_{ij} \quad (2.60)$$

Substituting Eqs. (2.7), (2.60) into Eq. (2.2)

$$\int_{{}_0V} {}^t S_{ij} \delta {}_0 \varepsilon_{ij} {}^0 dV + \int_{{}_0V} {}_0 S_{ij} \delta {}_0 \varepsilon_{ij} {}^0 dV = {}^{t+\Delta t} R \quad (2.61)$$

The strain increment components can be separated into linear and nonlinear parts

$${}_0 \varepsilon_{ij} = {}_0 e_{ij} + {}_0 n_{ij} \quad (2.62)$$

where

$${}_0 e_{ij} = 1/2 ({}_0 u_{i,j} + {}_0 u_{j,i} + {}^t u_{k,i} {}_0 u_{k,j} + {}_0 u_{k,i} {}^t u_{k,j}) \quad (2.63)$$

$${}^0n_{ij} = 1/2 {}^0u_{k,i} {}^0u_{k,j} \quad (2.64)$$

$${}^0u_{i,j} = \partial u_i / \partial {}^0x_j \quad (2.65)$$

$${}^t u_{i,j} = \partial {}^t u_i / \partial {}^0x_j \quad (2.66)$$

${}^0n_{ij}$  is the nonlinear part of strain increment  ${}^0\varepsilon_{ij}$ ,  ${}^0u_{i,j}$  is the derivative of displacement increment with respect to local coordinate  ${}^0x_j$ ,  ${}^t u_{i,j}$  is the derivative of displacement component in configuration at time  $t$  with respect to local coordinate  ${}^0x_j$ .

The constitutive law is [10]

$${}^0S_{ij} = {}^0C_{ijrs} {}^0\varepsilon_{rs} \quad (2.67)$$

where  ${}^0C_{ijrs}$  is the component of incremental material property tensor at time 0 referred to the configuration at time 0.

Using Eqs. (2.62), (2.67), Eq. (2.61) can be transformed to

$$\begin{aligned} \int_{0V} {}^0C_{ijrs} {}^0\varepsilon_{rs} \delta {}^0\varepsilon_{ij} {}^0dv + \int_{0V} {}^t S_{ij} \delta {}^0n_{ij} {}^0dv \\ = {}^{t+\Delta t} R - \int_{0V} {}^t S_{ij} \delta {}^0e_{ij} {}^0dv \end{aligned} \quad (2.68)$$

Eq. (2.68) is nonlinear in the incremental displacements  $u_i$ , and it can be linearized by using the approximations

$$\delta {}^0\varepsilon_{ij} = \delta {}^0e_{ij} \quad (2.69)$$

$${}^0S_{ij} = {}^0C_{ijrs} {}^0e_{rs} \quad (2.70)$$

Then Eq. (2.68) becomes

$$\begin{aligned} \int_{0_V} {}^0C_{ijrs} {}^0e_{rs} \delta {}^0e_{ij} {}^0dV + \int_{0_V} {}^tS_{ij} \delta {}^0n_{ij} {}^0dV \\ = {}^{t+\Delta t}R - \int_{0_V} {}^tS_{ij} \delta {}^0e_{ij} {}^0dV \end{aligned} \quad (2.71)$$

which is linear equation in the incremental displacements  $u_i$ , corresponding to the local coordinate system.

For three dimensional beam element with small deformation and uniaxial state of strain (i.e.  ${}^0\varepsilon_{11}$  only), in which torsion is treated independently, Eq. (2.71) can be specialized as

$$\begin{aligned} E \int_{0_V} {}^0e_{11} \delta {}^0e_{11} {}^0dV + \int_{0_V} {}^t\sigma \delta {}^0n_{11} {}^0dV \\ = {}^{t+\Delta t}R - \int_{0_V} {}^t\sigma \delta {}^0e_{11} {}^0dV \end{aligned} \quad (2.72)$$

where

$${}^t\sigma = {}^tS_{11} \quad (2.73)$$

is the 2nd Piola-Kirchhoff stress.

### 2.3.2 Incremental Strain

The incremental displacement components along the centroid axis of elements are interpolated as

$$\begin{pmatrix} {}^0u_1 \\ {}^0u_2 \\ {}^0u_3 \end{pmatrix} = \begin{pmatrix} {}^0h_1 & 0 & 0 & 0 & 0 & 0 & | & {}^0h_7 & 0 & 0 & 0 & 0 & 0 \\ 0 & {}^0h_2 & 0 & 0 & 0 & {}^0h_6 & | & 0 & {}^0h_8 & 0 & 0 & 0 & {}^0h_{12} \\ 0 & 0 & {}^0h_3 & 0 & {}^0h_5 & 0 & | & 0 & 0 & {}^0h_9 & 0 & {}^0h_{11} & 0 \end{pmatrix} \begin{pmatrix} \Delta^0d_1 \\ \Delta^0d_2 \\ \cdot \\ \cdot \\ \cdot \\ \Delta^0d_{11} \\ \Delta^0d_{12} \end{pmatrix} \quad (2.74)$$

where

${}^0u_i$  = increment in displacement component of element from  $t$  to  $t+\Delta t$  measured in the local axes  ${}^0x_i$ ;  $i=1,2,3$

${}^0h_k$  = finite element interpolation function corresponding to  $\Delta^0d_k$ ,  $k = 1$  to  $12$

$\Delta^0d_k$  = increment in nodal displacement component of element from  $t$  to  $t+\Delta t$  measured in the local axes  ${}^0x_i$

Neglecting the 2nd order terms of  ${}^t u_{1,1}$ ,  ${}^0 u_{1,1}$ , from Eq. (2.62) the incremental uniaxial strain along the beam element for small deformation is

$${}^0\varepsilon_{11} = {}^0e_{11} + {}^0n_{11} \quad (2.75)$$



where

$$\begin{aligned}
 {}_0^e{}_{11} = & \underbrace{{}_0^u{}_{1,1}}_{\text{due to axial force}} + \underbrace{{}_0^t{}_{u2,1} {}_0^u{}_{2,1} + {}_0^t{}_{u3,1} {}_0^u{}_{3,1}}_{\text{due to initial displacement}} \\
 & \underbrace{- y {}_0^u{}_{2,11} - z {}_0^u{}_{3,11}}_{\text{due to bending}}
 \end{aligned} \tag{2.76}$$

$${}_0^n{}_{11} = 1/2 ({}_0^u{}_{2,1})^2 + 1/2 ({}_0^u{}_{3,1})^2 \tag{2.77}$$

and

$${}_0^t{}_{u_i,j} = \partial^t u_i / \partial^0 x_j \tag{2.78}$$

$${}_0^u{}_{i,jj} = \partial^2 u_i / \partial^0 x_j^2 \tag{2.79}$$

### 2.3.3 Incremental Equilibrium Equation

Expressing Eqs. (2.75), (2.76), (2.77) in matrix form introduced by Wood and Schrefler [6]

$${}_0^\varepsilon = {}_0^e + {}_0^n \tag{2.80}$$

where

$${}_0^e = {}_0^L{}^T {}_0^\theta \tag{2.81}$$

$${}_0^n = 1/2 {}_0^\theta{}^T H {}_0^\theta \tag{2.82}$$

$${}_0L^T = [ 1 \quad {}_0^t u_{2,1} \quad {}_0^t u_{3,1} \quad -y \quad -z ] \quad (2.83)$$

$${}_0^\theta = [ {}_0^u u_{1,1} \quad {}_0^u u_{2,1} \quad {}_0^u u_{3,1} \quad {}_0^u u_{2,11} \quad {}_0^u u_{3,11} ]^T \quad (2.84)$$

${}_0L^T$  is the row vector defining linear strains  ${}_0e$  from displacement gradients,  ${}_0^\theta$  is the column vector of displacement gradient contributing to the strain  ${}_0\varepsilon$ ,  $H$  is defined as in Eq. (2.34).

Employing Eq. (2.74) into Eq. (2.84), we express  $\theta$  in terms of the nodal incremental displacements as

$${}_0^\theta = {}_0^t B_{NL} \Delta^0 d \quad (2.85)$$

where

$${}_0^t B_{NL} = \begin{pmatrix} {}_0^h h_{1,1} & 0 & 0 & 0 & 0 & 0 & {}_0^h h_{7,1} & 0 & 0 & 0 & 0 & 0 \\ 0 & {}_0^h h_{2,1} & 0 & 0 & 0 & {}_0^h h_{6,1} & 0 & {}_0^h h_{8,1} & 0 & 0 & 0 & {}_0^h h_{12,1} \\ 0 & 0 & {}_0^h h_{3,1} & 0 & {}_0^h h_{5,1} & 0 & 0 & 0 & {}_0^h h_{9,1} & 0 & {}_0^h h_{11,1} & 0 \\ 0 & {}_0^h h_{2,11} & 0 & 0 & 0 & {}_0^h h_{6,11} & 0 & {}_0^h h_{8,11} & 0 & 0 & 0 & {}_0^h h_{12,11} \\ 0 & 0 & {}_0^h h_{3,11} & 0 & {}_0^h h_{5,11} & 0 & 0 & 0 & {}_0^h h_{9,11} & 0 & {}_0^h h_{11,11} & 0 \end{pmatrix} \quad (2.86)$$

in which

$${}_0^h h_{i,j} = \partial_0^h h_i / \partial_0^x x_j \quad (2.87)$$

$${}^0h_{i,jj} = \partial_0^2 h_i / \partial_0^2 x_j^2 \quad (2.88)$$

Eq. (2.72) in matrix form is

$$\varepsilon \int_{\text{ov}} \delta_0 e^T \partial_0 e \, \text{ov} + \int_{\text{ov}} \delta_0 n^T t_{\sigma} \, \text{ov} = t^{\Delta} t_R - \int_{\text{ov}} \delta_0 e^T t_{\sigma} \, \text{ov} \quad (2.89)$$

Taking the variation of Eq. (2.81) and using Eq. (2.85) we obtain

$$\delta_0 e = t_{B_L} \delta \Delta^0 d \quad (2.90)$$

where

$$t_{B_L} = {}_0L^T t_{B_{NL}} = t_{B_{LO}} + t_{B_{L1}} \quad (2.91a)$$

and

$$t_{B_{LO}} = \begin{bmatrix} {}_0h_{1,1} & -y \, {}_0h_{2,11} & -z \, {}_0h_{3,11} & 0 & -z \, {}_0h_{5,11} & -y \, {}_0h_{6,11} \\ \vdots & {}_0h_{7,1} & -y \, {}_0h_{8,11} & -z \, {}_0h_{9,11} & 0 & -z \, {}_0h_{11,11} & -y \, {}_0h_{12,11} \end{bmatrix} \quad (2.91b)$$

$$t_{B_{L1}} = \begin{bmatrix} 0 & t_{u_{2,1}} \, {}_0h_{2,1} & t_{u_{3,1}} \, {}_0h_{3,1} & 0 & t_{u_{3,1}} \, {}_0h_{5,1} & t_{u_{2,1}} \, {}_0h_{6,1} \\ \vdots & 0 & t_{u_{2,1}} \, {}_0h_{8,1} & t_{u_{3,1}} \, {}_0h_{9,1} & 0 & t_{u_{3,1}} \, {}_0h_{11,1} & t_{u_{2,1}} \, {}_0h_{12,1} \end{bmatrix} \quad (2.91c)$$

Taking the variation of Eq. (2.82) and from Eq. (2.85) we have

$$\begin{aligned}
\delta_{0^{\theta}} \eta &= \frac{1}{2} (\delta_{0^{\theta}}^T H_{0^{\theta}} + {}_{0^{\theta}}^T H \delta_{0^{\theta}}) \\
&= {}_{0^{\theta}}^T H \delta_{0^{\theta}} \\
&= {}_{0^{\theta}}^T H {}_{0^{\text{B}_{\text{NL}}}}^t \delta \Delta^{\text{od}}
\end{aligned} \tag{2.92}$$

Substituting Eqs. (2.81), (2.85), (2.90), (2.92) into Eq. (2.89), we have

$$\int_{0V} {}_{0^{\text{B}_{\text{L}}}}^t \text{E} {}_{0^{\text{B}_{\text{L}}}}^t \Delta^{\text{od}} \text{odV} + \int_{0V} {}_{0^{\text{B}_{\text{NL}}}}^t \text{H} {}_{0^{\text{B}_{\text{NL}}}}^t \Delta^{\text{od}} \text{odV} = {}^{t+\Delta t} t_r - \int_{0V} {}_{0^{\text{B}_{\text{L}}}}^t \text{t}_{0^{\sigma}} \text{odV} \tag{2.93}$$

Let

$${}_{0^{\text{S}}}^t \text{t}_{0^{\sigma}} \text{H} = \begin{pmatrix} 0 & 0 & 0 & 0 & 0 \\ 0 & t_{0^{\sigma}} & 0 & 0 & 0 \\ 0 & 0 & t_{0^{\sigma}} & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 \end{pmatrix} \tag{2.94}$$

which is the 2nd Piola-Kirchhoff stress matrix.

Substituting Eq. (2.94) into Eq. (2.93), we obtain the incremental equilibrium equation

$$({}_{0^{\text{k}_{\text{L}}}}^t + {}_{0^{\text{k}_{\text{NL}}}}^t) \Delta^{\text{O}_d^{\text{k}}} = {}^{t+\Delta t} t_r - {}^{t+\Delta t} t_{0^{\text{f}^{\text{k}-1}}} \tag{2.95}$$

where

$${}^t_{0K}{}^k_L = \int_{0_V} {}^t_{0B}{}^T E {}^t_{0B}{}^k_L {}^0 dV \quad (2.96)$$

$${}^t_{0K}{}^k_{NL} = \int_{0_V} {}^t_{0B}{}^T {}^t_{0S} {}^t_{0B}{}^k_{NL} {}^0 dV \quad (2.97)$$

$\Delta^0_d{}^k$  = vector of incremental nodal displacement at time  $t$  in iteration  $k$ ; i.e.

$${}^{t+\Delta t}{}_{0d}{}^k = {}^{t+\Delta t}{}_{0d}{}^{k-1} + \Delta^0_d{}^k \quad (2.98)$$

${}^{t+\Delta t}{}_r$  = vector of externally applied element nodal loads at time  $t+\Delta t$

$${}^{t+\Delta t}{}_{0f}{}^{k-1} = \int_{0_V} {}^t_{0B}{}^T {}^t_{0\sigma} {}^0 dV \quad (2.99)$$

For the simplicity of presentation, Fig. 2.5 shows the local element displacements in T.L. formulation for a plane frame.

The finite element matrices in Eq. (2.95) are transformed to the global coordinate system for element  $i$

$${}^t_{0K}{}^i = {}^t_{0\Lambda}{}^T {}^t_{0K} {}^t_{0\Lambda} \quad (2.100)$$

$${}^t_{0F}{}^i = {}^t_{0\Lambda}{}^T {}^t_{0f} \quad (2.101)$$

$$\Delta^0_D = {}^t_{0\Lambda}{}^T \Delta^0_d \quad (\text{see sec. 2.3.4}) \quad (2.102)$$

in which

$${}^t_{0K}{}^i = {}^t_{0K}{}^L + {}^t_{0K}{}^{NL} \quad (2.103)$$

$t_{0\Lambda}$  = transformation matrix (see sec. 2.3.4)

Employing the member code technique [28], the incremental equilibrium equation in T.L. formulation of the whole structure is

$$t_{0K} \Delta q^k = t_{0Q}^{t+\Delta t} - t_{0F}^{t+\Delta t, k-1} \quad (2.104)$$

#### 2.3.4 Transformation Matrix

Similar to Eqs. (2.56) to (2.59), we have (Fig. 2.3 to Fig. 2.5)

$$t_{0d} = t_{0\Lambda} t_{0D} \quad (2.105)$$

$$t_{0F}^i = t_{0\Lambda}^T t_{0f} \quad (2.106)$$

where

$$t_{0\Lambda} = \begin{pmatrix} t_{0\lambda} & 0 & 0 & 0 \\ 0 & t_{0\lambda} & 0 & 0 \\ 0 & 0 & t_{0\lambda} & 0 \\ 0 & 0 & 0 & t_{0\lambda} \end{pmatrix} \quad (2.107)$$

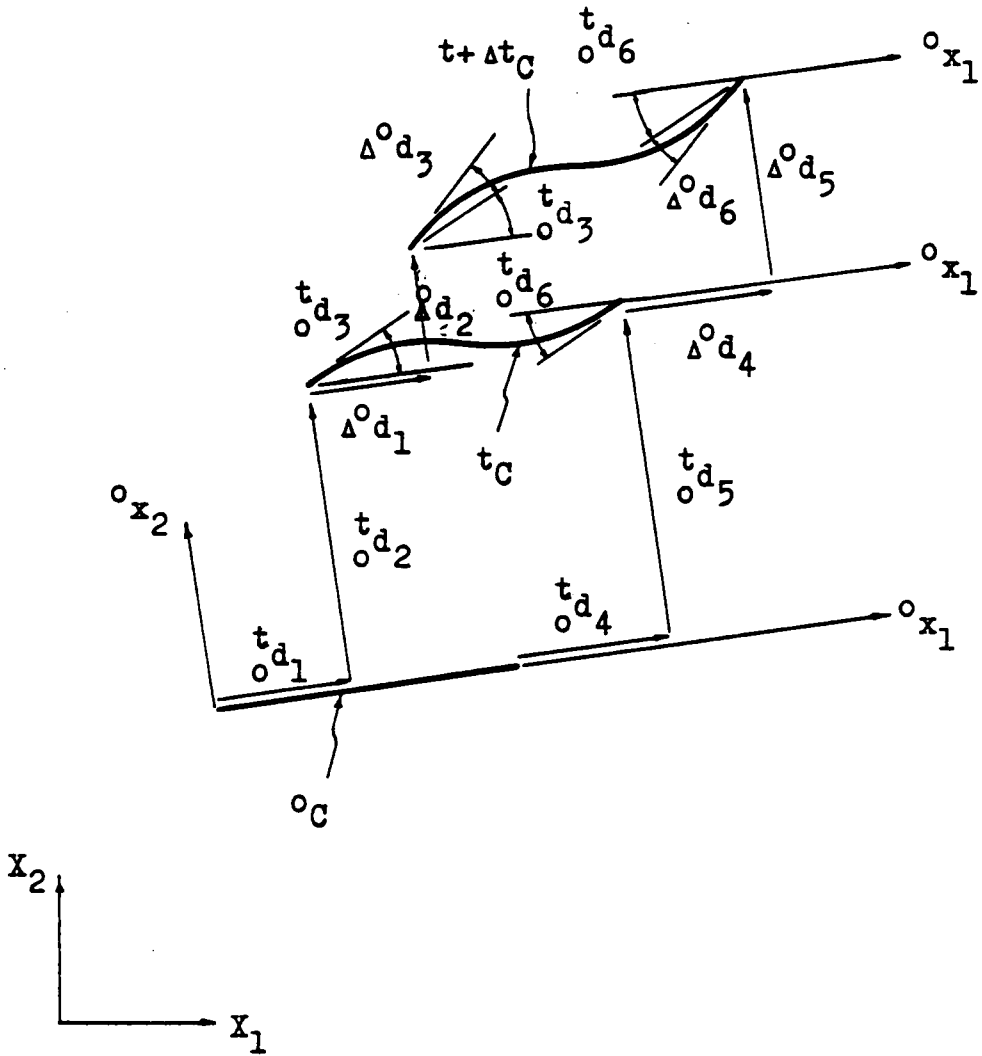


Fig. 2.5 Local Element Displacements in T.L. Formulation for Plane Frame

and  ${}^t_0\lambda$  is the orthogonal orientation matrix of element in configuration at time  $t$  referred to configuration at time  $0$ . Because all variables refer to the initial undeformed configuration  ${}^0C$  in the T.L. formulation, the transformation matrix  ${}^t_0A$  in Eq. (2.107) always not change from time  $0$  to  $t$  and  $t$  to  $t+\Delta t$ , i.e.

$${}^t_0\lambda = {}^0_0\lambda \quad (2.108)$$

The initial orientation matrix of space element in initial configuration  ${}^0C$  is [28]

$${}^0_0\lambda = \begin{pmatrix} {}^0c_1 & {}^0c_2 & {}^0c_3 \\ -\frac{{}^0c_1{}^0c_2}{o_1} \cos\phi - \frac{{}^0c_3}{o_1} \sin\phi & o_1 \cos\phi & -\frac{{}^0c_2{}^0c_3}{o_1} \cos\phi + \frac{{}^0c_1}{o_1} \sin\phi \\ \frac{{}^0c_1{}^0c_2}{o_1} \sin\phi - \frac{{}^0c_3}{o_1} \cos\phi & -o_1 \sin\phi & \frac{{}^0c_2{}^0c_3}{o_1} \sin\phi + \frac{{}^0c_1}{o_1} \cos\phi \end{pmatrix} \quad (2.109)$$

where

${}^0c_i$  = direction cosines of the local  $x_1$ -axis at time  $0$  with respect to the global coordinate system  $X_i$ -axis;  $i = 1, 2, 3$

$\phi$  = roll angle between the local  $x_2$ -axis and global  $X_2$ -axis



$${}^0\ell = ({}^0c_1^2 + {}^0c_3^2)^{1/2} \quad (2.110)$$

Eqs. (2.105), (2.106) yield the increments of the local element displacements and global element forces in T.L. formulation

$$\Delta^0d = {}^0\Lambda \Delta^0D \quad (2.111)$$

$$\Delta^0F^i = {}^0\Lambda^T \Delta^0f \quad (2.112)$$

where  $\Delta^0f$  and  $\Delta^0F^i$  are the vectors of the local and global incremental element forces from time  $t$  to  $t+\Delta t$  respectively, referred to configuration at time 0.

#### 2.4 CONVECTED COORDINATE FORMULATION

In the convected coordinate formulation developed by Belytschko [5], [17], all variables refer to the new incremented configuration at time  $t+\Delta t$ . The convected coordinate system means that each element is associated with a rigid cartesian coordinate system that rotates and translates with the element but does not deform with the element. Because the coordinate systems corresponding to the configurations at time  $t$  and  $t+\Delta t$  are independent of each other, an incremented concept in this formulation cannot be directly applied [18]. Hence, in the convected coordinate formulation

the displacements of each element at time  $t+\Delta t$  are decomposed into rigid body displacements and deformation displacements as shown in Fig. 2.6.

$$\begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} \mathbf{d} = \begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} \mathbf{d}^{\text{rig}} + \begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} \mathbf{d}^{\text{def}} \quad (2.113)$$

where

$$\begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} \mathbf{d} = \begin{bmatrix} t+\Delta t_{d_1} & t+\Delta t_{d_2} & t+\Delta t_{d_3} & \dots & t+\Delta t_{d_{10}} & t+\Delta t_{d_{11}} & t+\Delta t_{d_{12}} \end{bmatrix}^T \quad (2.114)$$

$$\begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} \mathbf{d}^{\text{rig}} = \begin{bmatrix} t+\Delta t_{d_1}^{\text{rig}} & t+\Delta t_{d_2}^{\text{rig}} & t+\Delta t_{d_3}^{\text{rig}} & \dots & t+\Delta t_{d_{10}}^{\text{rig}} & t+\Delta t_{d_{11}}^{\text{rig}} & t+\Delta t_{d_{12}}^{\text{rig}} \end{bmatrix}^T \quad (2.115)$$

$$\begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} \mathbf{d}^{\text{def}} = \begin{bmatrix} t+\Delta t_{d_1}^{\text{def}} & t+\Delta t_{d_2}^{\text{def}} & t+\Delta t_{d_3}^{\text{def}} & \dots & t+\Delta t_{d_{10}}^{\text{def}} & t+\Delta t_{d_{11}}^{\text{def}} & t+\Delta t_{d_{12}}^{\text{def}} \end{bmatrix}^T \quad (2.116)$$

The displacements in each element  $u$  can be similarly decomposed into rigid body displacements  $u^{\text{rig}}$  and deformation displacements  $u^{\text{def}}$ , i.e.

$$t+\Delta t_{\mathbf{u}} = t+\Delta t_{\mathbf{u}}^{\text{rig}} + t+\Delta t_{\mathbf{u}}^{\text{def}} \quad (2.117)$$

where  $u$  is measured from the initial position at time 0. For our problem small deformation and uniaxial state of strain are assumed; thus for space frame element  $u$  is  $3 \times 1$  column matrix and can be represented by interpolation functions  $h$  so that

$\gamma$  = rigid body rotation of element

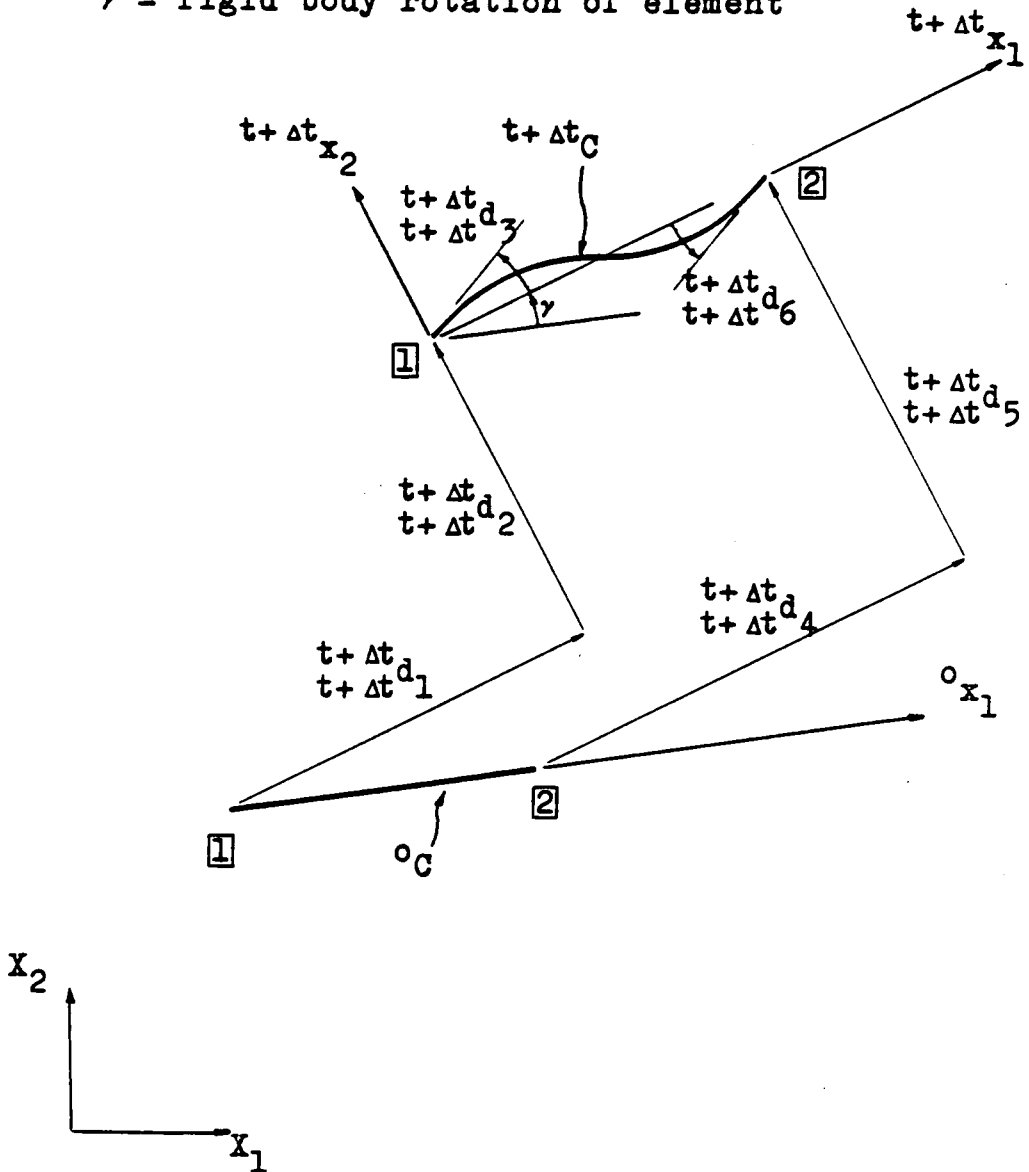


Fig. 2.6 Local Element Displacements in Convected Coordinate Formulation for Plane Frame

$$\begin{bmatrix} t+\Delta t_{u_1} \\ t+\Delta t_{u_2} \\ t+\Delta t_{u_3} \end{bmatrix} = \begin{bmatrix} h_1 & 0 & 0 & 0 & 0 & 0 & | & h_7 & 0 & 0 & 0 & 0 & 0 \\ 0 & h_2 & 0 & 0 & 0 & h_6 & | & 0 & h_8 & 0 & 0 & 0 & h_{12} \\ 0 & 0 & h_3 & 0 & h_5 & 0 & | & 0 & 0 & h_9 & 0 & h_{11} & 0 \end{bmatrix} \begin{bmatrix} t+\Delta t_{d_1} \\ t+\Delta t_{d_1} \\ t+\Delta t_{d_2} \\ t+\Delta t_{d_2} \\ \vdots \\ t+\Delta t_{d_{11}} \\ t+\Delta t_{d_{11}} \\ t+\Delta t_{d_{12}} \\ t+\Delta t_{d_{12}} \end{bmatrix}$$

or

$$t+\Delta t_u = h \begin{matrix} t+\Delta t_d \\ t+\Delta t_d \end{matrix} \quad (2.118)$$

Substituting Eq. (2.113) into Eq. (2.118) and from Eq. (2.117) we obtain

$$t+\Delta t_{u}^{rig} = h \begin{matrix} t+\Delta t_d^{rig} \\ t+\Delta t_d^{rig} \end{matrix} \quad (2.119)$$

and

$$t+\Delta t_{u}^{def} = h \begin{matrix} t+\Delta t_d^{def} \\ t+\Delta t_d^{def} \end{matrix} \quad (2.120)$$

In this formulation the global element displacements at time  $t+\Delta t$  are transformed to the local element displacements by the transformation (Figs. 2.3 and 2.4)

$$\begin{matrix} t+\Delta t_d \\ t+\Delta t_d \end{matrix} = \begin{matrix} t+\Delta t_A & t+\Delta t_D \\ t+\Delta t_A & t+\Delta t_D \end{matrix} \quad (2.121)$$

From the Appendix of reference [5], the relationships between the strains and deformation displacements in the convected coordinate system of each element are linear so that

$$\frac{t+\Delta t}{t+\Delta t} \varepsilon = B_L \frac{t+\Delta t}{t+\Delta t} d^{def} \quad (2.122)$$

where  $B_L$  is the linear strain displacement transformation matrix and can be obtained from the Eq. (2.120) and appropriate linear strain displacement equations.

As shown in the appendix of reference [5], consider two successive configurations at time  $t$  and  $t+\Delta t$ , the global internal nodal forces of each element can be derived from the principle of virtual work and it yields

$$\frac{t+\Delta t}{t+\Delta t} F = \frac{t+\Delta t}{t+\Delta t} \Lambda^T \int_V B_L^T \frac{t+\Delta t}{t+\Delta t} \sigma \frac{t+\Delta t}{t+\Delta t} dV \quad (2.123)$$

$$\frac{t}{t} F = \frac{t}{t} \Lambda^T \int_V B_L^T \frac{t}{t} \sigma \frac{t}{t} dV \quad (2.124)$$

in which  $\frac{t}{t} \sigma$  and  $\frac{t+\Delta t}{t+\Delta t} \sigma$  are the axial Cauchy stresses at time  $t$  and  $t+\Delta t$  respectively,  $V$  is the element volume which the change of volume can be neglected in the volume integration for small deformations,  $\Lambda$  is the transformation matrix from the global coordinate system to the convected coordinate system, defining

$$\Delta \frac{t}{t} \Lambda = \frac{t+\Delta t}{t+\Delta t} \Lambda - \frac{t}{t} \Lambda \quad (2.125)$$

$$\Delta \frac{t}{t} \sigma = \frac{t+\Delta t}{t+\Delta t} \sigma - \frac{t}{t} \sigma \quad (2.126)$$

and

$$\Delta^t_F = {}^{t+\Delta t}_F - {}^t_F \quad (2.127)$$

From Eqs. (2.123) to (2.127), we obtain

$$\Delta^t_F = {}^t_\Lambda \int_V B_L^T \Delta^t_\sigma dV + \Delta^t_\Lambda \int_V B_L^T {}^{t+\Delta t}_\sigma dV \quad (2.128)$$

To establish the global tangent stiffness matrix for an element, we must express  $\Delta^t_\sigma$  and  $\Delta^t_\Lambda$  in Eq. (2.128) in terms of the nodal incremental displacements  $\Delta^t_D$  defined as

$$\Delta^t_D = {}^{t+\Delta t}_D - {}^t_D \quad (2.129)$$

Because in this convected coordinate system the nodal element displacements between the configurations at time  $t$  and  $t+\Delta t$  are independent of each other, the nodal incremental displacements in Eq. (2.129) are not vector quantity. Hence, it is inappropriate to express  $\Delta^t_\sigma$  and  $\Delta^t_\Lambda$  in Eq. (2.128) in terms of  $\Delta^t_D$ . Therefore, using the finite element model in the convected coordinate formulation to establish the local element tangent stiffness matrix is inconvenient.

## 2.5 COMPARISON OF U.L. AND T.L. FORMULATIONS

The main difference between the U.L. and T.L. formulations lies in the coordinate system referred to formulate the incremental equilibrium equations; the former refers to the current deformed configuration  ${}^t C$ ; the latter refers to the initial undeformed configuration  ${}^0 C$ .

Regardless of which formulation used, we will show that the incremental equilibrium Eqs. (2.55) and (2.104) are identical in both formulations [9]. From Eqs. (2.50) and (2.100) the global element tangent stiffness matrices are

$${}^t K^i = \begin{matrix} {}^t \Lambda^T & {}^t k & {}^t \Lambda \\ {}^t \Lambda & {}^t k & {}^t \Lambda \end{matrix} \quad \text{for U.L. formulation} \quad (2.130)$$

$${}^0 K^i = \begin{matrix} {}^0 \Lambda^T & {}^0 k & {}^0 \Lambda \\ {}^0 \Lambda & {}^0 k & {}^0 \Lambda \end{matrix} \quad \text{for T.L. formulation} \quad (2.131)$$

in which

$\begin{matrix} {}^t \Lambda & {}^t \Lambda \\ {}^t \Lambda & {}^0 \Lambda \end{matrix}$  = transformation matrix between the local coordinate axes in configuration at time  $t$  and the global coordinate axes

Let

$$\begin{matrix} {}^t \Lambda & \\ {}^t \Lambda & \end{matrix} = R \begin{matrix} {}^t \Lambda \\ {}^0 \Lambda \end{matrix} \quad (2.132)$$

where  $R$  is the transformation matrix from the local coordinates  ${}^0x_i$  to  ${}^tx_i$  in space (Fig. 2.1); it has been derived by Bathe [9] and will not be described here.

Substituting Eq. (2.132) into (2.130)

$$\begin{matrix} t \\ t \end{matrix} K^i = \begin{matrix} t \\ 0 \end{matrix} \Lambda^T \begin{matrix} t \\ 0 \end{matrix} k \begin{matrix} t \\ 0 \end{matrix} \Lambda \quad (2.133)$$

where

$$\begin{matrix} t \\ 0 \end{matrix} k = R^T \begin{matrix} t \\ t \end{matrix} k R \quad (2.134)$$

From Eqs. (2.131) and (2.133), we have

$$\begin{matrix} t \\ t \end{matrix} K^i = \begin{matrix} t \\ 0 \end{matrix} K^i \quad (2.135)$$

Therefore, the global element tangent stiffness matrices  $K^i$  are identical in both formulations.

As a result of Eq. (2.135), the assembled system tangent stiffness matrices  $K$  in Eqs. (2.55) and (2.104) are identical in both formulations.

$$\begin{matrix} t \\ t \end{matrix} K = \begin{matrix} t \\ 0 \end{matrix} K \quad (2.136)$$

Similarly, from Eqs. (2.51), (2.101) the global element forces are

$$\begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} F^i{}^{k-1} = \begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} \Lambda^{k-1}{}^T \begin{matrix} t+\Delta t \\ t+\Delta t \end{matrix} f^{k-1}$$



for U.L. formulation (2.137)

$${}_{0F}^{t+\Delta t} \mathbf{i}^{k-1} = {}_{0A}^{t+\Delta t} \mathbf{k}^{-1T} {}_{0f}^{t+\Delta t} \mathbf{k}^{-1}$$

for T.L. formulation (2.138)

Substituting Eq. (2.132) into Eq. (2.137)

$$\frac{{}_{0F}^{t+\Delta t} \mathbf{i}^{k-1}}{t+\Delta t} = \frac{{}_{0A}^{t+\Delta t} \mathbf{k}^{-1T}}{t+\Delta t} \frac{{}_{0f}^{t+\Delta t} \mathbf{k}^{-1}}{t+\Delta t} \quad (2.139)$$

where

$$\frac{{}_{0f}^{t+\Delta t} \mathbf{k}^{-1}}{t+\Delta t} = \mathbf{R}^T \frac{{}_{0f}^{t+\Delta t} \mathbf{k}^{-1}}{t+\Delta t} \quad (2.140)$$

From Eqs. (2.138), (2.139), we have

$$\frac{{}_{0F}^{t+\Delta t} \mathbf{i}^{k-1}}{t+\Delta t} = \frac{{}_{0F}^{t+\Delta t} \mathbf{i}^{k-1}}{t+\Delta t} \quad (2.141)$$

Therefore, the global element force vectors are identical in both formulations.

As a result of Eq. (2.141), the vectors of equivalent nodal element forces  $F$  in Eqs. (2.55), (2.104) are identical in both formulations.

$$\frac{{}_{0F}^{t+\Delta t} \mathbf{k}^{-1}}{t+\Delta t} = \frac{{}_{0F}^{t+\Delta t} \mathbf{k}^{-1}}{t+\Delta t} \quad (2.142)$$

From Eqs. (2.136), (2.142) we may conclude that the incremental equilibrium Eqs. (2.55) and (2.104) are identical in both formulations.

As a result of Eqs. (2.55) and (2.104), we are led to the conclusion that

The main advantages of the U.L. formulation are:

1. Since the rotation referred to the current configuration  ${}^t_C$  is infinitesimal, it can be treated as vector in space frame [2], [18].
2. The linear incremental strains in Eq. (2.12), referred to the current configuration  ${}^t_C$ , not include the initial displacement effect; it results that the linear strain incremental stiffness matrix  ${}^t_K_L$  in Eq. (2.55) is not including the initial displacement effect.
3. Because the strains, referred to the current configuration  ${}^t_C$ , are so infinitesimal that sometimes we can neglect the nonlinear part of incremental strains in Eq. (2.11); it results that even the nonlinear strain incremental stiffness matrix  ${}^t_K_{NL}$  in Eq. (2.55) can be omitted. In fact, as we pointed out earlier, the stiffness used to solve Eq. (2.55) need not be exact.

4. Items 2 and 3 result that simple and approximate formulations of the U.L. in Eq. (2.55) may be developed, which make the U.L. formulation very efficient and clearly superior to the T.L. formulation.
5. Especially more suitable for the large displacement but small strain problems which are very common in practice for many types of problems [36].
6. If the relative member deformations are small enough that we can directly employ the beam-column model which uses convected coordinate system to formulate the tangent stiffness matrix and nodal equivalent force in Eq. (2.55), in addition to finite element model.

The disadvantages of the U.L. formulation are:

1. The transformation matrix  ${}^t_{t+\Delta t}\Lambda$  in Eq. (2.54), referred to the current configuration  ${}^tC$ , must be updated at each time step (iteration).
2. The stresses at time  $t+\Delta t$  as in Eq. (2.5)

$${}^{t+\Delta t}S_{ij} = {}^t\tau_{ij} + {}^tS_{ij}$$

where the stress increments  ${}^tS_{ij}$  from time  $t$  to  $t+\Delta t$  are the 2nd Piola-Kirchhoff stresses referred to time  $t$ , which must be transformed into Cauchy stresses.

Therefore, it is slightly more complicated than in T.L. formulation to compute the stresses [36].

3. Mass matrix would be updated at each iteration which leads to complexity in the U.L. formulation to dynamic problems.
4. As to economic consideration, items 1 to 3 result in more computational effort than T.L. formulation.

On the other hand, the advantages of the T.L. formulation are:

1. The transformation matrix  ${}^t_0\Lambda$  in Eq. (2.107) remains unchanged throughout each iteration as in Eq. (2.108).
2. The use of a unique type of the 2nd Piola-Kirchhoff stresses in Eq. (2.7).
3. Because the mass matrix would then be constant in the extension to dynamic problems throughout each iteration, it leads to simplification in the T.L. formulation [12], [3].
4. Items 1 to 3 result in less computational effort than U.L. formulation which saves computer time and money.

The disadvantages of the T.L. formulation are:

1. Because the rotation referred to the undeformed configuration  ${}^0C$  is finite, it cannot be treated as vector in space problems; therefore, the merit of matrix operation will be lost [3], [18].
2. The linear incremental strains in Eq. (2.63), referred to the initial configuration  ${}^0C$ , include the initial displacement effect; it results that the linear strain incremental stiffness matrix  ${}^t_0K_L$  in Eq. (2.104) include the initial displacement effect.
3. Because the strains referred to the initial configuration  ${}^0C$  is finite, as displacements become larger and larger, nonlinear term of incremental strains in Eq. (2.62) are significant; in other words, the nonlinear strain incremental stiffness matrix  ${}^t_0K_{NL}$  in Eq. (2.104) must be taken into account to obtain the exact stiffness for large displacements.
4. Employing the beam-column model, which uses convected coordinate system, in T.L. formulation, the local element stiffness matrix and force vector must be transformed to the initial coordinate system.

In conclusion, no matter which formulation we use, the same structural stiffness matrices and nodal equivalent ele-

ment forces should be obtained. Therefore, the solutions using different formulations must be the same, if the same number of elements are employed [9]. Whether to use the U.L. or the T.L. formulations depends largely on the program design and the practical problems. In beam analysis the U.L. formulation is more effective than the T.L. formulation in which the additional  ${}^t_0B_{L1}$  matrix must be evaluated as in Eq. (2.91c). Based on this concept, in next chapters we will only employ the U.L. formulation to solve the incremental equilibrium equation as shown in Eq. (2.55).

## Chapter III

### DEFORMATION DISPLACEMENTS OF SPACE FRAME ELEMENT

#### 3.1 INTRODUCTION

The element displacements are decomposed into rigid body displacements and deformation displacements. In this chapter the deformation displacements are derived for a space frame element with large displacements and rotations of nodes and members. However, the strain in each element is assumed to be small; that is, the element deformations are restricted to be small. Four types of coordinate systems are employed to derive the element deformation displacements.

The beam element is assumed to be straight and so slender that shear deformations can be neglected in comparison with bending deformations. It is assumed that the cross section of the beam remains plane and normal to the centroidal axis during deformation. The cross section is doubly symmetric, and the torsional stiffness will be treated independently from bending and axial stiffnesses. However, interaction of axial force and bending is considered [21].

### 3.2 COORDINATE SYSTEM OF SPACE FRAME

To derive the element deformation displacements, four types of coordinate systems are defined (Fig. 3.1):

1. The structural global coordinate system  $X_1, X_2, X_3$ .
2. The nodal coordinate system  $y_1, y_2, y_3$ ; which is rigidly connected to each node. The initial directions of the  $y_1, y_2, y_3$ -axes are chosen to coincide with the global axes.
3. The element convected coordinate system  $x_1, x_2, x_3$ ; which is associated with each element. The  $x_1$ -axis always moves with the beam element and passes through the centroids of the two end sections of the element; in the initial state the  $x_2, x_3$ -axes coincide with the principal axes of the cross section, and in the deformed state they are taken to be the average of the rotations of the two ends about the  $x_1$ -axis.
4. The end cross section (body [5],[17]) coordinate system  $\bar{x}_1, \bar{x}_2, \bar{x}_3$ ; which rotates with the end cross section. The  $\bar{x}_1$ -axis is tangent to the deformed axis of the member, and the  $\bar{x}_2, \bar{x}_3$ -axes coincide with the principal axes of the cross section.



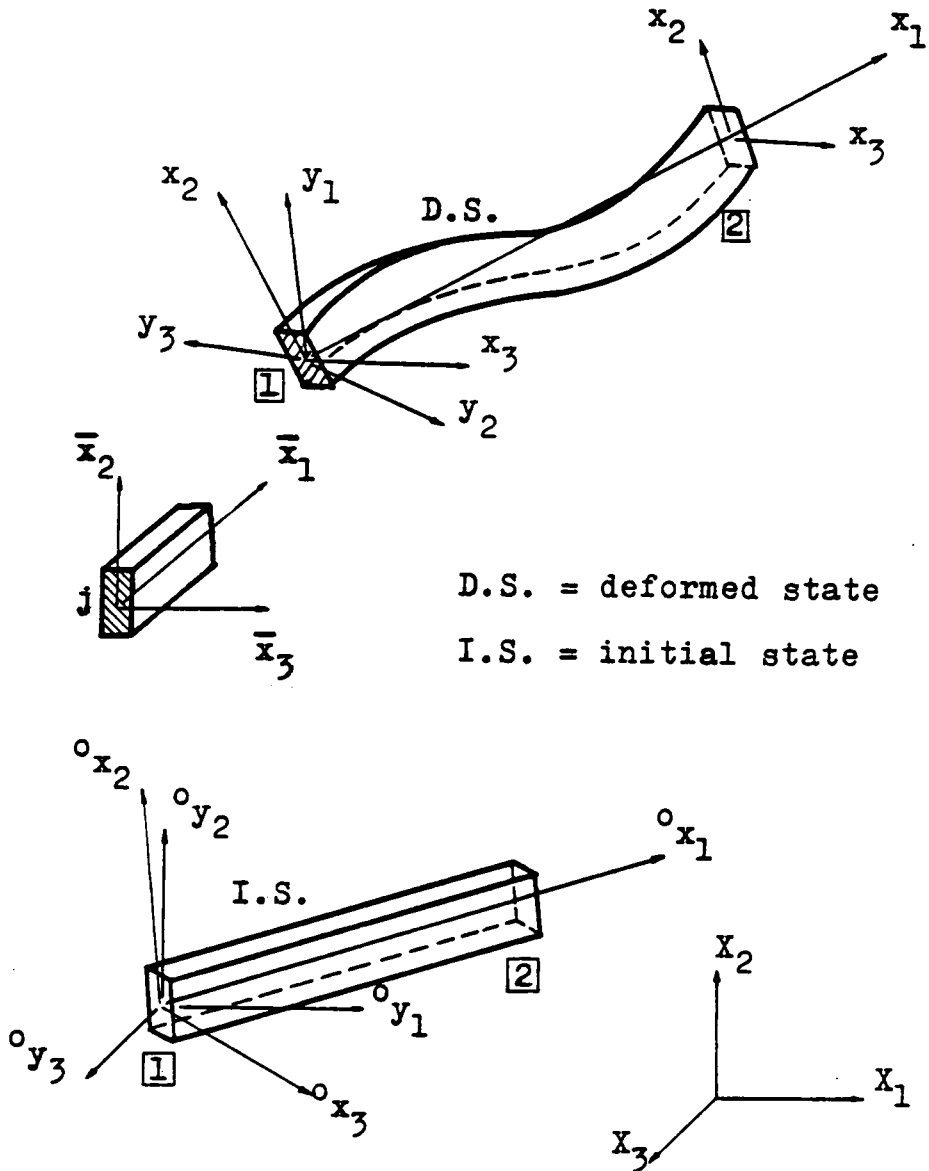


Fig. 3.1 Coordinate System of Space Frame

### 3.3 NODAL DISPLACEMENTS

The orientation of each deformed node  $j$  can be described in terms of the direction cosines of three node axes  $y_1, y_2, y_3$  relative to the global coordinates:

$${}^t \begin{pmatrix} y_1 \\ y_2 \\ y_3 \end{pmatrix}_j = \begin{pmatrix} \alpha_{11} & \alpha_{12} & \alpha_{13} \\ \alpha_{21} & \alpha_{22} & \alpha_{23} \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix}_j \begin{pmatrix} X_1 \\ X_2 \\ X_3 \end{pmatrix} \quad (3.1)$$

or

$${}^t y_j = {}^t \alpha_j X$$

where  ${}^t \alpha_j$  is the nodal orientation matrix at time  $t$ . The orthogonal nodal orientation matrix can be partitioned as

$$\alpha_j = [i_1 \ i_2 \ i_3]^T \quad (3.2)$$

where

$$i_\ell = [\alpha_{\ell 1} \ \alpha_{\ell 2} \ \alpha_{\ell 3}] ; \quad \ell=1,2,3 \quad (3.3)$$

and

$$\alpha_{\ell m} = \cos \phi_{\ell m} ; \quad m=1,2,3 \quad (3.4)$$

$\alpha_{\ell m}$  is the direction cosine and  $\phi_{\ell m}$  is the direction angle between the nodal  $y_{\ell}$ -axis and the global  $X_m$ -axis.

In analyzing large displacements of a space frame, the deformed configuration of each node  $j$  can be represented in terms of a translation vector  $[D_{j1} \ D_{j2} \ D_{j3}]^T$  relative to the global coordinate system and an node orientation matrix  $\alpha_j$ .

Consider the increments of nodal displacements which are assumed to be small during a load step. The incremental displacement vector of node  $j$  is

$$\Delta U_j = \begin{bmatrix} \Delta U_{jt} \\ \text{-----} \\ \Delta U_{jr} \end{bmatrix} \quad (3.5)$$

in which  $\Delta U_{jt}$ ,  $\Delta U_{jr}$  are the translational and rotational incremental vectors of node  $j$  in the global coordinate system:

$$\Delta U_{jt} = [ \Delta D_{j1} \ \Delta D_{j2} \ \Delta D_{j3} ]^T \quad (3.6)$$

$$\Delta U_{jr} = [ \Delta D_{j4} \ \Delta D_{j5} \ \Delta D_{j6} ]^T \quad (3.7)$$

where  $\Delta D_{j1}$ ,  $\Delta D_{j2}$ ,  $\Delta D_{j3}$  are the incremental deflections at node  $j$  and  $\Delta D_{j4}$ ,  $\Delta D_{j5}$ ,  $\Delta D_{j6}$  are the incremental rotations at node  $j$ .

According to appendix A, the node transformation from the configuration  ${}^t C$  to the configuration  ${}^{t+\Delta t} C$  is (Fig. 3.2)

$${}^{t+\Delta t} y_j = (I + {}^t R_j) {}^t y_j \quad (3.8)$$

where  $I$  is identity matrix and  ${}^t R_j$  is the rotation matrix at node  $j$ , defined as

$${}^t R_j = \begin{bmatrix} 0 & \Delta^t \omega_{j3} & -\Delta^t \omega_{j2} \\ -\Delta^t \omega_{j3} & 0 & \Delta^t \omega_{j1} \\ \Delta^t \omega_{j2} & -\Delta^t \omega_{j1} & 0 \end{bmatrix} \quad (3.9)$$

where  $\Delta \omega_{ji}$  is the incremental nodal rotation at node  $j$  about  ${}^t y_i$ -axis (Fig. 3.2),  $i=1, 2, 3$ ; defining

$$\Delta^t \omega_j = \begin{bmatrix} \Delta^t \omega_{j1} \\ \Delta^t \omega_{j2} \\ \Delta^t \omega_{j3} \end{bmatrix} \quad (3.10)$$

where  $\Delta^t \omega_j$  can be obtained from Eq. (3.1) as

$$\Delta^t \omega_j = {}^t \alpha_j \Delta^t U_{jr} \quad (3.11)$$

Substituting Eq. (3.1) into Eq. (3.8) yields

$${}^{t+\Delta t} y_j = (I + {}^t R_j) {}^t \alpha_j X \quad (3.12)$$

$$= \underbrace{{}^t \alpha_j X}_{{}^t y_j} + \underbrace{\Delta^t \alpha_j X}_{\Delta^t y_j} \quad (3.13)$$

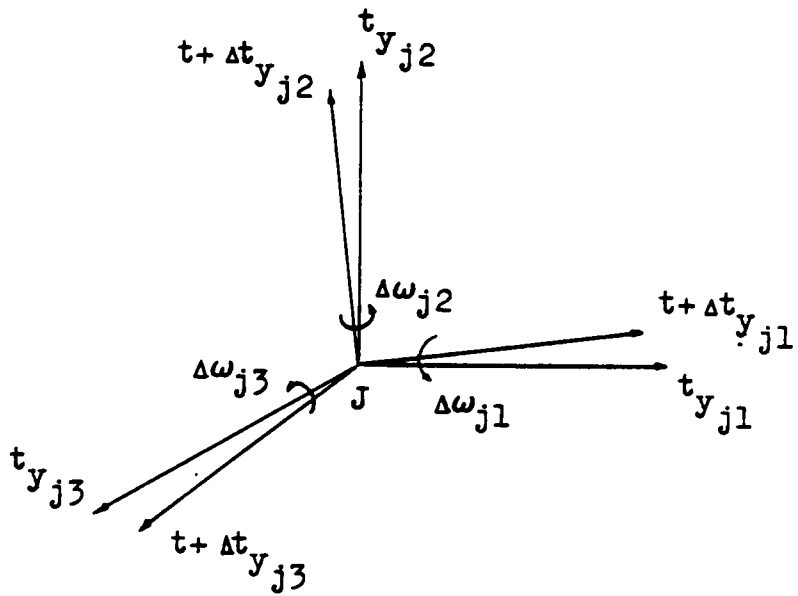


Fig. 3.2 Incremental Nodal Rotations

or

$${}^{t+\Delta t}y_j = {}^{t+\Delta t}\alpha_j X \quad (3.14)$$

where

$${}^{t+\Delta t}\alpha_j = {}^t\alpha_j + \Delta^t\alpha_j \quad (3.15)$$

$\Delta^t\alpha_j$  is the change of the node orientation matrix due to incremental rotations of node  $j$ , defined as

$$\Delta^t\alpha_j = {}^tR_j {}^t\alpha_j \quad (3.16)$$

### 3.4 ELEMENT DEFORMATION DISPLACEMENTS

The relative deformation displacements of an element are referred to the rigid-convected (corotational) coordinate system  $x_1, x_2, x_3$ . Each discrete element is referred to a cartesian coordinate system that rotates and translates with the element but does not deform with the element (Fig. 3.1).

The element orientation is defined by the element orientation matrix,  $\lambda$ . The rows of  $\lambda$  are the direction cosines of  $x_1, x_2, x_3$ -axes relative to the global coordinates  $X_1, X_2, X_3$ , respectively.

$${}^t \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} \lambda_{11} & \lambda_{12} & \lambda_{13} \\ \lambda_{21} & \lambda_{22} & \lambda_{23} \\ \lambda_{31} & \lambda_{32} & \lambda_{33} \end{pmatrix} \begin{pmatrix} X_1 \\ X_2 \\ X_3 \end{pmatrix} \quad (3.17)$$

or

$${}^t x = {}^t \lambda X$$

The orthogonal element orientation matrix can be partitioned as

$$\lambda = \begin{pmatrix} j_1 \\ j_2 \\ j_3 \end{pmatrix} \quad (3.18)$$

where

$$j_\ell = [ c_1 \ c_2 \ c_3 ] ; \quad \ell=1 \quad (3.19)$$

$$j_\ell = [ \lambda_{\ell 1} \ \lambda_{\ell 2} \ \lambda_{\ell 3} ] ; \quad \ell=2,3 \quad (3.20)$$

$c_i$  = direction cosines of the local  $x_1$ -axis in any configuration C with respect to the global coordinates  $X_i$

and

$$\lambda_{\ell m} = \cos \omega_{\ell m} ; \quad m=1,2,3 \quad (3.21)$$

$\lambda_{\ell m}$  is the direction cosine and  $\omega_{\ell m}$  is the direction angle between the local element  $x_{\ell}$ -axis and the global  $X_m$ -axis.

The element orientation matrix of a space element depends on the orientations of the principal plane of the element defined by the local  $x_1, x_2$ -axes. The initial orientation matrix of the space element in configuration  ${}^0C$  is defined in Eq. (2.109).

In the formulation developed here, no restrictions are made on the rotations and translations of the node. However, the relative deformations of the members are small such that we can apply the conventional beam-column theory to the member-force deformation relations.

The deformation displacements of a space frame element is represented by the end angles  $e_{13}, e_{23}, e_{12}, e_{22}$ ; total twist  $\phi_t$ , and relative axial displacement  $u$  (Fig. 3.3).



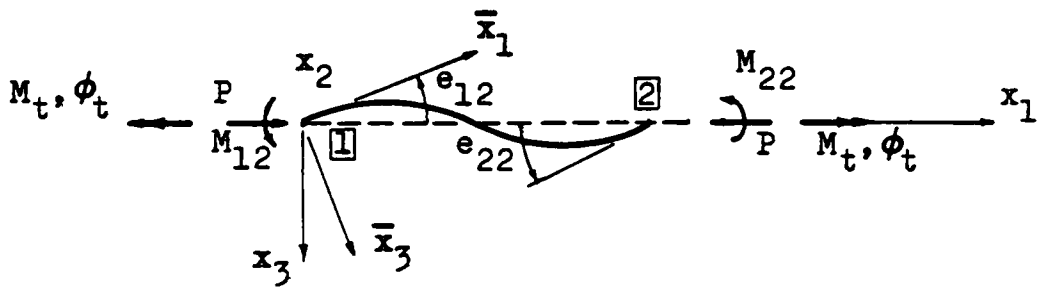
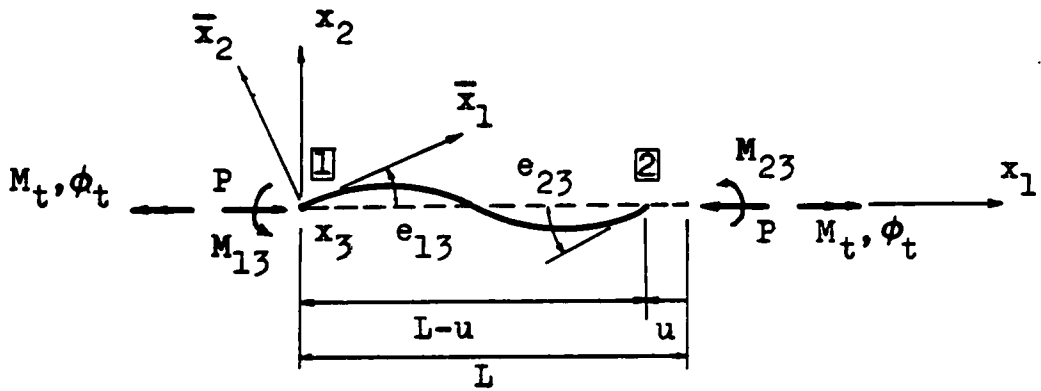


Fig. 3.3 Element Deformation Displacements and Associated Forces

Consider the increments of nodal displacements and element deformations which are assumed to be small during a load step. The element transformation from the configuration  ${}^t_C$  to the configuration  ${}^{t+\Delta t}_C$  is

$${}^{t+\Delta t}_x = (I + {}^t_R) {}^t_x \quad (3.22)$$

where

$${}^t_R = \begin{bmatrix} 0 & \Delta\psi_3 & -\Delta\psi_2 \\ -\Delta\psi_3 & 0 & \Delta\psi_1 \\ \Delta\psi_2 & -\Delta\psi_1 & 0 \end{bmatrix} \quad (3.23)$$

$$\Delta\psi_1 = (\Delta d_{10} - \Delta d_4)/2 \quad (3.24)$$

$$\Delta\psi_2 = -(\Delta d_9 - \Delta d_3)/l_e \quad (3.25)$$

$$\Delta\psi_3 = (\Delta d_8 - \Delta d_2)/l_e \quad (3.26)$$

and the local element displacements are numbered as shown in Fig. 2.3. Substituting Eq. (3.17) into Eq. (3.22) yield

$$\begin{aligned} {}^{t+\Delta t}_x &= (I + {}^t_R) {}^t_\lambda X \\ &= \underbrace{{}^t_\lambda X}_{{}^t_x} + \underbrace{\Delta {}^t_\lambda X}_{\Delta {}^t_x} \end{aligned} \quad (3.27)$$

or

$${}^{t+\Delta t}x = {}^{t+\Delta t}t_\lambda X \quad (3.28)$$

where

$${}^{t+\Delta t}t_\lambda = t_\lambda + \Delta t_\lambda \quad (3.29)$$

$\Delta t_\lambda$  is the change of element orientation matrix due to incremental rotations of element, defined as

$$\Delta t_\lambda = t_R t_\lambda \quad (3.30)$$

The nodal coordinates  $y_1, y_2, y_3$  are initially parallel to the global coordinates  $X_1, X_2, X_3$ . Thus, from Eqs. (3.1), (3.17) one obtains

$${}^0y_j = X = {}^0t_\lambda T {}^0x \quad (3.31)$$

where left superscript '0' refers to the initial undeformed configuration  ${}^0C$ .

Because the member-end sections are rigidly connected to the respective nodes, Eq. (3.31) may be written in the deformed configuration  ${}^{t+\Delta t}C$  [4] as

$${}^{t+\Delta t}y_j = {}^0t_\lambda T {}^{t+\Delta t}\bar{x}_j \quad (3.32)$$

where  ${}^{t+\Delta t}\bar{x}_j$  is the coordinate vector of end section at node  $j$  in the  ${}^{t+\Delta t}C$  configuration.

From Eqs. (3.32), (3.14), (3.17) one obtains

$${}^{t+\Delta t}\bar{x}_j = {}^{0\lambda}{}^{t+\Delta t}y_j = {}^{0\lambda}{}^{t+\Delta t}a_j X = [{}^{0\lambda}{}^{t+\Delta t}a_j \quad {}^{t+\Delta t}a_j^T] {}^{t+\Delta t}x \quad (3.33)$$

Since deformations are small

$$\sin e_{jn} = e_{jn}, \quad \cos e_{jn} = 1 \quad (3.34)$$

where  $e_{jn}$  is the relative element rotation at node  $j$  about  $x_n$ -axis.

From Figs. 3.1 and 3.3 one obtains

$${}^{t+\Delta t}\bar{x}_j = \begin{pmatrix} 1 & \cos(90^\circ - e_{j3}) & \cos(90^\circ + e_{j2}) \\ \cos(90^\circ + e_{j3}) & 1 & \cos(90^\circ \pm \phi_t/2) \\ \cos(90^\circ - e_{j2}) & \cos(90^\circ \mp \phi_t/2) & 1 \end{pmatrix} {}^{t+\Delta t}x \quad (3.35)$$

where  $j$  is 1 for the left end and 2 for the right end;  $e_{j2}$ ,  $e_{j3}$  are the relative end rotations at node  $j$  about the  $x_2$ ,  $x_3$ -axes respectively,  $\phi_t$  is the relative angle of twist of the element ends about the  $x_1$ -axis, and the upper and lower signs apply to nodes 1 and 2, respectively.

Substituting Eq. (3.34) into Eq. (3.35) yields

$${}^{t+\Delta t} \bar{x}_j = \begin{pmatrix} 1 & e_{j3} & -e_{j2} \\ -e_{j3} & 1 & \mp \phi_t/2 \\ e_{j2} & \pm \frac{\phi_t}{2} & 1 \end{pmatrix} {}^{t+\Delta t} x \quad (3.36)$$

Comparison of Eqs. (3.33) and (3.36) yields

$${}^{t+\Delta t} e_{j2} = [{}^0 \lambda \quad {}^{t+\Delta t} \alpha_j \quad {}^{t+\Delta t} \lambda^T]_{31} = - [{}^0 \lambda \quad {}^{t+\Delta t} \alpha_j \quad {}^{t+\Delta t} \lambda^T]_{13} \quad (3.37)$$

$${}^{t+\Delta t} e_{j3} = - [{}^0 \lambda \quad {}^{t+\Delta t} \alpha_j \quad {}^{t+\Delta t} \lambda^T]_{21} = [{}^0 \lambda \quad {}^{t+\Delta t} \alpha_j \quad {}^{t+\Delta t} \lambda^T]_{12} \quad (3.38)$$

$${}^{t+\Delta t} \phi_t/2 = \mp [{}^0 \lambda \quad {}^{t+\Delta t} \alpha_j \quad {}^{t+\Delta t} \lambda^T]_{23} = \pm [{}^0 \lambda \quad {}^{t+\Delta t} \alpha_j \quad {}^{t+\Delta t} \lambda^T]_{32} \quad (3.39)$$

where  $[ \ ]_{mn}$  denotes the entry in the  $m^{\text{th}}$  row and  $n^{\text{th}}$  column of the matrix.

From Eqs. (3.1) and (3.32) we obtain

$${}^{t+\Delta t} \bar{x}_j = {}^{t+\Delta t} p_j X \quad (3.40)$$

where

$${}^{t+\Delta t} p_j = {}^0 \lambda \quad {}^{t+\Delta t} \alpha_j \quad (3.41)$$

${}^{t+\Delta t}p_j$ ,  ${}^{t+\Delta t}\alpha_j$  are the end section and node orientation matrices at node  $j$  in the configuration at time  $t+\Delta t$ , respectively. The rows of  $p_j$  represent the direction cosines of the normal and principal directions of the corresponding end section at joint  $j$  with respect to the global coordinates  $X_1, X_2, X_3$ .

The deformation displacements can be expressed directly from Eqs. (3.37), (3.38), (3.39), (3.17) and (2.109) as

$${}^{t+\Delta t}e_{j2} = \begin{bmatrix} \frac{{}^0c_3}{{}_0e} & 0 & \frac{{}^0c_1}{{}_0e} \end{bmatrix} {}^{t+\Delta t}\alpha_j \begin{bmatrix} {}^{t+\Delta t}c_1 \\ {}^{t+\Delta t}c_2 \\ {}^{t+\Delta t}c_3 \end{bmatrix} \quad (3.42)$$

$${}^{t+\Delta t}e_{j3} = - \begin{bmatrix} \frac{{}^0c_1 {}^0c_2}{{}_0e} & \frac{{}^0c_2 {}^0c_3}{{}_0e} \end{bmatrix} {}^{t+\Delta t}\alpha_j \begin{bmatrix} {}^{t+\Delta t}c_1 \\ {}^{t+\Delta t}c_2 \\ {}^{t+\Delta t}c_3 \end{bmatrix} \quad (3.43)$$

$${}^{t+\Delta t}\phi_t = 2 \begin{bmatrix} \frac{{}^0c_1 {}^0c_2}{{}_0e} & \frac{{}^0c_2 {}^0c_3}{{}_0e} \end{bmatrix} {}^{t+\Delta t}\alpha_j \begin{bmatrix} {}^{t+\Delta t}\lambda_{31} \\ {}^{t+\Delta t}\lambda_{32} \\ {}^{t+\Delta t}\lambda_{33} \end{bmatrix} \quad (3.44)$$

and

$$u = {}^0L - {}^{t+\Delta t}L \quad (3.45)$$

Chapter IV  
FINITE ELEMENT MODEL

4.1 INTRODUCTION

To solve the incremental equilibrium equation in the U.L. formulation we must to evaluate the finite element matrices in Eq. (2.55). The local element secant stiffness matrix of the plane frame element, developed by Katzenberger [48], is extended to form the local element secant stiffness matrix of the space frame element in which the torsional forces are treated independently and obtained by linear theory.

For a straight small strain beam of constant cross section in the convected coordinate system, Hermitian interpolation functions are employed to interpolate the transverse bending displacements, and Lagrange interpolation functions are used to interpolate the axial and torsional displacements [9],[33],[34].

For convenience, the left superscripts and subscripts are not shown in this formulation.

## 4.2 INTERPOLATION FUNCTIONS FOR INCREMENTAL DISPLACEMENTS

The natural coordinate system in Fig. 4.1 is used. The interpolation matrix of Eq. (2.24) contains the entries [9], [33], [34]:

$$\begin{aligned}
 h_1 &= 1 - \xi \\
 h_2 &= 1 - 3\xi^2 + 2\xi^3 \\
 h_3 &= 1 - 3\xi^2 + 2\xi^3 \\
 h_5 &= -L\xi + 2L\xi^2 - L\xi^3 \\
 h_6 &= L\xi - 2L\xi^2 + L\xi^3 \\
 h_7 &= \xi \\
 h_8 &= 3\xi^2 - 2\xi^3 \\
 h_9 &= 3\xi^2 - 2\xi^3 \\
 h_{11} &= L\xi^2 - L\xi^3 \\
 h_{12} &= -L\xi^2 + L\xi^3
 \end{aligned} \tag{4.1}$$

where

$$\xi = x_1 / L \tag{4.2}$$

Eqs. (4.1) yield the 1st and 2nd derivatives of the interpolation functions with respect to  $\xi$ -axis:

$$\begin{aligned}
 h_{1,1} &= -1 \\
 h_{2,1} &= -6\xi + 6\xi^2 \\
 h_{3,1} &= -6\xi + 6\xi^2 \\
 h_{5,1} &= -L + 4L\xi - 3L\xi^2
 \end{aligned}$$



$$\begin{aligned}
h_{6,1} &= L - 4L\xi + 3L\xi^2 \\
h_{7,1} &= 1 \\
h_{8,1} &= 6\xi - 6\xi^2 \\
h_{9,1} &= 6\xi - 6\xi^2 \\
h_{11,1} &= 2L\xi - 3L\xi^2 \\
h_{12,1} &= -2L\xi + 3L\xi^2
\end{aligned}
\tag{4.3}$$

and

$$\begin{aligned}
h_{2,11} &= -6 + 12\xi \\
h_{3,11} &= -6 + 12\xi \\
h_{5,11} &= 4L - 6L\xi \\
h_{6,11} &= -4L + 6L\xi \\
h_{8,11} &= 6 - 12\xi \\
h_{9,11} &= 6 - 12\xi \\
h_{11,11} &= 2L - 6L\xi \\
h_{12,11} &= -2L + 6L\xi
\end{aligned}
\tag{4.4}$$

Using Eqs. (4.1), (4.3), and (4.4), the linear and nonlinear strain-displacement transformation matrices that are required to evaluate the tangent stiffness matrix and nodal force vector of an element can directly be evaluated.

$\xi = \text{natural coordinate} = x_1 / L$

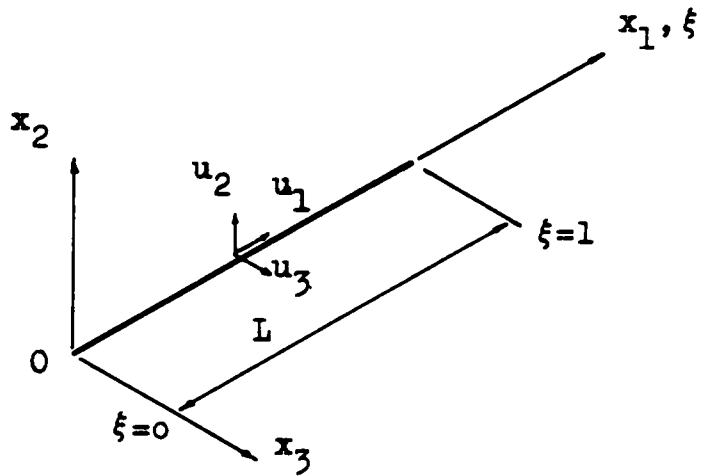


Fig. 4.1 Local and Natural Coordinate Systems

### 4.3 LINEAR STRAIN INCREMENTAL STIFFNESS MATRIX

Employing Eqs. (4.3), (4.4) and substituting Eq. (2.49) into Eq. (2.45) one obtains the linear strain incremental stiffness matrix of the local element model with the coordinate system in Fig. 2.3 [28]

$$k_L = \int_V B_L^T E B_L dV =$$

1	2	3	4	5	6	7	8	9	10	11	12	
$\frac{EA}{L}$	0	0	0	0	0	$-\frac{EA}{L}$	0	0	0	0	0	1
	$\frac{12EI_3}{L^3}$	0	0	0	$\frac{6EI_3}{L^2}$	0	$-\frac{12EI_3}{L^3}$	0	0	0	$\frac{6EI_3}{L^2}$	2
		$\frac{12EI_2}{L^3}$	0	$-\frac{6EI_2}{L^2}$	0	0	0	$-\frac{12EI_2}{L^3}$	0	$-\frac{6EI_2}{L^2}$	0	3
			$\frac{GJ}{L}$	0	0	0	0	0	$-\frac{GJ}{L}$	0	0	4
				$\frac{4EI_2}{L}$	0	0	0	$\frac{6EI_2}{L^2}$	0	$\frac{2EI_2}{L}$	0	5
					$\frac{4EI_3}{L}$	0	$-\frac{6EI_3}{L^2}$	0	0	0	$\frac{2EI_3}{L}$	6
							$\frac{EA}{L}$	0	0	0	0	7
								$\frac{12EI_3}{L^3}$	0	0	$-\frac{6EI_3}{L^2}$	8
symmetric									$\frac{12EI_2}{L^3}$	0	$\frac{6EI_2}{L^2}$	9
										$\frac{GJ}{L}$	0	10
											$\frac{4EI_2}{L}$	11
											$\frac{4EI_3}{L}$	12

(4.5)

where

A = area of cross section

L = undeformed element length

$I_n$  = moment of inertia about the  $x_n$ -axis

$G$  = shear modulus of elasticity

$J$  = polar moment of inertia

#### 4.4 NONLINEAR STRAIN INCREMENTAL STIFFNESS MATRIX

Similarly, employing Eqs. (4.3), (4.4) and substituting Eqs. (2.36), (2.43) into Eq. (2.46) one obtains the nonlinear strain incremental stiffness matrix of the local element model with coordinate system in Fig. 2.3 [18]

$$k_{NL} = \int_V B_{NL}^T \tau B_{NL} dV = T$$

1	2	3	4	5	6	7	8	9	10	11	12	
0	0	0	0	0	0	0	0	0	0	0	0	1
$\frac{6}{5L}$	0	0	0	0	$\frac{1}{10}$	0	$\frac{-6}{5L}$	0	0	0	$\frac{1}{10}$	2
	$\frac{6}{5L}$	0	$\frac{-1}{10}$	0	0	0	0	$\frac{-6}{5L}$	0	$\frac{-1}{10}$	0	3
		0	0	0	0	0	0	0	0	0	0	4
			$\frac{2L}{15}$	0	0	0	0	$\frac{1}{10}$	0	$\frac{-1}{30}$	0	5
				$\frac{2L}{15}$	0	0	$\frac{-1}{10}$	0	0	0	$\frac{-1}{30}$	6
						0	0	0	0	0	0	7
symmetric							$\frac{6}{5L}$	0	0	0	$\frac{-1}{10}$	8
							$\frac{6}{5L}$	0	$\frac{1}{10}$	0	0	9
								0	0	0	0	10
									$\frac{2L}{15}$	0	0	11
										$\frac{2L}{15}$	0	12

(4.6)

where  $T = \sigma A$  is positive in tension.

The nonlinear strain incremental stiffness matrix  $k_{NL}$  is independent of elastic properties. It depends only on the element's geometry, displacement field, and the current state of stress level.

#### 4.5 LOCAL STRAIN INCREMENTAL STIFFNESS MATRIX

The local strain incremental stiffness matrix in Eq. (2.53) is

$$k = k_L + k_{NL} = \begin{array}{c} \left[ \begin{array}{cccccc|cccccc} 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 & 9 & 10 & 11 & 12 \\ T_1 & 0 & 0 & 0 & 0 & 0 & -T_1 & 0 & 0 & 0 & 0 & 0 \\ & T_2 & 0 & 0 & 0 & T_3 & 0 & -T_2 & 0 & 0 & 0 & T_3 \\ & & T_4 & 0 & T_5 & 0 & 0 & 0 & -T_4 & 0 & T_5 & 0 \\ & & & T_6 & 0 & 0 & 0 & 0 & 0 & -T_6 & 0 & 0 \\ & & & & T_7 & 0 & 0 & 0 & -T_5 & 0 & T_8 & 0 \\ & & & & & T_9 & 0 & -T_3 & 0 & 0 & 0 & T_{10} \\ \hline & & & & & & T_1 & 0 & 0 & 0 & 0 & 0 \\ & & & & & & & T_2 & 0 & 0 & 0 & -T_3 \\ & & & & & & & & T_4 & 0 & -T_5 & 0 \\ & & & & & & & & & T_6 & 0 & 0 \\ & & & & & & & & & & T_7 & 0 \\ & & & & & & & & & & & T_9 \end{array} \right] \begin{array}{l} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \\ 9 \\ 10 \\ 11 \\ 12 \end{array} \end{array} \right. \quad (4.7)$$

where

$$\begin{aligned}
T_1 &= EA / L \\
T_2 &= 12EI_3 / L^3 + 6T / 5L \\
T_3 &= 6EI_3 / L^2 + T / 10 \\
T_4 &= 12EI_2 / L^3 + 6T / 5L \\
T_5 &= -6EI_2 / L^2 - T / 10 \\
T_6 &= GJ / L \\
T_7 &= 4EI_2 / L + 2TL / 15 \\
T_8 &= 2EI_2 / L - TL / 30 \\
T_9 &= 4EI_3 / L + 2TL / 15 \\
T_{10} &= 2EI_3 / L - TL / 30
\end{aligned} \tag{4.8}$$

In Eqs. (4.8) the first terms are contributed by  $k_L$  and the second terms by  $k_{NL}$ .

#### 4.6 GLOBAL STRAIN INCREMENTAL STIFFNESS MATRIX

The global strain incremental stiffness matrix for element  $i$  may be obtained by the standard transformation in Eq. (2.50) [28]

$$K^i = \Lambda^T k \Lambda \tag{4.9}$$

where  $\Lambda$  is the transformation matrix defined in Eq. (2.54). Substituting Eqs. (4.7), (3.17) into Eq. (4.9), one can express  $K^i$  in terms of coefficient functions  $g_i$ ,  $i = 1, 2, \dots, 42$ . The following index matrix locates  $g_i$  in the stiffness matrix storing the subscripts and negative signs of  $g_i$  [28].

$$\text{INDEX} = \left[ \begin{array}{cccccc|cccccc} 1 & 2 & 3 & 4 & 5 & 6 & -1 & -2 & -3 & 7 & 8 & 9 \\ 2 & 10 & 11 & 12 & 13 & 14 & -2 & -10 & -11 & 15 & 16 & 17 \\ 3 & 11 & 18 & 19 & 20 & 21 & -3 & -11 & -18 & 22 & 23 & 24 \\ 4 & 12 & 19 & 25 & 26 & 27 & -4 & -12 & -19 & 28 & 29 & 30 \\ 5 & 13 & 20 & 26 & 31 & 32 & -5 & -13 & -20 & 29 & 33 & 34 \\ 6 & 14 & 21 & 27 & 32 & 35 & -6 & -14 & -21 & 30 & 34 & 36 \\ \hline -1 & -2 & -3 & -4 & -5 & -6 & 1 & 2 & 3 & -7 & -8 & -9 \\ -2 & -10 & -11 & -12 & -13 & -14 & 2 & 10 & 11 & -15 & -16 & -17 \\ -3 & -11 & -18 & -19 & -20 & -21 & 3 & 11 & 18 & -22 & -23 & -24 \\ 7 & 15 & 22 & 28 & 29 & 30 & -7 & -15 & -22 & 37 & 38 & 39 \\ 8 & 16 & 23 & 29 & 33 & 34 & -8 & -16 & -23 & 38 & 40 & 41 \\ 9 & 17 & 24 & 30 & 34 & 36 & -9 & -17 & -24 & 39 & 41 & 42 \end{array} \right]$$

(4.10)

the coefficient functions are

$$g_1 = T_1 c^2_1 + T_2 \lambda^2_{21} + T_4 \lambda^2_{31}$$

$$g_2 = T_1 c_1 c_2 + T_2 \lambda_{21} \lambda_{22} + T_4 \lambda_{31} \lambda_{32} \tag{4.11}$$

$$g_3 = T_1 c_1 c_3 + T_2 \lambda_{21} \lambda_{23} + T_4 \lambda_{31} \lambda_{33}$$

$$\begin{aligned}
g_4 &= T_3 \lambda_{21} \lambda_{31} + T_5 \lambda_{31} \lambda_{21} \\
g_5 &= T_3 \lambda_{21} \lambda_{32} + T_5 \lambda_{31} \lambda_{22} \\
g_6 &= T_3 \lambda_{21} \lambda_{33} + T_5 \lambda_{31} \lambda_{23} \\
g_7 &= T_3 \lambda_{21} \lambda_{31} + T_5 \lambda_{31} \lambda_{21} \\
g_8 &= T_3 \lambda_{21} \lambda_{32} + T_5 \lambda_{31} \lambda_{22} \\
g_9 &= T_3 \lambda_{21} \lambda_{33} + T_5 \lambda_{31} \lambda_{23} \\
g_{10} &= T_1 c_2^2 + T_2 \lambda_{22}^2 + T_4 \lambda_{32}^2 \\
g_{11} &= T_1 c_2 c_3 + T_2 \lambda_{22} \lambda_{23} + T_4 \lambda_{32} \lambda_{33} \\
g_{12} &= T_3 \lambda_{22} \lambda_{31} + T_5 \lambda_{32} \lambda_{21} \\
g_{13} &= T_3 \lambda_{22} \lambda_{32} + T_5 \lambda_{32} \lambda_{22} \\
g_{14} &= T_3 \lambda_{22} \lambda_{33} + T_5 \lambda_{32} \lambda_{23} \\
g_{15} &= T_3 \lambda_{22} \lambda_{31} + T_5 \lambda_{32} \lambda_{21} \\
g_{16} &= T_3 \lambda_{22} \lambda_{32} + T_5 \lambda_{32} \lambda_{22} \\
g_{17} &= T_3 \lambda_{22} \lambda_{33} + T_5 \lambda_{32} \lambda_{23} \\
g_{18} &= T_1 c_3^2 + T_2 \lambda_{23}^2 + T_4 \lambda_{33}^2 \\
g_{19} &= T_3 \lambda_{23} \lambda_{31} + T_5 \lambda_{33} \lambda_{21} \\
g_{20} &= T_3 \lambda_{23} \lambda_{32} + T_5 \lambda_{33} \lambda_{22} \\
g_{21} &= T_3 \lambda_{23} \lambda_{33} + T_5 \lambda_{33} \lambda_{23} \\
g_{22} &= T_3 \lambda_{23} \lambda_{31} + T_5 \lambda_{33} \lambda_{21}
\end{aligned} \tag{4.11}$$



$$\begin{aligned}
g_{23} &= T_3 \lambda_{23} \lambda_{32} + T_5 \lambda_{33} \lambda_{22} \\
g_{24} &= T_3 \lambda_{23} \lambda_{33} + T_5 \lambda_{33} \lambda_{23} \\
g_{25} &= T_6 c_1^2 + T_9 \lambda_{31}^2 + T_7 \lambda_{21}^2 \\
g_{26} &= T_6 c_1 c_2 + T_9 \lambda_{31} \lambda_{32} + T_7 \lambda_{21} \lambda_{22} \\
g_{27} &= T_6 c_1 c_3 + T_9 \lambda_{31} \lambda_{33} + T_7 \lambda_{21} \lambda_{23} \\
g_{28} &= -T_6 c_1^2 + T_{10} \lambda_{31}^2 + T_8 \lambda_{21}^2 \\
g_{29} &= -T_6 c_1 c_2 + T_{10} \lambda_{31} \lambda_{32} + T_8 \lambda_{21} \lambda_{22} \\
g_{30} &= -T_6 c_1 c_3 + T_{10} \lambda_{31} \lambda_{33} + T_8 \lambda_{21} \lambda_{23} \\
g_{31} &= T_6 c_2^2 + T_9 \lambda_{32}^2 + T_7 \lambda_{22}^2 \\
g_{32} &= T_6 c_2 c_3 + T_9 \lambda_{32} \lambda_{33} + T_7 \lambda_{22} \lambda_{23} \\
g_{33} &= -T_6 c_2^2 + T_{10} \lambda_{32}^2 + T_8 \lambda_{22}^2 \\
g_{34} &= -T_6 c_2 c_3 + T_{10} \lambda_{32} \lambda_{33} + T_8 \lambda_{22} \lambda_{23} \\
g_{35} &= T_6 c_3^2 + T_9 \lambda_{33}^2 + T_7 \lambda_{23}^2 \\
g_{36} &= -T_6 c_3^2 + T_{10} \lambda_{33}^2 + T_8 \lambda_{23}^2 \\
g_{37} &= T_6 c_1^2 + T_9 \lambda_{31}^2 + T_7 \lambda_{21}^2 \\
g_{38} &= T_6 c_1 c_2 + T_9 \lambda_{31} \lambda_{32} + T_7 \lambda_{21} \lambda_{22} \\
g_{39} &= T_6 c_1 c_3 + T_9 \lambda_{31} \lambda_{33} + T_7 \lambda_{21} \lambda_{23} \\
g_{40} &= T_6 c_2^2 + T_9 \lambda_{32}^2 + T_7 \lambda_{22}^2 \\
g_{41} &= T_6 c_2 c_3 + T_9 \lambda_{32} \lambda_{33} + T_7 \lambda_{22} \lambda_{23} \\
g_{42} &= T_6 c_3^2 + T_9 \lambda_{33}^2 + T_7 \lambda_{23}^2
\end{aligned} \tag{4.11}$$

#### 4.7 LOCAL ELEMENT FORCES

The relative deformation displacements and forces can be defined by Figs. 4.2 and 4.3, corresponding to the local  $x_1-x_2$  and  $x_1-x_3$  planes, respectively.

The deformation forces in the  $x_1-x_2$  plane (Fig. 4.2) can be obtained by the equilibrium equation

$$p' = \hat{k}'_s e' \quad (4.12)$$

where

$$p' = [ T \quad M_{13} \quad M_{23} ]^T \quad (4.13)$$

$$e' = [ e_1 \quad e_{13} \quad e_{23} ]^T \quad (4.14)$$

$p'$  is the deformation force vector in the  $x_1-x_2$  plane,  $\hat{k}'_s$  is the local secant stiffness matrix of the element in the  $x_1-x_2$  plane,  $e'$  is the deformation displacement vector in the  $x_1-x_2$  plane,  $T$  is the axial force (positive in tension),  $e_{mn}$  are the relative end rotations (Figs. 4.2 and 4.3),  $M_{mn}$  are the bending moments corresponding to  $e_{mn}$ , and  $e_1$  is the relative axial elongation.

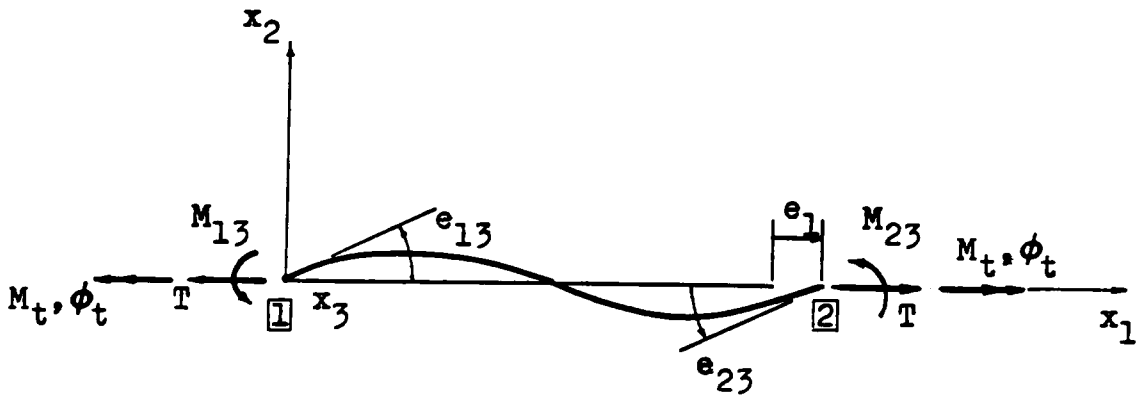


Fig. 4.2 Deformation Displacements and Forces  
in  $x_1$ - $x_2$  Plane

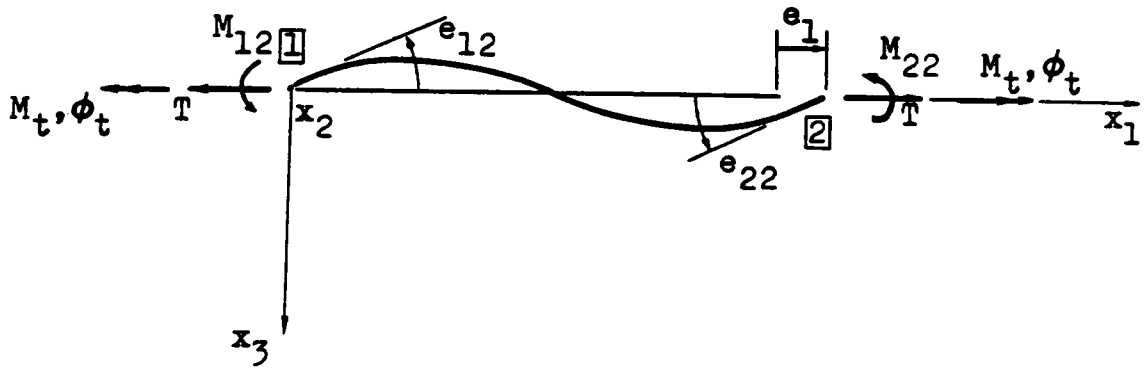


Fig. 4.3 Deformation Displacements and Forces  
in  $x_1$ - $x_3$  Plane

Katzenberger [48] derived the local element secant stiffness matrix and the deformation forces for the plane frame element. The symmetric form of the secant stiffness matrix in the  $x_1$ - $x_2$  plane is

$$\hat{k}_s = \begin{matrix} & \begin{matrix} 1 & & 2 & & & & 3 \end{matrix} \\ \begin{matrix} 1 \\ 2 \\ 3 \end{matrix} & \left[ \begin{array}{cccccc} \frac{EA}{L} & & & & & & \\ & \frac{EA}{60} (4e_{13} - e_{23}) & & & & & \\ & & & & & & \frac{EA}{60} (4e_{23} - e_{13}) \\ \frac{4EI_3}{L} + \frac{EA}{15} e_1 + \frac{EAL}{420} (12e_{13}^2 - 3e_{13}e_{23} + e_{23}^2) & & & & & & \frac{2EI_3}{L} - \frac{EA}{60} e_1 + \frac{EAL}{840} (-3e_{13}^2 + 4e_{13}e_{23} - 3e_{23}^2) \\ \text{symm} & & & & & & \frac{4EI_3}{L} + \frac{EA}{15} e_1 + \frac{EAL}{420} (e_{13}^2 - 3e_{13}e_{23} + 12e_{23}^2) \end{array} \right] \end{matrix} \begin{matrix} 1 \\ 2 \\ 3 \end{matrix} \quad (4.15)$$

Similarly, the deformation forces in the  $x_1$ - $x_3$  plane (Fig. 4.3) are

$$p'' = \hat{k}_s'' e'' \quad (4.16)$$

where

$$p'' = [ T \quad M_{12} \quad M_{22} ]^T \quad (4.17)$$

$$e'' = [ e_1 \quad e_{12} \quad e_{22} ]^T \quad (4.18)$$

$p''$  is the deformation force vector in  $x_1$ - $x_3$  plane,  $\hat{k}_s''$  is the local secant stiffness matrix of the element in the  $x_1$ - $x_3$  plane, and  $e''$  is the deformation displacement vector in the  $x_1$ - $x_3$  plane.

Comparing Fig. 4.2 with Fig. 4.3, we obtain the local element secant stiffness matrix in  $x_1-x_3$  plane, similar to Eq. (4.15):

$$\hat{k}_s = \begin{bmatrix} 1 & & & & \\ & 4 & & & \\ & & & & 5 \\ & & & & & & & & & \\ \text{symm} & & & & & & & & & \end{bmatrix} \begin{matrix} \frac{EA}{L} \\ \frac{EA}{60}(4e_{12} - e_{22}) \\ \frac{4EI_2}{L} + \frac{EA}{15}e_1 + \frac{EAL}{420}(12e_{12}^2 - 3e_{12}e_{22} + e_{22}^2) \\ \frac{2EI_2}{L} - \frac{EA}{60}e_1 + \frac{EAL}{840}(-3e_{12}^2 + 4e_{12}e_{22} - 3e_{22}^2) \\ \frac{4EI_2}{L} + \frac{EA}{15}e_1 + \frac{EAL}{420}(e_{12}^2 - 3e_{12}e_{22} + 12e_{22}^2) \end{matrix} \begin{matrix} 1 \\ 4 \\ 5 \end{matrix} \quad (4.19)$$

Combining Eqs. (4.12) and (4.16), one obtains the deformation forces of the space frame element

$$p = \hat{k}_s e \quad (4.20)$$

where

$$p = [ T \quad M_{13} \quad M_{23} \quad M_{12} \quad M_{22} ]^T \quad (4.21)$$

$$e = [ e_1 \quad e_{13} \quad e_{23} \quad e_{12} \quad e_{22} ]^T \quad (4.22)$$

$p$ ,  $e$  are the deformation force and displacement vectors of the space frame element, respectively.

The relative element forces and displacements are as shown in Fig. 4.4. The superposition of Eqs. (4.15) and (4.19) in accordance with the sequence in Eqs. (4.21) and

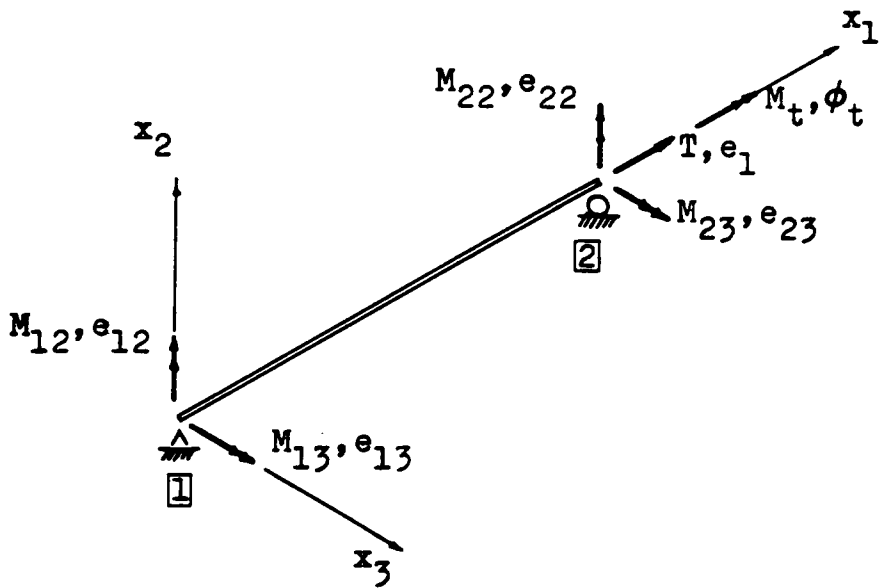


Fig. 4.4 Relative Element Forces and Displacements  
in Space Frame Element

(4.22) yields the local element secant stiffness matrix of space frame

$$\hat{k}_s = \begin{array}{c} \begin{array}{ccc} 1 & & 2 & & 3 \\ \left[ \begin{array}{ccc} EA/L & & \frac{EA}{60}(4e_{13} - e_{23}) & & \frac{EA}{60}(4e_{23} - e_{13}) \\ \frac{EA}{60}(4e_{13} - e_{23}) & \frac{4EI_3}{L} + \frac{EA}{15} e_1 + \frac{EAL}{420}(12e_{13}^2 - 3e_{13}e_{23} + e_{23}^2) & \frac{2EI_3}{L} - \frac{EA}{60}e_1 + \frac{EAL}{840}(-3e_{13}^2 + 4e_{13}e_{23} - 3e_{23}^2) \\ \frac{EA}{60}(4e_{23} - e_{13}) & \frac{2EI_3}{L} - \frac{EA}{60}e_1 + \frac{EAL}{840}(-3e_{13}^2 + 4e_{13}e_{23} - 3e_{23}^2) & \frac{4EI_3}{L} + \frac{EA}{15}e_1 + \frac{EAL}{420}(e_{13}^2 - 3e_{13}e_{23} + 12e_{23}^2) \\ EA/60(4e_{12} - e_{22}) & 0 & 0 \\ EA/60(4e_{22} - e_{12}) & 0 & 0 \end{array} \right] \\ \\ 4 & & 5 \\ \left[ \begin{array}{cc} \frac{EA}{60}(4e_{12} - e_{22}) & \frac{EA}{60}(4e_{22} - e_{12}) \\ 0 & 0 \\ 0 & 0 \\ \frac{4EI_2}{L} + \frac{EA}{15}e_1 + \frac{EAL}{420}(12e_{12}^2 - 3e_{12}e_{22} + e_{22}^2) & \frac{2EI_2}{L} - \frac{EA}{60}e_1 + \frac{EAL}{840}(-3e_{12}^2 + 4e_{12}e_{22} - 3e_{22}^2) \\ \frac{2EI_2}{L} - \frac{EA}{60}e_1 + \frac{EAL}{840}(-3e_{12}^2 + 4e_{12}e_{22} - 3e_{22}^2) & \frac{4EI_2}{L} + \frac{EA}{15}e_1 + \frac{EAL}{420}(e_{12}^2 - 3e_{12}e_{22} + 12e_{22}^2) \end{array} \right] \end{array} \end{array} \quad (4.23)$$

From equilibrium conditions one obtains the local element forces corresponding to Fig. 2.3:

$$f_1 = -T$$



$$\begin{aligned}
f_2 &= (M_{13} + M_{23}) / L \\
f_3 &= -(M_{12} + M_{22}) / L \\
f_5 &= M_{12} \\
f_6 &= M_{13} \\
f_7 &= -f_1 \\
f_8 &= -f_2 \\
f_9 &= -f_3 \\
f_{11} &= M_{22} \\
f_{12} &= M_{23}
\end{aligned} \tag{4.24}$$

The torsional forces are treated independently and can be obtained from the linear theory [28]

$$\begin{aligned}
f_4 &= -(GJ/L) \phi_t \\
f_{10} &= -f_4
\end{aligned} \tag{4.25}$$

where  $\phi_t$  is the relative rotation of element ends about the element axis.

#### 4.8 GLOBAL ELEMENT FORCES

The global element forces for element  $i$  can be obtained from Eq. (2.51):

$$F^i = \Lambda^T f \tag{4.26}$$

where  $\Lambda$  is defined in Eq. (3.17). The components of  $F^i$  are

$$F_1 = c_1 f_1 + \lambda_{21} f_2 + \lambda_{31} f_3$$

$$F_2 = c_2 f_1 + \lambda_{22} f_2 + \lambda_{32} f_3$$

$$F_3 = c_3 f_1 + \lambda_{23} f_2 + \lambda_{33} f_3$$

$$F_4 = c_1 f_4 + \lambda_{21} f_5 + \lambda_{31} f_6$$

$$F_5 = c_2 f_4 + \lambda_{22} f_5 + \lambda_{32} f_6$$

$$F_6 = c_3 f_4 + \lambda_{23} f_5 + \lambda_{33} f_6$$

$$F_7 = c_1 f_7 + \lambda_{21} f_8 + \lambda_{31} f_9$$

$$F_8 = c_2 f_7 + \lambda_{22} f_8 + \lambda_{32} f_9$$

$$F_9 = c_3 f_7 + \lambda_{23} f_8 + \lambda_{33} f_9$$

$$F_{10} = c_1 f_{10} + \lambda_{21} f_{11} + \lambda_{31} f_{12}$$

$$F_{11} = c_2 f_{10} + \lambda_{22} f_{11} + \lambda_{32} f_{12}$$

$$F_{12} = c_3 f_{10} + \lambda_{23} f_{11} + \lambda_{33} f_{12}$$

(4.27)

Chapter V  
BEAM-COLUMN MODEL

5.1 INTRODUCTION

This chapter presents the three-dimensional beam-column model in U.L. formulation [2],[4], which is based on the conventional beam-column theory [23]. Because using the actual solution to the differential equation, the beam-column model can trace equilibrium path accurately. In this model the behavior of the element is referred to the convected coordinate system, and then a transformation is applied from a local to a global coordinate system.

The rotations and translations of the nodes are considered to be arbitrarily large, but the relative deformations of the element are assumed to be small such that the conventional beam-column theory can be applied. In the beam-column model, the effect of length shortening due to bending is considered, and the external loads are supposed to be applied at the nodes only.

## 5.2 ELEMENT END FORCE-DEFORMATION RELATIONS

The relationships between element-end forces and deformations (Fig. 5.1) based on the conventional beam-column theory are [2],[3],[23]:

$$M_{1n} = EI_n / L (c_{1n}e_{1n} + c_{2n}e_{2n}) \quad (5.1)$$

$$M_{2n} = EI_n / L (c_{2n}e_{1n} + c_{1n}e_{2n}) \quad (5.2)$$

$$M_t = (GJ / L) \phi_t \quad (5.3)$$

$$P = EA (u / L - c_{b2} - c_{b3}) \quad (5.4)$$

in which

$e_{mn}$  = relative end rotations (Fig. 5.1); the first subscript refers to the node where the angle is measured (1 for left node and 2 for right node); the second subscript indicates the axis about which the rotation takes place

$M_{mn}$  = bending moment corresponding to  $e_{mn}$

$M_t$  = torque

$P$  = axial force, positive in compression

$c_{mn}$  = stability functions [1],[4] associated with bending moments about the  $x_n$ -axis (Appendix B)

$L$  = element length

$u$  = relative axial displacement

$c_{bn}$  is length correction factor resulting from the flexural deformations about the  $x_n$ -axis [1],[2] defined as

$$c_{bn} = b_{1n} (e_{1n} + e_{2n})^2 + b_{2n} (e_{1n} - e_{2n})^2 \quad (5.5)$$

$b_{1n}$ ,  $b_{2n}$  are the bowing functions given by

$$b_{1n} = (c_{1n} + c_{2n})(c_{2n} - 2) / 8 \pi^2 p_n$$

$$b_{2n} = c_{2n} / 8(c_{1n} + c_{2n}) \quad (5.6)$$

in which

$$p_n = P / P_{En} \quad (5.7)$$

$p_n$  is the dimensionless axial force parameter and

$$P_{En} = \pi^2 E I_n / L^2 \quad (5.8)$$

Thus

$$p_n = PL^2 / \pi^2 E I_n \quad (5.9)$$

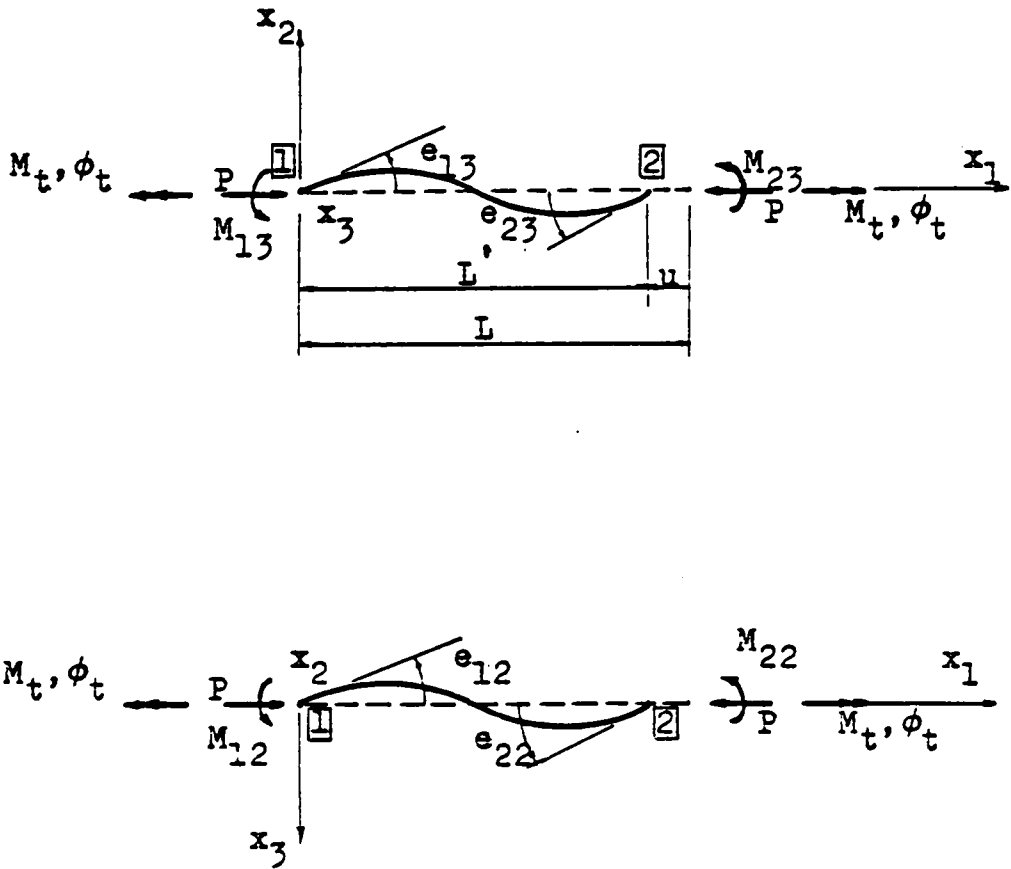


Fig. 5.1 Relative Member End Deformations and Associated Forces in Space Frame

### 5.3 TANGENT STIFFNESS MATRIX FOR RELATIVE DEFORMATIONS

The relation between incremental vectors, the relative end forces and deformations is

$$\Delta f_{bc} = \hat{k} \Delta e \quad (5.10)$$

where

$$\begin{aligned} \Delta f_{bc} &= [ \Delta M_{13} \quad \Delta M_{23} \quad \Delta M_{12} \quad \Delta M_{22} \quad \Delta M_t \quad \Delta PL ]^T \\ \Delta e &= [ \Delta e_{13} \quad \Delta e_{23} \quad \Delta e_{12} \quad \Delta e_{22} \quad \Delta \phi_t \quad \Delta u/L ]^T \end{aligned} \quad (5.11)$$

$\hat{k}$  is the tangent stiffness matrix for relative deformations of the beam-column model and has been derived by Oran [2]:

$$\hat{k} = \frac{EI}{L} \begin{pmatrix} \epsilon_3^c c_{13}^+ \frac{G_{13}^2}{\pi^2 H} & \epsilon_3^c c_{23}^+ \frac{G_{13} G_{23}}{\pi^2 H} & 0 & 0 & 0 & \frac{G_{13}}{H} \\ \epsilon_3^c c_{23}^+ \frac{G_{13} G_{23}}{\pi^2 H} & \epsilon_3^c c_{13}^+ \frac{G_{23}^2}{\pi^2 H} & 0 & 0 & 0 & \frac{G_{23}}{H} \\ 0 & 0 & \epsilon_2^c c_{12}^+ \frac{G_{12}^2}{\pi^2 H} & \epsilon_2^c c_{22}^+ \frac{G_{12} G_{22}}{\pi^2 H} & 0 & \frac{G_{12}}{H} \\ 0 & 0 & \epsilon_2^c c_{22}^+ \frac{G_{12} G_{22}}{\pi^2 H} & \epsilon_2^c c_{12}^+ \frac{G_{22}^2}{\pi^2 H} & 0 & \frac{G_{22}}{H} \\ 0 & 0 & 0 & 0 & \eta & 0 \\ \frac{G_{13}}{H} & \frac{G_{23}}{H} & \frac{G_{12}}{H} & \frac{G_{22}}{H} & 0 & \frac{\pi^2}{H} \end{pmatrix} \quad (5.12)$$

where

$$G_{1n} = c'_{1n} e_{1n} + c'_{2n} e_{2n} \quad (5.13)$$

$$G_{2n} = c'_{2n} e_{1n} + c'_{1n} e_{2n} \quad (5.14)$$

$$H = \pi^2 / \zeta^2 + \sum_{n=1,2} 1 / \xi_n [b'_{1n} (e_{1n} + e_{2n})^2 + b'_{2n} (e_{1n} - e_{2n})^2] \quad (5.15)$$

$$\eta = GJ/EI \quad (5.16)$$

$$\zeta = L / (I/A)^{1/2} \quad (5.17)$$

$$\xi_n = I_n / I \quad (5.18)$$

$I$  is the reference moment of inertia, a prime superscript on  $c_{mn}$  or  $b_{mn}$  indicates one differentiation with respect to  $P_n$ .

#### 5.4 LOCAL ELEMENT TANGENT STIFFNESS MATRIX

From Fig. 2.3 and Fig. 5.2 the relationship between local element forces  $f$  and relative local element forces  $f_{bc}$  is

$$f = C^T f_{bc} \quad (5.19)$$

where

$$f = [ f_1 \ f_2 \ f_3 \ \dots \ f_{12} ]^T \quad (5.20)$$



$$f_{bc} = [ M_{13} \ M_{23} \ M_{12} \ M_{22} \ M_t \ PL ]^T \quad (5.21)$$

and C is the local instantaneous transformation matrix defined as

$$C = \begin{pmatrix} 0 & 1/L' & 0 & 0 & 0 & 1 & 0 & -1/L' & 0 & 0 & 0 & 0 \\ 0 & 1/L' & 0 & 0 & 0 & 0 & 0 & -1/L' & 0 & 0 & 0 & 1 \\ 0 & 0 & -1/L' & 0 & 1 & 0 & 0 & 0 & 1/L' & 0 & 0 & 0 \\ 0 & 0 & -1/L' & 0 & 0 & 0 & 0 & 0 & 1/L' & 0 & 1 & 0 \\ 0 & 0 & 0 & -1 & 0 & 0 & 0 & 0 & 0 & 1 & 0 & 0 \\ 1/L' & 0 & 0 & 0 & 0 & 0 & -1/L' & 0 & 0 & 0 & 0 & 0 \end{pmatrix} \quad (5.22)$$

$L'$  is the deformed length of element given by

$$L' = L(1+\delta) \quad (5.23)$$

The length correction factor  $\delta$  is a function of  $d$ , since  $\delta$  is small in comparison with unity,  $L'$  may be approximated by  $L$ ; i.e., one can set  $\delta = 0$ . However,  $\Delta\delta$  which is a function of  $\Delta d$  must be considered to determine  $\Delta C$  [2], [29].

The incremental form of Eq. (5.19) is

$$\Delta f = C^T \Delta f_{bc} + \Delta C^T f_{bc} \quad (5.24)$$

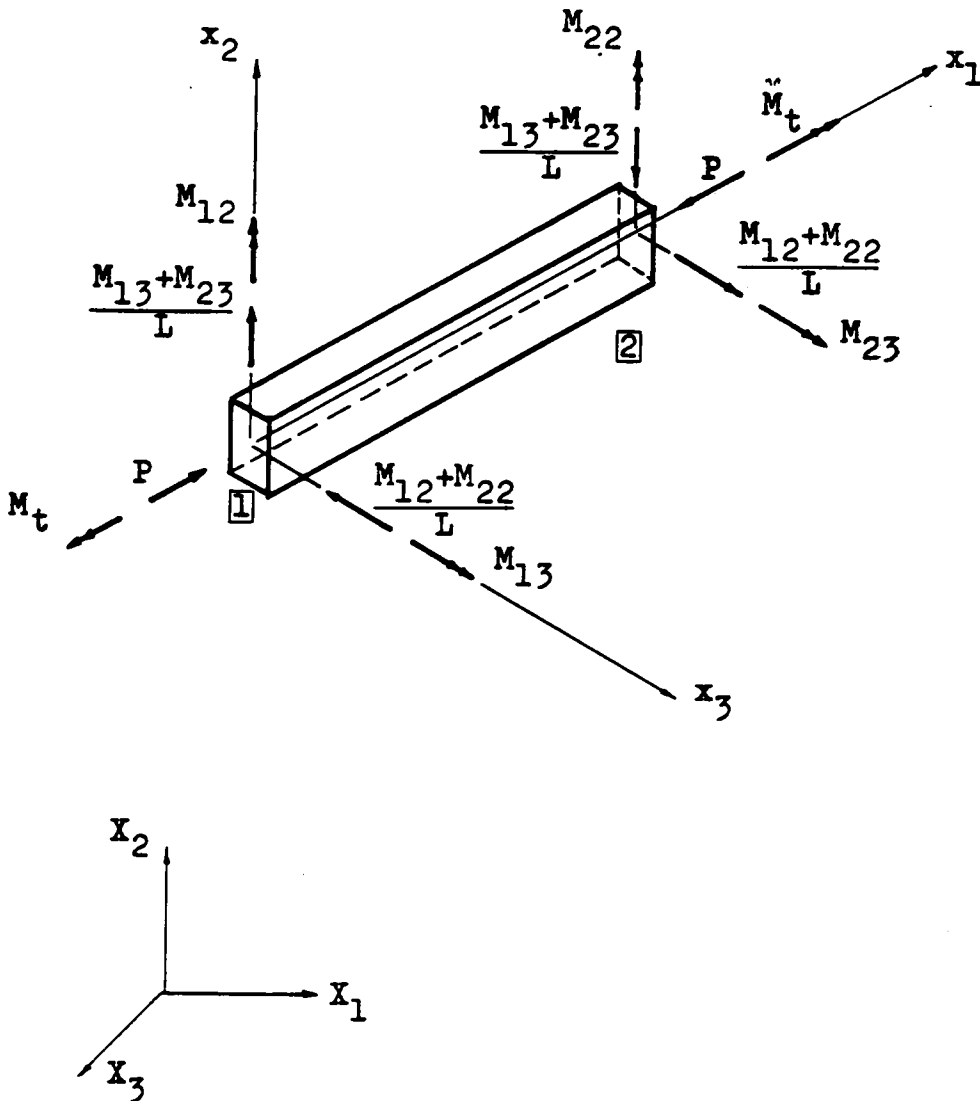


Fig. 5.2 Basic Local Element Forces Associated with Oran's Beam-Column Model in Space Frame

where  $\Delta C$  is the change in  $C$  resulting from  $\Delta d$  (see Appendix C)

By contragredience from Eq. (5.19)

$$\Delta e = C \Delta d \quad (5.25)$$

Substituting Eqs. (5.10), (5.25), (C.1) into Eq. (5.24), one obtain

$$\Delta f = k \Delta d \quad (5.26)$$

where  $k$  is the local tangent stiffness matrix of element defined as

$$k = k_L + k_{NL} \quad (5.27)$$

where

$$k_L = C^T \hat{k} C \quad (5.28)$$

$$k_{NL} = \sum_{i=1}^6 f_{bc_i} \bar{g}^{(i)} \quad (5.29)$$

$k_L$ ,  $k_{NL}$  are the linear and nonlinear tangent stiffness matrix respectively,  $\bar{g}^{(i)}$  are the geometric matarices defined in appendix C.

Substituting Eqs. (5.12), (5.22) into Eq. (5.28), one can express the  $k_L$  in terms of coefficient functions  $g'_i$ ,  $i = 1, \dots, 20$ :

$$\mathbf{k}_L = \begin{array}{c} \left[ \begin{array}{cccccc|cccccc} 1 & 2 & 3 & 4 & 5 & 6 & 7 & 8 & 9 & 10 & 11 & 12 \\ g'_1 & g'_2 & g'_3 & 0 & g'_4 & g'_5 & -g'_1 & -g'_2 & -g'_3 & 0 & g'_6 & g'_7 \\ g'_2 & g'_8 & 0 & 0 & 0 & g'_9 & -g'_2 & -g'_8 & 0 & 0 & 0 & g'_{10} \\ g'_3 & 0 & g'_{11} & 0 & g'_{12} & 0 & -g'_3 & 0 & -g'_{11} & 0 & g'_{13} & 0 \\ 0 & 0 & 0 & g'_{14} & 0 & 0 & 0 & 0 & 0 & -g'_{14} & 0 & 0 \\ g'_4 & 0 & g'_{12} & 0 & g'_{15} & 0 & -g'_4 & 0 & -g'_{12} & 0 & g'_{16} & 0 \\ g'_5 & g'_9 & 0 & 0 & 0 & g'_{17} & -g'_5 & -g'_9 & 0 & 0 & 0 & g'_{18} \\ \hline -g'_1 & -g'_2 & -g'_3 & 0 & -g'_4 & -g'_5 & g'_1 & g'_2 & g'_3 & 0 & -g'_6 & -g'_7 \\ -g'_2 & -g'_8 & 0 & 0 & 0 & -g'_9 & g'_2 & g'_8 & 0 & 0 & 0 & -g'_{10} \\ -g'_3 & 0 & -g'_{11} & 0 & -g'_{12} & 0 & g'_3 & 0 & g'_{11} & 0 & -g'_{13} & 0 \\ 0 & 0 & 0 & -g'_{14} & 0 & 0 & 0 & 0 & 0 & g'_{14} & 0 & 0 \\ g'_6 & 0 & g'_{13} & 0 & g'_{16} & 0 & -g'_6 & 0 & -g'_{13} & 0 & g'_{19} & 0 \\ g'_7 & g'_{10} & 0 & 0 & 0 & g'_{18} & -g'_7 & -g'_{10} & 0 & 0 & 0 & g'_{20} \end{array} \right] \begin{array}{l} 1 \\ 2 \\ 3 \\ 4 \\ 5 \\ 6 \\ 7 \\ 8 \\ 9 \\ 10 \\ 11 \\ 12 \end{array} \end{array}$$

(5.30)

where

$$g'_1 = \alpha ( \pi^2 / L^2 H )$$

$$g'_2 = \alpha ( G_{13} + G_{23} ) / L^2 H$$

$$g'_3 = \alpha ( -G_{12} - G_{22} ) / L^2 H$$

$$g'_4 = \alpha ( G_{12} / L H )$$

$$g'_5 = \alpha ( G_{13} / L H )$$

$$g'_6 = \alpha ( G_{22} / L H )$$

$$g'_7 = \alpha ( G_{23} / L H )$$

$$g'_8 = \alpha [ 2(\xi_3 c_{13} + \xi_3 c_{23}) / L^2 + (G_{13} + G_{23})^2 / L^2 \pi^2 H ]$$

$$g'_9 = \alpha [ \xi_3 (c_{13} + c_{23}) / L + (G_{13}^2 + G_{13} G_{23}) / L \pi^2 H ]$$

$$g'_{10} = \alpha [ \xi_3 (c_{13} + c_{23}) / L + (G_{13} G_{23} + G_{23}^2) / L \pi^2 H ]$$

$$g'_{11} = \alpha [ 2(\xi_2 c_{12} + \xi_2 c_{22}) / L^2 + (G_{12} + G_{22})^2 / L^2 \pi^2 H ]$$

$$g'_{12} = \alpha [ -\xi_2 (c_{12} + c_{22}) / L - (G_{12}^2 + G_{12} G_{22}) / L \pi^2 H ]$$

$$g'_{13} = \alpha [ -\xi_2 (c_{12} + c_{22}) / L - (G_{12} G_{22} + G_{22}^2) / L \pi^2 H ]$$

$$g'_{14} = \alpha \eta$$

$$g'_{15} = \alpha ( \xi_2 c_{12} + G_{12}^2 / \pi^2 H )$$

$$g'_{16} = \alpha ( \xi_2 c_{22} + G_{12} G_{22} / \pi^2 H )$$

$$g'_{17} = \alpha ( \xi_3 c_{13} + G_{13}^2 / \pi^2 H )$$

$$g'_{18} = \alpha ( \xi_3 c_{23} + G_{13} G_{23} / \pi^2 H )$$

$$g'_{19} = \alpha ( \xi_2 c_{12} + G_{22}^2 / \pi^2 H )$$

$$g'_{20} = \alpha ( \xi_3 c_{13} + G_{23}^2 / \pi^2 H )$$

and

$$\alpha = EI/L$$

(5.31)

### 5.5 GLOBAL ELEMENT TANGENT STIFFNESS MATRIX

From Eq. (2.59)

$$\Delta F = \Lambda^T \Delta f \quad (5.32)$$

Substituting Eqs. (5.26), (2.58) into Eq. (5.32) yields

$$\Delta F = K^i \Delta D \quad (5.33)$$

where

$$K^i = \Lambda^T k \Lambda \quad (5.34)$$

is the global strain incremental stiffness matrix of element  $i$ . Substituting Eqs. (5.27), (5.28), (5.29) into Eq. (5.34) yields

$$K^i = K_L^i + K_{NL}^i \quad (5.35)$$

where

$$K_L^i = \Lambda^T k_L \Lambda \quad (5.36)$$

$$K_{NL}^i = \Lambda^T k_{NL} \Lambda \quad (5.37)$$

Substituting Eq. (5.28) into Eq. (5.37) yields

$$K_{NL}^i = \sum_{i=1}^6 f_{bc_i} g^{(i)} \quad (5.38)$$

$$= (M_{13} + M_{23})g^{(1)} + (M_{12} + M_{22})g^{(3)} + PL g^{(6)} \quad (5.39)$$

where

$$g^{(i)} = \Lambda^T \bar{g}^{(i)} \Lambda \quad (5.40)$$

$\bar{g}^{(i)}$  and  $g^{(i)}$  are the geometric matrices defined in the appendices C and D respectively.  $K_L^i$ ,  $K_{NL}^i$  are the linear and nonlinear tangent stiffness matrices in the global coordinates respectively,  $f_{bc_i}$  is the  $i^{\text{th}}$  component of  $f_{bc}$ .

Sustituting Eqs. (2.54), (5.30) into Eq. (5.36), one can express  $K_L^i$ , similar to Eq. (4.10), in terms of coefficient functions  $g_i$ ,  $i = 1, 2, \dots, 42$ . The index matrix of  $g_i$  is [28]

$$\text{INDEX} = \left[ \begin{array}{cccccc|cccccc} 1 & 2 & 3 & 4 & 5 & 6 & -1 & -2 & -3 & 7 & 8 & 9 \\ 2 & 10 & 11 & 12 & 13 & 14 & -2 & -10 & -11 & 15 & 16 & 17 \\ 3 & 11 & 18 & 19 & 20 & 21 & -3 & -11 & -18 & 22 & 23 & 24 \\ 4 & 12 & 19 & 25 & 26 & 27 & -4 & -12 & -19 & 28 & 29 & 30 \\ 5 & 13 & 20 & 26 & 31 & 32 & -5 & -13 & -20 & 29 & 33 & 34 \\ 6 & 14 & 21 & 27 & 32 & 35 & -6 & -14 & -21 & 30 & 34 & 36 \\ \hline -1 & -2 & -3 & -4 & -5 & -6 & 1 & 2 & 3 & -7 & -8 & -9 \\ -2 & -10 & -11 & -12 & -13 & -14 & 2 & 10 & 11 & -15 & -16 & -17 \\ -3 & -11 & -18 & -19 & -20 & -21 & 3 & 11 & 18 & -22 & -23 & -24 \\ 7 & 15 & 22 & 28 & 29 & 30 & -7 & -15 & -22 & 37 & 38 & 39 \\ 8 & 16 & 23 & 29 & 33 & 34 & -8 & -16 & -23 & 38 & 40 & 41 \\ 9 & 17 & 24 & 30 & 34 & 36 & -9 & -17 & -24 & 39 & 41 & 42 \end{array} \right] \quad (5.41)$$

which is consistent with Eq. (4.10) and the coefficient functions are

$$\begin{aligned}
 g_1 &= \alpha(c_1^2 g_1' + 2c_1 \lambda_{21} g_2' + 2c_1 \lambda_{31} g_3' + \lambda_{21}^2 g_8' + \lambda_{31}^2 g_{11}') \\
 g_2 &= \alpha[c_1 c_2 g_1' + (c_1 \lambda_{22} + c_2 \lambda_{21})g_2' + (c_1 \lambda_{32} + c_2 \lambda_{31})g_3' \\
 &\quad + \lambda_{21} \lambda_{22} g_8' + \lambda_{31} \lambda_{32} g_{11}'] \\
 g_3 &= \alpha[c_1 c_3 g_1' + (c_1 \lambda_{23} + c_3 \lambda_{21})g_2' + (c_1 \lambda_{33} + c_3 \lambda_{31})g_3' \\
 &\quad + \lambda_{21} \lambda_{23} g_8' + \lambda_{31} \lambda_{33} g_{11}'] \tag{5.42}
 \end{aligned}$$

$$\begin{aligned}
 g_4 &= \alpha(c_1 \lambda_{21} g_4' + c_1 \lambda_{31} g_5' + \lambda_{21} \lambda_{31} g_9' + \lambda_{31} \lambda_{21} g_{12}') \\
 g_5 &= \alpha(c_1 \lambda_{22} g_4' + c_1 \lambda_{32} g_5' + \lambda_{21} \lambda_{32} g_9' + \lambda_{31} \lambda_{22} g_{12}') \\
 g_6 &= \alpha(c_1 \lambda_{23} g_4' + c_1 \lambda_{33} g_5' + \lambda_{21} \lambda_{33} g_9' + \lambda_{31} \lambda_{23} g_{12}') \\
 g_7 &= \alpha(c_1 \lambda_{21} g_6' + c_1 \lambda_{31} g_7' + \lambda_{21} \lambda_{31} g_{10}' + \lambda_{31} \lambda_{21} g_{13}') \\
 g_8 &= \alpha(c_1 \lambda_{22} g_6' + c_1 \lambda_{32} g_7' + \lambda_{21} \lambda_{32} g_{10}' + \lambda_{31} \lambda_{22} g_{13}') \\
 g_9 &= \alpha(c_1 \lambda_{23} g_6' + c_1 \lambda_{33} g_7' + \lambda_{21} \lambda_{33} g_{10}' + \lambda_{31} \lambda_{23} g_{13}') \\
 g_{10} &= \alpha(c_2^2 g_1' + 2c_2 \lambda_{22} g_2' + 2c_2 \lambda_{32} g_3' + \lambda_{22}^2 g_8' + \lambda_{32}^2 g_{11}') \\
 g_{11} &= \alpha[c_2 c_3 g_1' + (c_2 \lambda_{23} + c_3 \lambda_{22})g_2' + (c_2 \lambda_{33} + c_3 \lambda_{32})g_3' \\
 &\quad + \lambda_{22} \lambda_{23} g_8' + \lambda_{32} \lambda_{33} g_{11}'] \\
 g_{12} &= \alpha(c_2 \lambda_{21} g_4' + c_2 \lambda_{31} g_5' + \lambda_{22} \lambda_{31} g_9' + \lambda_{32} \lambda_{21} g_{12}')
 \end{aligned}$$



$$\begin{aligned}
g_{13} &= \alpha(c_2 \lambda_{22} g'_4 + c_2 \lambda_{32} g'_5 + \lambda_{22} \lambda_{32} g'_9 + \lambda_{32} \lambda_{22} g'_{12}) \\
g_{14} &= \alpha(c_2 \lambda_{23} g'_4 + c_2 \lambda_{33} g'_5 + \lambda_{22} \lambda_{33} g'_9 + \lambda_{32} \lambda_{23} g'_{12}) \\
g_{15} &= \alpha(c_2 \lambda_{21} g'_6 + c_2 \lambda_{31} g'_7 + \lambda_{22} \lambda_{31} g'_{10} + \lambda_{32} \lambda_{21} g'_{13}) \\
g_{16} &= \alpha(c_2 \lambda_{22} g'_6 + c_2 \lambda_{32} g'_7 + \lambda_{22} \lambda_{32} g'_{10} + \lambda_{32} \lambda_{22} g'_{13}) \\
g_{17} &= \alpha(c_2 \lambda_{23} g'_6 + c_2 \lambda_{33} g'_7 + \lambda_{22} \lambda_{33} g'_{10} + \lambda_{32} \lambda_{23} g'_{13}) \\
g_{18} &= \alpha(c_3^2 g'_1 + 2c_3 \lambda_{23} g'_2 + 2c_3 \lambda_{33} g'_3 + \lambda_{23}^2 g'_8 + \lambda_{33}^2 g'_{11}) \\
g_{19} &= \alpha(c_3 \lambda_{21} g'_4 + c_3 \lambda_{31} g'_5 + \lambda_{23} \lambda_{31} g'_9 + \lambda_{33} \lambda_{21} g'_{12}) \\
g_{20} &= \alpha(c_3 \lambda_{22} g'_4 + c_3 \lambda_{32} g'_5 + \lambda_{23} \lambda_{32} g'_9 + \lambda_{33} \lambda_{22} g'_{12}) \\
g_{21} &= \alpha(c_3 \lambda_{23} g'_4 + c_3 \lambda_{33} g'_5 + \lambda_{23} \lambda_{33} g'_9 + \lambda_{33} \lambda_{23} g'_{12}) \\
g_{22} &= \alpha(c_3 \lambda_{21} g'_6 + c_3 \lambda_{31} g'_7 + \lambda_{23} \lambda_{31} g'_{10} + \lambda_{33} \lambda_{21} g'_{13}) \\
g_{23} &= \alpha(c_3 \lambda_{22} g'_6 + c_3 \lambda_{32} g'_7 + \lambda_{23} \lambda_{32} g'_{10} + \lambda_{33} \lambda_{22} g'_{13}) \\
g_{24} &= \alpha(c_3 \lambda_{23} g'_6 + c_3 \lambda_{33} g'_7 + \lambda_{23} \lambda_{33} g'_{10} + \lambda_{33} \lambda_{23} g'_{13}) \\
g_{25} &= \alpha(c_1^2 g'_{14} + \lambda_{21}^2 g'_{15} + \lambda_{31}^2 g'_{17}) \\
g_{26} &= \alpha(c_1 c_2 g'_{14} + \lambda_{21} \lambda_{22} g'_{15} + \lambda_{31} \lambda_{32} g'_{17}) \\
g_{27} &= \alpha(c_1 c_3 g'_{14} + \lambda_{21} \lambda_{23} g'_{15} + \lambda_{31} \lambda_{33} g'_{17}) \\
g_{28} &= \alpha(-c_1^2 g'_{14} + \lambda_{21}^2 g'_{16} + \lambda_{31}^2 g'_{18}) \\
g_{29} &= \alpha(-c_1 c_2 g'_{14} + \lambda_{21} \lambda_{22} g'_{16} + \lambda_{31} \lambda_{32} g'_{18}) \\
g_{30} &= \alpha(-c_1 c_3 g'_{14} + \lambda_{21} \lambda_{23} g'_{16} + \lambda_{31} \lambda_{33} g'_{18}) \\
g_{31} &= \alpha(c_2^2 g'_{14} + \lambda_{22}^2 g'_{15} + \lambda_{32}^2 g'_{17})
\end{aligned} \tag{5.42}$$

$$\begin{aligned}
g_{32} &= \alpha(c_2 c_3 g'_{14} + \lambda_{22} \lambda_{23} g'_{15} + \lambda_{32} \lambda_{33} g'_{17}) \\
g_{33} &= \alpha(-c_2^2 g'_{14} + \lambda_{22}^2 g'_{16} + \lambda_{32}^2 g'_{18}) \\
g_{34} &= \alpha(-c_2 c_3 g'_{14} + \lambda_{22} \lambda_{23} g'_{16} + \lambda_{32} \lambda_{33} g'_{18}) \\
g_{35} &= \alpha(c_3^2 g'_{14} + \lambda_{23}^2 g'_{15} + \lambda_{33}^2 g'_{17}) \\
g_{36} &= \alpha(-c_3^2 g'_{14} + \lambda_{23}^2 g'_{16} + \lambda_{33}^2 g'_{18}) \\
g_{37} &= \alpha(c_1^2 g'_{14} + \lambda_{21}^2 g'_{19} + \lambda_{31}^2 g'_{20}) \\
g_{38} &= \alpha(c_1 c_2 g'_{14} + \lambda_{21} \lambda_{22} g'_{19} + \lambda_{31} \lambda_{32} g'_{20}) \\
g_{39} &= \alpha(c_1 c_3 g'_{14} + \lambda_{21} \lambda_{23} g'_{19} + \lambda_{31} \lambda_{33} g'_{20}) \\
g_{40} &= \alpha(c_2^2 g'_{14} + \lambda_{22}^2 g'_{19} + \lambda_{32}^2 g'_{20}) \\
g_{41} &= \alpha(c_2 c_3 g'_{14} + \lambda_{22} \lambda_{23} g'_{19} + \lambda_{32} \lambda_{33} g'_{20}) \\
g_{42} &= \alpha(c_3^2 g'_{14} + \lambda_{23}^2 g'_{19} + \lambda_{33}^2 g'_{20})
\end{aligned} \tag{5.42}$$

where

$$\alpha = EI/L$$

Chapter VI  
SOLUTION ALGORITHMS

6.1 INTRODUCTION

The response of a nonlinear structure may be solved approximately for incremental nodal displacements by taking a series of linear steps. Many different solution schemes have been proposed to solve the nonlinear equilibrium Eqs. (2.55) and (2.104). Here the numerical solutions are obtained by applying either the Newton-Raphson method or the modified Riks/Wempner method to the nonlinear equilibrium equation. The former, efficient for low convergence tolerance [16], is popularly used to trace nonlinear prebuckling paths of structures; however, it cannot trace the response beyond the limit point. The latter has been especially proposed to overcome this problem and can trace nonlinear post critical response [27]. Both solution techniques are described briefly in this chapter.

## 6.2 NEWTON-RAPHSON METHOD

The Newton-Raphson method for solving the nonlinear equilibrium equations (2.55) and (2.104) may be stated as [43],[45]

$$K^{k-1} \Delta q^k = R^{k-1} \quad k=1,2,\dots,n \quad (6.1)$$

in which

$$K^{k-1} = \sum_{i=1}^{NE} K^{(i)k-1} \quad (6.2)$$

$$R^{k-1} = t+\Delta t Q - F^{k-1} \quad (6.3)$$

$$q^k = q^{k-1} + \Delta q^k \quad (6.4)$$

$q^k$  is the  $k^{\text{th}}$  trial solution corresponding to a given load level  $t+\Delta t Q$ ,  $K^{k-1}$  is the structural tangent stiffness matrix corresponding to  $q^{k-1}$ ,  $K^{(i)k-1}$  is the generalized tangent stiffness matrix of element  $i$  corresponding to  $q^{k-1}$ ,  $F^{k-1}$  is the equilibrating nodal force vector corresponding to  $q^{k-1}$ ,  $R^{k-1}$  is the unbalanced force vector corresponding to  $q^{k-1}$ .

The procedure, illustrated for a one-degree-of-freedom system in Fig. 6.1, is as follows. For a given new load level  $t+\Delta t Q$

1. Establish system tangent stiffness matrix  $K^{k-1}$ .
2. Evaluate the equilibrating nodal force vector  $F^{k-1}$ .
3. Compute the vector of unbalanced nodal force  $R^{k-1}$  using Eq. (6.3).
4. Solve Eq. (6.1) for the vector of incremental nodal displacements  $\Delta q^k$ .
5. Update the vector of nodal displacements by Eq. (6.4).
6. Test for convergence.
7. If the process has not converged return to step 1. Otherwise increment the load vector and seek a new solution.

To reduce the amount of computations per iteration, the Newton-Raphson method is often modified by using the same tangent stiffness matrix during several iterative cycles (Fig. 6.2). It is updated only at every load level or only when the convergence rate becomes poor. This method requires more steps to reach a new equilibrium point. The choice between both methods depends on the closeness of the initial vector to the true solution.

Using Newton-Raphson method in the neighborhood of a limit point, the tangent stiffness matrix approaches singulari-

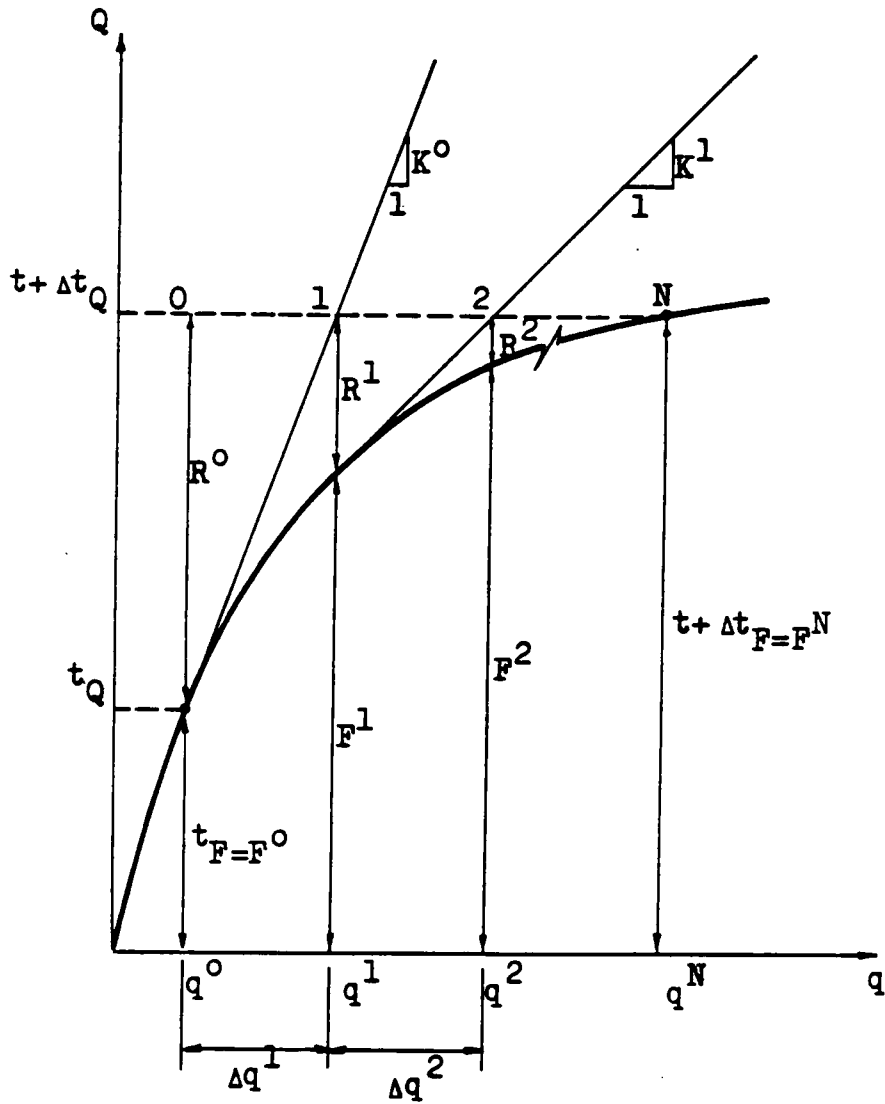


Fig. 6.1 Newton-Raphson Iteration

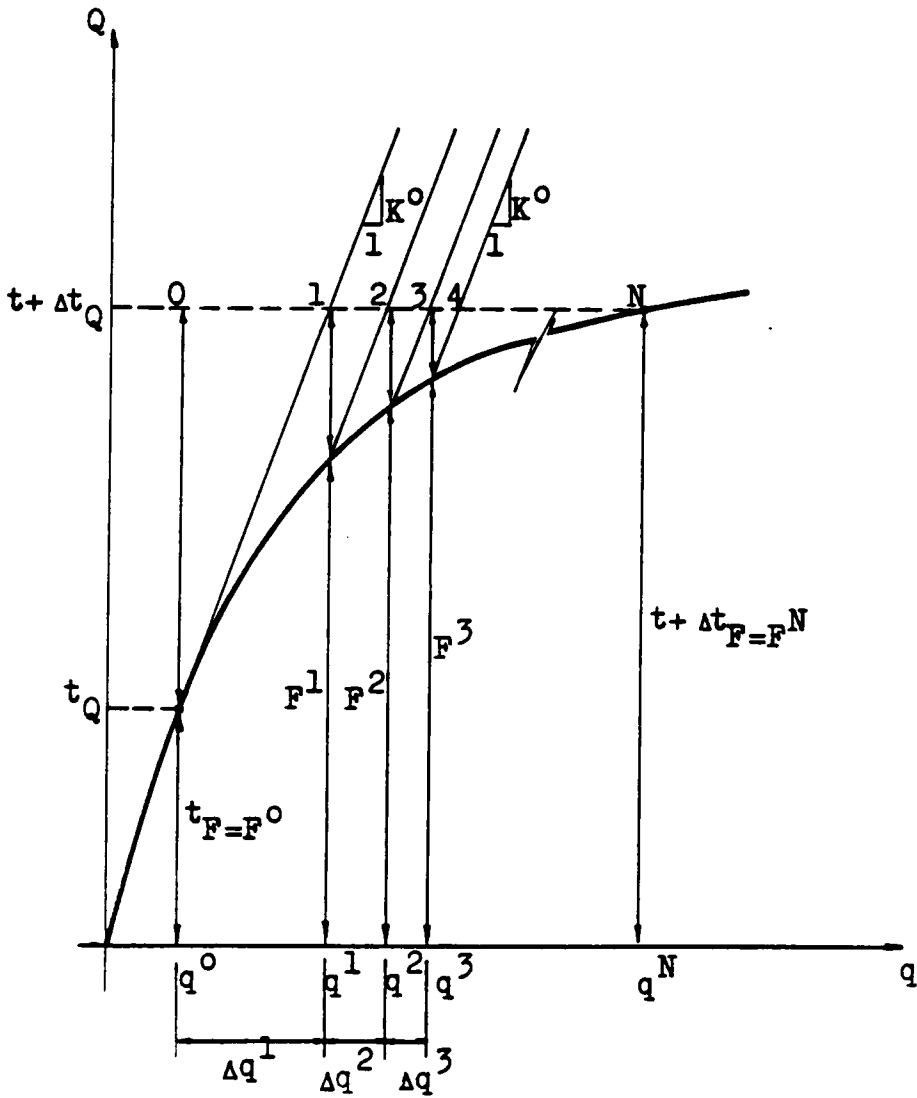


Fig. 6.2 Modified Newton-Raphson Iteration

ty resulting in an increasing number of iterations and smaller and smaller load step; finally the solution diverges. Therefore, it cannot trace the response beyond the limit point. To solve this problem, the modified Riks/Wempner method is recommended [27].

### 6.3 MODIFIED RIKS/WEMPNER METHOD

The theoretical development was recently summarized in reference [27]. Therefore, details are omitted here. The basic idea of the Riks/Wempner method is to choose a generalized arc length to facilitate the procedure of seeking a new solution [45],[46]. For a given basic load incremental  $\Delta\lambda^1$ , the generalized arc length  $\Delta S$  can be computed from the constraint equation. Then the iteration path follows a normal plane to the tangent, and the new equilibrium point will be the intersection of the normal plane with the equilibrium path. Iteration along a circle (or sphere in space) will not be considered here [27],[47]. In this section the notation of reference [27] is adopted.

The algorithm of the modified Riks/Wempner method is briefly summarized as follows (Figs. 6.3 and 6.4):

#### For the First Step



1. Choose a basic load increment  $\Delta\lambda^1$
2. Establish the system tangent stiffness matrix  $K^0$  in the current configuration  $q^0$
3. Solve equation

$$K^0 \Delta q^{1I} = \bar{Q} \quad \text{for } \Delta q^{1I} \quad (6.5)$$

where  $\bar{Q}$  is the vector of constant load distribution

4. Compute the generalized arc length  $\Delta S$  from the constraint equation

$$\Delta S = \Delta\lambda^1 (\Delta q^{1I} \cdot \Delta q^{1I} + 1)^{1/2} \quad (6.6)$$

or scale  $\Delta S$  to control the number of future iterations

$$\Delta \hat{S} = \Delta S (\hat{I}/I)^{1/2} \quad (6.7)$$

where  $\hat{I}$  is the desired number of iterations and  $I$  is the required number of iterations in the previous step.

5. Compute the incremental nodal displacements

$$\Delta q^1 = \Delta\lambda^1 \Delta q^{1I} \quad (6.8)$$

6. Update the nodal displacements and the load parameter

$$q^1 = q^0 + \Delta q^1 \quad (6.9)$$

$$\lambda^1 = \lambda^0 + \Delta\lambda^1 \quad (6.10)$$

During the  $k^{\text{th}}$  iteration on normal plane,  $k=2,3,\dots$

7. Compute the nodal force vector  $F^{k-1}$  corresponding to  $q^{k-1}$

8. Update the load vector

$$Q^{k-1} = \lambda^{k-1} \bar{Q} \quad (6.11)$$

9. Compute the unbalanced force vector

$$R^{k-1} = Q^{k-1} - F^{k-1} \quad (6.12)$$

10. Update the system tangent stiffness matrix  $K^{k-1}$  if desirable

11. Solve for  $\Delta q^{kI}$  and  $\Delta q^{kII}$  from the two sets of equilibrium equations

$$K^{k-1} \Delta q^{kI} = \bar{Q} \quad (6.13)$$

$$K^{k-1} \Delta q^{kII} = R^{k-1} \quad (6.14)$$

12. Compute the incremental load parameter

$$\Delta \lambda^k = -(\Delta q^1 \cdot \Delta q^{kII}) / (\Delta q^1 \cdot \Delta q^{kI} + \Delta \lambda^1) \quad (6.15)$$

13. Compute the incremental nodal displacement vector

$$\Delta q^k = \Delta \lambda^k \Delta q^{kI} + \Delta q^{kII} \quad (6.16)$$

14. Update the nodal displacement vector and the load parameter

$$q^k = q^{k-1} + \Delta q^k \quad (6.17)$$

$$\lambda^k = \lambda^{k-1} + \Delta \lambda^k \quad (6.18)$$

15. Repeat steps 7 to 14 until process has converged  
 16. Start a new step by returning to step 2.

In conclusion, the modified Riks/Wempner method has been successfully applied to a variety problems [27],[43].

#### 6.4 CONVERGENCE CRITERIA

The incremental solution at the end of each iteration should be checked to see whether it has converged within preset tolerances. Two displacement criteria and one force criterion are discussed in this section.

Cook [33] adopted the displacement criterion based on the infinity vector norm:

$$\| \Delta q^k \|_{\infty} / \| q^k \|_{\infty} \leq \text{CPDC} \quad (6.19)$$

where  $\Delta q^k$  is the maximum incremental displacement,  $q^k$  is the maximum total displacement of same type, and CPDC is the

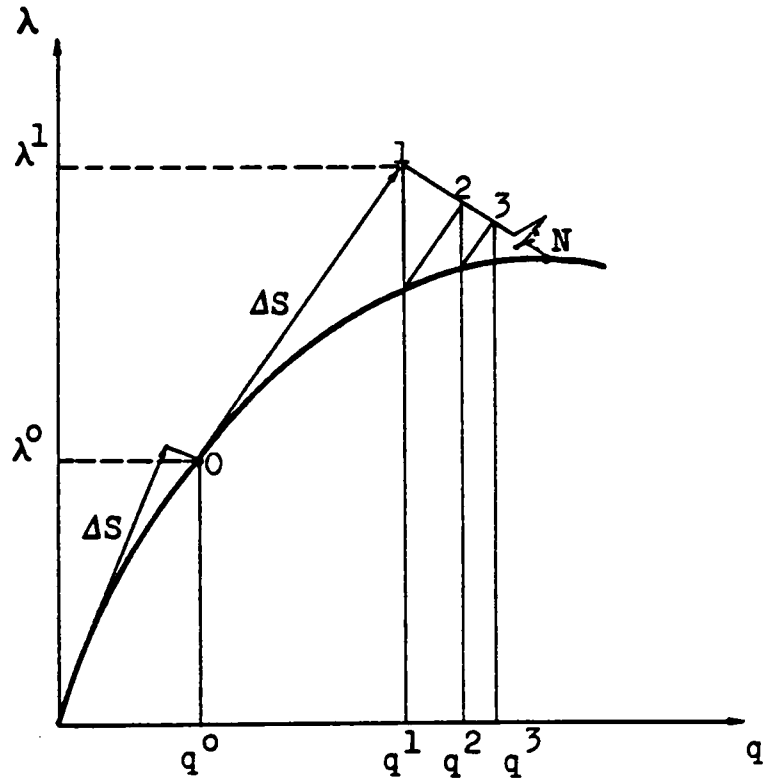


Fig. 6.3 Modified Riks/Wempner Iteration  
Along the Normal Plane

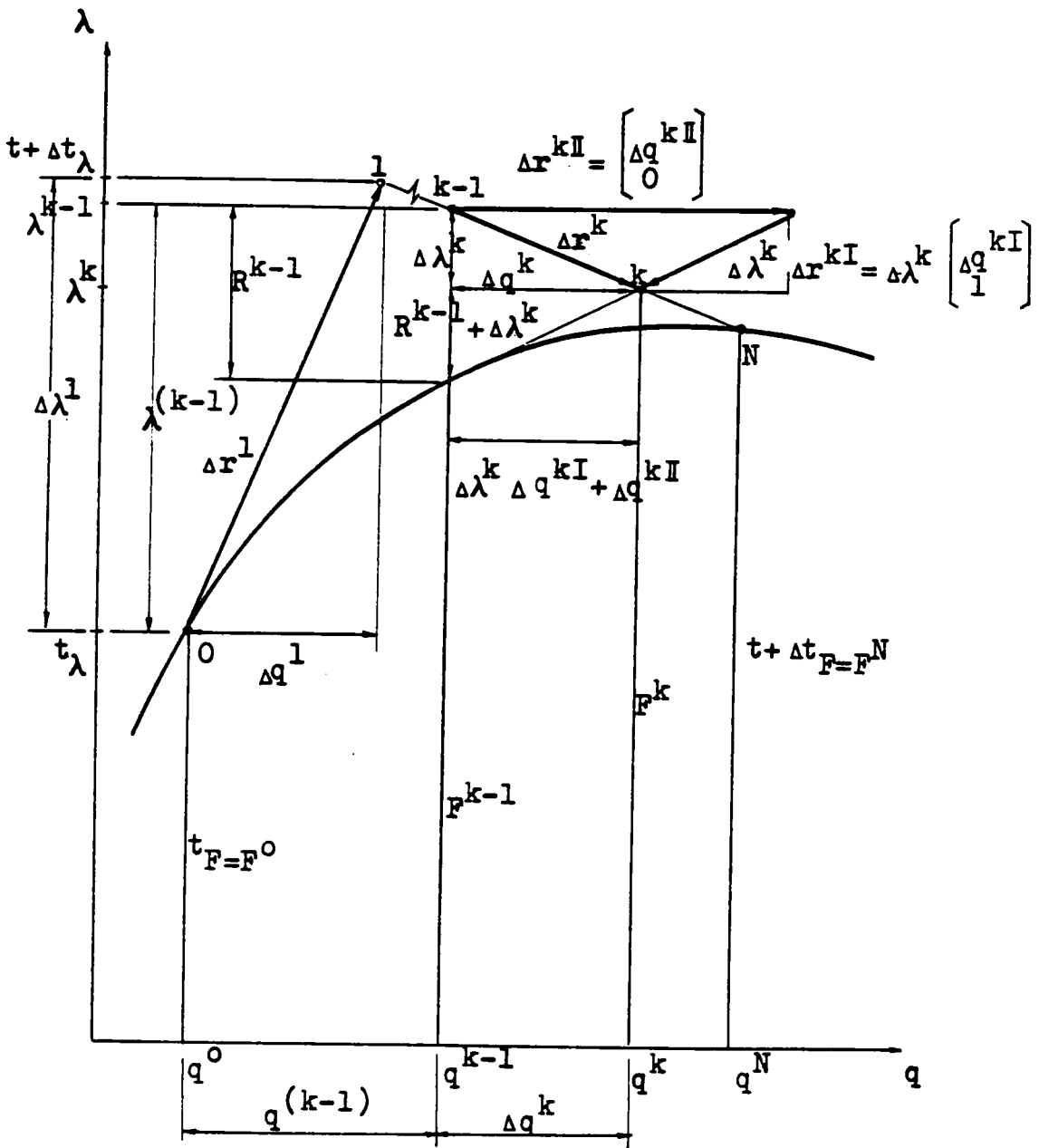


Fig. 6.4 Iteration Procedure along Normal Plane with Updating K

displacement convergence parameter adopted by Cook with the following range:

$$10^{-6} \leq \text{CPDC} \leq 10^{-2} \quad (6.20)$$

Bathe [38],[49] recommends the displacement criterion based on the Euclidean vector norm:

$$\| \Delta \mathbf{q}^k \|_2 / \| \mathbf{q}^k \|_2 \leq \text{CPDB} \quad (6.21)$$

where CPDB is the displacement convergence parameter adopted by Bathe with the value

$$\text{CPDB} = 0.001 \quad (6.22)$$

A third convergence criterion could be defined as the ratio of the norm of the residual load to the norm of the original load increment [38],[49]:

$$\| \mathbf{t}^{t+\Delta t} \mathbf{Q} - \mathbf{t}^{t+\Delta t} \mathbf{F}^k \|_2 / \| \mathbf{t}^{t+\Delta t} \mathbf{Q} - \mathbf{F}^0 \|_2 \leq \text{CPF} \quad (6.23)$$

where CPF is the convergence parameter for the unbalanced force criterion with a convergence tolerance of 0.1.

## Chapter VII

### SAMPLE ANALYSIS

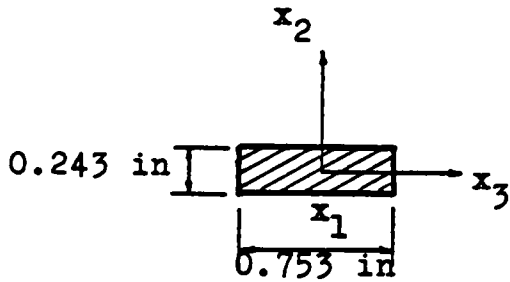
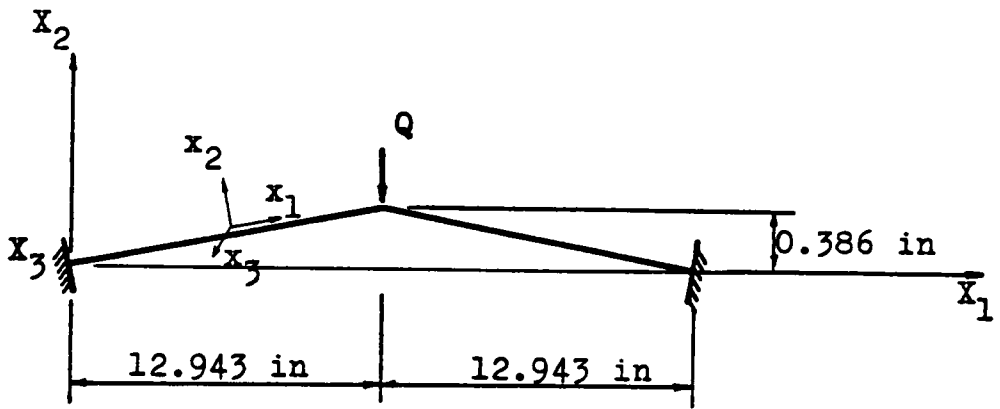
#### 7.1 INTRODUCTION

The U.L formulation described in Chapter 2 has been employed in the computer program. Four examples were investigated in this study and comparisons between the finite element model and the beam-column model were made. All of the test structures were treated as space frames and analyzed by the modified Riks/Wempner method.

#### 7.2 EXAMPLE 1: WILLIAMS' TOGGLE FRAME

Williams [24], Wood and Zienkiewicz [20], and Papadrakakis [21] investigated this problem. Fig. 7.1 shows the configuration and properties.

The finite element solution is presented in Fig. 7.2 to compare the effects of mesh refinement. One, four, six, and twelve elements per member are used. The responses of the different meshes are in close agreement up to near the limit point. Beyond this region, the single element is too stiff, but four elements are adequate for modeling the behavior of the toggle frame.



$$A = 0.183 \text{ in}^2$$

$$E = 10,300. \text{ ksi}$$

$$I_3 = 0.00090039 \text{ in}^4$$

Fig. 7.1 William's Toggle Frame



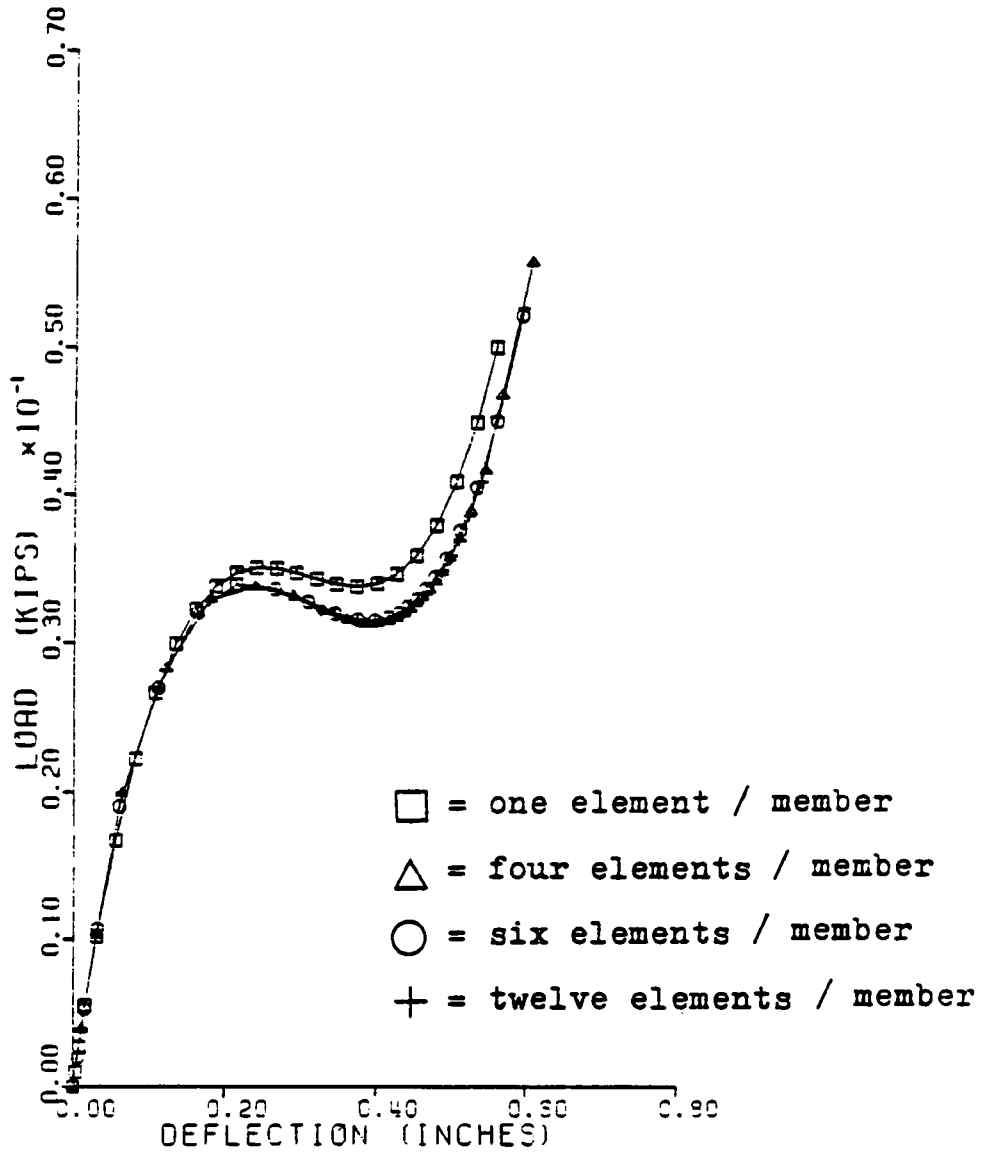


Fig. 7.2 Finite Element Model

In Fig. 7.3 a comparison of twelve finite elements/member to one beam-column element/member is presented. It shows that the accuracy of one beam-column element is comparable to that of twelve finite elements. The execution times for both elements are as follows:

one finite element	:	0.43 sec
four finite elements	:	1.12 sec
six finite elements	:	1.37 sec
twelve finite elements	:	2.40 sec
one beam-column element:		0.49 sec

### 7.3 EXAMPLE 2 : THREE DIMENSIONAL CANTILEVER BEAM OF A 45-DEGREE BEND

A tip loaded cantilever beam of a 45-degree bend, undergoing large displacements (Fig. 7.4), was investigated by Bathe [9]. The beam lies in the  $X_1$ - $X_3$  plane. The average radius of the bend is 100 inch. The concentrated end load is applied in the  $X_2$ -direction.

Using both models, the tip deflections in the  $X_2$ -direction are shown in Figs. 7.5 and 7.6 respectively. Fig. 7.5 presents the finite element solution in which the

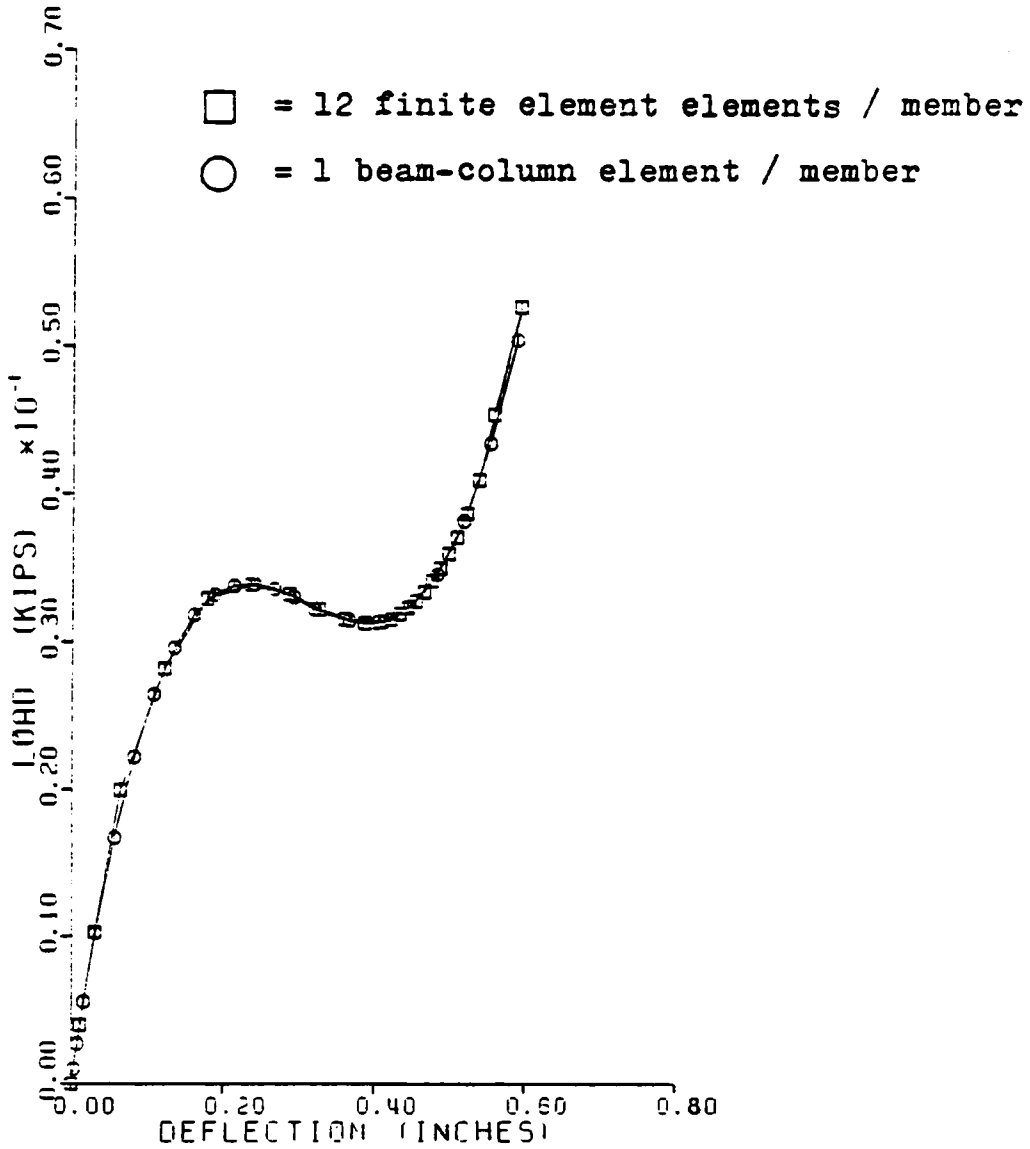
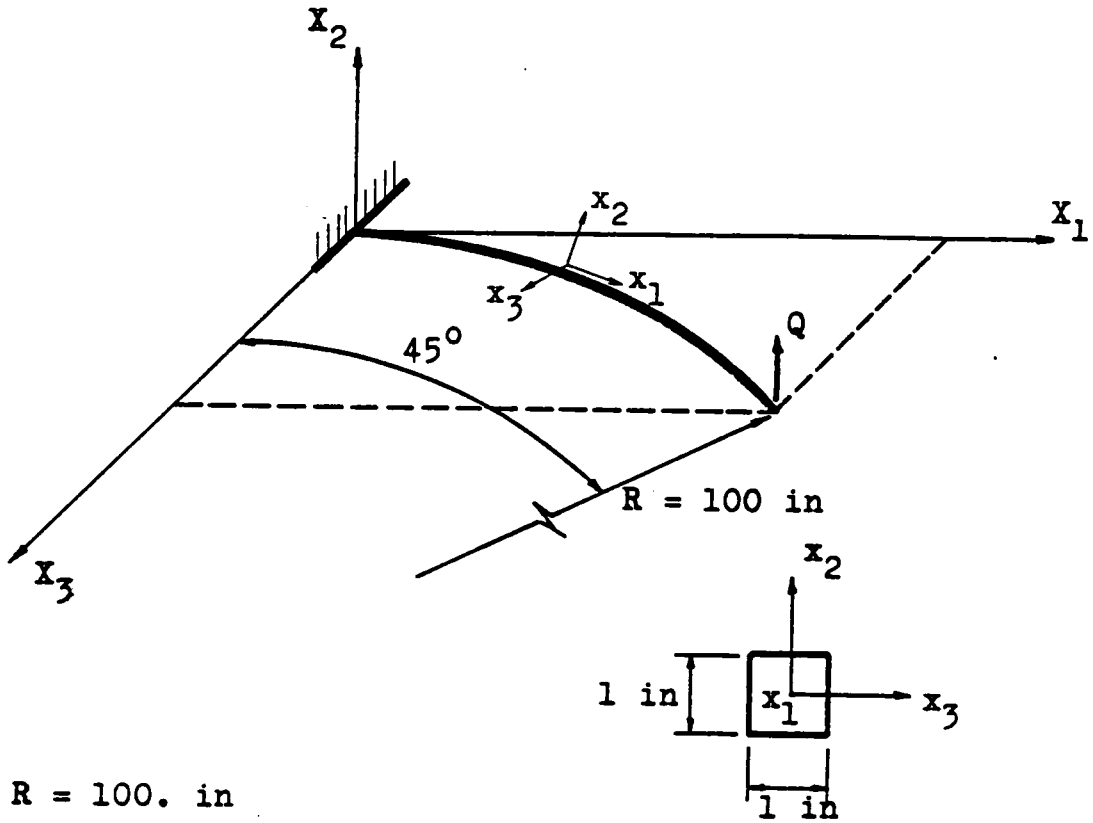


Fig. 7.3 Comparison of Models



$$R = 100. \text{ in}$$

$$E = 10^7 \text{ psi}$$

$$I_2 = I_3 = 0.0833333 \text{ in}^4$$

$$\nu = 0.$$

$$G = \frac{E}{2(1+\nu)}$$

$$J = 0.141a^4 = 0.141 \text{ in}^4$$

beam cross section

Fig. 7.4 Three Dimensional Cantilever Beam of  $45^\circ$  Bend

effects of different mesh refinements are compared; it shows that two and four elements fail (solution procedure breaks down), but eight elements are adequate to model the behavior of the curved beam. Fig. 7.6 presents the beam-column solution obtained using four and eight elements; it gives similar results for both meshes. Fig. 7.7 presents a comparison of four beam-column elements to eight finite elements; the response predicted by eight finite elements is slightly stiffer and less accurate than the response predicted by four beam-column elements.

The result obtained using four beam-column elements is also compared to the Bathe's solution [9] using eight finite elements (Fig. 7.8); it can be seen that the beam-column solution produces a similar level of accuracy using a much coarser mesh. The deflected shapes of the cantilever beam at various load levels are depicted in Fig. 7.9. The execution times for both models are as follows:

four elements:

finite element model: 5.66 sec

beam-column model: 7.19 sec

eight elements:

finite element model: 37.27 sec

beam-column model: 37.35 sec

- △ = two elements  
○ = four elements  
□ = eight elements

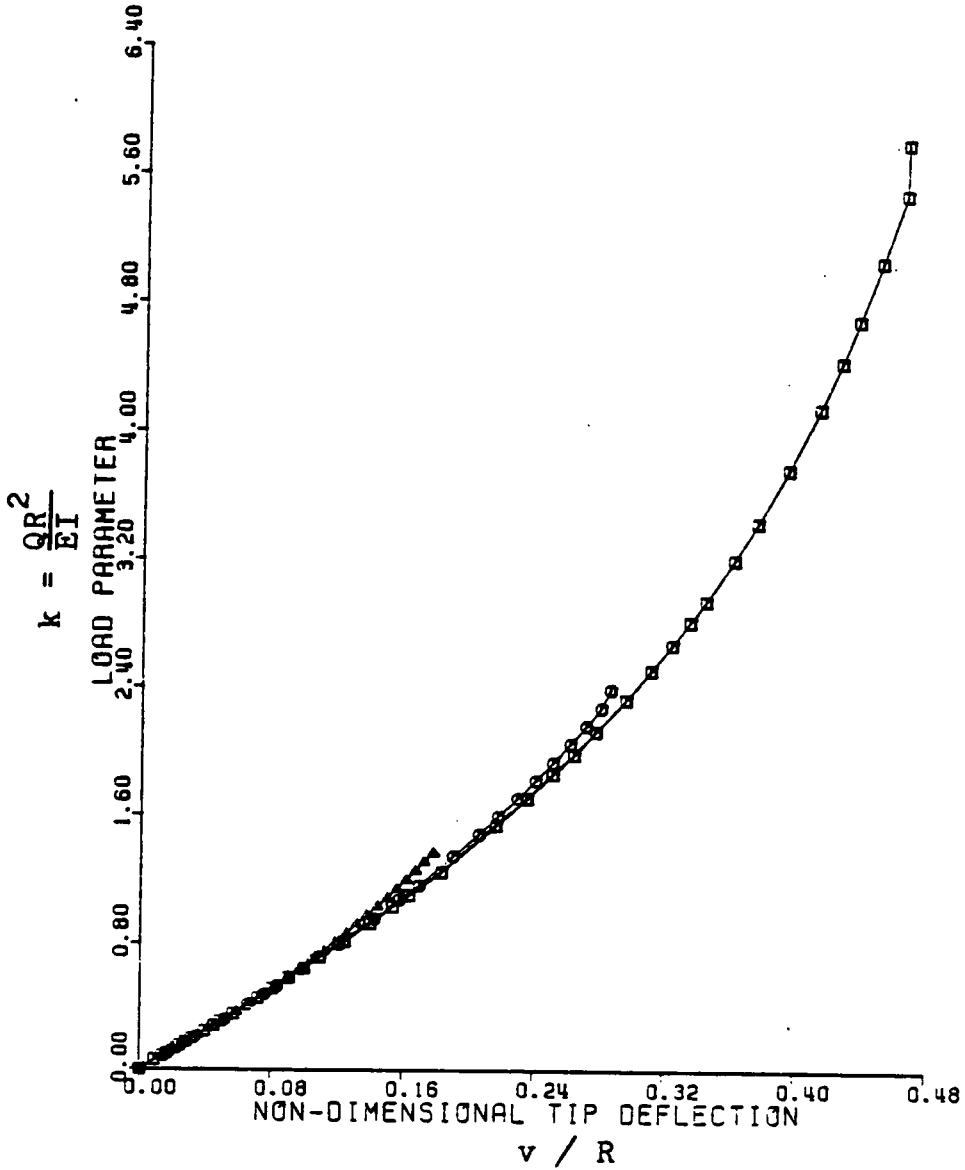


Fig. 7.5 Finite Element Model

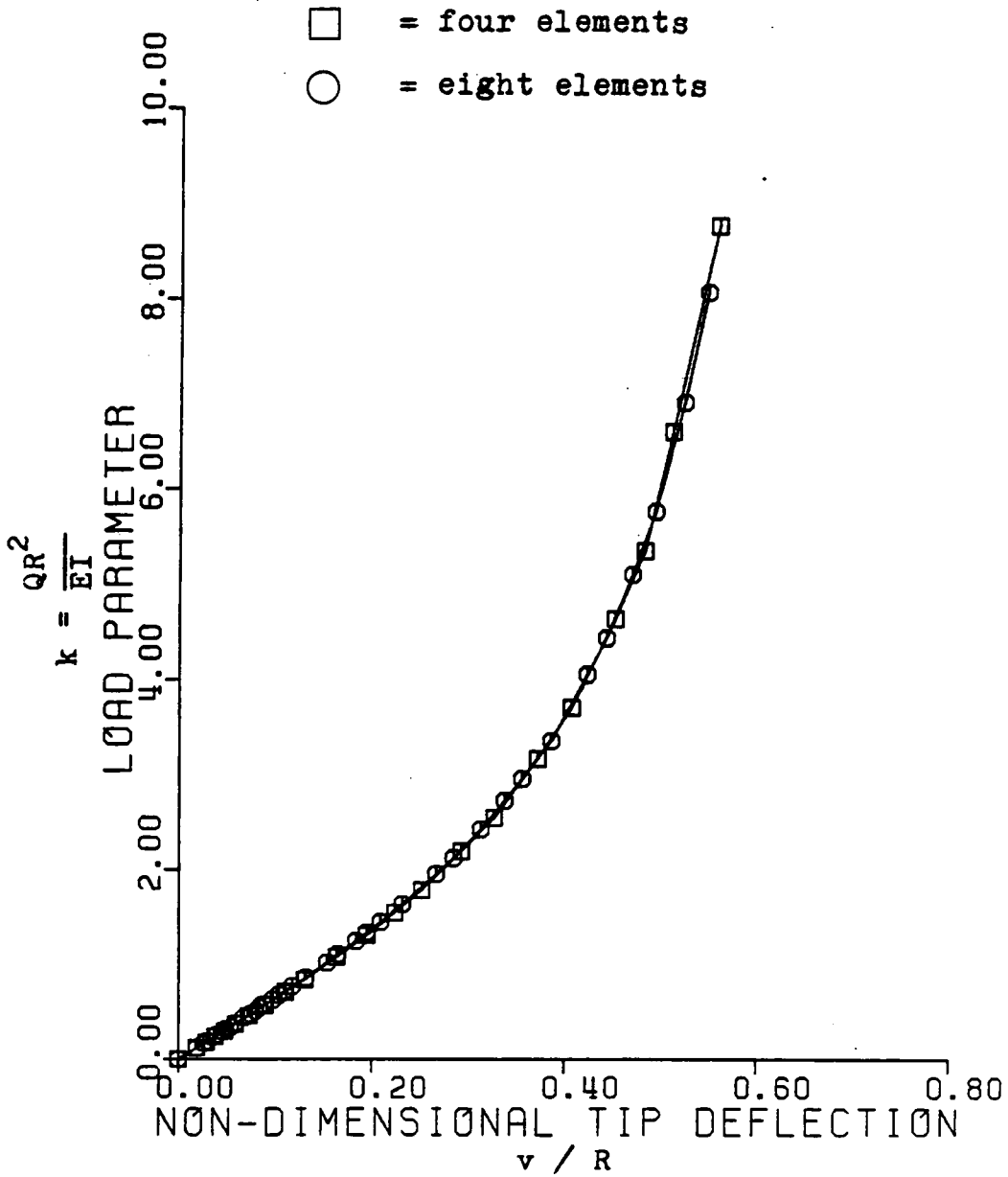


Fig. 7.6 Beam-Column Model

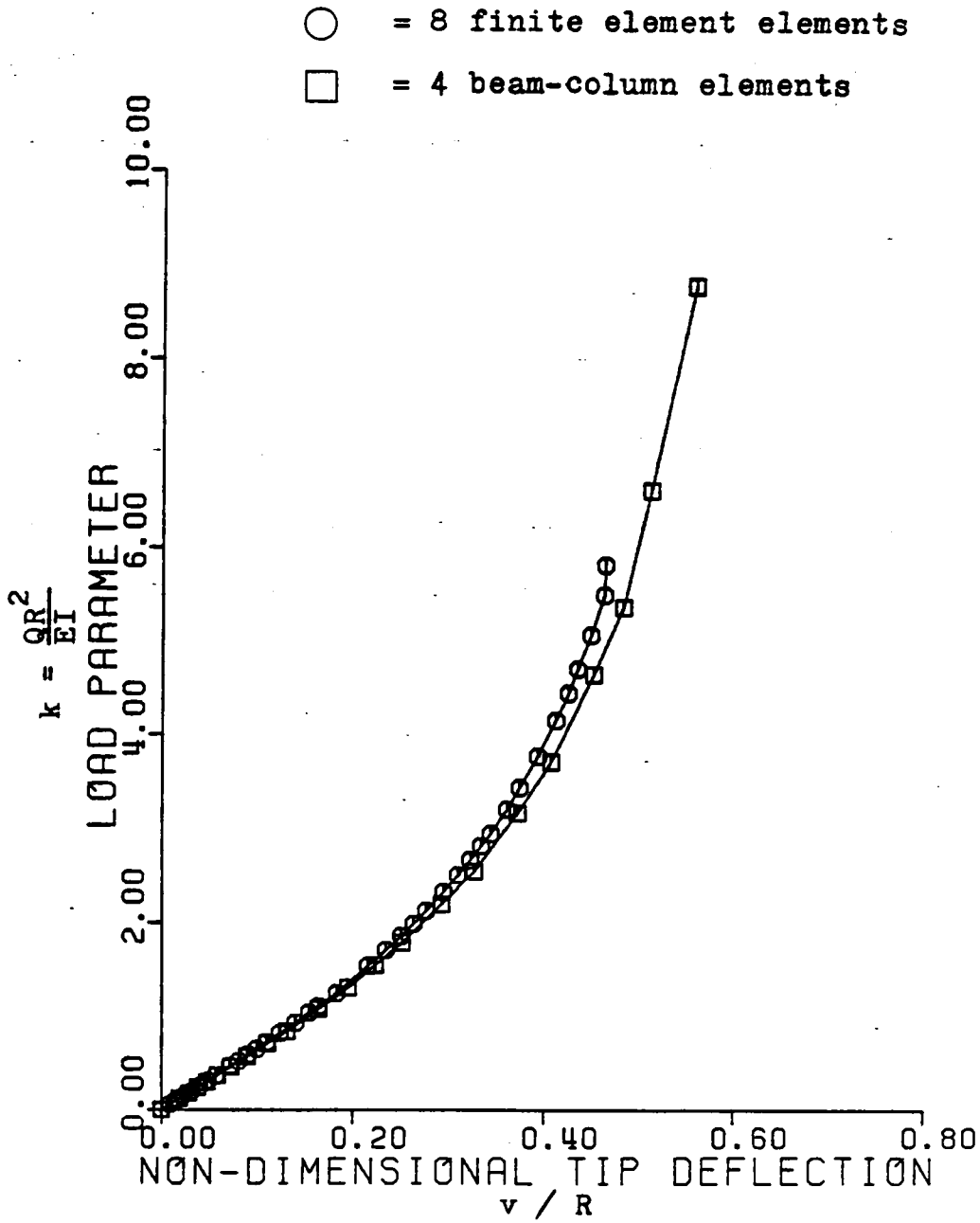


Fig. 7.7 Comparison of Models



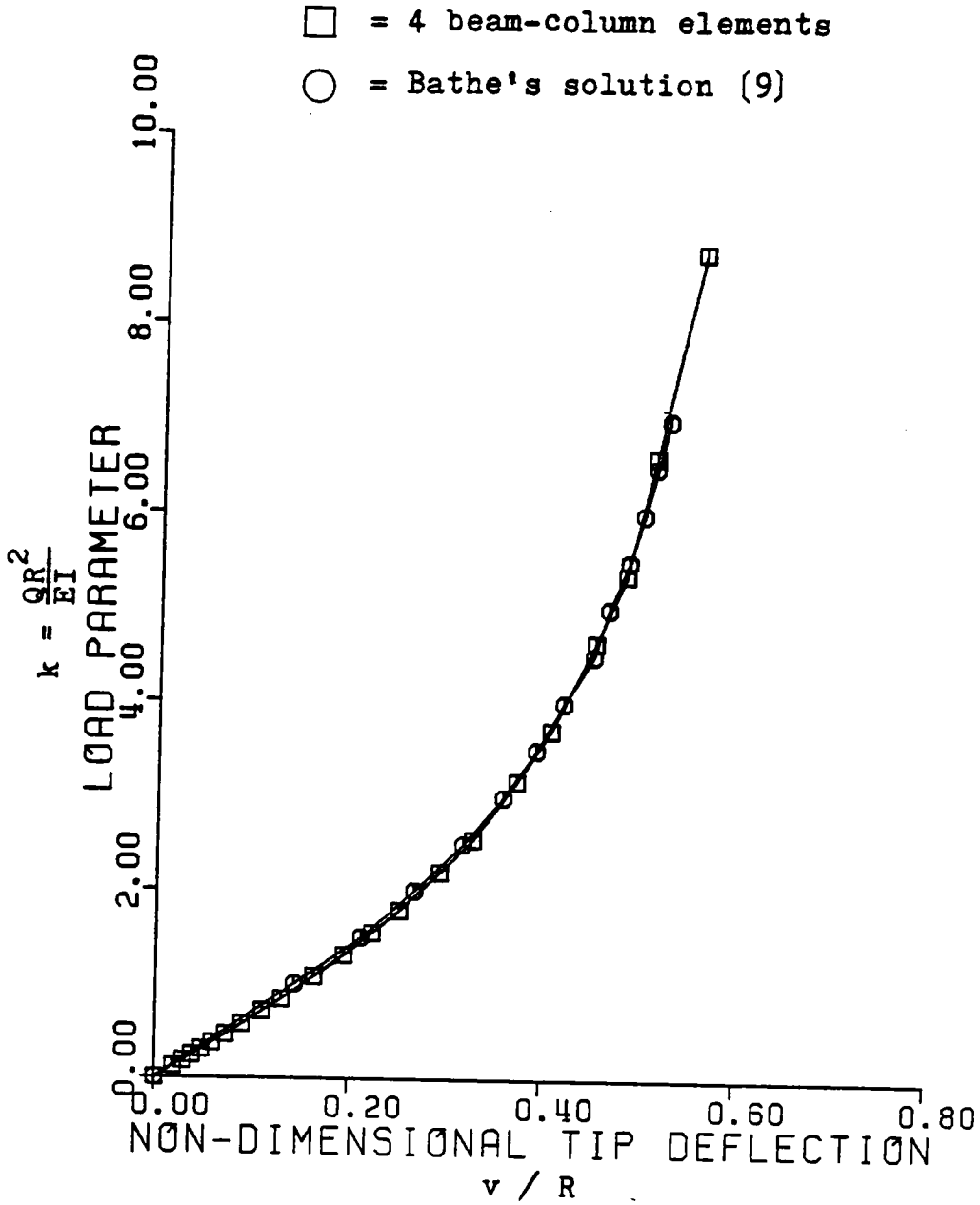
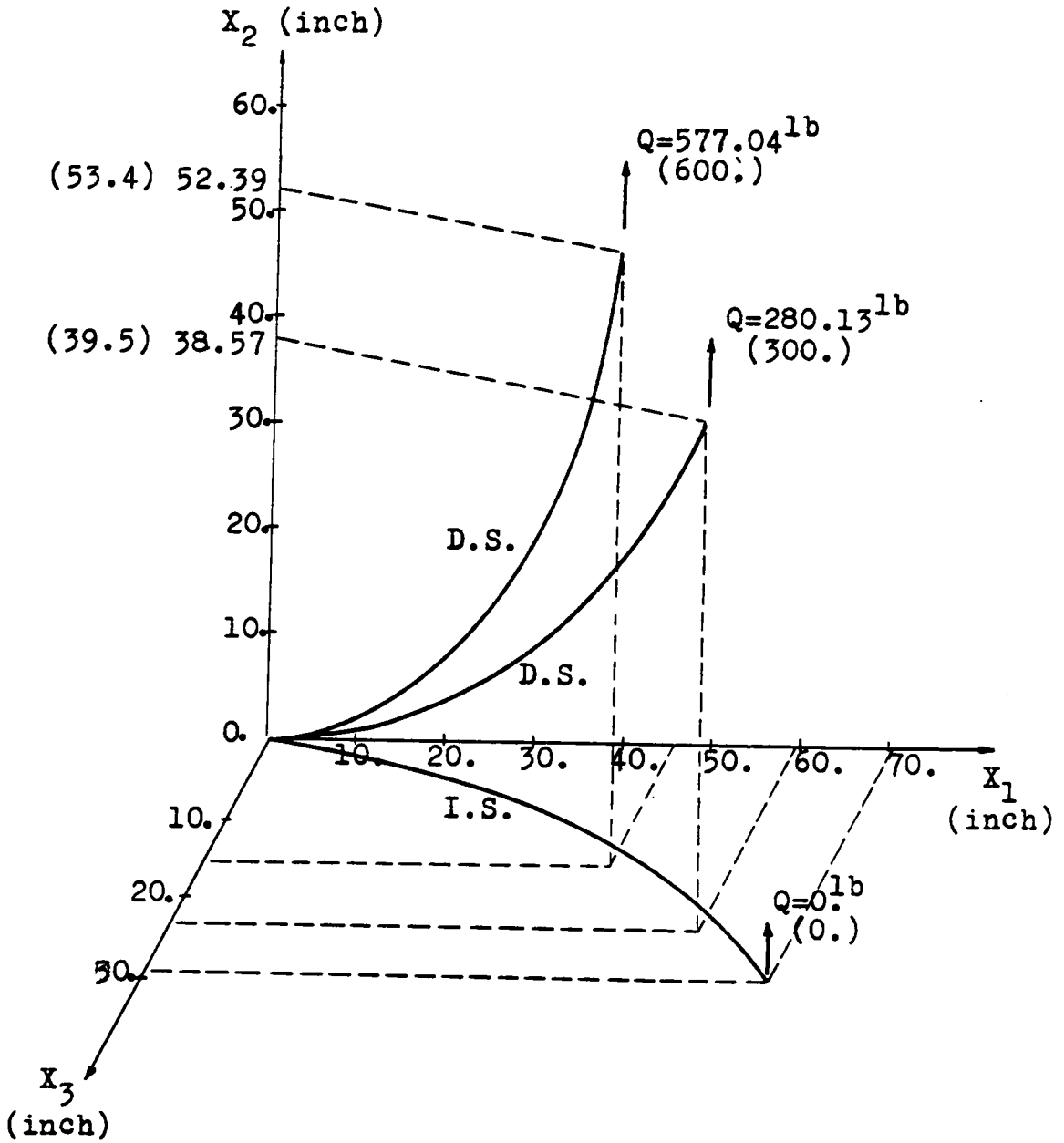


Fig. 7.8 Load Deflection Curves



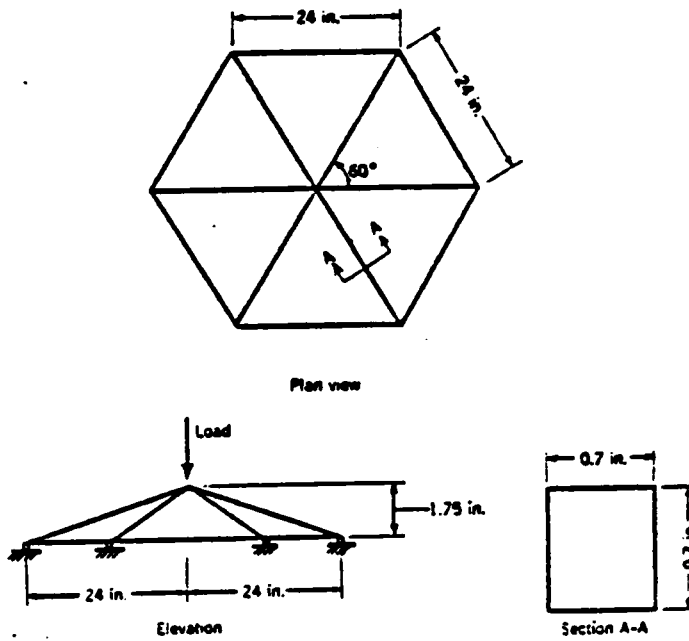
( ) = Bathe's results [9]

Fig. 7.9 Deflected Shapes of a 45° Circular Bend using the Beam-Column Model

#### 7.4 EXAMPLE 3: 12 MEMBER MODEL FRAME

The model frame shown in Fig. 7.10 was analyzed by Connor [16], Papadrakakis [21], and Chu [39], one and two elements per member are employed in this study. Fig. 7.11 shows the finite element solution; it can be seen that one element cannot accurately represent the behavior of the structure. Fig. 7.12 presents the beam-column solution; it gives similar results for both meshes. In Fig. 7.13 a comparison of equilibrium paths up to the limit point of the finite element model with the beam-column model is presented; both paths agree quite closely. The numerical solutions of others workers [16],[21] and [39] are given in Fig. 7.14. In Fig. 7.15 a comparison of Papadrakakis solution with the two beam-column elements/member solution is presented; the results are in close agreement.

Combining Fig. 7.12,7.13 and 7.15, it is seen that the beam-column solution obtained using one element per member give satisfactory result in which two elements of finite element solution are needed, to predict the structural response.



$$E = 439,800. \text{ psi}$$

$$G = 159,000. \text{ psi}$$

$$J = 0.141 a^4 = 0.0344382 \text{ in}^4$$

Fig. 7.10 12 Member Model Frame

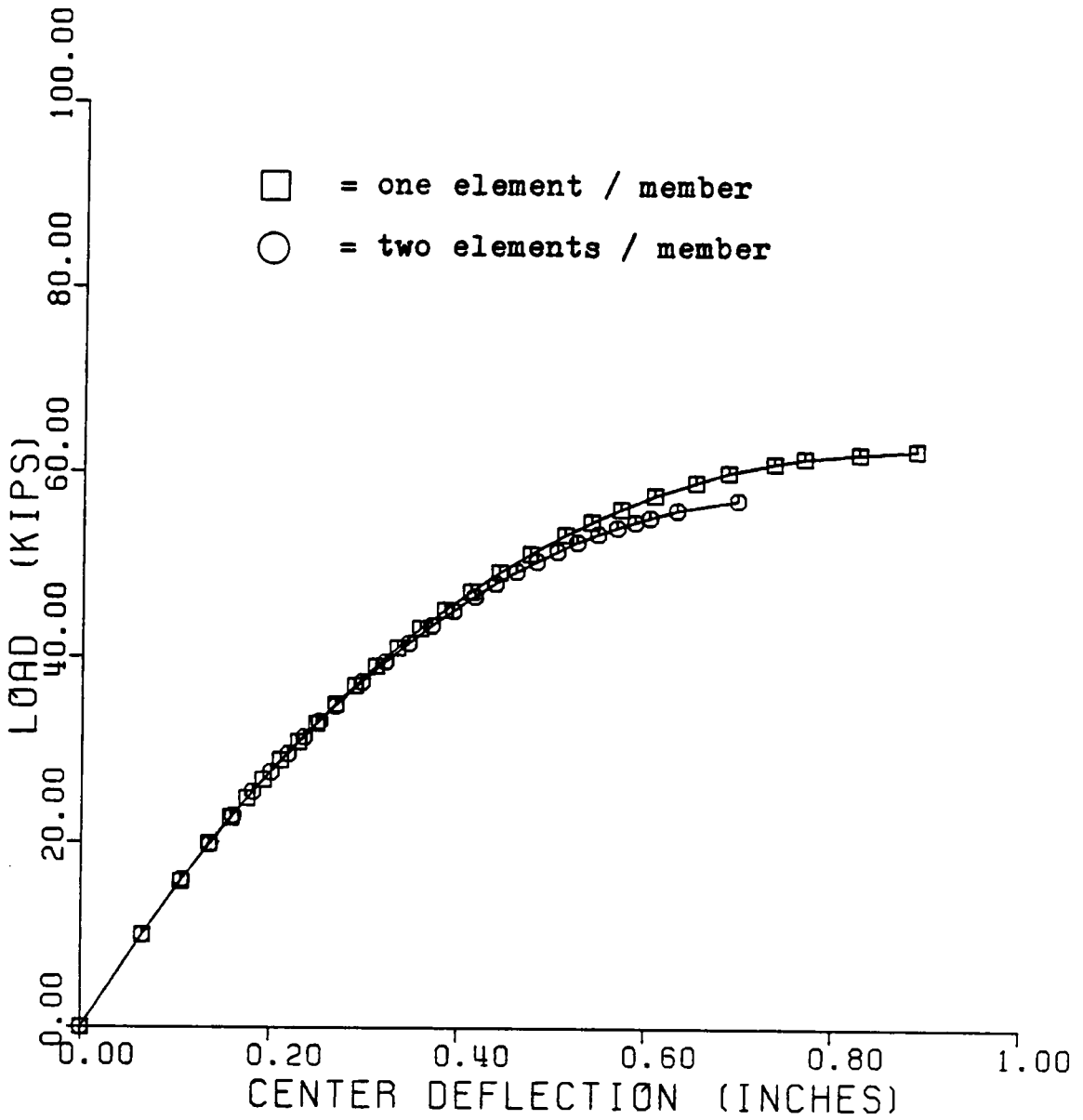


Fig. 7.11 Finite Element Model

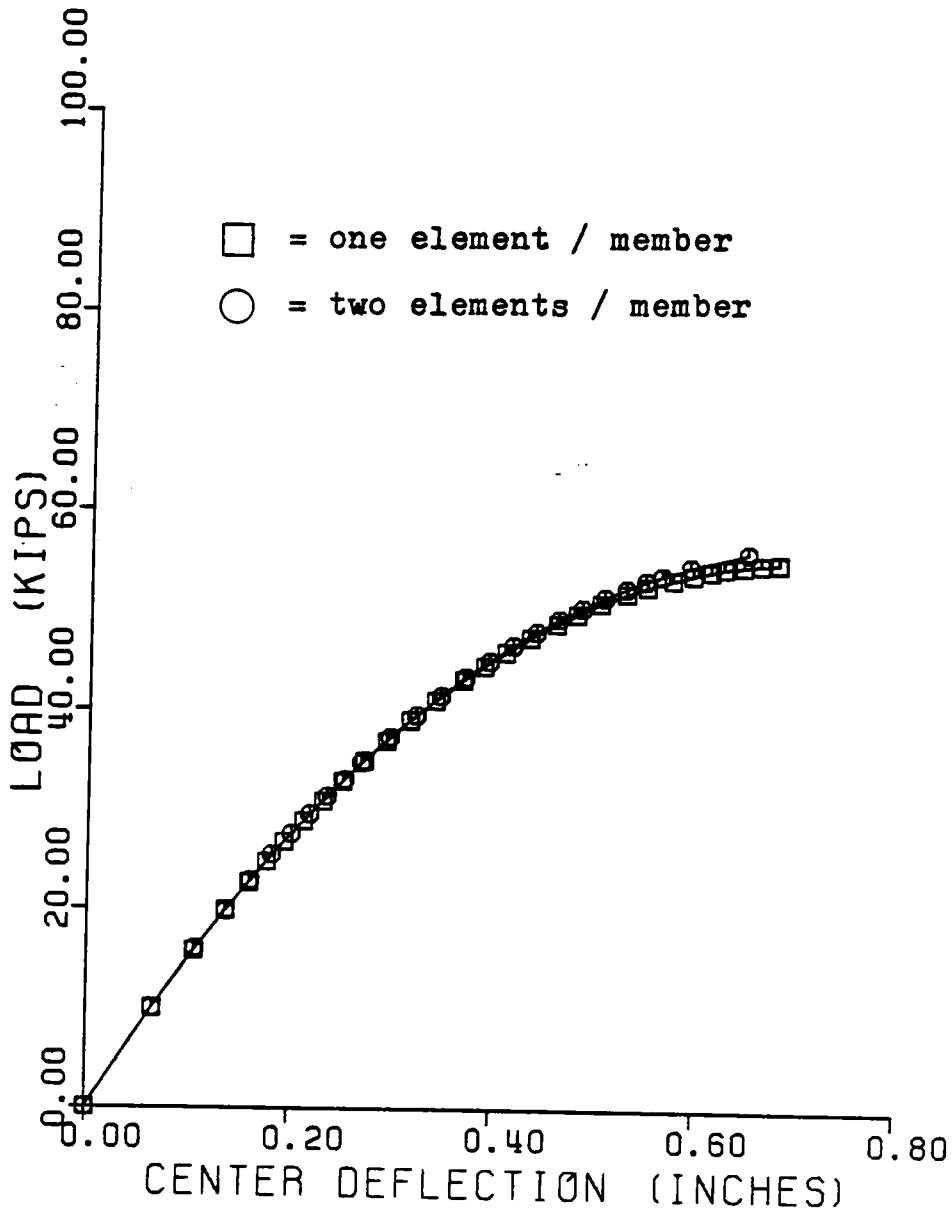


Fig. 7.12 Beam-Column Model

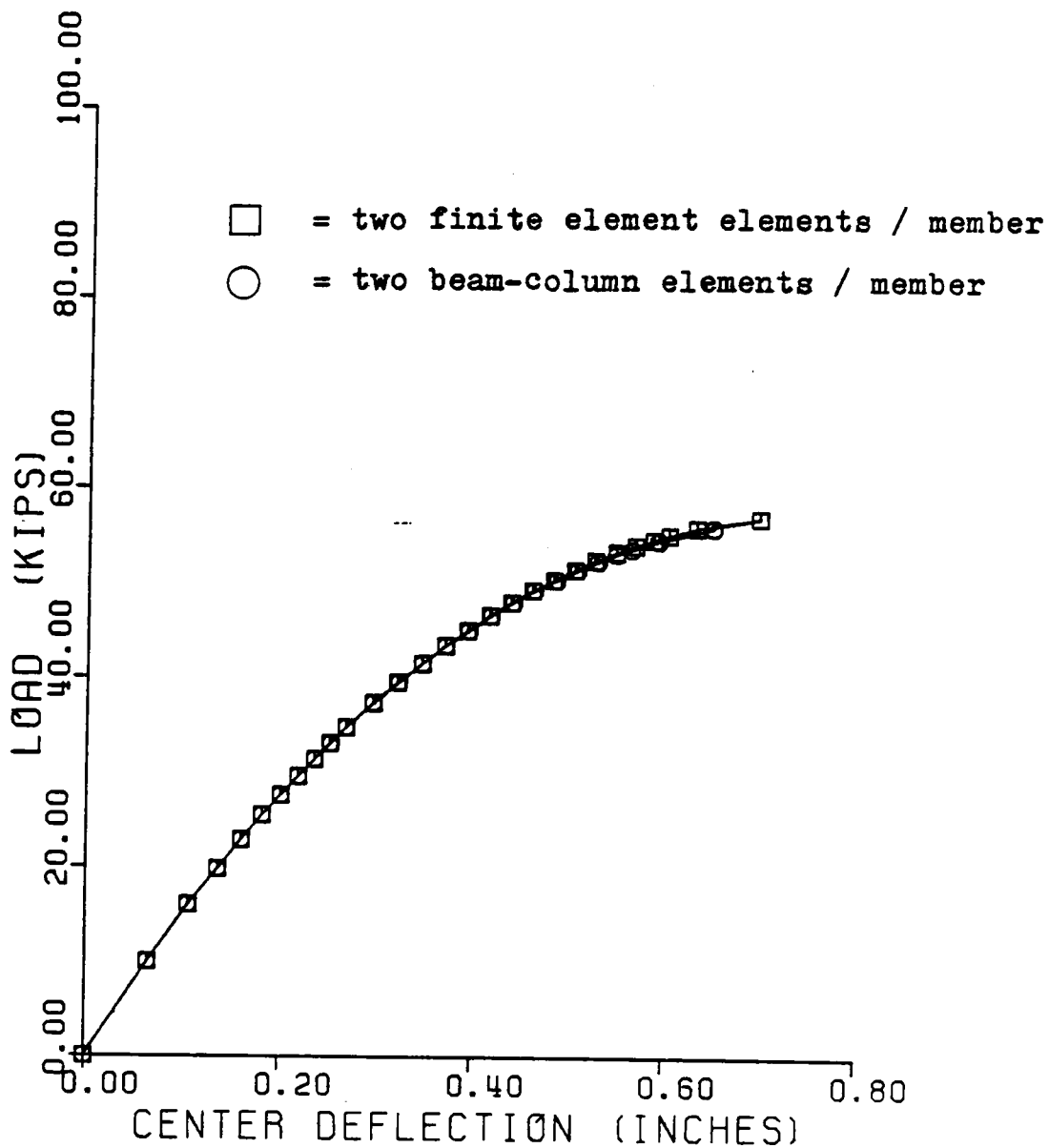


Fig. 7.13 Comparison of Models

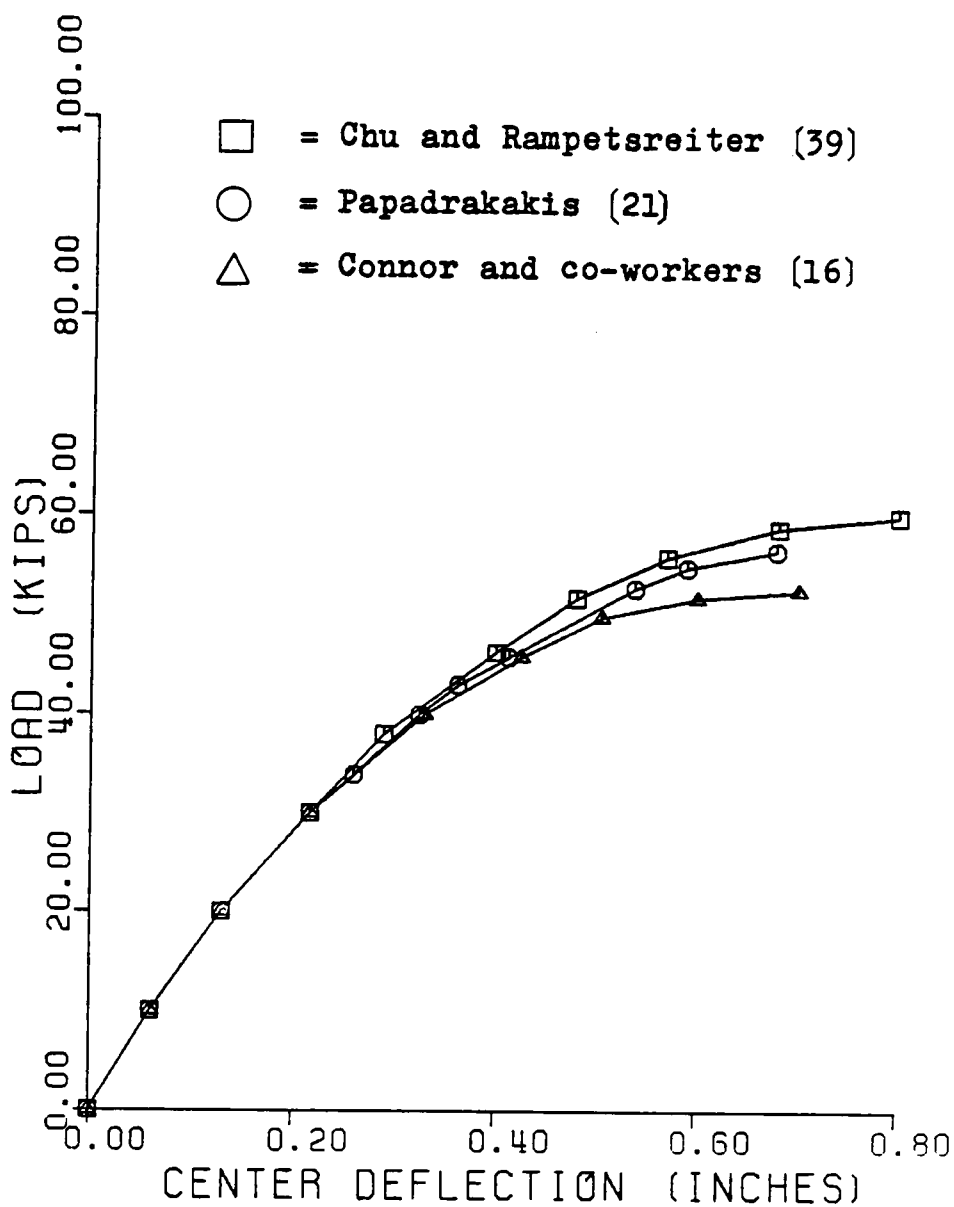


Fig. 7.14 Load Deflection Curves



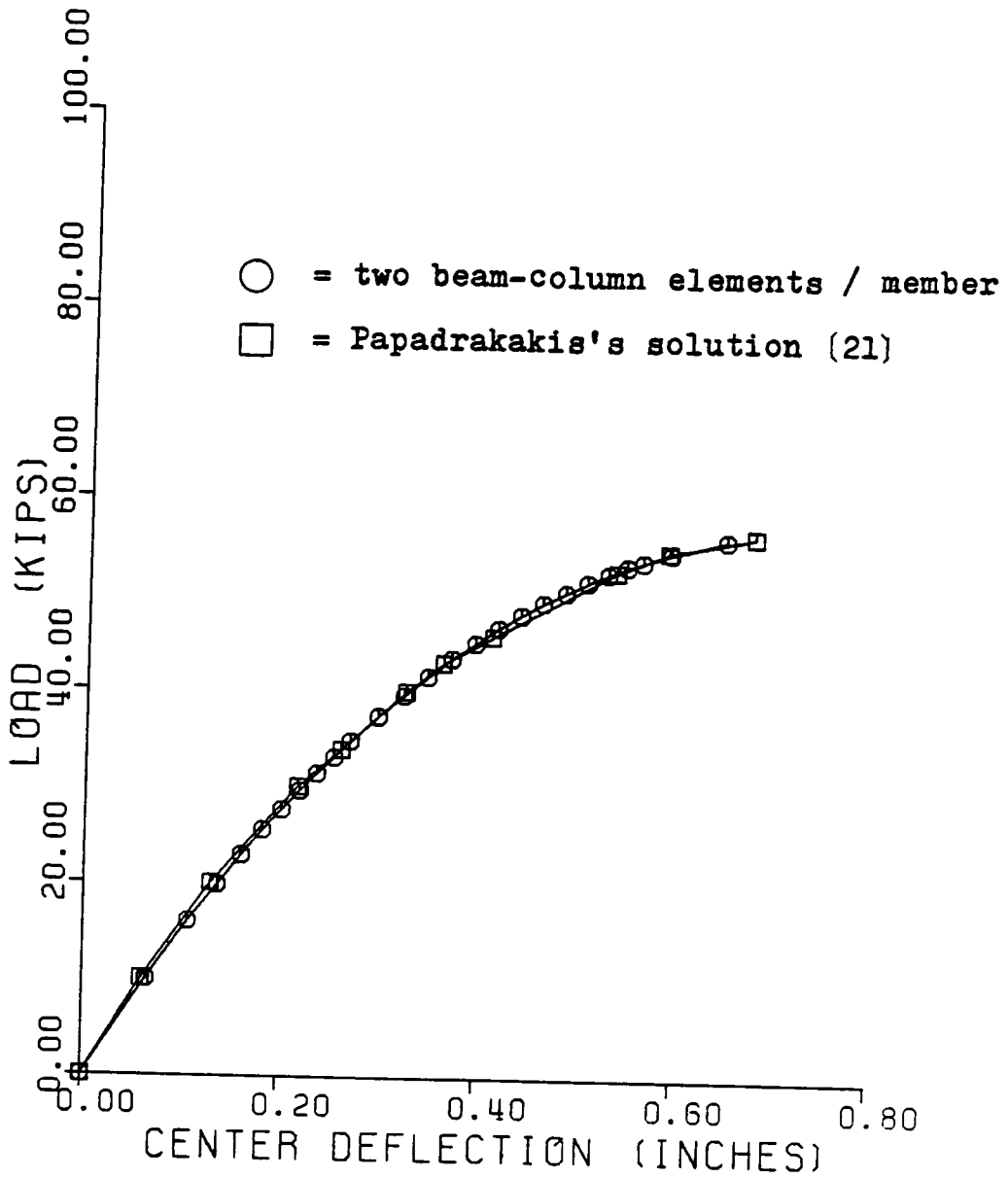


Fig. 7.15 Load Deflection Curves

Employing two elements per member, the buckling load obtained by the beam-column model is 56.14 lb and by the finite element model is 57.24 lb. The execution time for both models are as follows:

one element:

finite element model: 10.55 sec

beam-column model: 14.55 sec

two elements:

finite element model: 71.53 sec

beam-column model: 75.59 sec

#### 7.5 EXAMPLE 4: RETICULATED DOME

Shallow truss domes were investigated by Holzer [27],[51], Hangai [52], Paradiso [53]. In this study a shallow frame dome designed in accordance with the Specification of the Aluminum Association [54] is considered. The geometry of the reticulated frame dome is given in Fig. 7.16. The dome is subjected to a single vertical load at the center.

Fig. 7.17 presents the finite element solution obtained using one and two elements per member; it is seen that single element is inadequate to model the structural behavior.

Fig. 7.18 presents the beam-column solution obtained using one and two elements per member; similar results for both meshes are obtained. Fig. 7.19 presents a comparison of two finite elements to one beam-column element; it shows again that fewer elements in the beam column model are often needed than in the finite element model for satisfactory accuracy.

In addition Fig. 7.19 shows a comparison of two extreme cases, the shallow truss dome, considered by Uliana [55], with the present frame dome, is presented; it shows that the behavior of the frame dome remains linear when the truss dome reaches the limit point due to a central, concentrated load. Clearly, the frame dome is much stronger than the truss dome. The execution time for both models are given as follows:

one element:

finite element model: 52.71 sec

beam-column model: 60.53 sec

two elements:

finite element model: 339.24 sec

beam-column model: 459.41 sec

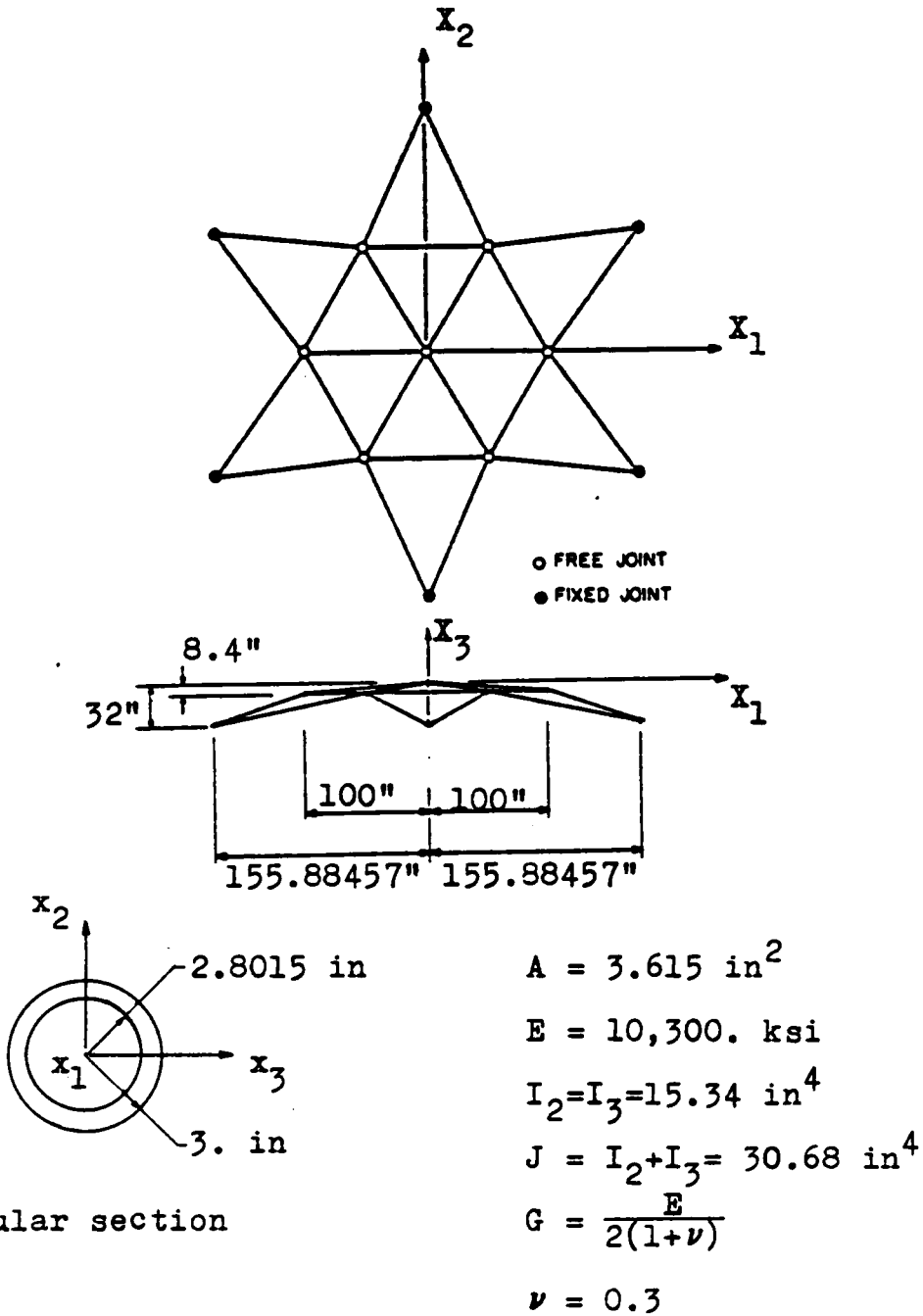


Fig. 7.16 Reticulated Dome

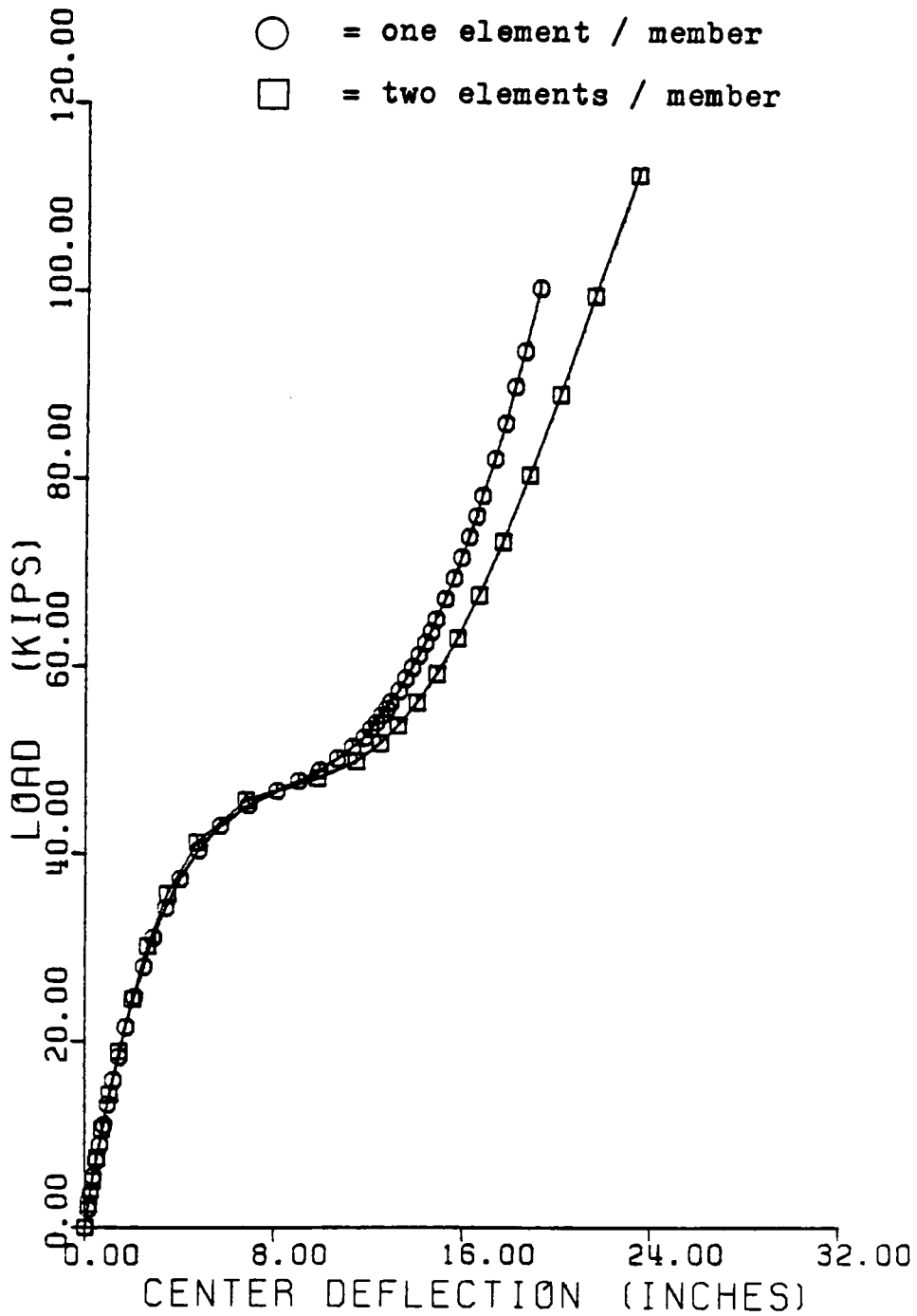


Fig. 7.17 Finite Element Model

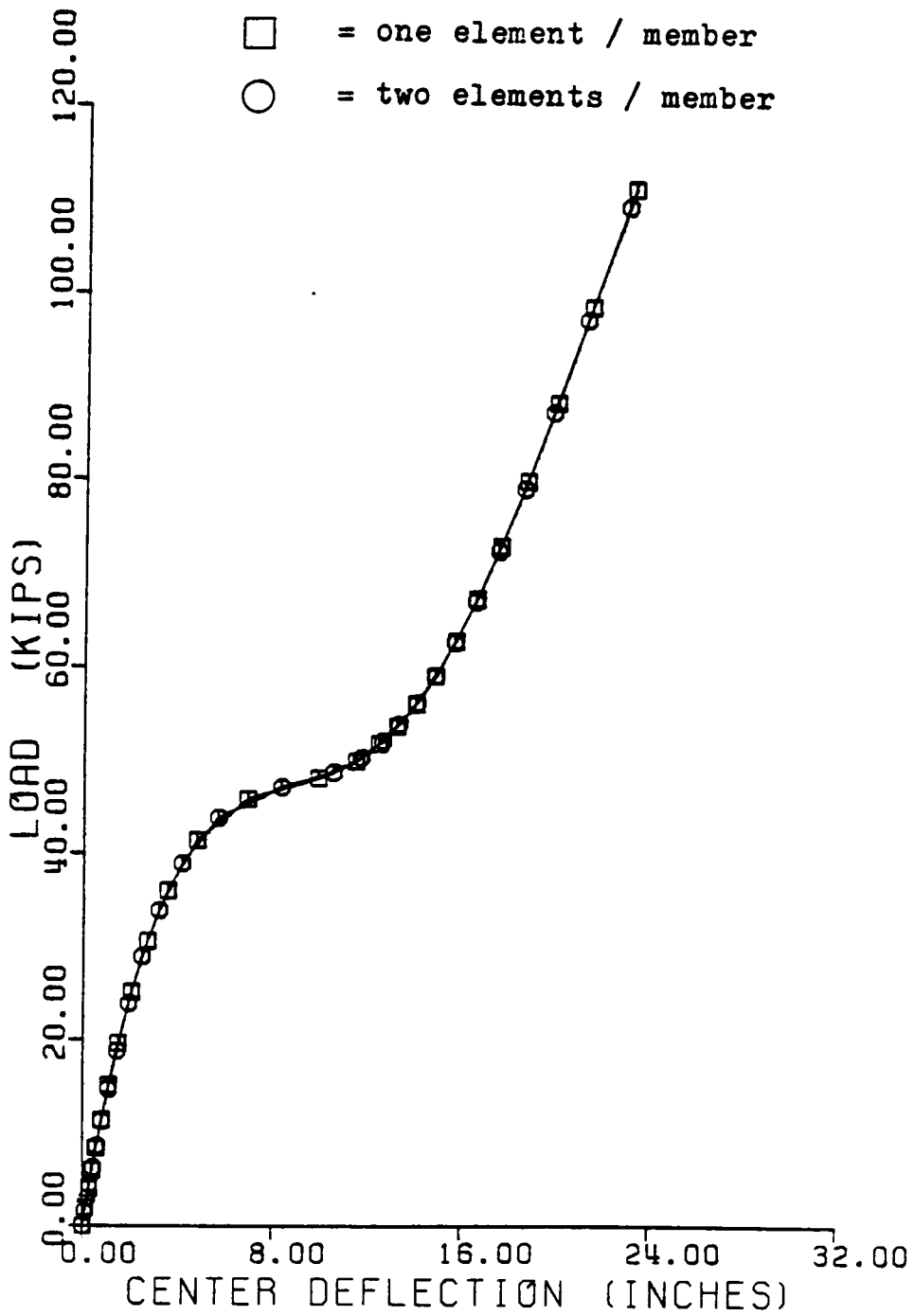


Fig. 7.18 Beam-Column Model

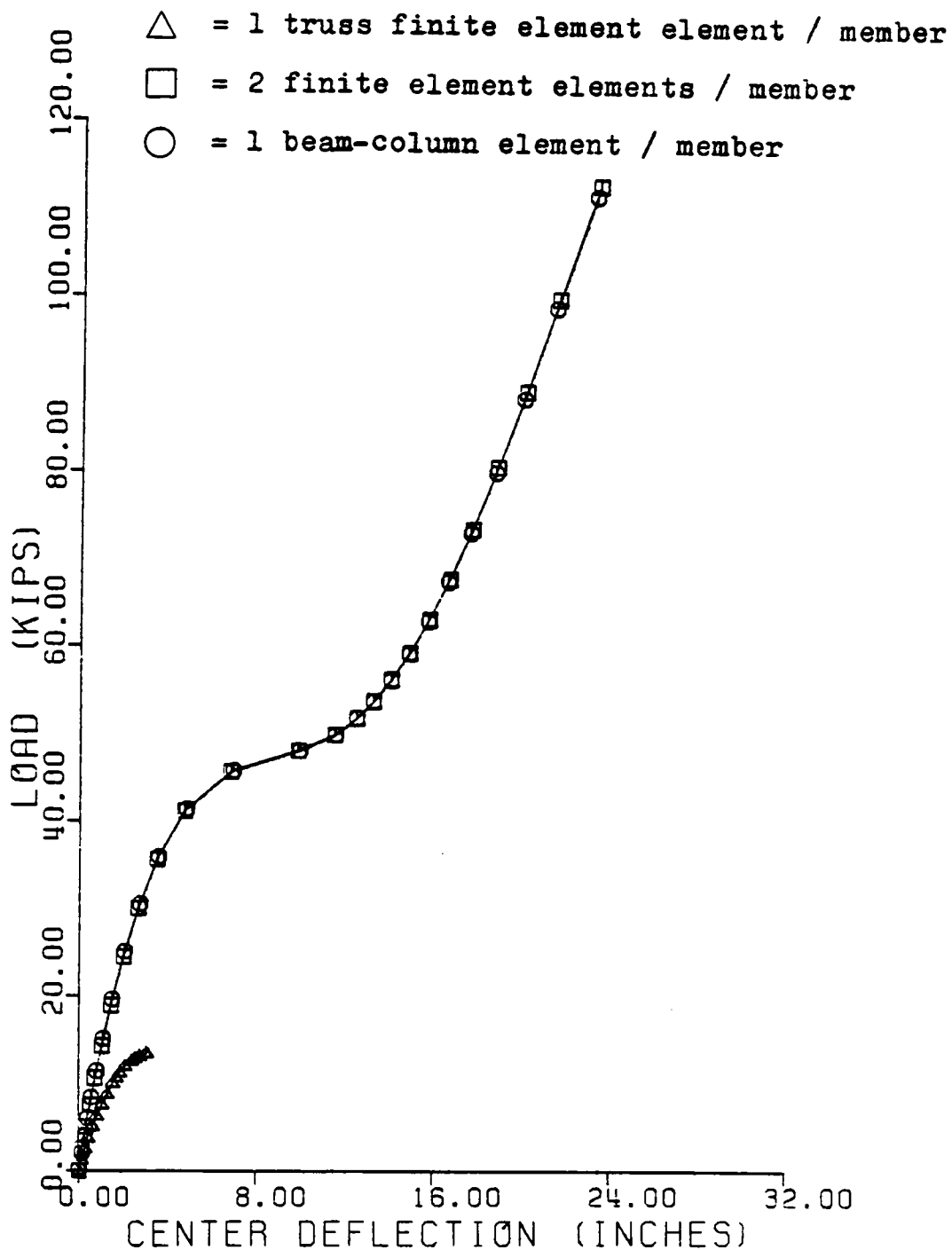


Fig. 7.19 Comparison of Frame Dome with Truss Dome

Chapter VIII  
PROGRAM DEVELOPMENT

8.1 INTRODUCTION

The program developed by Butler [29], for the geometrically nonlinear static analysis of plane frames, has been modified and extended to the GNSSF (Geometrically Nonlinear Static Analysis of Space Frames) program using U.L. formulation as in appendix F. The modified program has the capability to analyze space frames with arbitrary large displacement. It is restricted to continuous, elastic structures, i.e., only geometrical nonlinearity is considered.

The GNSSF computer program is written in the WATFIV version of FORTRAN code having the following options:

1. Newton-Raphson method or modified Riks/Wempner method
2. beam-column model or finite element model

Subroutines FORCES and STIFF can be approached by using either the beam-column model or the finite element model. Subroutines FACTOR, REDUCE, and BACSUB, adapted from Bathe [38] and Subroutines SOLVE, TEST, DISPLC, DISPLB, UNBALF, UPDATE, DOTPRD, adapted from Butler [29], will not be described here.



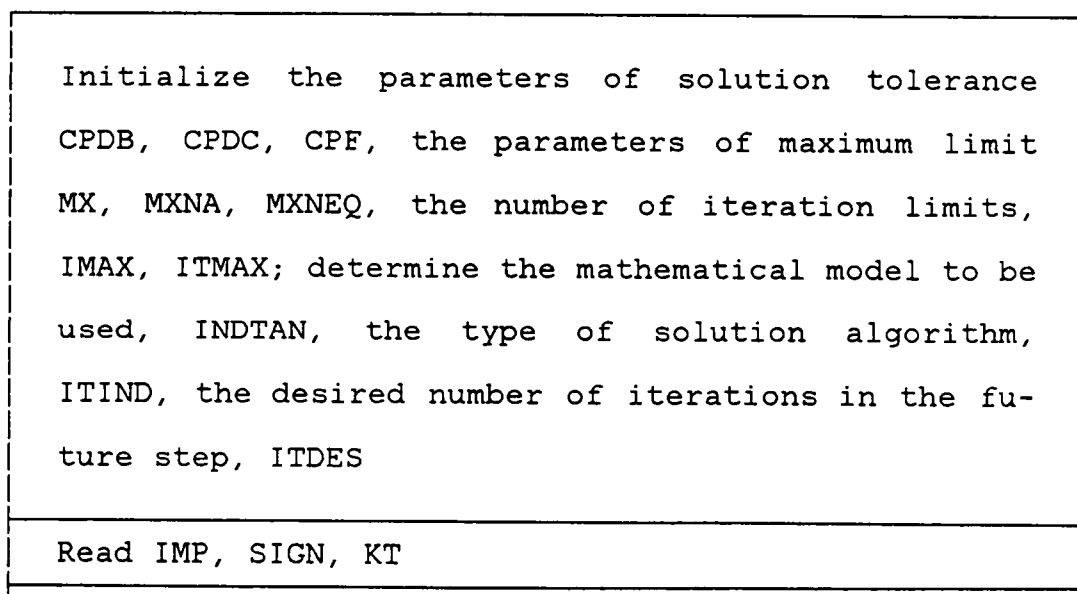
## 8.2 PROGRAM STRUCTURE

Fig. 8.1 presents the tree chart of the GNSSF computer program. The MAIN program and each subroutine are described by Nassi-Schneiderman (N-S) diagrams in this section [28].

### 8.2.1 Main Program

Function: Initialize and read in control parameters; select the mathematical model, INDTAN; determine the desired number of iterations in the future step; call DATA; initialize the joint and element orientation matrices to the undeformed state; determine the type of solution algorithm, call either NEWRAP or RIKWEM.

N-S diagram:



Echo input data	
Call DATA	
Initialize F, FP, FPI, D, DD, DDO, Z, to zero; let Q=QJ	
Initialize the joint orientation matrix at time zero to identity martrix	
Calculate the components of the initial orientation matrix of element I	
If ITIND=0	
then	else
Call NEWRAP	Call RIKWEM

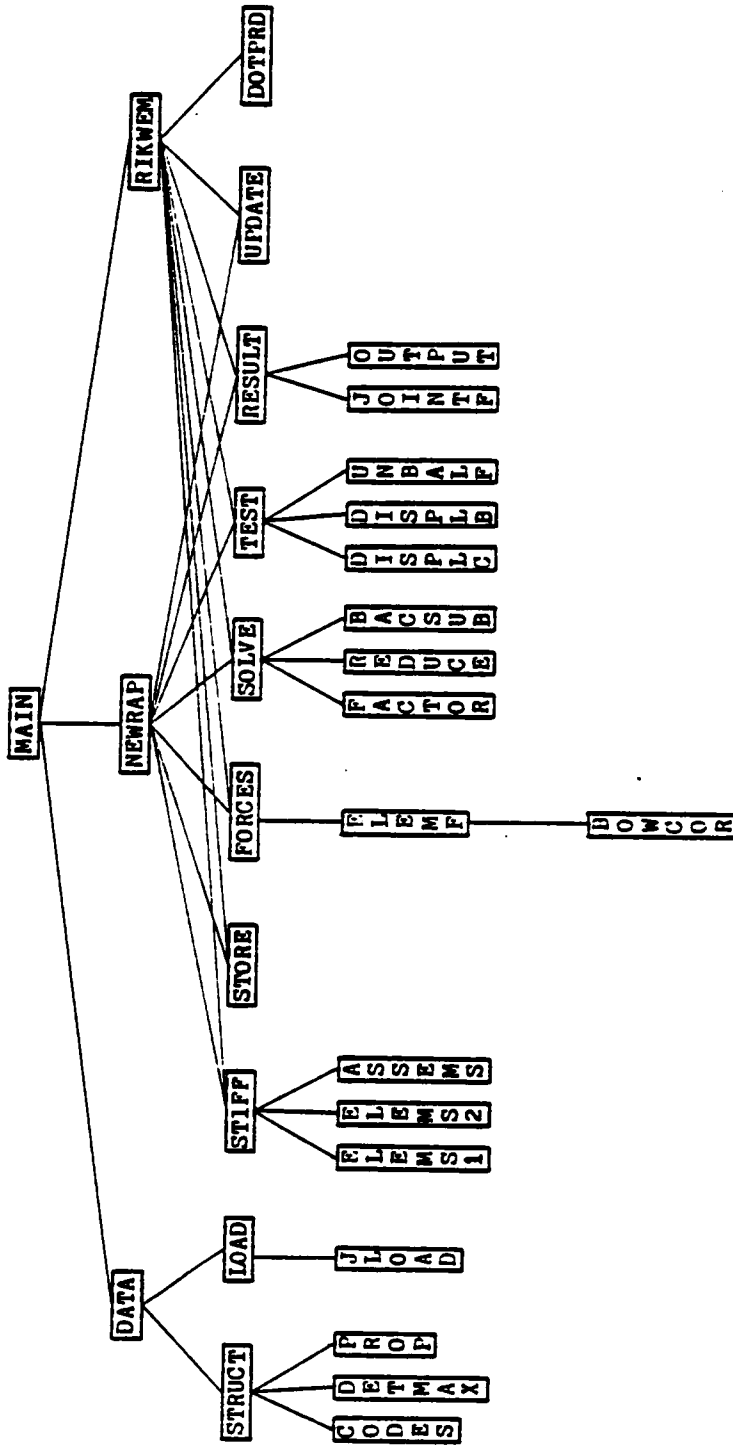


Fig. 8.1 Program Structure

8.2.2 Subroutine DATA

Function: Read and echo NE, NJ; call STRUCT and LOAD.

Input arguments: MX, MXNA, MXNEQ

Output arguments: MINC, JCODE, MCODE, NEQ, MAXA, NKT, X,  
 AREA, ZI2, ZI3, ZJ, EMOD, GMOD, ZJ, NE,  
 NJ, Q, C01, C02, C03, C1PI, C2PI, C3PI,  
 C1, C2, C3, ELENGO, ELENG

N-S diagram:

Read and echo NE and NJ	
If $NE \leq MX$ and $NJ \leq MX$	
then	else
Call STRUCT	Print error message
Call LOAD	

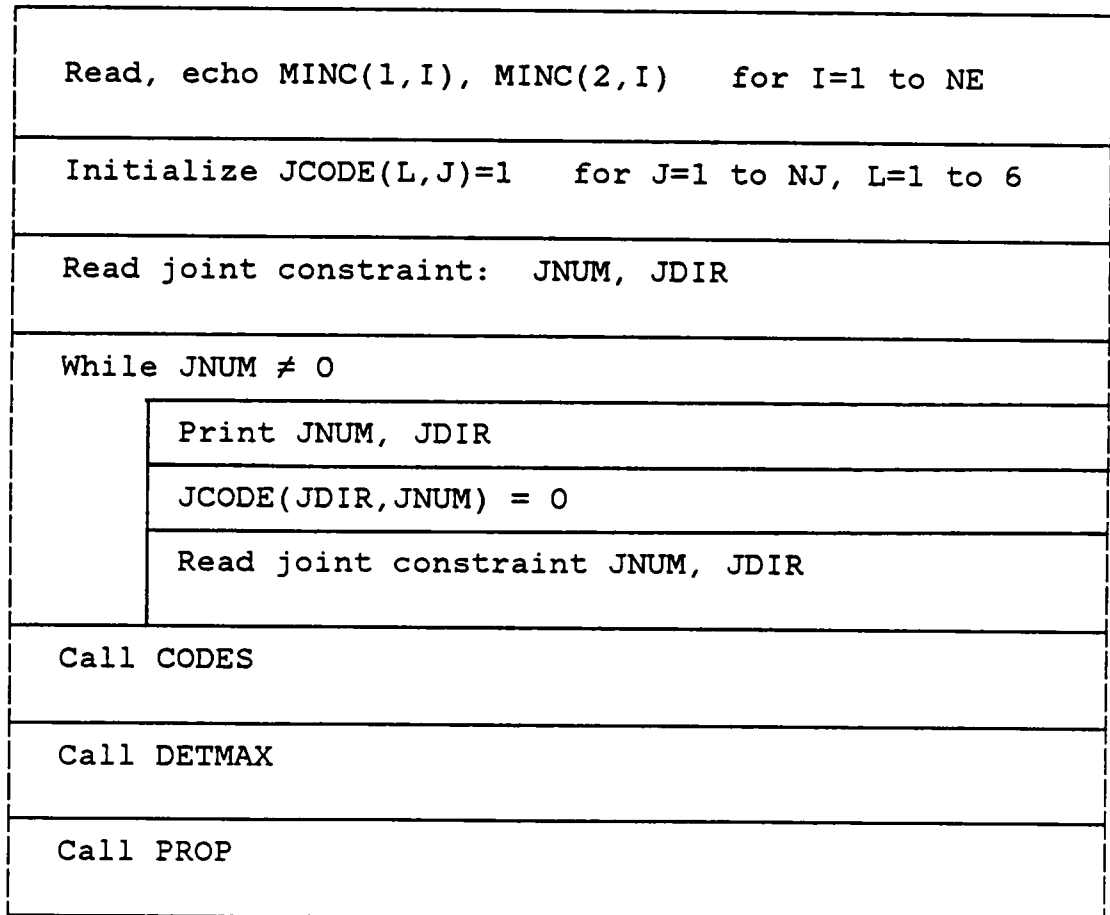
8.2.3 Subroutine STRUCT

Function: Read echo, and process the structural data.

Input arguments: NE, NJ, MXNEQ

Output arguments: MINC, JCODE, MCODE, NEQ, MAXA, NKT, X,  
 AREA, ZI2, ZI3, ZJ, EMOD, GMOD, ELENGO,  
 ELENG, CO1, CO2, CO3, C1PI, C2PI, C3PI,  
 C1, C2, C3

N-S diagram:



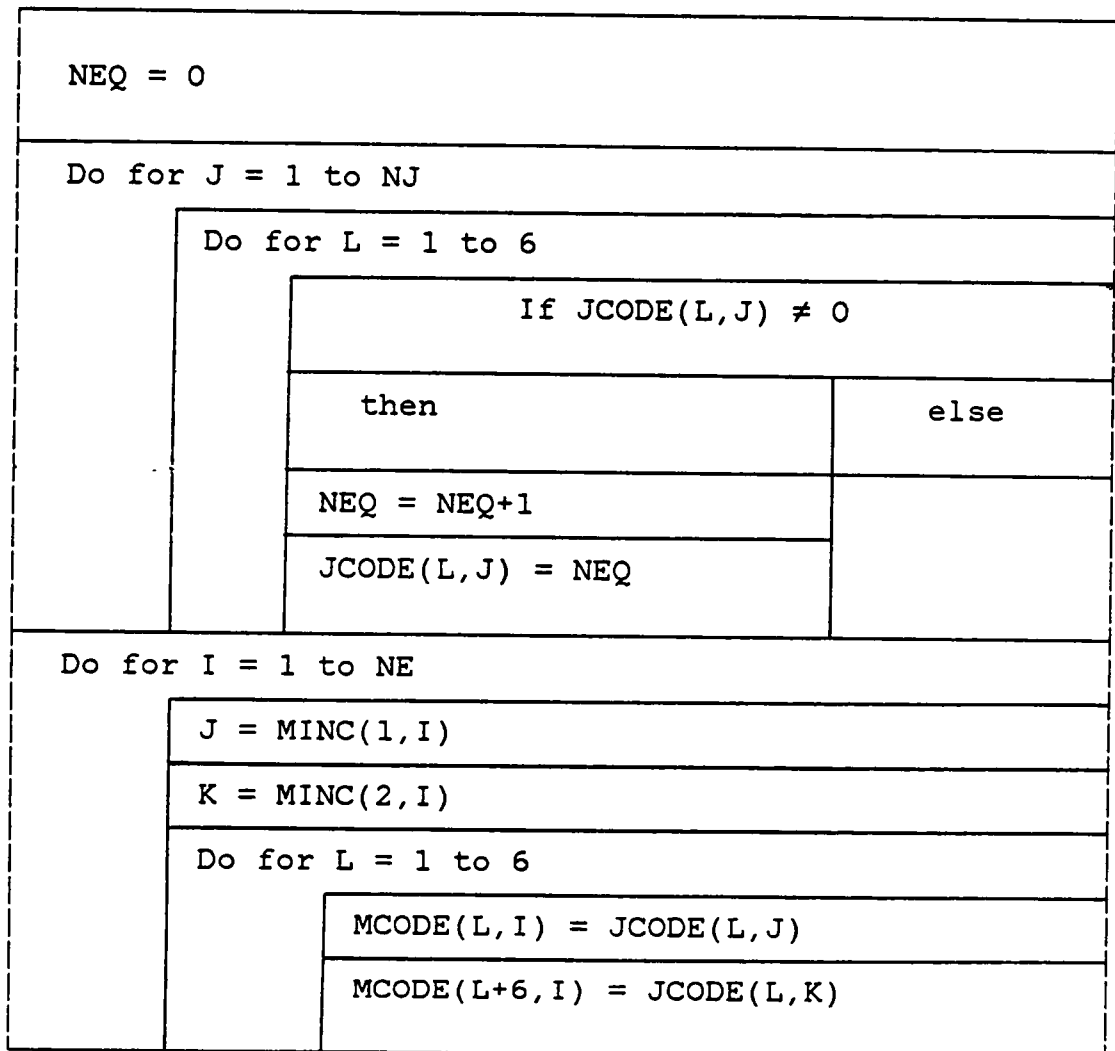
8.2.4 Subroutine CODES

Function: Generate the JCODE and MCODE.

Input arguments: JCODE, MINC, NE, NJ, MXNEQ

Output arguments: MCODE, JCODE, NEQ

N-S diagram:



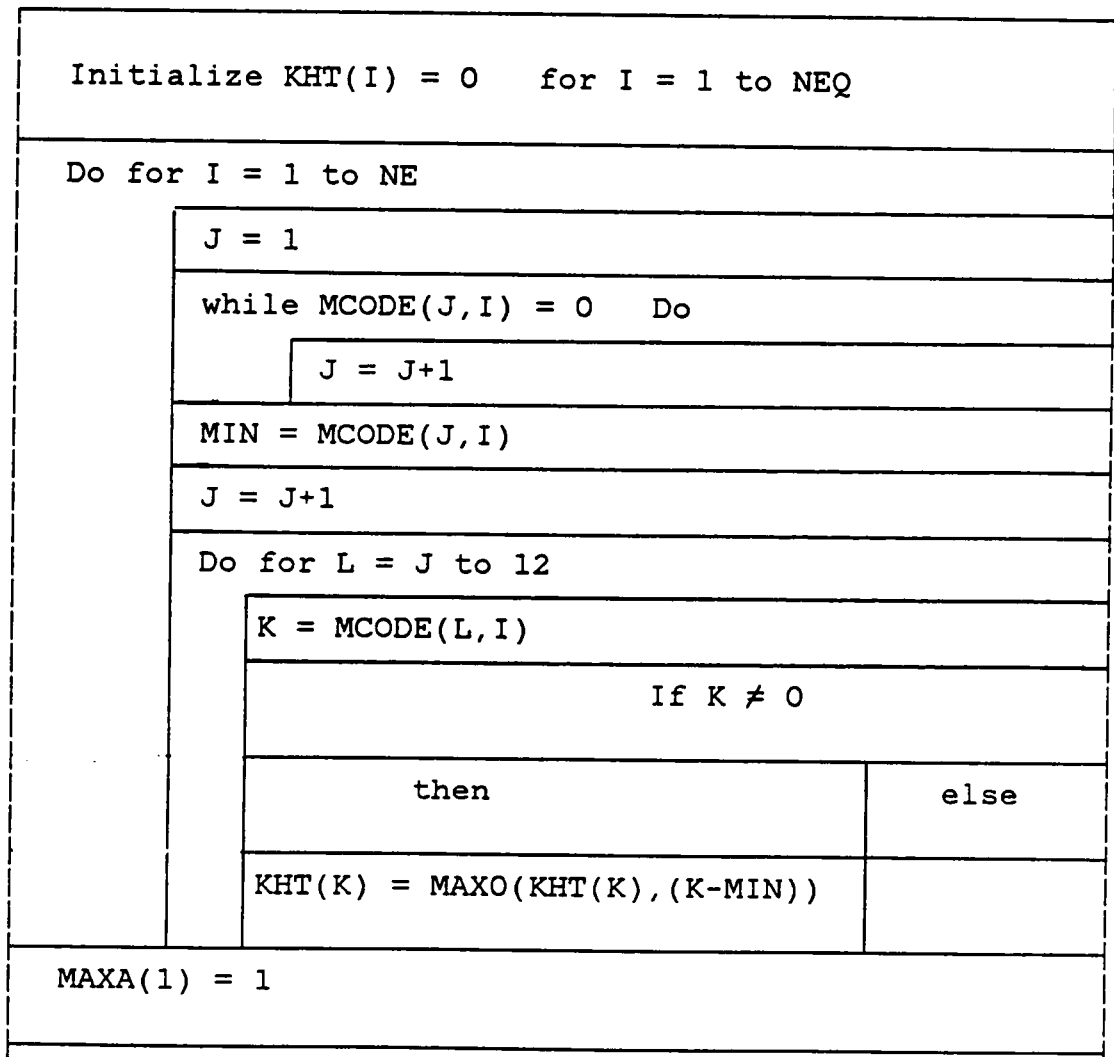
8.2.5 Subroutine DETMAX

Function: To calculate column heights,  $KHT(NEQ)$ ; addresses of diagonal elements in banded matrix  $K$  (column heights known),  $MAXA(NEQ+1)$ ; and the number of elements below skyline of matrix  $K$ ,  $NKT$ .

Input arguments:  $NEQ$ ,  $NE$ ,  $MCODE$

Output arguments:  $MAXA$ ,  $NKT$

N-S diagram:



```
DO for I = 1 to NEQ
```

```
    MAXA(I+1) = MAXA(I) + KHT(I)+ 1
```

```
NKT = MAXA(NEQ+1) - 1
```



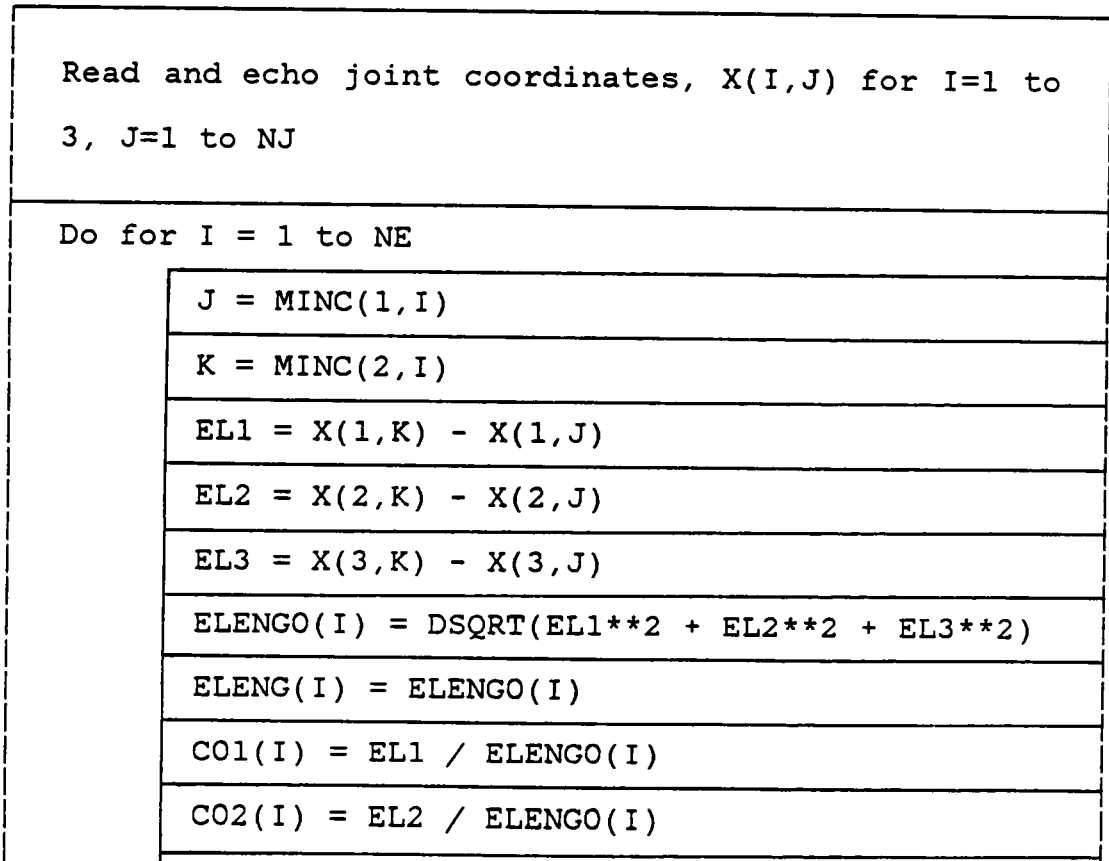
8.2.6 Subroutine PROP

Function: Read and echo the joint coordinates, X; compute and echo the undeformed element length and direction cosines, ELENGO, CO1, CO2, CO3; store ELENGO, CO1, CO2, CO3 in the previous iteration state, ELENG, C1PI, C2PI, C3PI; read and print element properties: AREA, ELENGO, EMOD, GMOD, ZI2, ZI3, ZJ.

Input arguments: NJ, NE, MINC

Output arguments: X, ELENGO, ELENG, CO1, CO2, CO3, C1PI, C2PI, C3PI, AREA, EMOD, GMOD, ZI2, ZI3, ZJ

N-S diagram



$CO3(I) = EL3 / ELENGO(I)$
$C1PI(I) = CO1(I)$
$C2PI(I) = CO2(I)$
$C3PI(I) = CO3(I)$
Read AREA(I), EMOD(I), GMOD(I), ZI2(I), ZI3(I), ZJ(I)
Echo AREA(I), ELENGO(I), EMOD(I), GMOD(I), ZI2(I), ZI3(I), ZJ(I)

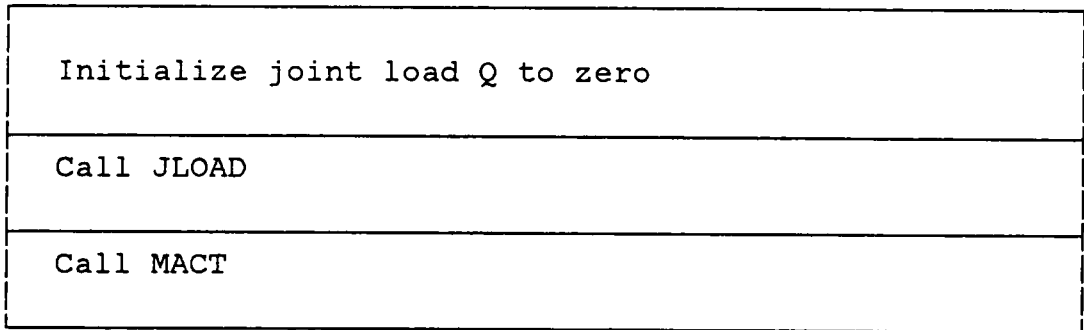
8.2.7 Subroutine LOAD

Function: Initialize to zero the joint load vector,  $Q$ ; Call  
JLOAD and MACT.

Input arguments: NEQ, JCODE

Output arguments:  $Q$

N-S diagram:



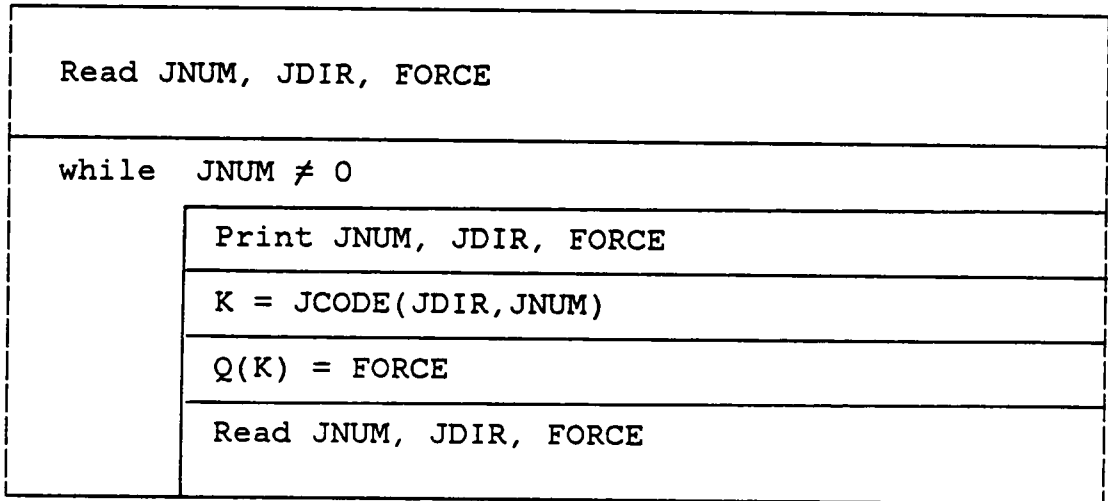
8.2.8 Subroutine JLOAD

Function: Read JNUM, JDIR, and FORCE; while JNUM = 0, print JNUM, JDIR, FORCE; store FORCE in Q, and read JNUM, JDIR, FORCE.

Input arguments: JCODE

Output arguments: Q

N-S diagram:



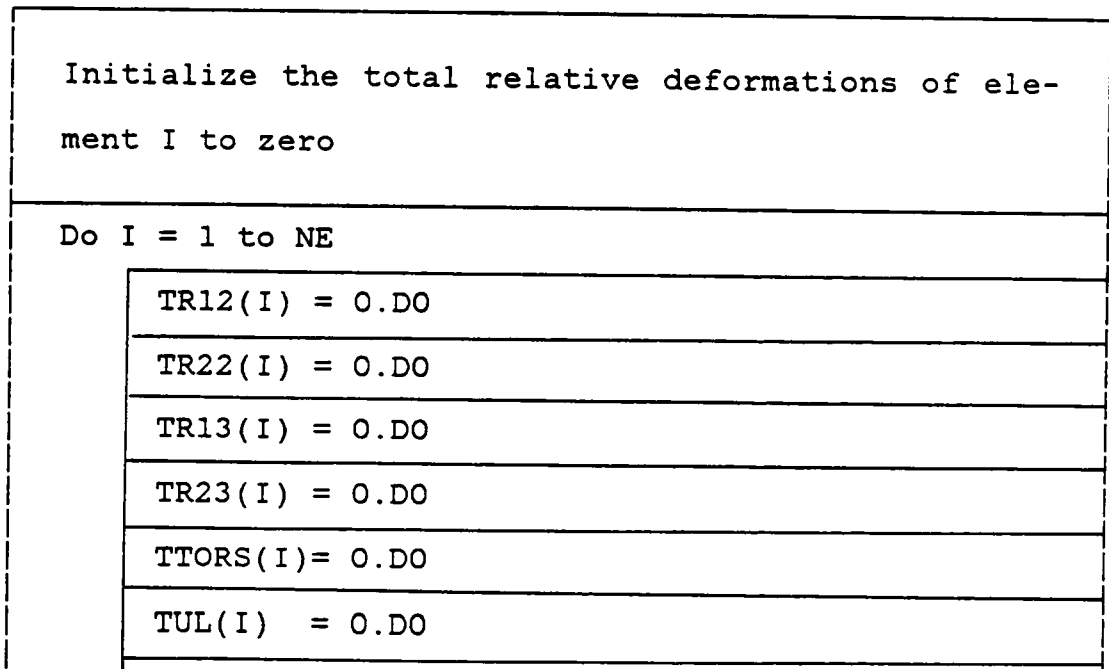
8.2.9 Subroutine NEWRAP

Function: Perform the Newton-Raphson or modified Newton-Raphson method to trace the prelimit equilibrium path. Call STIFF, STORE, SOLVE, FORCES, TEST, RESULT, and UPDATE.

Input arguments: AREA, C01, C02, C03, C1PI, C2PI, C3PI, C1, C2, C3, ELENGO, ELENG, EMOD, GMOD, JCODE, MAXA, MCODE, MINC, ITIND, INDTAN, KT, IMP, Q, QJ, CPDB, CPDC, CPF, DD, DDO, DQQR, FQR, IMAX, ITMAX, NE, NJ, NEQ, NKT, NUPD, QI, QIMAX, DQI, NPRINT, IELS, AP, OLAMD

Output arguments: none

N-S diagram:



TCB2(I) = 0.DO	
TCB3(I) = 0.DO	
TU(I) = 0.DO	
NN = 1	
Call FORCES	
Call UPDATE	
While QI $\leq$ QIMAX	
NC = NUPD	
ITCT = 0	
ICI = 1	
While ICI $\neq$ 0 and ITCT $\leq$ ITMAX	
If NC $\geq$ NUPD	
then	else
Call STIFF	
NC = 0	
Call STORE	
Call SOLVE	
Do for I = 1 to NEQ	
D(I) = D(I)+DD(I)	
Call FORCES	
Call TEST	
NC = NC+1	

ITCT = ITCT+1	
Do for I = 1 to NEQ	
FPI(I) = F(I)	
Do for I = 1 to NEQ	
FP(I) = F(I)	
Call RESULT	
QI = QI+DQI	
If ICI ≠ 0	
then	else
Print 'last solution is not converged'	
Stop	
Call UPDATE	

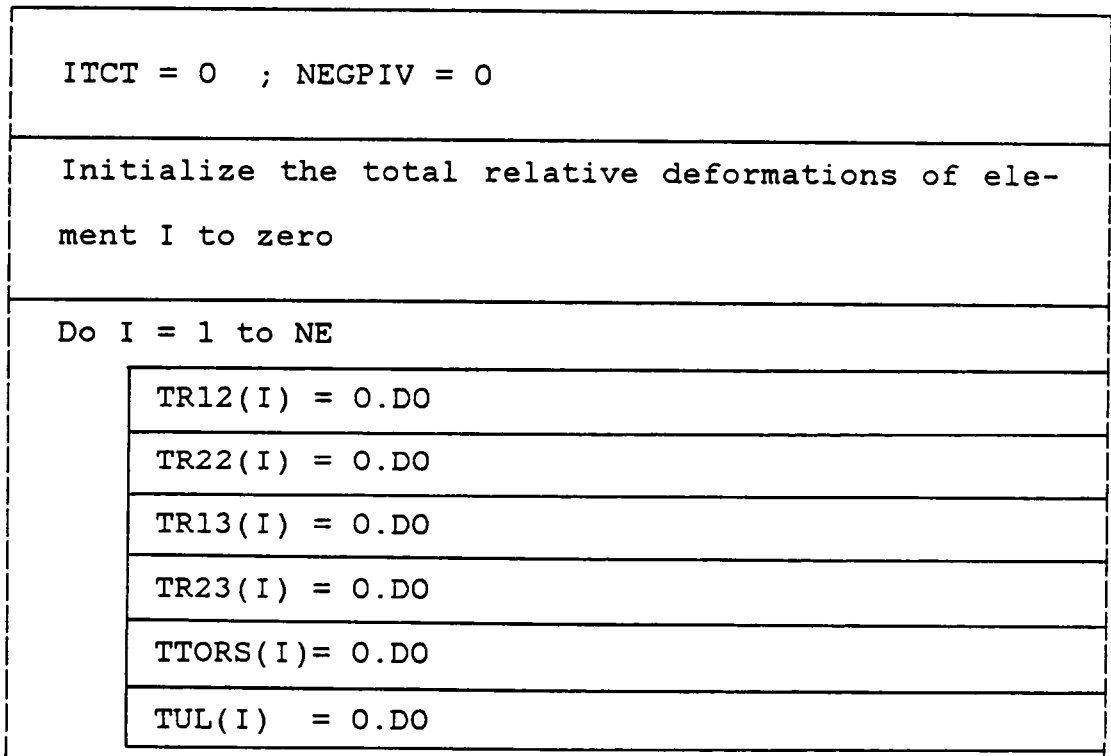
8.2.10 Subroutine RIKWEM

Function: Perform the Riks/Wempner or modified Riks/Wempner iteration on the normal plane. Call STIFF, STORE, SOLVE, FORCES, TEST, RESULT, and UPDATE.

Input arguments: AREA, CO1, CO2, CO3, C1PI, C2PI, C3PI, C1, C2, C3, ELENCO, ELENG, EMOD, GMOD, JCODE, MAXA, MCODE, MINC, ITIND, INDIAN, KT, IMP, Q, QJ, CPDB, CPDC, CPF, DD, DDO, DQQR, FQR, IMAX, ITMAX, NE, NJ, NEQ, NKT, NUPD, QI, QIMAX, DQI, NPRINT, SIGN, IELS, AP, OLAMD

Output arguments: none

N-S diagram:





TCB2(I) = 0.DO	
TCB3(I) = 0.DO	
TU(I) = 0.DO	
NN = 1	
Call FORCES	
While QI ≤ QIMAX and ITCT ≤ ITMAX	
Call STIFF	
Call STORE	
Call SOLVE : SKT DDO1 = Q for DDO1	
If ITCT = 0	
then	else
DS=DQI*[DOTPRD (DDO1.DDO1)+1] <sup>1/2</sup>	DQI=SIGN(NEGPIV+1)*DS/ [DOTPRD(DDO1.DDO1)+1] <sup>1/2</sup>
DQI1 = DQI	
Do for I = 1 to NEQ	
DDO(I) = DQI*DDO1(I)	
D(I) = D(I)+DDO(I)	
QI = QI+DQI	
NN = 1	
Call FORCES	
ICI = 1 ; NC = NUPD ; IT = 0	
While ICI ≠ 0 and IT ≤ IMAX	
Call UPDATE	

If $NC \geq NUPD$	
then	else
Call STIFF	
NC = 0	
Call STORE	
Call SOLVE: SKT DD1 = Q for DD1	
Call SOLVE: SKT DD2 = QT-F for DD2	
DQI = -DOTPRD(DDO.DD2)/(DOTPRD(DDO.DD1)+DQI1)	
Do for I = 1 to NEQ	
DD(I) = DQI*DD1(I)+DD2(I)	
D(I) = D(I) + DD(I)	
NN = 2	
Call FORCES	
QI = QI + DQI	
Call TEST	
NC = NC + 1 ; IT = IT + 1	
ITCT = ITCT + 1	
Do for I = 1 to NEQ	
FP(I) = F(I)	
Call RESULT	
If ICI $\neq$ 0	
then	else

Print 'last solution is not converged'
Stop
$DS = DS * (ITDES/IT)^{1/2}$

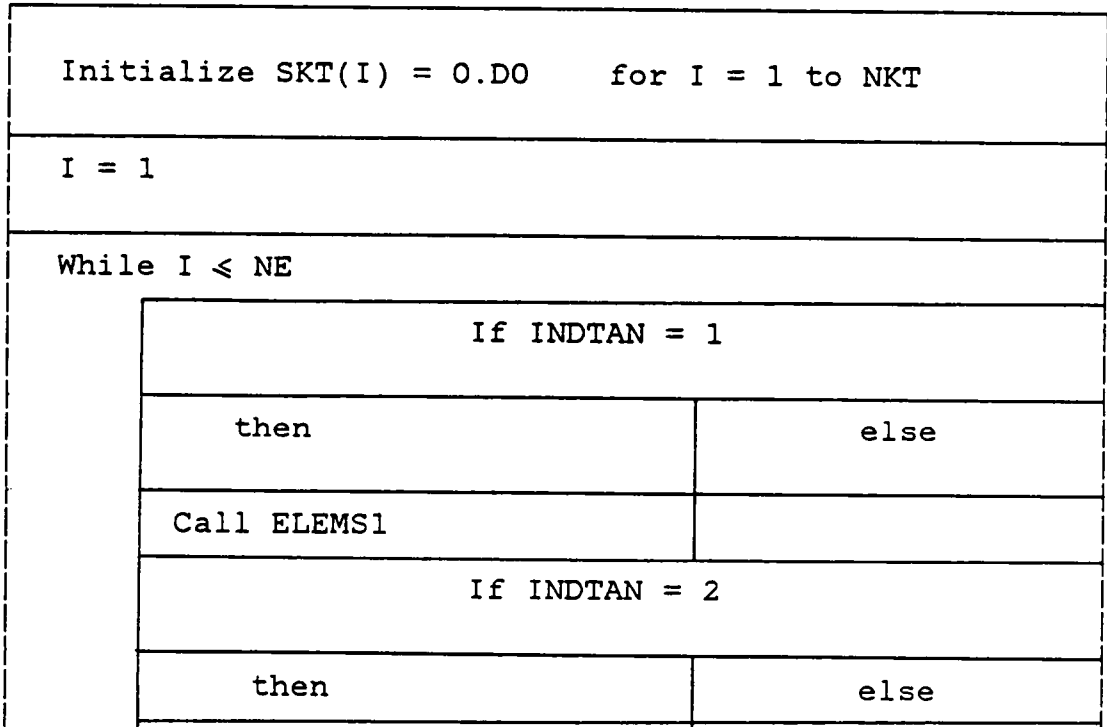
8.2.11 Subroutine STIFF

Function: Initialize the system tangent stiffness matrix, SKT, to zero; for each element call either ELEMS1 for INDTAN = 1, or ELEMS2 for INDTAN = 2, and call ASSEMS.

Input arguments: AREA, BP12, BP22, BP13, BP23, C1, C2, C3, ELENGO, EMOD, GMOD, MAXA, MCODE, TR12, TR22, TR13, TR23, SP12, SP22, SP13, SP23, ST12, ST22, ST13, ST23, RM32, RM33, QL, TU, INDTAN, KT, ZI2, ZI3, ZJ, IELS, NE, NKT, NPRINT, OLAMD

Output arguments: SKT

N-S diagram:



	Call ELEMS2	
	Call ASSEMS	
	I = I + 1	
If NPRINT = 3		
then		else
Print SKT(K) for K = 1 to NKT		

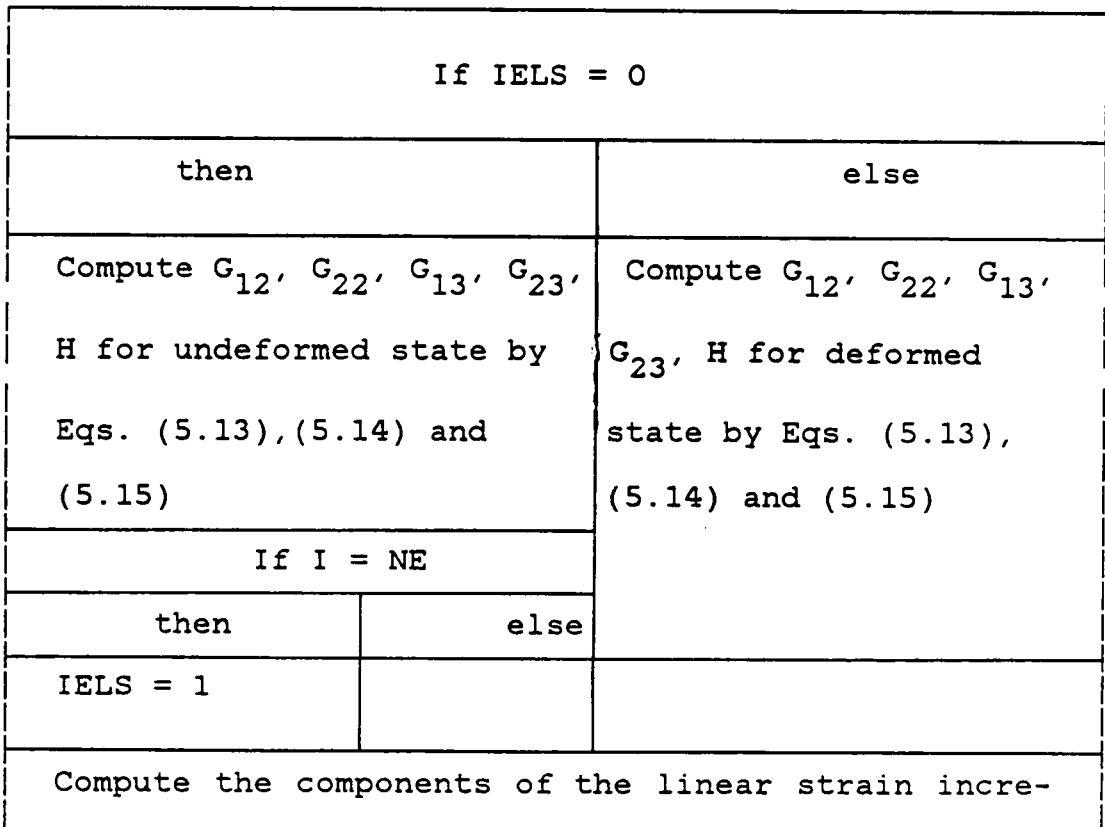
8.2.12 Subroutine ELEMS1

Function: Compute the components of the element tangent stiffness matrix for the beam-column model.

Input arguments: AREA, BP12, BP22, BP13, BP23, C1, C2, C3, EMOD, GMOD, TR12, TR22, TR13, TR23, SP12, SP22, SP13, SP23, ST12, ST22, ST13, ST23, ZI2, ZI3, ZJ, I, IELS, NE, RM32, RM33, QL, NPRINT, KT, OLAMD, ELENGO

Output arguments: G

N-S diagram:



<p>mental stiffness matrix <math>k_L</math> in the local coordinate system <math>g'(1)</math> to <math>g'(20)</math>, by Eq. (5.31)</p>	
<p>Compute the components of the linear strain incremental stiffness matrix <math>K_L</math> in the global coordinate system, <math>g(1)</math> to <math>g(42)</math>, by Eqs.(5.41) and (5.42)</p>	
<p>If <math>KT(2) &gt; 0</math></p>	
<p>then</p>	<p>else</p>
<p>Add the contribution of the nonlinear strain incremental stiffness matrix in the global coordinate system, <math>K_{NL}</math>, to <math>G(1), G(2), G(3), G(10), G(11), G(18)</math> by Eq. (5.38)</p>	
<p>If <math>NPRINT = 3</math></p>	
<p>then</p>	<p>else</p>
<p>Print <math>G(KI)</math> of element I, for <math>KI = 1</math> to 42</p>	

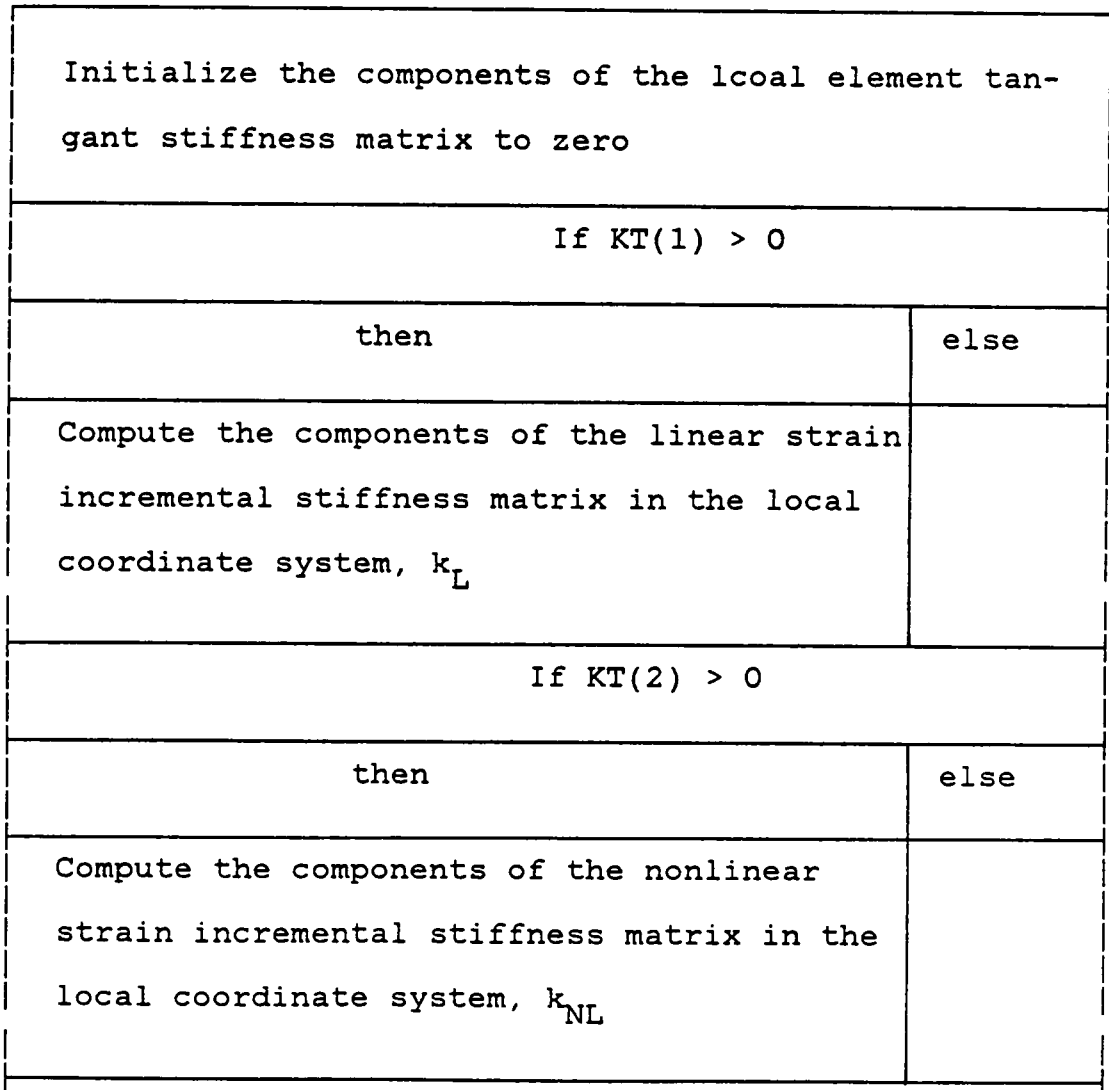
8.2.13 Subroutine ELEMS2

Function: Compute the components of the element tangent stiffness matrix for the finite element model.

Input arguments: AREA, C1, C2, C3, ELENGO, EMOD, GMOD, I, ZI2, ZI3, ZJ, KT, QL, OLAMD

Output arguments: G

N-S diagram:





Compute the components of the global element tangent stiffness matrix,  $K^i$

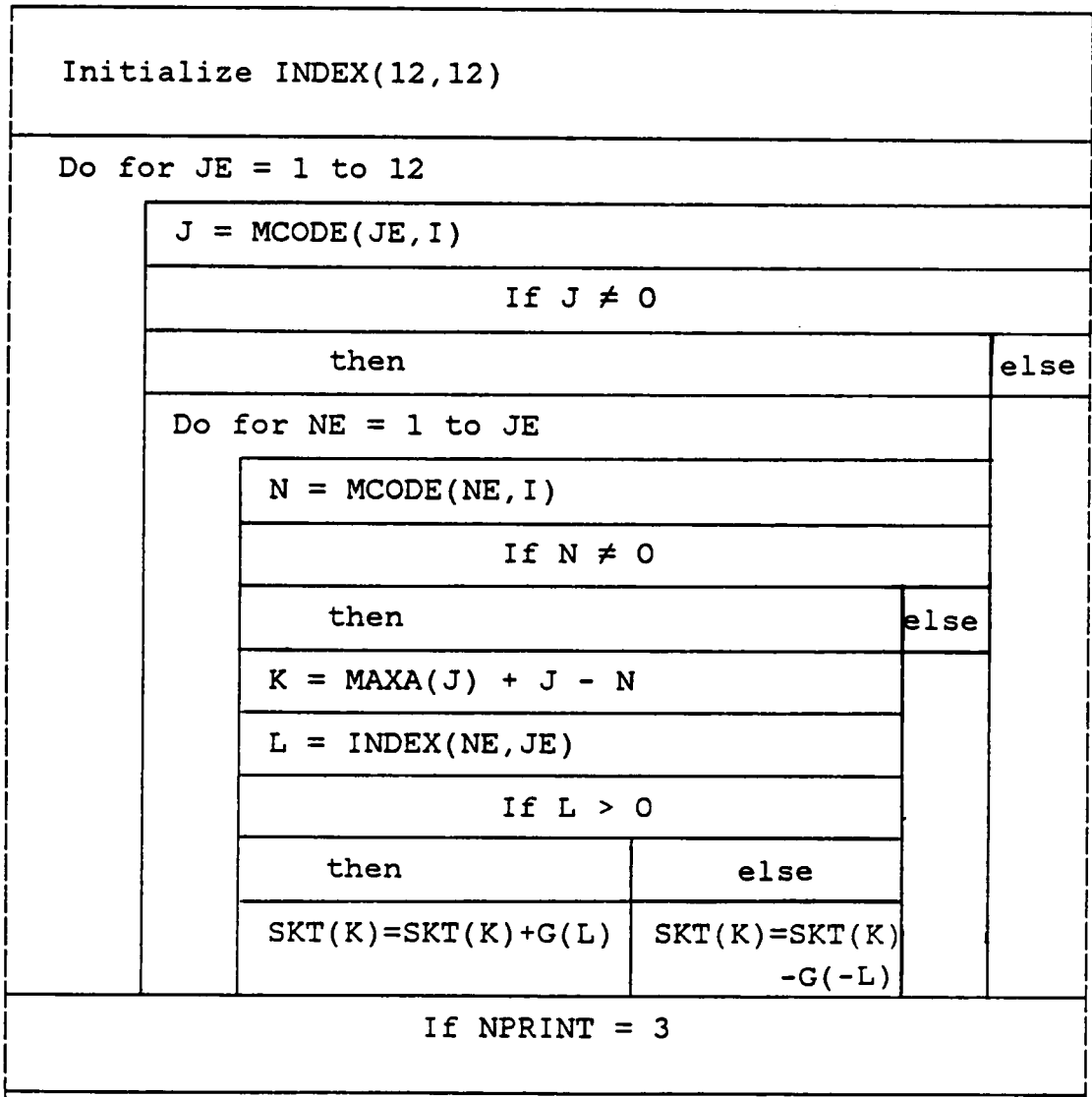
8.2.14 Subroutine ASSEMS

Function: Initialize INDEX; assign the global stiffness coefficients,  $G(L)$ , of element to the system stiffness column matrix, SKT, by INDEX, MAXA.

Input arguments: G, I, MAXA, MCODE, NKT, NPRINT

Output arguments: SKT

N-S diagram:



then	else
Print SKT(KI) from element I, for KI = 1 to NKT	

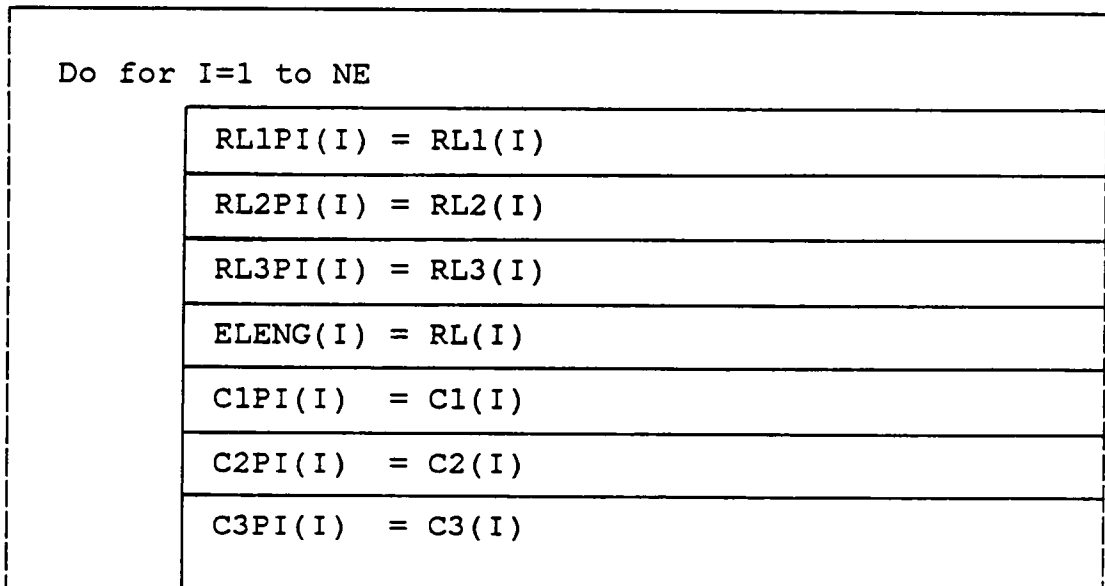
### 8.2.15 Subroutine STORE

Function: Store the previous iteration length and direction cosines of element I.

Input arguments: NE, RL1, RL2, RL3, RL, C1, C2, C3

Output arguments: RL1PI, RL2PI, RL3PI, ELENG, C1PI,  
C2PI, C3PI

N-S diagram:



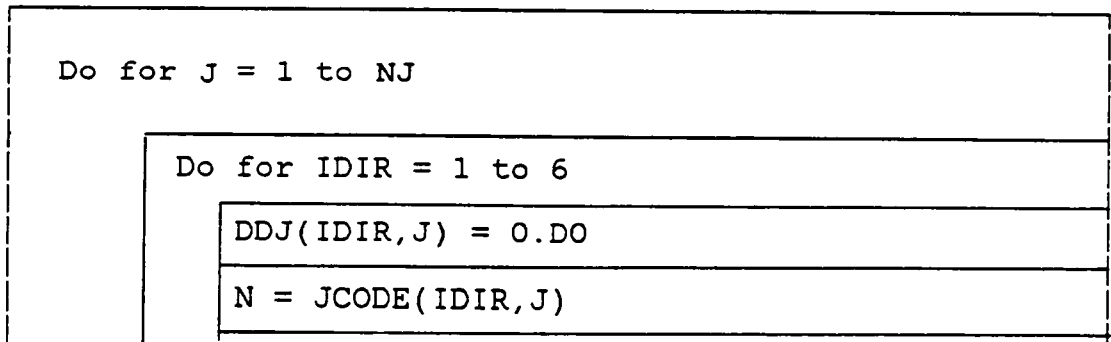
8.2.16 Subroutine FORCES

Function: Generate the incremental joint displacement matrix DDJ; calculate the incremental nodal rotations at joint J and yield the components of the rotation matrix of joint J; for each iteration employing the last joint orientation matrix, AP, and rotation matrix of joint J, RJ, to generate the new joint orientation matrix; initialize the vector of the equivalent nodal element force, F, to zero; for each element call ELEMF.

Input arguments: AREA, C01, C02, C03, C1PI, C2PI, C3PI,  
 OL, DD, DDO, NN, ELENCO, ELENG, EMOD, GMOD,  
 MCODE, ZI2, ZI3, ZJ, DQQR, FQR, IMAX, NE,  
 NJ, NEQ, INDTAN, AP, NPRINT, OLAMD

Output arguments: BP12, BP22, BP13, BP23, C1, C2, C3, F, FG,  
 FL, TR12, TR22, TR13, TR23, TTORS, SP12,  
 SP22, SP13, SP23, ST12, ST22, ST13, ST23,  
 RM32, RM33, QL, OLAMD

N-S diagram:



If $N \neq 0$	
then	else
If $NN=1$ and $ITIND=1$	
then	else
DDJ(IDIR, J) = DDO(N)	DDJ(IDIR, J) =DD(N)
Calculate the incremental nodal rotations at joint J, OME1, OME2, OME3, by Eq. (3.11)	
Yield the incremental nodal rotation matrix at joint J, RJ, by Eq. (3.9)	
Generate the current joint orientation matrix, AP, by Eq. (3.12)	
Do for L = 1 to 3, M = 1 to 3	
APNEW(L, M, J) = 0.DO	
Do for K = 1 to 3	
APNEW(L, M, J) = RJ(L, K)*AP(K, M, J)+APNEW(L, M, J)	
Do for L = 1 to 3, K = 1 to 3	
AP(L, K, J) = APNEW (L, K, J)	
Do for I = 1 to NEQ	
F(I) = 0.DO	
I = 1	
While I ≤ NE	
Call ELEMF	

	$I = I + 1$
--	-------------

8.2.17 Subroutine ELEM

Function: Employing the MCODE to get the global incremental displacements and direction cosines of element I; generate the increments of the local displacements and the components of the rotation matrix of element I; obtain the current element orientation matrix; employing the MINC to calculate the total relative deformations of element I; compute the local internal element forces by either beam-column model or finite element model, due to the deformations; transform them to the global element forces, and compute the equivalent nodal element forces.

Input arguments: AREA, C1PI, C2PI, C3PI, DD, DDO, NN,  
 ELENGO, ELENG, EMOD, GMOD, MCODE, ZI2, ZI3,  
 ZJ, DQQR, FQR, I, IMAX, INDTAN, OLAMD, C01,  
 C02, C03, OL, AP

Output arguments: BP12, BP22, BP13, BP23, C1, C2, C3, F,  
 FL, FG, TR12, TR22, TR13, TR23, TTORS,  
 SP12, SP22, SP13, SP23, ST12, ST22, ST13,  
 ST23, RM32, RM33, QL, OLAMD

N-S diagram:

```
Do for J = 1 to 12
```

```
  DE(J) = 0.DO
```

K = MCODE(J,I)	
If K ≠ 0	
then	else
If NN=1 and ITIND=1	
then	else
DE(J)=DDO(K)	DE(J)=DD(K)
RL1PI(I) = ELENG(I)*C1PI(I)	
RL2PI(I) = ELENG(I)*C2PI(I)	
RL3PI(I) = ELENG(I)*C3PI(I)	
RL1(I) = RL1PI(I)+DE(7)-DE(1)	
RL2(I) = RL2PI(I)+DE(8)-DE(2)	
RL3(I) = RL3PI(I)+DE(9)-DE(3)	
RL(I) = DSQRT(RL1(I)**2+RL2(I)**2+RL3(I)**2)	
C1(I)=RL1(I)/RL(I)	
C2(I)=RL2(I)/RL(I)	
C3(I)=RL3(I)/RL(I)	
Calculate the local incremental element displacements by Eq. (2.56)	



Yield the incremental element rotation matrix,  $R$ ,  
by Eqs. (3.23) to (3.25)

Generate the current element orientation matrix,  
OLANEW, by Eq. (3.26)

Do for L=1 to 3, J=1 to 3

    OLANEW(L,J,I) = 0.DO

        Do for K=1 to 3

            OLANEW(L,J,I)=R(L,K,I)\*OLAMD(K,J,I)  
                                  +OLANEW(L,J,I)

Do for L=1 to 3, J=1 to 3

    OLAMD(L,J,I)=OLANEW(L,J,I)

TUL = (ELENGO(I) - RL(I)) / ELENGO(I)

L = MINC(1,I)

K = MINC(2,I)

Compute the total relative element deformations,  
TR12, TR13, TR22, TR23, TTORS

    If INDTAN = 1

        then

        Call BOWCOR  
        Compute RM12, RM22, RM13,  
        RM23, FL(1,I) by Eqs.

        else

        Compute the RM12, RM22,  
        RM13, RM23, FL(1,I)  
        using the secant stiff-

(5.1), (5.2) and (5.4)	ness matrix by Eq. (4.20)
Compute the remaining local element forces	
Transform the local element forces, FL, to the global element forces, FG, by Eq. (4.27)	
Do for J = 1 to 12	
K = MCODE(J, I)	
If K ≠ 0	
then	else
Compute the equivalent nodal element forces F(K)=F(K)+FG(J, I)	

### 8.2.18 Subroutine BOWCOR

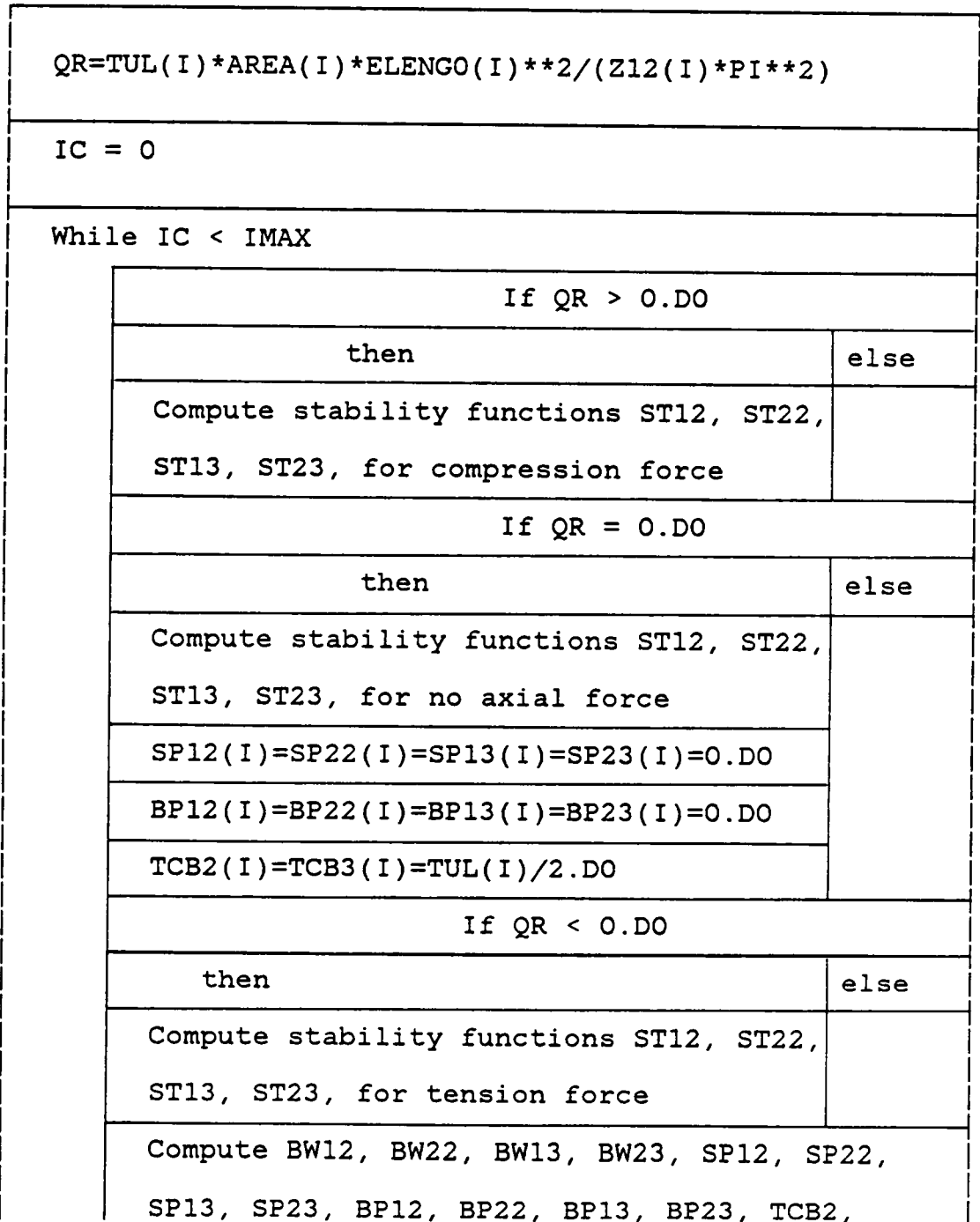
Function: Calculate stability functions, ST12, ST22, ST13, ST23; bowing functions, BW12, BW22, BW13, BW23; length correction factors for the bowing actions, TCB2, TCB3; axial force, QPRINT.

Input arguments: AREA, ELENGO, EMOD, GMOD, ZI2, ZI3, ZJ, I, TUL,

Output arguments: BP12, BP22, BP13, BP23, TR12, TR22, TR13,

TR23, SP12, SP22, SP13, SP23, ST12, ST22,  
 ST13, ST23, TCB2, TCB3

N-S Diagram:



TCB3, TCB2P, TCB3P	
Compute FQRC and DQ (ref. [29])	
$QR = QR + DQ$	
$DQQRC = DABS(DQ/QR)$	
$IC = IC + 1$	
If $DQ/QR \leq DQQR$ and $FQRC \leq FQR$	
then	else
$IC = IMAX + 1$	
If $DQ/QR > DQQR$ or $FQRC > FQR$	
then	else
Stop	

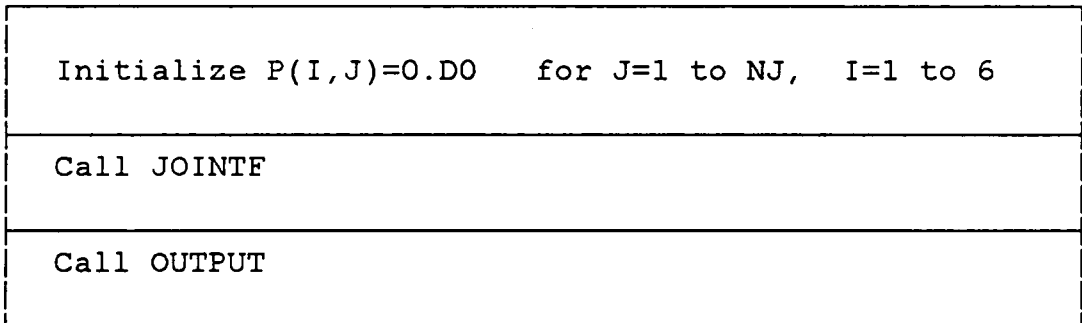
8.2.19 Subroutine RESULT

Function: Initialize the joint force matrix, P, to zero;  
call JOINTF and OUTPUT.

Input arguments: NJ, NE, MINC, JCODE, FG, D, NPRINT,  
QI, QT, NEQ, FL, IMP

Output arguments: none

N-S diagram:



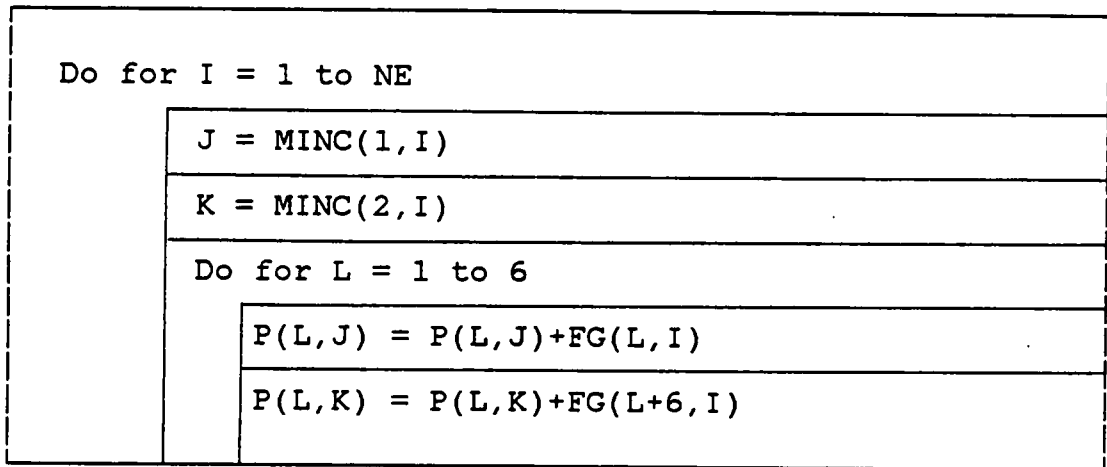
8.2.20 Subroutine JOINTF

Function: Assign the global element forces, FG, to the joint force matrix, P, via MINC.

Input arguments: NE, MINC, FG

Output arguments: P

N-S diagram:



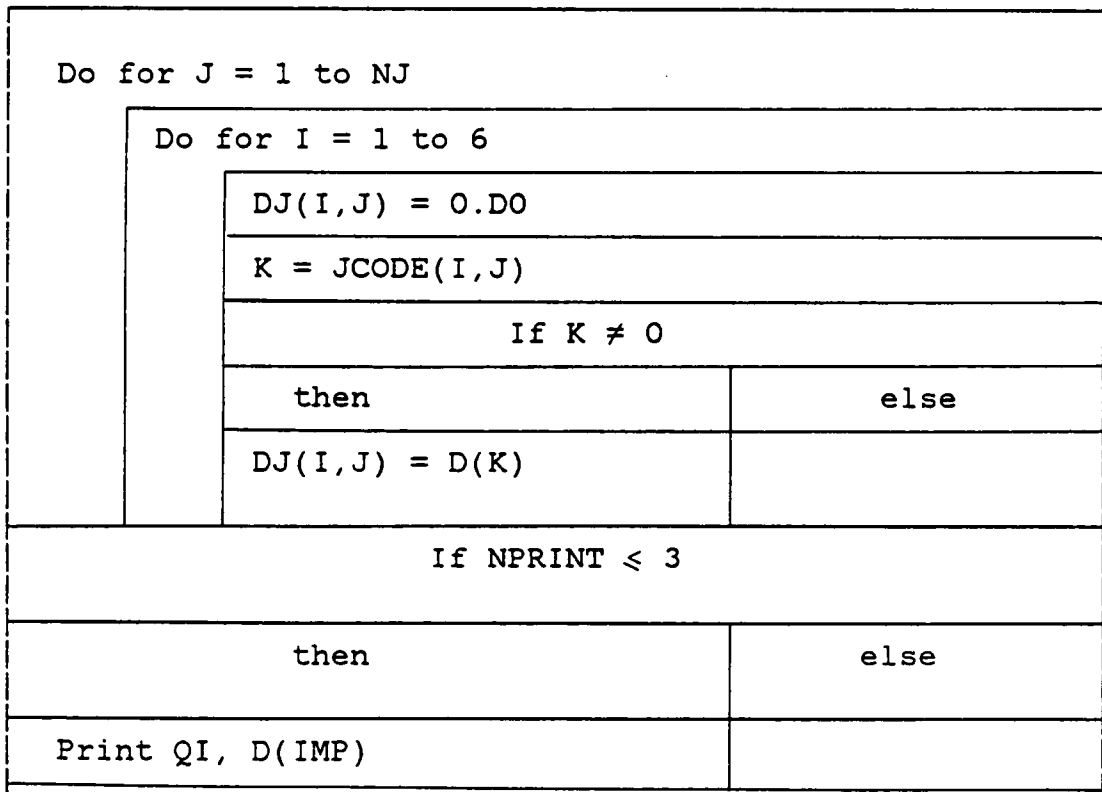
8.2.21 Subroutine OUTPUT

Function: Assign the generalized displacements, D, to the joint displacements, DJ, via JCODE; print the loading parameter, QI, generalized displacements, D; print the current generalized external force vector, QT; print the global joint displacements (including joint constraints), DJ; print the local element forces, FL; print the joint forces, P.

Input arguments: NJ, JCODE, D, NPRINT, QI, QT, NEQ, FL,  
NE, P, IMP

Output arguments: none

N-S diagram:



If NPRINT = 4	
then	else
Print QI	
Print QT(I) for I = 1 to NEQ	
Print DJ(I,J) for I=1 to 6, J=1 to NJ	
Print FL(I,J) for I=1 to 12, J=1 to NE	
Print P(I,J) for I=1 to 6, J=1 to NJ	



## Chapter IX

### CONCLUSION

#### 9.1 CONCLUSION

Updated Lagrangian and a total Lagrangian formulations of a geometrically nonlinear three-dimensional beam element, undergoing large displacements with small strain, have been developed. In the T.L. formulation the incremental linear strains  ${}_0e_{ij}$  (Eqs. (2.63) and (2.76)) contain an initial displacement effect that leads to a more complex strain-displacement matrix than in the U.L. formulation. The main difference between the U.L. and the T.L. formulations is that in the T.L. formulation the interpolation functions in Eq. (2.74) are obtained by referring the displacements to the initial configuration and the  ${}_0^tB_{L1}$  matrix (Eq. 2.91c) is included in the calculations.

It was been shown in Chapter 2 that two formulations yield identical element stiffness matrices and vectors of nodal element forces. In fact, if the appropriate constitutive relations are employed, same results are obtained. In this study for three dimensional beam element with small deformations and uniaxial strain, the same Young's modulus

used in the T.L. and U.L. formulations yields practically the same results.

The choice of using either the U.L. or the T.L. formulation depends on their numerical efficiency, i.e., economy. Which formulation is most effective depends on the program design and the actual structures.

The finite element model and the beam-column model have been compared for space frame problems. Both the linear and nonlinear strain incremental stiffness matrices must be considered in both models in order to predict the response accurately and to detect instability of large displacement problems. The finite element model is simple in theory, but quite a number of elements are often needed for satisfactory accuracy. The beam-column model uses the conventional beam-column theory to determine the member force-deformation relations so that fewer elements are often needed than in the case of a finite element model to give the same accuracy.

Because the stability functions  $c_{mn}$  in the beam-column model depend on whether the axial force is compressive, zero, or tensile, the axial force must be determined by an iterative procedure (i.e. Newton method). Consequently, more

computational effort is required in the beam-column model than in the finite element model for the same mesh.

From the sample analyses (chapter 7) it is seen that both element models are reliable for studying the behavior of space frames. In conclusion the following observations are presented:

The only disadvantage of the finite element model is that a fine mesh is often needed for satisfactory accuracy. The main advantages of the finite element model are:

1. The formulations are simple without recourse to complex differential equations.
2. It can be extended to the elements undergoing large deformations in which the cross-sectional area would be updated.
3. The material properties can be adjusted at each element integration point. Combinations of material nonlinearity and geometric nonlinearity is particularly simple if the increment of the material properties may be obtained. The operations required in the solution of problems of material and geometric nonlinearity are similar.
4. The external loads can be configuration-dependent, i.e., nonconservative forces, like follower forces can be involved [48].

On the other hand, the main advantage of the beam-column model is that it yields higher accuracy than the finite element model for the same mesh. However, the disadvantages of the beam-column model are:

1. The formulations are complicated by the fact that the stability coefficients (appendix B) are functions of the axial force which must be obtained by an iterative procedure.
2. The basic element force-deformation relations (sec. 5.2) are derived from the conventional beam-column theory [23]; thus the relative deformations of the element are limited to be small.
3. Since the formulations are based on elastic behavior (sec. 5.2), extensions to material nonlinearity is impractical.
4. The external loads are assumed to be configuration-independent, i.e., conservative forces, and applied at the nodes only.

## 9.2 SUGGESTIONS FOR FUTURE DEVELOPMENT

Some interesting extensions of the present study are suggested:

1. Implement the T.L. formulation in the computer program using both models by transformation of the interpolation functions [9].
2. Use numerical integration, that is, Newton-Cotes or Gauss integration [38], to evaluate the vectors of nodal element forces.
3. Include the material nonlinearity in the finite element model.
4. The displacements in the U.L. formulation referred to the current coordinates at time  $t$  are usually small enough that  $K_{NL}$  can be omitted [33]. The effects of neglecting the nonlinear strain incremental stiffness matrix in the U.L. formulation can be studied.
5. Different convergence criteria will effect the accuracy and computer cost. The effects of using different convergence criteria in nonlinear analysis are of interest.
6. A comparative study of different numerical solution techniques in nonlinear analysis.

7. The final goal is to develop a nonlinear dynamic analysis program to analyze space frames under earthquake loads.

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## Appendix A

### JOINT ORIENTATION MATRIX $\alpha_j$ FOR SMALL JOINT ROTATIONS

In this appendix Eq. (3.9) will be derived for small joint rotation. A small change in the orientation of the joint can be represented by rotation components  $\rho_1, \rho_2, \rho_3$  in the global coordinates  $X_1, X_2, X_3$  (Fig. A.1). Let

$$\rho = \begin{bmatrix} \rho_1 \\ \rho_2 \\ \rho_3 \end{bmatrix} \quad (\text{A.1})$$

To derive Eq. (3.9) we can formulate the joint orientation matrix in three steps:

First, consider Fig. A.2 which shows a small joint rotation  $\rho_2$  about  $X_2$ -axis, similar to Eq. (3.1)

$$b = \alpha_{j2} X \quad (\text{A.2})$$

where

$$\alpha_{j2} = \begin{bmatrix} \cos \rho_2 & \cos 90^\circ & \cos(90^\circ + \rho_2) \\ \cos 90^\circ & \cos 0^\circ & \cos 90^\circ \\ \cos(90^\circ - \rho_2) & \cos 90^\circ & \cos \rho_2 \end{bmatrix} \quad (\text{A.3})$$

For infinitesimal  $\rho_2$ :

$$\rho_2 \approx 0, \quad \cos \rho_2 \approx 1, \quad \sin \rho_2 \approx \rho_2 \quad (\text{A.4})$$

Substituting Eq. (A.4) into (A.3), we have

$$\alpha_{j2} = \begin{bmatrix} 1 & 0 & -\rho_2 \\ 0 & 1 & 0 \\ \rho_2 & 0 & 1 \end{bmatrix} \quad (\text{A.5})$$

Second, we rotate a small joint rotation  $\rho_3$  about  $b_3$ -axis (Fig. A.3). In a similar manner we have

$$a = \alpha_{j3} b \quad (\text{A.6})$$

where

$$\alpha_{j3} = \begin{bmatrix} 1 & \rho_3 & 0 \\ -\rho_3 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \quad (\text{A.7})$$

Third, we rotate a small joint rotation  $\rho_1$  about  $a_1$ -axis (Fig. A.4). Similarly

$$y = \alpha_{j1} a \quad (\text{A.8})$$

where

$$\alpha_{j1} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & \rho_1 \\ 0 & -\rho_1 & 1 \end{pmatrix} \quad (\text{A.9})$$

Combining Eqs. (A.2), (A.6) and (A.8), we obtain

$$y = \alpha_j X \quad (\text{A.10})$$

where

$$\alpha_j = \alpha_{j1} \alpha_{j3} \alpha_{j2} = \begin{pmatrix} 1 & \rho_3 & -\rho_2 \\ -\rho_3 + \rho_1 \rho_2 & 1 & \rho_1 + \rho_2 \rho_3 \\ \rho_2 + \rho_1 \rho_3 & -\rho_1 & 1 - \rho_1 \rho_2 \rho_3 \end{pmatrix} \quad (\text{A.11})$$

Neglecting the higher order terms, we obtain Eq. (3.9)

$$\alpha_j = \begin{pmatrix} 1 & \rho_3 & -\rho_2 \\ -\rho_3 & 1 & \rho_1 \\ \rho_2 & -\rho_1 & 1 \end{pmatrix} \quad (\text{A.12})$$

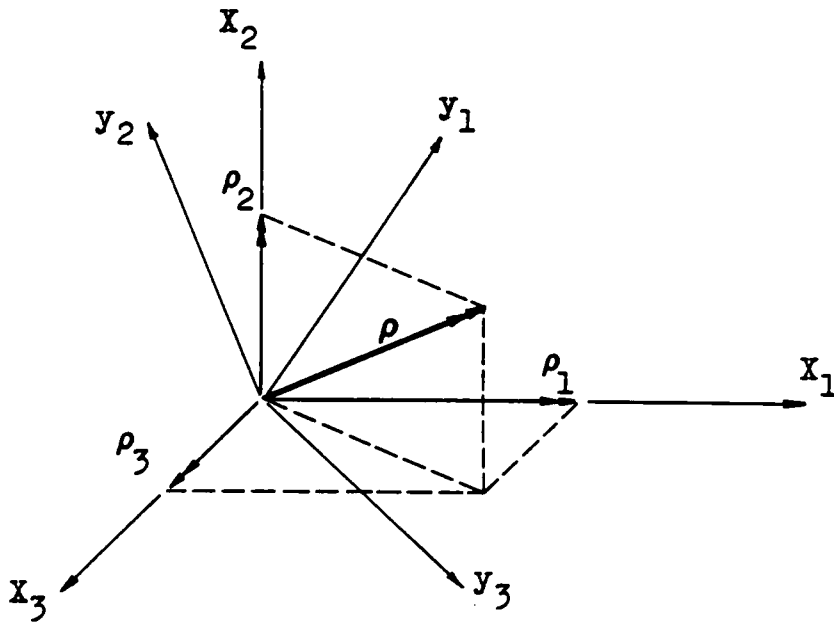


Fig. A.1 Small Joint Rotation



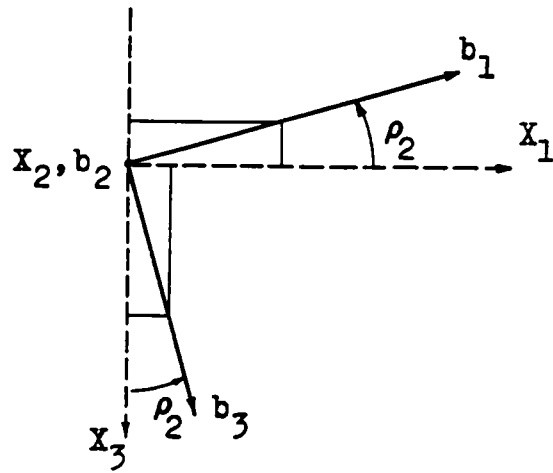


Fig. A.2 Small joint rotation  $\rho_2$  about  $X_2$ -axis

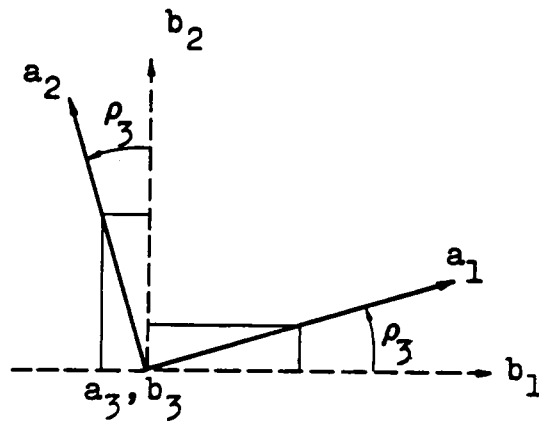


Fig. A.3 Small Joint rotation  $\rho_3$  about  $b_3$ -axis

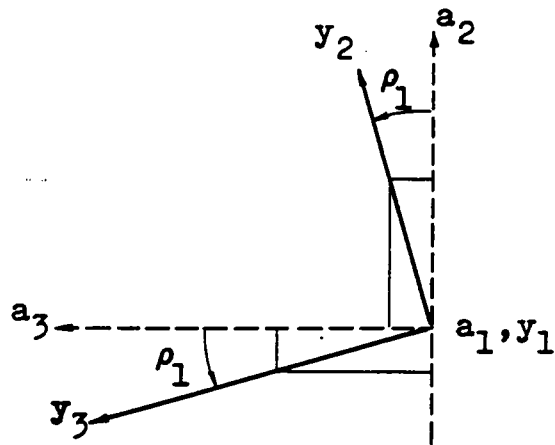


Fig. A.4 Small joint rotation  $\rho_1$  about  $a_1$ -axis

Appendix B  
STABILITY FUNCTIONS

1. Compression ( $p_n > 0$ )

$$c_{1n} = \frac{(k_n^t L) [\sin(k_n^t L) - (k_n^t L) \cos(k_n^t L)]}{2[1 - \cos(k_n^t L)] - (k_n^t L) \sin(k_n^t L)} \quad (B.1)$$

$$c_{2n} = \frac{(k_n^t L) [(k_n^t L) - \sin(k_n^t L)]}{2[1 - \cos(k_n^t L)] - (k_n^t L) \sin(k_n^t L)} \quad (B.2)$$

with

$$k_n = (P/EI_n)^{1/2} \quad (B.3)$$

2. No Axial Force ( $p_n = 0$ )

$$c_{1n} = 4 \quad (B.4)$$

$$c_{2n} = 2 \quad (B.5)$$

3. Tension ( $p_n < 0$ )

$$c_{1n} = \frac{k_n^t L [(k_n^t L) \cosh(k_n^t L) - \sinh(k_n^t L)]}{2[(1 - \cosh(k_n^t L))] + (k_n^t L) \sinh(k_n^t L)} \quad (B.6)$$

$$c_{2n} = \frac{(k_n^t L) [\sinh(k_n^t L) - (k_n^t L)]}{2[(1 - \cosh(k_n^t L))] + (k_n^t L) \sinh(k_n^t L)} \quad (B.7)$$

with

$$k_n = (-P/EI_n)^{1/2} \quad (\text{B.8})$$

## Appendix C

### GEOMETRIC MATRICES $\bar{g}^{(i)}$

The geometric matrices  $\bar{g}^{(i)}$  can be defined as [2], [29]

$$\Delta C^T^{(i)} = \bar{g}^{(i)} \Delta d \quad \text{for } i = 1 \text{ to } 6 \quad (C.1)$$

in which

$$\Delta C^T^{(i)} = \text{the } i^{\text{th}} \text{ column of } \Delta C^T$$

such that

$$\bar{g}^{(1)} = \bar{g}^{(2)} = 1/L^2 \begin{bmatrix} \delta^{(1)} & -\delta^{(1)} \\ -\delta^{(1)} & \delta^{(1)} \end{bmatrix}$$

$$\bar{g}^{(3)} = \bar{g}^{(4)} = 1/L^2 \begin{bmatrix} \delta^{(3)} & -\delta^{(3)} \\ -\delta^{(3)} & \delta^{(3)} \end{bmatrix} \quad (C.2)$$

$$\bar{g}^{(5)} = 0$$

$$\bar{g}^{(6)} = 1/L^2 \begin{bmatrix} \delta^{(6)} & -\delta^{(6)} \\ -\delta^{(6)} & \delta^{(6)} \end{bmatrix}$$

in which the size of matrices  $\delta^{(1)}$ ,  $\delta^{(3)}$ ,  $\delta^{(6)}$  are 6x6, the values of their element are defined as:

$$\delta_{12}^{(1)} = \delta_{21}^{(1)} = 1, \text{ other } \delta_{ij}^{(1)} = 0 \quad (\text{C.3})$$

$$\delta_{13}^{(3)} = \delta_{31}^{(3)} = -1, \text{ other } \delta_{ij}^{(3)} = 0 \quad (\text{C.4})$$

$$\delta_{22}^{(6)} = \delta_{33}^{(6)} = -1, \text{ other } \delta_{ij}^{(6)} = 0 \quad (\text{C.5})$$

Appendix D

$$g^{(i)} = \Lambda^T \bar{g}^{(i)} \Lambda$$

Sustituting Eqs. (2.54) and appendix C into Eq. (5.40), one obtains  $g^{(i)}$ , corresponding to the coordinate system in Fig. 2.4:

$$g^{(1)} = g^{(2)} = \frac{1}{L^2} \left[ \begin{array}{c|c} \underline{a(1)} & -\underline{a(1)} \\ \hline -\underline{a(1)} & \underline{a(1)} \end{array} \right] \quad (D.1)$$

where

$$a^{(1)} = \left[ \begin{array}{ccc|c} 2c_1^{\lambda 21} & c_2^{\lambda 21} + c_1^{\lambda 22} & c_3^{\lambda 21} + c_1^{\lambda 23} & \\ c_1^{\lambda 22} + c_2^{\lambda 21} & 2c_2^{\lambda 22} & c_3^{\lambda 22} + c_2^{\lambda 23} & \underline{0} \\ c_1^{\lambda 23} + c_3^{\lambda 21} & c_2^{\lambda 23} + c_3^{\lambda 22} & 2c_3^{\lambda 23} & \\ \hline & \underline{0} & & \underline{0} \end{array} \right] \quad (D.2)$$

$\underline{0}$  is a null matrix.

$$g^{(3)} = g^{(4)} = \frac{1}{L^2} \left[ \begin{array}{c|c} \underline{b(3)} & -\underline{b(3)} \\ \hline -\underline{b(3)} & \underline{b(3)} \end{array} \right] \quad (D.3)$$



where

$$b^{(3)} = \left( \begin{array}{ccc|c} -2c_1^{\lambda 31} & -c_2^{\lambda 31} - c_1^{\lambda 32} & -c_3^{\lambda 31} - c_1^{\lambda 33} & \\ -c_1^{\lambda 32} - c_2^{\lambda 31} & -2c_2^{\lambda 32} & -c_3^{\lambda 32} - c_2^{\lambda 33} & 0 \\ -c_1^{\lambda 33} - c_3^{\lambda 31} & -c_2^{\lambda 33} - c_3^{\lambda 32} & -2c_3^{\lambda 33} & \\ \hline & & & \\ & 0 & & 0 \end{array} \right) \quad (D.4)$$

$$g^{(5)} = 0 \quad (D.5)$$

$$g^{(6)} = \frac{1}{L^2} \left( \begin{array}{c} -i^{(6)} \\ -i^{(6)} \\ -i^{(6)} \\ -i^{(6)} \end{array} \right) \quad (D.6)$$

where

$$i^{(6)} = \left( \begin{array}{ccc|c} -\lambda_{21}^{\lambda 31} & -\lambda_{21}^{\lambda 22} - \lambda_{31}^{\lambda 32} & -\lambda_{21}^{\lambda 23} - \lambda_{31}^{\lambda 33} & \\ -\lambda_{22}^{\lambda 21} - \lambda_{32}^{\lambda 31} & -\lambda_{22}^{\lambda 32} & -\lambda_{22}^{\lambda 23} - \lambda_{32}^{\lambda 33} & 0 \\ -\lambda_{23}^{\lambda 21} - \lambda_{33}^{\lambda 31} & -\lambda_{23}^{\lambda 22} - \lambda_{33}^{\lambda 32} & -\lambda_{23}^{\lambda 33} & \\ \hline & & & \\ & 0 & & 0 \end{array} \right) \quad (D.7)$$

## Appendix E

### NOTATION

$A$  = area of cross section

$${}^t_{0B_L}, {}^t_{B_L} =$$

linear strain-displacement matrix at time  $t$  referred to configuration at time  $0$ ,  $t$

$${}^t_{0B_{NL}}, {}^t_{B_{NL}} =$$

nonlinear strain-displacement matrix at time  $t$  referred to configuration at time  $0$ ,  $t$

$b_{mn}$  = bowing functions

$$= (c_{1n} + c_{2n})(c_{2n} - 2) / 8\pi^2 q_n \quad \text{for } m = 1$$

$$= c_{2n} / 8(c_{1n} + c_{2n}) \quad \text{for } m = 2$$

$C$  = local transformation matrix defined in Eq. (5.22)

$${}^0C_{ijrs}, {}^tC_{ijrs} =$$

component of incremental material property tensor at time  $t$  referred to configuration at time  $0$ ,  $t$

$${}^0C, {}^tC, {}^{t+\Delta t}C = \text{element configuration at time } 0, t, t+\Delta t$$

CPDB = displacement convergence parameter adopted by Bathe

CPDC = displacement convergence parameter adopted by Cook

CPF = convergence parameter for the unbalanced force criterion

$c_{bn}$  = length correction factors for bending action

$$= b_{1n} (e_{1n} + e_{2n})^2 + b_{2n} (e_{1n} - e_{2n})^2$$

$${}^0c_i =$$

direction cosines of the local  $x_i$ -axis at time 0 with respect to the global coordinate system  $X_i$ -axis;  
 $i=1,2,3$

$c_{mn} =$

stability functions corresponding to  $e_{mn}$  defined in appendix B [1], [4]

$D_{jk} =$

component of node displacement at node  $j$  in the global coordinate system,  $k=1$  to 6

${}^t_0D, {}^t_tD, {}^{t+\Delta t}_tD =$

vector of nodal element displacements at time  $t, t, t+\Delta t$  referred to configuration at time 0,  $t, t+\Delta t$  in global coordinate system

$\Delta D_{ji} =$

incremental component of nodal displacement at node  $j$  in the global coordinate system;  $i=1$  to 6

$\Delta^0D, \Delta^tD =$

incremental nodal displacement vector of element from  $t$  to  $t+\Delta t$  measured in global axis  ${}^0X_i, {}^tX_i$

${}^t_0d, {}^t_t d, {}^{t+\Delta t}_t d =$

vector of nodal element displacements at time  $t, t, t+\Delta t$  referred to configuration at time 0,  $t, t+\Delta t$  in local coordinate system

$\Delta^0d, \Delta^t d =$

incremental nodal displacement vector of element from  $t$  to  $t+\Delta t$  measured in local axis  ${}^0x_i, {}^t x_i$

$$\Delta^0 d_k, \Delta^t d_k =$$

increment in nodal displacement component of element from  $t$  to  $t+\Delta t$  measured in local axis  ${}^0x_i, {}^t x_i, k=1$  to  $12$

$E$  = Young's modulus of elasticity

$e$  = vector of deformation displacements in space

$e_1$  = relative axial elongation of element

$$e_1', e_2', e_3' =$$

relative deformation displacements of element for plane frame

$$e_{mn} =$$

member relative end rotations: the first subscript  $m=1,2$  referred to the nodal point (1 for left and 2 for right); the second subscript  $n=2,3$  indicates the member axis,  $x_n$ , about which the rotation takes place (Fig. 3.3)

$$e', e'' =$$

vector of deformation displacements in  $x_1-x_2$  and  $x_1-x_3$  plane

$${}^0 e, {}^t e =$$

linear part of incremental strain vector (Green-Lagrange)  ${}^0 \varepsilon, {}^t \varepsilon$

$${}^0 e_{ij}, {}^t e_{ij} = \text{linear part of strain increment } {}^0 \varepsilon_{ij}, {}^t \varepsilon_{ij}$$

$${}^t_{0F}{}^i, {}^t_F{}^i =$$

vector of global end forces for element  $i$  at time  $t$  referred to configuration at time  $0, t$

$${}^{t+\Delta t}_{0F}{}^{k-1}, {}^{t+\Delta t}_F{}^{k-1} =$$

vector of nodal equivalent element force at time  $t+\Delta t$  referred to configuration at time  $0, t+\Delta t$ , during  $(k-1)$ th iteration (see Figs. 6.2 and 6.4)

$$\Delta^0 F^i, \Delta^t F^i =$$

vector of global incremental end forces for element  $i$  from time  $t$  to  $t+\Delta t$  referred to configuration at time  $0, t$

$$f_{bc} =$$

vector of relative local element forces in beam-column model

$f_i, F_i$  = component of local and global nodal element force

$${}^t_{0f}, {}^t_f, {}^{t+\Delta t}_f =$$

vector of local nodal element forces at time  $t, t, t+\Delta t$  referred to configuration at time  $0, t, t+\Delta t$

$${}^{t+\Delta t}_{0f}{}^{k-1}, {}^{t+\Delta t}_f{}^{k-1} =$$

vector of nodal equivalent element forces corresponding to  ${}^{t+\Delta t}_{0d}{}^{k-1}, {}^{t+\Delta t}_d{}^{k-1}$

$$\Delta^0 f, \Delta^t f =$$

vector of local incremental element forces from time  $t$  to  $t+\Delta t$  referred to configuration at time  $0, t$

$G$  = shear modulus of elasticity

$GJ$  = torsional rigidity

$g_i$  =

coefficient function defined in Eqs. (4.11) and (5.42),  
 $i = 1$  to  $42$

$$g^{(i)} = \Lambda^T g^{(i)} \Lambda$$

$g_i$  = coefficient functions defined in Eq. (5.31)

$g^{(i)}$  = geometric matrix defined in Eq. (D.2)

$H$  = symmetric matrix containing entries of 1 and 0

$h_k$  =

interpolation function corresponding to nodal direction  
 $k$ ,  $k = 1$  to  $12$ .

$${}^0 h_{i,j}' \quad {}^t h_{i,j} =$$

derivative of component of interpolation function  $h_i$   
with respect to coordinate  ${}^0 x_j$ ,  ${}^t x_j$

$${}^0 h_{i,jj}' \quad {}^t h_{i,jj} =$$

2nd order derivative of component of interpolation  
function  $h_i$  with respect to coordinate  ${}^0 x_j$ ,  ${}^t x_j$

$I$  = identity matrix or reference moment of inertia

$I_n$  = moment of inertia about the  $x_n$ -axis

INDEX =

location matrix of the components in the stiffness ma-  
trix, which stores the subscripts and negative signs of  
these components [28]

$J$  = polar moment of inertia (torsional stiffness)

$${}^t K, \quad {}^0 K =$$

structural strain incremental stiffness matrix at time  $t$  referred to configuration at time  $0, t$

$${}^t_0K^i, {}^tK^i =$$

strain incremental stiffness matrix at time  $t$  referred to configuration at time  $0, t$  for element  $i$  in global coordinate system

$$K_q^{(i)k-1} =$$

generalized tangent stiffness matrix of element  $i$  corresponding to  $q^{k-1}$

$${}^t_0k, {}^t_k =$$

element strain incremental stiffness matrix at time  $t$  referred to configuration at time  $0, t$ , in local coordinate system

$${}^t_0k_L, {}^t_kL =$$

linear strain incremental stiffness matrix at time  $t$  referred to configuration at time  $0, t$ , in local coordinate system

$${}^t_0k_{NL}, {}^t_kNL =$$

nonlinear strain incremental stiffness matrix at time  $t$  referred to configuration at time  $0, t$ , in local coordinate system

$$\hat{k} =$$

local tangent stiffness matrix of element for relative deformations in space

$$\hat{k}_s =$$

local secant stiffness matrix of element for relative deformations in space

$$\hat{k}'_s, \hat{k}''_s =$$

local secant stiffness matrix of element for relative deformations in  $x_1-x_2, x_1-x_3$  plane

$L$  = undeformed element length

$L'$  = new deformed element length

$${}^0L, {}_tL =$$

column vector defining linear strain  ${}^0e, {}_te$  from displacement gradients

${}^0L, {}^{t+\Delta t}L$  = element length at time 0,  $t+\Delta t$

$${}^0\epsilon = ({}^0c_1^2 + {}^0c_3^2)^{1/2}$$

$M_{mn}$  = bending moment corresponding to  $e_{mn}$

$M_t$  = torque

$P$  = axial compression force

$P_{En}$  = classical Euler buckling load =  $\pi^2 EI_n / L^2$

$p =$

vector of deformation forces in space defined in Eq.

(4.21)

$p_n$  = dimensionless axial force parameter =  $P / P_{En}$

${}^{t+\Delta t}p_j$  = end section matrix at node  $j$  at time  $t+\Delta t$

$p', p'' =$

vector of deformation forces in  $x_1-x_2$  and  $x_1-x_3$  plane

${}^{t+\Delta t}Q =$



vector of given new load level in configuration at time  $t+\Delta t$

$q^{k-1}$  = vector of nodal displacements at  $k-1$  iteration

$\Delta q^k =$

vector of incremental nodal displacements at  $k^{\text{th}}$  iteration in configuration at time  $t$

$R =$

transformation matrix from local coordinates  ${}^0x_i$  to  ${}^tx_i$  in space

$R^{k-1}$  = unbalanced force vector corresponding to  $q^{k-1}$

${}^tR$  = rotation matrix defined in Eq. (3.23)

${}^tR_j$  = rotation matrix at node  $j$  at time  $t$

${}^{t+\Delta t}R$  = external virtual work at time  $t+\Delta t$

${}^{t+\Delta t}r =$

vector of externally applied element nodal loads at time  $t+\Delta t$  in local coordinate system

${}^{t+\Delta t}{}_tS_{ij} =$

component of 2nd Piola-Kirchhoff stress tensor at time  $t+\Delta t$  referred to configuration at time  $t$

${}^tS_{ij}, {}^{t+\Delta t}{}_0S_{ij} =$

component of 2nd Piola-Kirchhoff stress tensor at time  $t, t+\Delta t$  referred to configuration at time  $0$

${}^0S_{ij}, {}^tS_{ij} =$

component of 2nd Piola-Kirchhoff stress increment at time  $t$

$T$  = axial force (positive in tension)

$T_i =$

entry of local strain incremental stiffness matrix defined in Eq. (4.8),  $i=1$  to 10

$\Delta U_j =$

increment vector of node displacements at node  $j$  in the global coordinate system

$\Delta U_{jr} =$

rotational increment vector of node displacements at node  $j$  in the global coordinate system

$\Delta U_{jt} =$

translational increment vector of node displacements at node  $j$  in the global coordinate system

$u$  = relative axial displacement

$u_i =$  increment in displacement component,  $u_i = {}^{t+\Delta t}u_i - {}^t u_i$

${}^t_0 u_{i,j} =$

derivative of displacement component at time  $t$  with respect to coordinate  ${}^0x_j$

${}^{t+\Delta t}u =$

displacement vector in each element at time  $t+\Delta t$  measured from initial position at time 0

${}^0u_i, {}^t u_i =$

increment in displacement component of element from  $t$  to  $t+\Delta t$  measured in local coordinate system at time 0,  $t$

$${}^t u_i, {}^{t+\Delta t} u_i =$$

displacement component from initial position at time 0 to configuration at time  $t, t+\Delta t$

$${}^0 u_{i,j}, {}^t u_{i,j} =$$

derivative of displacement increment with respect to coordinate  ${}^0 x_j, {}^t x_j$

$${}^0 u_{i,jj}, {}^t u_{i,jj} =$$

2nd order derivative of displacement increment with respect to coordinate  ${}^0 x_j, {}^t x_j$

$${}^0 V, {}^t V, {}^{t+\Delta t} V = \text{volume at time } 0, t, t+\Delta t$$

$${}^0 X_i, {}^t X_i, {}^{t+\Delta t} X_i =$$

global coordinate system at 0,  $t, t+\Delta t$  respectively

$$x_i =$$

end cross section (body) coordinate system of element associated with each node,  $i=1,2,3$

$${}^0 x_i, {}^t x_i, {}^{t+\Delta t} x_i =$$

local coordinate system at time 0,  $t, t+\Delta t$  respectively

$$y_i =$$

node coordinate system of element which is rigidly connected to each node,  $i=1,2,3$

$$\Delta^t y_j =$$

increment vector of node coordinate vector of node  $j$  at time  $t$

$y, z =$

distance from centroid axis in local  $x_2, x_3$ -axes direction

$\alpha = EI / L$

$t_{\alpha j}, t_{\alpha j}^{t+\Delta t} =$

orthogonal node orientation matrix at node  $j$  at time  $t, t+\Delta t$

$\Delta_{\alpha j}^t =$

incremental node orientation matrix of node  $j$  at time  $t$

$\delta =$  'variation in' or length correction factor

$0^\varepsilon, t^\varepsilon =$

incremental Green-Lagrange strain vector referred to configuration at time  $0, t$

$t_{\varepsilon ij}^{t+\Delta t} =$

component of Green-Lagrange strain tensor at time  $t+\Delta t$  referred to configuration at time  $t$  (i.e. using displacements from the configuration at time  $t$  to the configuration at time  $t+\Delta t$ )

$0_{\varepsilon ij}^t, 0_{\varepsilon ij}^{t+\Delta t} =$

component of Green-Lagrange strain tensor at time  $t+\Delta t, t$ , referred to configuration at time  $0$ .

$0_{\varepsilon ij}^t, t_{\varepsilon ij}^\varepsilon =$

component of strain increment tensor (Green-Lagrange)  
referred to configuration at time 0, t

$$\eta = GJ / EI$$

$${}^0\eta, {}^t\eta =$$

nonlinear part of incremental strain vector (Green-Lagrange)  ${}^0\varepsilon, {}^t\varepsilon$

$${}^0\eta_{ij}, {}^t\eta_{ij} =$$

nonlinear part of strain increment  ${}^0\varepsilon_{ij}, {}^t\varepsilon_{ij}$

$\gamma$  = rigid body rotation of element for plane frame

$${}^t_{0\Lambda}, {}^t_{t\Lambda}, {}^t_{t+\Delta t\Lambda} =$$

coordinate transformation matrix at time 0, t, t+ $\Delta t$  defined in Eqs. (2.54) and (2.107)

$\lambda_{ij}$  = component of element orientation matrix

$${}^t_{0\lambda}, {}^t_{t\lambda} =$$

element orientation matrix at time 0, t defined in Eqs. (2.109) and (3.17)

$$\Delta^t_{\lambda} =$$

change of element orientation matrix due to incremental element rotations

$$\omega_{\ell m} =$$

direction angle between the local element  $x_{\ell}$ -axis and the global  $X_m$ -axis

$\Delta^t_{\omega_j}$  = increment vector of nodal rotations at node j

$\phi$  = roll angle between the local  $x_2$ -axis and global  $X_2$ -axis

$$\phi_{\ell m} =$$

direction angle between the local node  $y_\ell$ -axis and the global  $X_m$ -axis

$\phi_t$  = member relative end torsion about  $x_1$ -axis

$\Delta\psi_i$  = incremental element rotation in local  $x_i$ -axis direction

$\rho$  = vector of small joint rotation

${}^t_0\sigma = 2\text{nd Piola-Kirchhoff stress} = {}^t_0S_{11}$

${}^t_\sigma, {}^{t+\Delta t}_\sigma = \text{axial Cauchy stress at time } t, t+\Delta t$

${}^t_\tau = \text{Cauchy stress matrix at time } t$

${}^t_{\tau_{ij}} = \text{component of Cauchy stress tensor at time } t = {}^t_{S_{ij}}$

${}^0_\theta, {}^t_\theta = \text{column vector of displacement gradient contributing to strain } {}^0_\epsilon, {}^t_\epsilon$

$\nu = \text{Poisson's ratio}$

$\xi = \text{natural coordinate} = x_1 / L$

$\xi_n = I_n / I$

$\zeta = L / (I/A)^{1/2}$

Appendix F  
PROGRAM LISTING

\*\*\*\*\*  
 \* GNSSF \*  
 \*\*\*\*\*

PARAMETERS  
 \*\*\*\*\*

\* TOLERANCES MAY BE DEFINED BY USER AT THE  
 \* BEGINNING OF THE PROGRAM. SET A CONVERGENCE  
 \* TOLERANCE GREATER THAN 1 TO NOT PERFORM THE  
 \* CORRESPONDING TEST.

CPDB CONVERGENCE TOLERANCE FOR 2-NORM DISPLACEMENT TEST

CPDC CONVERGENCE TOLERANCE FOR INFINITY-NORM DISPLACEMENT TEST

CPF CONVERGENCE TOLERANCE FOR UNBALANCED FORCE TEST

DQOR CONVERGENCE TOLERANCE FOR ITERATION IN BOWCOR

IMAX MAX NO. OF ITERATIONS PERMITTED IN BOWCOR

IMP DEGREE OF FREEDOM FOR WHICH RESULTS ARE PRINTED FOR NPRINT = 0, 1, 2

ITMAX MAX NO. OF ITERATIONS PERMITTED FOR NEWTON-RAPHSON OR MODIFIED RIKS/WEMPNER METHOD

MX MAX NUMBER OF ELEMENTS OR NODES

MXNEQ MAX NUMBER OF EQUATIONS (D.O.F.)

NPRINT PRINT INDICATOR = 0 FOR EQUILIBRIUM PATH PLOT DATA (FOR ONE D.O.F. ) (NO WORDS ARE PRINTED)  
 = 1 FOR EQUILIBRIUM PATH PLOT DATA WITH TANGENTS  
 = 2 FOR EQUILIBRIUM PATH RESULTS INCLUDING EXPLANATIONS OF RESULTS  
 = 3 FOR FULL DEBUGGING OUTPUT  
 = 4 FOR FULL EQUILIBRIUM CONFIGURATION OUTPUT

NUPD NUMBER OF ITERATIONS BETWEEN UPDATING TANGENT STIFFNESS MATRIX. NUPD= 1 FOR UPDATING AT EVERY ITERATION. (MUST BE .GE. 1 )

VARIABLES  
 \*\*\*\*\*

AREA(MX) CROSS SECTIONAL AREA OF ELEMENT I

AP(3,3,MX) CURRENT NODAL ORIENTATION MATRIX

APNEW(3,3,MX) NEW NODAL ORIENTATION MATRIX

BP12(MX),BP22(MX),BP13(MX),BP23(MX) FIRST DERIVATIVE OF BW12,BW22,BW13,BW23 RESPECTIVELY WITH RESPECT TO QR

BW12,BW22,BW13,BW23 BOWING FUNCTIONS

C CONVERGENCE RATIO TO BE COMPARED TO CONVERGENCE PARAMETERS

C01(MX),C02(MX),C03(MX) INITIAL DIRECTION COSINES OF ELEMENT I

C1PI(MX),C2PI(MX),C3PI(MX) CURRENT DIRECTION COSINES OF ELEMENT I

C1(MX),C2(MX),C3(MX) NEW DIRECTION COSINES OF ELEMENT I

CD,CN DENOMINATOR, NUMERATOR OF C

CPR,CPT CONVERGENCE RATIOS FOR ROTATION,TRANSLATION TO BE COMPARED TO CONV. PARAMETER



```

C CR,CT MAXIMUM CHANGE IN ROTATION, TRANSLATION
C D(MXNEQ) TOTAL GLOBAL DISPLACEMENT VECTOR
C DD(MXNEQ) CHANGE IN D ("DELTA-D")
C DDO(MXNEQ) DD FOR FIRST STEP OF MOD. R/W ITERATION TO A
C NEW EQUILIBRIUM POINT
C DDO1(MXNEQ) FOR MOD. R/W METHOD: DDO CONTRIBUTION
C (DDO=DQI*DDO1)
C DD1(MXNEQ) FOR N-R METHOD: DD FOR FIRST ITERATION
C FOR MOD. R/W METHOD: DD CONTRIBUTION
C (DD=DQI*DD1+DD2)
C DD2(MXNEQ) FOR MOD. R/W METHOD: DD CONTRIBUTION
C (DD=DQI*DD1+DD2)
C DE(12) GLOBAL DISPLACEMENTS OF ELEMENT BEING CONSIDERED
C DJ(6,MX) GLOBAL JOINT DISPLACEMENT MATRIX
C DDJ(6,MX) INCREMENTAL GLOBAL NODAL DISPLACEMENT MATRIX
C DL(12) LOCAL DISPLACEMENTS OF ELEMENT BEING CONSIDERED
C DOT1(N),DOT2(N) VECTORS FOR WHICH DOT PRODUCT IS OBTAINED IN
C DOTPRD
C DQ CHANGE IN QR USED TO CONVERGE TO QR IN BOWCOR
C DQI CHANGE IN LOAD INCREMENT
C ELENGO(MX) INITIAL LENGTH OF ELEMENT I
C ELENG(MX) CURRENT LENGTH OF ELEMENT I
C EMOD(MX) MODULUS OF ELASTICITY OF ELEMENT I
C F(MXNEQ) ELEMENT FORCE VECTOR CAUSED BY DEFORMATIONS
C FG(12,MX) GLOBAL ELEMENT FOCE MATRIX
C F(L,I) = LTH GLOBAL FORCE OF MEMBER I
C FL(12,MX) LOCAL ELEMENT FORCE MATRIX
C F(L,I)= LTH LOCAL FORCE OF ELEMENT I
C FP(MXNE) F FROM PREVIOUS LOAD INCREMENT
C FPI(MXNE) F FROM PREVIOUS N-R ITERATION
C G(42) GLOBAL ELEMENT STIFFNESS COEFFICIENTS
C G12,G22,G13,G23 INTERMEDIATE FUNCTIONS USED TO COMPUTE SKT FOR
C BEAM-COLUMN MODEL
C GP1-GP20 INTERMEDIATE FUNCTION USED TO COMPUTE SKT
C GMOD(MX) SHEAR MODULUS OF ELASTICITY OF ELEMENT I
C H INTERMEDIATE FUNCTION USED TO COMPUTE SKT
C IC ITERATION COUNTER IN BOWCOR
C ICI CONVERGENCE INDICATOR
C = 0 AFTER TEST FOR CONVERGED
C .NE.0 AFTER TEST FOR NOT CONVERGED
C N-R OR MOD. R/W ITERATION PROCEEDS AS LONG AS
C ICI DOES NOT = 0 (UNTIL MAX NO. ITERATIONS)
C ICI IS SET = 0 AT THE BEGINNING OF TEST SO
C THAT SOLUTION IS ASSUMED TO BE CONVERGED UNTIL
C IT IS PROVEN OTHERWISE. ICI IS ONLY CHANGED
C FROM ITS ZERO VALUE IF A CONVERGENCE TEST
C IS FAILED. AFTER TEST THE VALUE OF ICI MAY BE
C PRINTED TO INDICATE WHICH TESTS WERE FAILED:
C ICI=ICI+1 IF DISPLC WAS FAILED
C ICI=ICI+10 IF DISPLB WAS FAILED
C ICI=ICI+100 IF UNBALF WAS FAILED
C INDEX(12,12) MATRIX DEFINING LOCATION OF G IN SKT
C SKT(I,J)=G(INDEX(I,J))
C THE SIGN OF INDEX(I,J) IS THE SIGN OF SKT(I,J)
C ITCT ITERATION COUNTER FOR N-R AND MOD. R/W ITERATIONS
C ITDES NUMBER OF ITERATIONS DESIRED FOR MOD. R/W METHOD
C TO OBTAIN NEXT EQUILIB. CONFIGURATION
C JCODE(6,MX) NODAL CODE MATRIX: JCODE(I,J)= THE D. O. F. NUMBER
C AT NODE J IN THE GLOBAL I-DIRECTION
C JDIR NODAL DIRECTION OF APPLIED NODAL FORCE
C JNUM NODAL NUMBER THAT FORCE IS APPLIED TO
C KHT(I) COLUMN HEIGHT OF FULL TANGENT STIFFNESS FOR
C DEGREE OF FREEDOM I
C KT(5) KT(I) IS GREATER THAN ZERO TO INCLUDE THE
C CORRESPONDING CONTRIBUTION TO SKT
C MAXA(I) STORES ADDRESSES OF DIAGONAL TERMS IN THE COLUMN

```

```

C      VECTOR REPRESENTATION OF THE STIFFNESS MATRIX
C      MCODE(12,MX) MEMBER CODE MATRIX:
C      MCODE(I,J) IS THE DEGREE OF FREEDOM NUMBER IN
C      THE ITH GLOBAL DIRECTION OF MEMBER J
C      MINC(2,MX) MEMBER INCIDENCE MATRIX = NODAL NUMBERS AT EACH
C      END OF MEMBER
C      NC COUNTER FOR NUMBER OF ITERATIONS BETWEEN UPDATING
C      SKT
C      NE NUMBER OF ELEMENTS
C      NEGPIV NUMBER OF NEGATIVE PIVOTS ENCOUNTERED IN THE
C      FACTORIZATION OF SKT = NO. OF NEGATIVE
C      EIGENVALUES
C      NEQ NUMBER OF EQUATIONS (D.O.F.)
C      NJ NUMBER OF NODES
C      OL(I) (C01(I)**2+C03(I)**2)**0.5 OF ELEMENT I
C      OLAMD(3,3,MX) CURRENT ELEMENT ORIENTATION MATRIX
C      OLANEW(3,3,MX) NEW ELEMENT ORIENTATION MATRIX
C      P(6,MX) NODAL FORCE MATRIX:
C      P(I,J)= FORCE AT NODE J IN DIRECTION I
C      Q(NEQ) APPLIED FORCE DISTRIBUTION (UNFACTORED)
C      QI MULTIPLIER OF Q FOR CURRENT LOAD LEVEL OR
C      INCREMENT (SOME REFERENCES CALL THIS LAMBDA)
C      QIMAX MAXIMUM ALLOWED LOAD LEVEL
C      QL AXIAL FORCE / ELENGO
C      QT(NEQ) APPLIED LOAD VECTOR FOR CURRENT LOAD INCREMENT
C      QR RATIO OF AXIAL FORCE TO EULER BUCKLING LOAD
C      R(3,3,MX) ROTATION MATRIX OF ELEMENT I
C      RJ(3,3) ROTATION MATRIX OF NODE J
C      RL(MX) NEW LENGTH OF ELEMENT I
C      RL1(MX),RL2(MX),RL3(MX)
C      PROJECTION OF NEW ELEMENT LENGTH ON TO THE
C      DIRECTION OF X1,X2,X3-AXIS
C      RL1PI(MX),RL2PI(MX),RL3PI(MX)
C      PROJECTION OF CURRENT ELEMENT LENGTH ON TO THE
C      DIRECTION OF X1,X2,X3-AXIS
C      RM12,RM22,RM13,RM23
C      MOMENTS AT A- AND B-END OF MEMBER ABOUT X2 AND
C      X3-AXIS RESPECTIVELY
C      RM32(MX) RM12+RM22 / ELENGO
C      RM33(MX) RM13+RM23 / ELENGO
C      RMR,RMT MAXIMUM ROTATION, TRANSLATION IN STRUCTURE
C      SIGN(NEGPIV+1) +1.DO FOR A LOADING REGION
C      -1.DO FOR AN UNLOADING REGION
C      SKT(I) SYSTEM TANGENT STIFFNESS MATRIX
C      SP12(MX),SP22(MX),SP13(MX),SP23(MX)
C      FIRST DERIVATIVE OF ST12,ST22,ST13,ST23 W.R.T. QR
C      ST12(MX),ST22(MX),ST13(MX),ST23(MX)
C      STABILITY FUNCTION CORRESPONDING TO TR12,TR22
C      TR13,TR23
C      TR12(MX),TR22(MX),TR13(MX),TR23(MX)
C      TOTAL ROTATION AT A- AND B-END OF MEMBER ABOUT
C      X2 AND X3-AXIS RESPECTIVELY
C      TTORS(I) TOTAL RELATIVE TORSION AT B-END ABOUT X1-AXIS
C      OF ELEMENT I
C      TCB2,TCB3 LENGTH CORRECTION FACTOR FOR BENDING ACTION
C      TCB2P,TCB3P DERIVATIVES OF TCB2,TCB3 WITH RESPECT TO QR
C      TUL(I) TOTAL AXIAL SHORTENING DIVIDED BY INITIAL LENGTH
C      OF MEMBER I (POS FOR SHORTENING)
C      U(I) INCREMENTAL AXIAL LENGTHENING OF MEMBER I
C      TU(I) TOTAL AXIAL LENGTHENING OF MEMBER I
C      X(1,J),X(2,J),X(3,J)
C      GLOBAL X1,X2,X3-COORDINATE OF NODE J
C      ZI2(I),ZI3(I) MOMENT OF INERTIA ABOUT LOCAL X2 AND
C      X3-AXIS OF ELEM. I
C      ZJ(I) POLAR MOMENT OF INERTIA (TORSIONAL STIFFNESS)
C      OF ELEMENT I

```

DATA CARDS  
 \*\*\*\*\*

\*NOTE: ALL VARIABLES USE STANDARD WATFIV DEFAULT TYPING;  
 I.E. IF THE FIRST LETTER OF THE VARIABLE NAME IS BETWEEN  
 I AND N INCLUSIVE, THE VARIABLE IS AN INTEGER.

QI,QIMAX,DQI

IMP

SIGN(I)

: :

: :

0.DO

KT(I),FOR I=1 TO 5

NE,NJ

MINC(1,I),MINC(2,I),FOR I=1 TO NE

NODAL NUMBER, CONSTRAINT GLOBAL DIRECTION

: :

(FOR ALL CONSTRAINTS)

0,0

X(1,J),X(2,J),X(3,J) FOR J=1 TO NJ

AREA(I),EMOD(I),GMOD(I)

: :

(FOR ALL MEMBERS)

ZI2(I),ZI3(I),ZJ(I)

: :

(FOR ALL MEMBERS)

JNUM , JDIR , FORCE

: :

(FOR ALL NODAL FORCES)

0,0,0.DO

\*\*\*\*\*

\*

MAIN

\*

\*\*\*\*\*

INITIALIZE PARAMETERS CPDB,CPDC,CPF,CPE,DP,IMAX,INDTAN,ITDES,

ITIND,ITMAX,MX,MXNA,MXNEQ,NUPD,NPRINT,IELS; READ QI,QIMAX,

DQI,IMP,SIGN,KT; INITIALIZE THE CONFIGURATION TO THE

UNDEFORMED STATE; CALL DATA AND EITHER NEWRAP OR RIKWEM.

IMPLICIT REAL\*8(A-H,O-Z)

DIMENSION AREA(60),BP12(60),BP22(60),BP13(60),BP23(60),

\* C1(60),C2(60),C3(60),D(360),DD(360),DD1(360),

\* DE(12),DL(12),DJ(6,60),ELENG(60),EMOD(60),

\* GMOD(60),F(360),FG(12,60),FL(12,60),FP(360),

\* FPI(360),G(42),INDEX(12,12),JCODE(6,60),KHT(361),

\* MAXA(361),MCODE(12,60),MINC(2,60),P(6,60),

\* Q(360),QT(360),R12(60),R22(60),R13(60),R23(60),

\* TORS(60),TR12(60),TR22(60),TR13(60),TR23(60),

\* TTORS(60),SKT(65160),SP12(60),SP22(60),SP13(60),

\* SP23(60),ST12(60),ST22(60),ST13(60),ST23(60),

\* ZI2(60),ZI3(60),ZJ(60),C1PI(60),C2PI(60),C3PI(60),

\* DD2(360),Z(360),SIGN(10),DOT1(360),DOT2(360),

\* DDO(360),DDO1(360),RM32(60),RM33(60),QL(60),

\* U(60),KT(5),RL1PI(60),RL2PI(60),RL3PI(60),RL1(60),

\* RL2(60),RL3(60),RL(60),TUL(60),TCB2(60),TCB3(60),

\* TU(60),DDJ(6,60),AP(3,3,60),APNEW(3,3,60),

\* RJ(3,3),OL(60),C01(60),C02(60),C03(60),R(3,3,60),

```

*          OLAMD(3,3,60),OLANEW(3,3,60),ELENGO(60),
*          X(3,60)
C      CPDB : ICI+10
        CPDB=1.D-03
C      1E-06 < CPDC < 1E-02
C      CPDC : ICI+1
C      CPDC=3.D-01
        CPDC=1.D-02
C      CPF : ICI+100
        CPF=1.D-01
        DQQR=1.D-03
        FQR=1.D-05
        IMAX=30
C      INDTAN = 1 FOR BEAM-COLUMN TANGENT STIFFNESS
C              = 2 FOR FINITE ELEMENT TANGENT STIFFNESS
        INDTAN=1
C      ITDES = NUMBER OF ITERATIONS DESIRED FOR CONVERGENCE IN
C              MODIFIED RIKS WEMPNER METHOD
        ITDES=3
C      ITIND = 0 FOR NEWTON RAPHSON ITERATION
C              = 1 FOR MODIFIED RIKS WEMPNER METHOD
        ITIND=1
        ITMAX=200
        MX=60
        MXNA=60
        MXNEQ=6*MX
        NUPD=300
C      NPRINT = 0 FOR EQUILIBRIUM PATH PLOT
C              1 FOR EQUILIBRIUM PATH PLOT WITH TANGENTS
C              2 FOR EQUILIBRIUM PATH RESULTS
C              3 FOR FULL DEBUGGING OUTPUT
C              4 FOR FULL FINAL OUTPUT
        NPRINT=0
        READ,QI,QIMAX,DQI
C      IMP = IMPORTANT DEGREE OF FREEDOM (D.O.F. FOR WHICH PRINTOUT
C              IS DESIRED)
        READ,IMP
        PRINT 100,NUPD,CPDB,CPDC,CPF,IMAX,ITMAX,QI,QIMAX,
*          DQI,ITDES,ITIND,INDTAN,IMP,NPRINT
100    FORMAT(' -NUPD = ',I4,' OCPDB = ',F10.6,' OCPDC = ',F10.6/
*          ' CPF = ',F10.6,' IMAX = ',I6,' ITMAX = ',I6/
*          ' OQI = ',F15.5,' OQIMAX = ',F15.5,' DQI = ',
*          F15.5,' OITDES = ',I6,' ITIND = ',I6,' INDTAN = ',
*          I3,' D.O.F. PRINTED = ',I3,' NPRINT = ',I3)
        PRINT 200
200    FORMAT(/' 0  I',10X,' SIGN(I)')
        I=1
        READ,SIGN(I)
        WHILE (SIGN(I) .NE. 0.D00) DO
            PRINT 300,I,SIGN(I)
300    FORMAT(' ',I3,5X,F10.1)
            I=I+1
            READ,SIGN(I)
        END WHILE
        READ,(KT(J),J=1,5)
        PRINT 400,(KT(J),J=1,5)
400    FORMAT(/'  KT=',5(3X,I4))
        CALL DATA(AREA,C1,C2,C3,ELENG,EMOD,GMOD,FL,JCODE,KHT,
*          MAXA,MCODE,MINC,Q,QT,X,ZI2,ZI3,ZJ,MX,MXNA,
*          MXNEQ,NE,NEQ,NJ,NKT,QI,C1PI,C2PI,C3PI,
*          C01,C02,C03,ELENGO)
        DO 10 I=1,NEQ
            F(I)=0.D00
            FP(I)=0.D00
            FPI(I)=0.D00
            D(I)=0.D00

```

```

      DD(I)=0. D00
      DD0(I)=0. D0
      Z(I)=0. D00
10  CONTINUE
      IELS=0

```

C  
C  
C

```

      INITIALIZE THE NODAL ORIENTATION MATRIX AT TIME 0 TO
      IDENTITY MATRIX.
      DO 35 J=1,NJ
        DO 30 K=1,3
          DO 20 L=1,3
            IF (K.EQ. L) THEN DO
              AP(K,L,J)=1. D0
            ELSE DO
              AP(K,L,J)=0. D0
            END IF
          20  CONTINUE
        30  CONTINUE
      35  CONTINUE

```

C  
C  
C

```

      CALCULATE THE COMPONENTS OF THE INITIAL ORIENTATION MATRIX
      OF ELEMENT I.
      DO 40 I=1,NE
        OL(I)=(C01(I)**2+C03(I)**2)**0.5
        OLAMD(1,1,I)=C01(I)
        OLAMD(1,2,I)=C02(I)
        OLAMD(1,3,I)=C03(I)
        OLAMD(2,1,I)=-C01(I)*C02(I)/OL(I)
        OLAMD(2,2,I)=OL(I)
        OLAMD(2,3,I)=-C02(I)*C03(I)/OL(I)
        OLAMD(3,1,I)=-C03(I)/OL(I)
        OLAMD(3,2,I)=0. D0
        OLAMD(3,3,I)=C01(I)/OL(I)
      40  CONTINUE

```

C

```

      IF (ITIND.EQ. 0) THEN DO
        CALL NEWRAP(AREA, BP12, BP22, BP13, BP23, C1, C2, C3, D, DD, DD1,
          * DE, ELENG, EMOD, GMOD, F, FG, FL, FP, FPI, G, JCODE,
          * MAXA, MCODE, QT, R12, R22, R13, R23, TORS, TR12, TR22,
          * TR13, TR23, TTORS, SKT, SP12, SP22, SP13, SP23, ST12,
          * ST22, ST13, ST23, ZI2, ZI3, ZJ, CPDB, CPDC, CPF,
          * ICI, IELS, IMAX, ITMAX, NE, NEQ, NJ, NKT, NUPD,
          * Q, Q1, C1PI, C2PI, C3PI, QIMAX, DOI, DJ, MINC,
          * P, ITIND, RM32, RM33, QL, TU, INDTAN, KT, IMP,
          * NPRINT, QJ, RL1PI, RL2PI, RL3PI, RL1, RL2, RL3, RL,
          * TUL, TCB2, TCB3, DDJ, AP, APNEW, RJ, OL, CO1, CO2,
          * CO3, DL, R, OLAMD, OLANEW, DQQR, FQR, ELENG0, DD0)
      ELSE DO
        QI=0. D0
        CALL RIKWEM(AREA, BP12, BP22, BP13, BP23, C1, C2, C3, D, DD, DD1,
          * DE, ELENG, EMOD, GMOD, F, FG, FL, FP, FPI, G, JCODE,
          * MAXA, MCODE, QT, R12, R22, R13, R23, TORS, TR12,
          * TR22, TR13, TR23, TTORS, SKT, SP12, SP22, SP13,
          * SP23, ST12, ST22, ST13, ST23, ZI2, ZI3, ZJ, CPDB,
          * CPDC, CPF, ICI, IELS, IMAX, ITMAX, NE, NEQ, NJ,
          * NKT, NUPD, Q, Q1, C1PI, C2PI, C3PI, DD2, Z, SIGN,
          * ITDES, QIMAX, DOI, DJ, MINC, P, DOT1, DOT2, ITIND,
          * DDO, DDO1, RM32, RM33, QL, TU, INDTAN, KT, IMP, NPRINT,
          * RL1PI, RL2PI, RL3PI, RL1, RL2, RL3, RL, TUL, TCB2, TCB3,
          * DDJ, AP, APNEW, RJ, OL, CO1, CO2, CO3, DL, R, OLAMD,
          * OLANEW, DQQR, FQR, ELENG0)
      END IF
      STOP
      END

```

C  
C

```

C
C
C *****
C *                               DATA                               *
C *****
SUBROUTINE DATA(AREA,C1,C2,C3,ELENG,EMOD,GMOD,FL,JCODE,
*           KHT,MAXA,MCODE,MINC,Q,QT,X,ZI2,ZI3,
*           ZJ,MX,MXNA,MXNEQ,NE,NEQ,NJ,NKT,QI,
*           C1PI,C2PI,C3PI,C01,C02,C03,ELENGO)
C READ NE,NJ; IF NE .LE. MX AND NJ .LE. MX, CALL STRUCT AND
C LOAD.
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION AREA(1),C1(1),C2(1),C3(1),ELENG(1),
*           EMOD(1),GMOD(1),FL(12,1),JCODE(6,1),
*           KHT(1),MAXA(1),MCODE(12,1),MINC(2,1),
*           Q(1),QT(1),X(3,1),ZI2(1),ZI3(1),ZJ(1),
*           C1PI(1),C2PI(1),C3PI(1),C01(1),C02(1),
*           C03(1),ELENGO(1)
  READ,NE,NJ
  PRINT 100,NE,NJ
100 FORMAT(' -NE = ',I4,7X,'NJ = ',I4)
  IF (NE .LE. MX AND NJ .LE. MX) THEN DO
    CALL STRUCT(AREA,ELENG,EMOD,GMOD,JCODE,KHT,MAXA,MCODE,
*           MINC,X,ZI2,ZI3,ZJ,MXNEQ,NE,NEQ,NJ,NKT,
*           C1PI,C2PI,C3PI,C01,C02,C03,ELENGO)
    CALL LOAD(C1,C2,C3,ELENG,EMOD,GMOD,FL,JCODE,
*           MCODE,Q,QT,ZI2,ZI3,ZJ,MXNA,NE,NEQ,QI)
  ELSE DO
    PRINT 200
200 FORMAT(' -***NE OR NJ EXCEEDS MX; REDIMENSION ARRAYS' )
    STOP
  END IF
  RETURN
  END
C
C
C *****
C *                               STRUCT                               *
C *****
SUBROUTINE STRUCT(AREA,ELENG,EMOD,GMOD,JCODE,KHT,MAXA,MCODE,
*           MINC,X,ZI2,ZI3,ZJ,MXNEQ,NE,NEQ,NJ,NKT,
*           C1PI,C2PI,C3PI,C01,C02,C03,ELENGO)
C READ AND ECHO, MINC; INITIALIZE THE ELEMENTS OF JCODE,
C TO UNITY; FOR EACH NODAL CONSTRAINT, READ AND ECHO, JNUM, JDIR,
C AND STORE A ZERO IN THE CORRESPONDING LOCATION OF JCODE;
C CALL CODES, DETMAX, AND PROP.
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION AREA(1),ELENG(1),EMOD(1),GMOD(1),JCODE(6,1),
*           KHT(1),MAXA(1),MCODE(12,1),MINC(2,1),X(3,1),
*           ZI2(1),ZI3(1),ZJ(1),C1PI(1),C2PI(1),C3PI(1),
*           C01(1),C02(1),C03(1),ELENGO(1)
  READ,(MINC(1,I),MINC(2,I),I=1,NE)
  PRINT 100
100 FORMAT('/'-' ',7X,'MINC' )
  DO 10 I=1,2
    PRINT 200,(MINC(I,J),J=1,NE)
200 FORMAT(' 0 ',22(2X,I4))
  10 CONTINUE
  DO 30 J=1,NJ
    DO 20 I=1,6
      JCODE(I,J)=1
    20 CONTINUE
  30 CONTINUE
  PRINT 300
300 FORMAT('/'-' ',10X,'NODAL NUMBER',10X,'CONSTRAINT DIRECTION' )
  READ,JNUM,JDIR

```

```

      WHILE (JNUM .NE. 0) DO
        PRINT 400, JNUM, JDIR
400    FORMAT(' 0', 11X, I4, 24X, I2)
        JCODE(JDIR, JNUM)=0
        READ, JNUM, JDIR
      END WHILE

```

C

```

      CALL CODES(JCODE, MCODE, MINC, MXNEQ, NE, NEQ, NJ)
      CALL DETMAX(KHT, MAXA, MCODE, NE, NEQ, NKT)
      CALL PROP(AREA, ELENG, EMOD, GMOD, MINC, X, ZI2, ZI3, ZJ,
*          NE, NJ, C1PI, C2PI, C3PI, CO1, CO2, CO3, ELENG0)
      RETURN
      END

```

C

C

C

C

C

C

C

```

*****
*

```

```

*          CODES          *
*****

```

```

*****

```

```

SUBROUTINE CODES(JCODE, MCODE, MINC, MXNEQ, NE, NEQ, NJ)

```

```

GENERATE JCODE; AND GENERATE MCODE, BY USING MINC AND JCODE.

```

```

IMPLICIT REAL*8(A-H, O-Z)

```

```

DIMENSION JCODE(6, 1), MCODE(12, 1), MINC(2, 1)

```

```

NEQ=0

```

```

DO 20 J=1, NJ

```

```

  DO 10 L=1, 6

```

```

    IF (JCODE(L, J) .NE. 0) THEN DO

```

```

      NEQ=NEQ+1

```

```

      JCODE(L, J)=NEQ

```

```

    END IF

```

```

10  CONTINUE

```

```

20  CONTINUE

```

```

IF (NEQ .GT. MXNEQ) THEN DO

```

```

  PRINT 100

```

```

100  FORMAT(' - *** NEQ EXCEDES MXNEQ ; REDIMENSION ARRAYS ***' )

```

```

  STOP

```

```

END IF

```

```

PRINT 200

```

```

200  FORMAT('/ -', 7X, ' JCODE' )

```

```

DO 30 I=1, 6

```

```

  PRINT 300, (JCODE(I, J), J=1, NJ)

```

```

300  FORMAT(' 0', 22(2X, I4))

```

```

30  CONTINUE

```

```

DO 50 I=1, NE

```

```

  J=MINC(1, I)

```

```

  K=MINC(2, I)

```

```

  DO 40 L=1, 6

```

```

    MCODE(L, I)=JCODE(L, J)

```

```

    MCODE(L+6, I)=JCODE(L, K)

```

```

40  CONTINUE

```

```

50  CONTINUE

```

```

PRINT 400

```

```

400  FORMAT('/ -', 7X, ' MCODE' )

```

```

DO 60 I=1, 12

```

```

  PRINT 500, (MCODE(I, J), J=1, NE)

```

```

500  FORMAT(' 0', 22(2X, I4))

```

```

60  CONTINUE

```

```

RETURN

```

```

END

```

C

C

C

C

C

C

```

*****
*

```

```

*          DETMAX          *
*****

```

```

*****

```

```

SUBROUTINE DETMAX(KHT, MAXA, MCODE, NE, NEQ, NKT)

```

```

C      CALCULATE COLUMN HEIGHTS, KHT; AND CALCULATE ADDRESSES OF
C      DIAGONAL ELEMENTS IN Banded MATRIX WHOSE COLUMN HEIGHTS
C      ARE KNOWN, MAXA.
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION KHT(1),MAXA(1),MCOde(12,1)
      DO 10 I=1,NEQ
        KHT(I)=0
10     CONTINUE
      DO 30 I=1,NE
        J=1
        WHILE( MCODE(J,I) .EQ. 0) DO
          J=J+1
        END WHILE
        MIN=MCODE(J,I)
        J=J+1
        DO 20 L=J,12
          K=MCODE(L,I)
          IF (K .NE. 0) THEN DO
            KHT(K)=MAX0(KHT(K),(K-MIN))
          END IF
20      CONTINUE
30     CONTINUE
      PRINT 100
100    FORMAT(// ' - ' ,5X, ' I ' ,10X, ' KHT(I) ' ,10X, ' MAXA(I) ' )
      MAXA(1)=1
      DO 40 I=1,NEQ
        PRINT 200, I, KHT(I), MAXA(I)
200    FORMAT( ' 0 ' ,1X, I5, 9X, I5, 11X, I5)
        MAXA(I+1)=MAXA(I)+KHT(I)+1
40     CONTINUE
      NKT=MAXA(NEQ+1)-1
      I=NEQ+1
      PRINT 300, I, MAXA(I), NKT
300    FORMAT( ' 0 ' ,1X, I5, 25X, I5//1X, ' NKT = ' , I5)
      RETURN
      END

```

```

C
C
C
C      *****
C      *                                     PROP                                     *
C      *****
C      SUBROUTINE PROP(AREA,ELENG,EMOD,GMOD,MINC,X,ZI2,ZI3,ZJ,
C      *      NE,NJ,C1PI,C2PI,C3PI,C01,C02,C03,ELENG0)
C      READ AND ECHO THE NODAL COORDINATES, X(I,J), AND ELEMENT
C      PROPERTIES. FOR EACH ELEMENT COMPUTE THE INITIAL ELEMENT
C      LENGTH AND DIRECTION COSINES.
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION AREA(1),ELENG(1),EMOD(1),GMOD(1),MINC(2,1),
      *      X(3,1),ZI2(1),ZI3(1),ZJ(1),C1PI(1),C2PI(1),
      *      C3PI(1),C01(1),C02(1),C03(1),ELENG0(1)
      READ,(X(1,J),X(2,J),X(3,J),J=1,NJ)
      PRINT 100
100    FORMAT(// ' - ' : ' GLOBAL NODAL COORDINATES' ,/,18X, ' X1 ' ,10X, ' X2 ' ,
      *      10X, ' X3 ' )
      DO 10 J=1,NJ
        PRINT 200, J, (X(I,J), I=1,3)
200    FORMAT( ' 0 ' , ' NODE ' , I3, 3(2X, F10.4))
10     CONTINUE
      PRINT 300
300    FORMAT(// ' - ' : ' ELEMENT PROPERTIES' /, 9X, ' NO. ' , 9X, ' AREA ' ,10X,
      *      ELENG' ,12X, ' EMOD ' ,12X, ' GMOD ' )
      DO 20 I=1,NE
        J=MINC(1,I)
        K=MINC(2,I)
        EL1=X(1,K)-X(1,J)

```



```

      EL2=X( 2,K)-X( 2,J)
      EL3=X( 3,K)-X( 3,J)
      ELENGO(I)=DSQRT( EL1**2+EL2**2+EL3**2)
      ELENG(I)=ELENGO(I)
      CO1(I)=EL1/ELENGO(I)
      CO2(I)=EL2/ELENGO(I)
      CO3(I)=EL3/ELENGO(I)
      C1PI(I)=CO1(I)
      C2PI(I)=CO2(I)
      C3PI(I)=CO3(I)
      READ,AREA(I),EMOD(I),GMOD(I)
      PRINT 400,I,AREA(I),ELENG(I),EMOD(I),GMOD(I)
400   FORMAT(' 0',6X,I4,7X,F8.3,3(2X,D15.7))
      20 CONTINUE

```

```

C
      PRINT 410
410   FORMAT(//,' ',8X,'NO.',12X,'ZI2',18X,'ZI3',17X,'ZJ')
      DO 430 I=1,NE
      READ,ZI2(I),ZI3(I),ZJ(I)
      PRINT 420,I,ZI2(I),ZI3(I),ZJ(I)
420   FORMAT(' ',6X,I4,3X,D15.7,8X,D15.7,6X,D15.7)
430   CONTINUE
      RETURN
      END

```

```

C
C
C
C
C
*****
*                               LOAD                               *
*****
SUBROUTINE LOAD(C1,C2,C3,ELENG,EMOD,GMOD,FL,JCODE,
*               MCODE,Q,QT,ZI2,ZI3,ZJ,MXNA,NE,
*               NEQ,QI)
C
C
      INITIALIZE THE NODAL LOAD VECTOR,Q,TO ZERO. CALL JLOAD AND
      MACT.
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION C1(1),C2(1),C3(1),ELENG(1),EMOD(1),GMOD(1),
*             FL(12,1),JCODE(6,1),MCODE(12,1),Q(1),
*             QT(1),ZI2(1),ZI3(1),ZJ(1)
      DO 10 K=1,NEQ
      Q(K)=0.D00
10    CONTINUE
      CALL JLOAD(JCODE,Q)
      RETURN
      END

```

```

C
C
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C
*****
*                               JLOAD                               *
*****
SUBROUTINE JLOAD(JCODE,Q)
      READ,ECHO,JNUM,JDIR,AND THE APPLIED FORCE, FORCE; STORE FORCE
      IN Q.
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION JCODE(6,1),Q(1),QT(1)
      READ,JNUM,JDIR,FORCE
      IF (JNUM.NE. 0) THEN DO
      PRINT 100
100   FORMAT(//'- ',3X,'NODAL NUMBER',10X,'GLOBAL DIRECTION',10X,
*           APPLIED FORCE')
      END IF
      WHILE (JNUM.NE. 0) DO
      PRINT 200,JNUM,JDIR,FORCE
200   FORMAT(' 0',6X,I4,22X,I2,13X,F16.5,/)
      K=JCODE(JDIR,JNUM)

```

```

      Q(K)=FORCE
      READ,JNUM,JDIR,FORCE
      END WHILE
      RETURN
      END

```

C  
C  
C  
C  
C

```

*****
*                               NEWRAP                               *
*****

```

```

SUBROUTINE NEWRAP( AREA, BP12, BP22, BP13, BP23, C1, C2, C3, D, DD, DD1,
*                DE, ELENG, EMOD, GMOD, F, FG, FL, FP, FPI, G, JCODE,
*                MAXA, MCODE, QT, R12, R22, R13, R23, TORS, TR12,
*                TR22, TR13, TR23, TTORS, SKT, SP12, SP22, SP13,
*                SP23, ST12, ST22, ST13, ST23, ZI2, ZI3, ZJ, CPDB,
*                CPDC, CPF, ICI, IELS, IMAX, ITMAX, NE, NEQ, NJ,
*                NKT, NUPD, Q, QI, C1PI, C2PI, C3PI, QIMAX, DQI, DJ,
*                MINC, P, ITIND, RM32, RM33, QL, TU, INDTAN, KT, IMP,
*                NPRINT, RL1PI, RL2PI, RL3PI, RL1, RL2, RL3, RL,
*                TUL, TCB2, TCB3, DDJ, AP, APNEW, RJ, OL, C01, C02,
*                C03, DL, R, OLAMD, OLANEW, DQOR, FOR, ELENGO, DDO)

```

C  
C

```

      EMPLOY NEWTON-RAPHSON ITERATION OR MODIFIED NEWTON-RAPHSON
      ITERATION TO ARRIVE AT THE EQUILIBRIUM POINT.
      IMPLICIT REAL*8(A-H,O-Z)

```

```

      DIMENSION AREA(1), BP12(1), BP22(1), BP13(1), BP23(1), C1(1), C2(1),
*            C3(1), D(1), DD(1), DD1(1), DE(1), ELENG(1), EMOD(1),
*            GMOD(1), F(1), FG(12,1), FL(12,1), FP(1), FPI(1), G(1),
*            JCODE(6,1), MAXA(1), MCODE(12,1), Q(1), QT(1), R12(1),
*            R22(1), R13(1), R23(1), TORS(1), TR12(1), TR22(1),
*            TR13(1), TR23(1), TTORS(1), SKT(1), SP12(1), SP22(1),
*            SP13(1), SP23(1), ST12(1), ST22(1), ST13(1), ST23(1),
*            ZI2(1), ZI3(1), ZJ(1), C1PI(1), C2PI(1), C3PI(1),
*            DJ(6,1), MINC(2,1), P(6,1), RM32(1), RM33(1), QL(1),
*            TU(1), KT(1), RL1PI(1), RL2PI(1), RL3PI(1), RL1(1),
*            RL2(1), RL3(1), RL(1), TUL(1), TCB2(1), TCB3(1),
*            DDJ(6,1), AP(3,3,60), APNEW(3,3,60), RJ(3,1),
*            OL(1), C01(1), C02(1), C03(1), DL(1), R(3,3,60),
*            OLAMD(3,3,60), OLANEW(3,3,60), ELENGO(1),
*            DDO(1)

```

C  
C  
C

```

      INITIALIZE THE TOTAL RELATIVE DEFORMATION OF ELEMENT I TO
      ZERO.

```

```

      DO 5 I=1, NE
        TR12(I)=0. DO
        TR22(I)=0. DO
        TR13(I)=0. DO
        TR23(I)=0. DO
        TTORS(I)=0. DO
        TUL(I)=0. DO
        TCB2(I)=0. DO
        TCB3(I)=0. DO
        TU(I)=0. DO

```

5 CONTINUE

C

```

      NN=1
      CALL FORCES( AREA, BP12, BP22, BP13, BP23, C1, C2, C3, DD, DE, ELENG,
*            EMOD, GMOD, F, FG, FL, MCODE, QT, R12, R22, R13, R23, TORS,
*            TR12, TR22, TR13, TR23, TTORS, SP12, SP22, SP13, SP23,
*            ST12, ST22, ST13, ST23, ZI2, ZI3, ZJ, IMAX, NE, NEQ,
*            Q, QI, C1PI, C2PI, C3PI, RM32, RM33, QL, TU, INDTAN,
*            NPRINT, RL1PI, RL2PI, RL3PI, RL1, RL2, RL3, RL,
*            TUL, TCB2, TCB3, DDJ, JCODE, AP, APNEW, RJ,
*            MINC, OL, C01, C02, C03, DL, R, OLAMD, OLANEW,
*            DQOR, FOR, ELENGO, NJ, DDO, ITIND, NN)
      CALL UPDATE( QT, Q, NEQ, QI)

```

```

WHILE( QI .LE. QIMAX) DO
  NC=NUPD
  ITCT=0
  ICI=1
  WHILE(ICI .NE. 0 .AND. ITCT .LE. ITMAX) DO
    IF (NC .GE. NUPD) THEN DO
      CALL STIFF(AREA,BP12,BP22,BP13,BP23,C1PI,C2PI,C3PI,
*          ELENG,EMOD,GMOD,G,MAXA,MCODE,TR12,TR22,
*          TR13,TR23,SKT,SP12,SP22,SP13,SP23,ST12,
*          ST22,ST13,ST23,ZI2,ZI3,ZJ,IELS,NE,NKT,
*          RM32,RM33,QL,TU,INDTAN,KT,NPRINT,RL1PI,
*          RL2PI,RL3PI,RL1,RL2,RL3,RL,C1,C2,C3,
          OLAMD,ELENGO)
      NC=0
    END IF
C
    CALL STORE(RL1PI,RL2PI,RL3PI,ELENG,RL1,RL2,RL3,RL,
*          C1PI,C2PI,C3PI,C1,C2,C3,NE)
    CALL SOLVE(DD,F,MAXA,QT,SKT,NC,NEQ,NEGPIV,NKT,NPRINT)
    DO 10 I=1,NEQ
      D(I)=D(I)+DD(I)
    CONTINUE
    IF (ITCT .EQ. 0) THEN DO
      DO 20 I=1,NEQ
        DD1(I)=DD(I)
      CONTINUE
    END IF
C
    CALL FORCES(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,DD,DE,
*          ELENG,EMOD,GMOD,F,FG,FL,MCODE,QT,R12,R22,
*          R13,R23,TORS,TR12,TR22,TR13,TR23,TTORS,SP12,
*          SP22,SP13,SP23,ST12,ST22,ST13,ST23,ZI2,ZI3,
*          ZJ,IMAX,NE,NEQ,Q,QI,C1PI,C2PI,C3PI,RM32,
*          RM33,QL,TU,INDTAN,NPRINT,RL1PI,RL2PI,RL3PI,
*          RL1,RL2,RL3,RL,TUL,TCB2,TCB3,DDJ,JCODE,
*          AP,APNEW,RJ,MINC,OL,C01,C02,C03,DL,R,OLAMD,
*          OLANEW,DQOR,FOR,ELENGO,NJ,DDO,ITIND,NN)
    CALL TEST(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,D,DD,DD1,DE,
*          ELENG,EMOD,GMOD,F,FG,FL,FP,FPI,JCODE,MCODE,
*          QT,R12,R22,R13,R23,TORS,SP12,SP22,SP13,
*          SP23,ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,CPDB,
*          CPDC,CPF,ICI,IMAX,NE,NEQ,NJ,Q,QI,C11,C12,
*          C13,ITIND,NPRINT)
    IF (NPRINT .EQ. 3) THEN DO
      PRINT 100,ITCT
      FORMAT(//'-',7X,'NEWTON-RAPHSON ITERATION ',I5//10X,
*          QT,20X,'F',20X,'D')
      DO 30 I=1,NEQ
        PRINT 200,QT(I),F(I),D(I)
        FORMAT('0',3(5X,D15.7))
      CONTINUE
    END IF
C
    NC=NC+1
    ITCT=ITCT+1
    DO 40 I=1,NEQ
      FPI(I)=F(I)
    CONTINUE
    END WHILE
    IF (NPRINT .EQ. 3) THEN DO
      PRINT 300,ICI,ITCT
      FORMAT(//'ICI=',I4,5X,'ITCT=',I4)
    END IF
    DO 50 I=1,NEQ
      FP(I)=F(I)
    CONTINUE

```

```

* CALL RESULT(C1,C2,C3,D,DJ,ELENG,EMOD,GMOD,FG,FL,
* JCÓDE,MCÓDE,MINC,P,Q,QT,ZI2,ZI3,ZJ,
* NE,NEQ,NJ,QI,IMP,NPRINT)
  QI=QI+DQI
  IF (ICI.NE. 0) THEN DO
    PRINT 400
400   FORMAT(' *** LAST SOLUTION IS NOT CONVERGED *** ')
    STOP
  END IF
  CALL UPDATE(QT,Q,NEQ,QI)
END WHILE
RETURN
END

```

C  
C  
C  
C  
C

```

*****
*                               RIKWEM                               *
*****
SUBROUTINE RIKWEM(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,D,DD,DD1,
* DE,ELENG,EMOD,GMOD,F,FG,FL,FP,FPI,G,JCÓDE,
* MAXA,MCÓDE,QT,R12,R22,R13,R23,TORS,TR12,
* TR22,TR13,TR23,TTORS,SKT,SP12,SP22,SP13,
* SP23,ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,CPDB,
* CPDC,CPF,ICI,IELS,IMAX,ITMAX,NE,NEQ,NJ,
* NKT,NUPD,Q,QI,C1PI,C2PI,C3PI,DD2,Z,SIGN,
* ITDES,QIMAX,DQI,DJ,MINC,P,DOT1,DOT2,ITIND,
* DDO,DD01,RM32,RM33,QL,TU,INDTAN,KT,IMP,
* NPRINT,RL1PI,RL2PI,RL3PI,RL1,RL2,RL3,
* RL,TUL,TCB2,TCB3,DDJ,AP,APNEW,RJ,OL,
* CO1,CO2,CO3,DL,R,OLAMD,OLANEW,DQQR,
* FQR,ELENGO)
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION AREA(1),BP12(1),BP22(1),BP13(1),BP23(1),C1(1),C2(1),
* C3(1),D(1),DD(1),DD1(1),DE(1),ELENG(1),EMOD(1),
* GMOD(1),F(1),FG(12,1),FL(12,1),FP(1),FPI(1),G(1),
* JCÓDE(6,1),MAXA(1),MCÓDE(12,1),Q(1),QT(1),R12(1),
* R22(1),R13(1),R23(1),TORS(1),TR12(1),TR22(1),
* TR13(1),TR23(1),TTORS(1),SKT(1),SP12(1),SP22(1),
* SP13(1),SP23(1),ST12(1),ST22(1),ST13(1),ST23(1),
* ZI2(1),ZI3(1),ZJ(1),C1PI(1),C2PI(1),C3PI(1),DD2(1),
* Z(1),SIGN(1),DJ(6,1),MINC(2,1),P(6,1),DOT1(1),
* DOT2(1),DDO(1),DD01(1),RM32(1),RM33(1),QL(1),
* TU(1),KT(1),RL1PI(1),RL2PI(1),RL3PI(1),RL1(1),
* RL2(1),RL3(1),RL(1),TUL(1),TCB2(1),TCB3(1),
* DDJ(6,1),AP(3,3,60),APNEW(3,3,60),RJ(3,1),OL(1),
* CO1(1),CO2(1),CO3(1),DL(1),R(3,3,60),
* OLAMD(3,3,60),OLANEW(3,3,60),ELENGO(1)
  ITCT=0
  NEGPIV=0

```

C  
C  
C

```

  INITIALIZE THE TOTAL RELATIVE DEFORMATION OF ELEMENT I
  TO ZERO.
  DO 5 I=1,NE
    TR12(I)=0. DO
    TR22(I)=0. DO
    TR13(I)=0. DO
    TR23(I)=0. DO
    TTORS(I)=0. DO
    TUL(I)=0. DO
    TCB2(I)=0. DO
    TCB3(I)=0. DO
    TU(I)=0. DO
  5 CONTINUE
C
  NN=1

```

```

CALL FORCES(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,DD,DE,ELENG,
*      EMOD,GMOD,F,FG,FL,MCODE,QT,R12,R22,R13,R23,TORS,
*      TR12,TR22,TR13,TR23,TTORS,SP12,SP22,SP13,SP23,
*      ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,IMAX,NE,NEQ,
*      Q,QI,C1PI,C2PI,C3PI,RM32,RM33,QL,TU,INDTAN,
*      NPRINT,RL1PI,RL2PI,RL3PI,RL1,RL2,RL3,RL,
*      TUL,TCB2,TCB3,DDJ,JCODE,AP,APNEW,RJ,
*      MINC,OL,C01,C02,C03,DL,R,OLAMD,OLANEW,
*      DQQR,FQR,ELENGO,NJ,DDO,ITIND,NN)

```

C

```

WHILE (QI .LE. QIMAX .AND. ITCT .LE. ITMAX) DO
CALL STIFF(AREA,BP12,BP22,BP13,BP23,C1PI,C2PI,C3PI,ELENG,
*      EMOD,GMOD,G,MAXA,MCODE,TR12,TR22,TR13,TR23,
*      SKT,SP12,SP22,SP13,SP23,ST12,ST22,ST13,
*      ST23,ZI2,ZI3,ZJ,IELS,NE,NKT,RM32,RM33,QL,
*      TU,INDTAN,KT,NPRINT,RL1PI,RL2PI,RL3PI,
*      RL1,RL2,RL3,RL,C1,C2,C3,OLAMD,ELENGO)
CALL STORE(RL1PI,RL2PI,RL3PI,ELENG,RL1,RL2,RL3,RL,
*      C1PI,C2PI,C3PI,C1,C2,C3,NE)
CALL SOLVE(DDO1,Z,MAXA,Q,SKT,0,NEQ,NEGPIV,NKT,NPRINT)
IF (ITCT .EQ. 0) THEN DO
DS=DQI*DSQRT(DOTPRD(DDO1,DDO1,NEQ)+1.DO)
ELSE DO
DQI=SIGN(NEGPIV+1)*DS/DSQRT(DOTPRD(DDO1,DDO1,NEQ)+1.DO)
IF (NPRINT .EQ. 3) THEN DO
250 PRINT 250,SIGN(NEGPIV+1),NEGPIV,DS,DQI
*      FORMAT(//,'SIGN=' ,D15.7,' NEGPIV=' ,I4/' DS=' ,D15.7,
*      DQI=' ,D15.7)
END IF
END IF
DQI1=DQI
DO 10 I=1,NEQ
DDO(I)=DQI*DDO1(I)
D(I)=D(I)+DDO(I)
10 CONTINUE
QI=QI+DQI
IF (NPRINT .EQ. 1) THEN DO
400 PRINT 400,QI,D(IMP)
*      FORMAT(' ',20X,F13.7/' ',20X,F14.8)
END IF
IF (NPRINT .EQ. 3) THEN DO
500 PRINT 500,QI,D(IMP)
*      FORMAT(//,'QI INIT = ',F13.7,' D INIT = ',F14.8//)
END IF

```

C

```

NN=1
CALL FORCES(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,DD,DE,ELENG,
*      EMOD,GMOD,F,FG,FL,MCODE,QT,R12,R22,R13,R23,
*      TORS,TR12,TR22,TR13,TR23,TTORS,SP12,SP22,SP13,
*      SP23,ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,IMAX,NE,
*      NEQ,Q,QI,C1PI,C2PI,C3PI,RM32,RM33,QL,TU,
*      INDTAN,NPRINT,RL1PI,RL2PI,RL3PI,RL1,RL2,RL3,
*      RL,TUL,TCB2,TCB3,DDJ,JCODE,AP,APNEW,RJ,
*      MINC,OL,C01,C02,C03,DL,R,OLAMD,OLANEW,
*      DQQR,FQR,ELENGO,NJ,DDO,ITIND,NN)

```

```

ICI=1
NC=NUPD
IT=0

```

C

```

WHILE (ICI .NE. 0 .AND. IT .LE. IMAX) DO
CALL UPDATE(QT,Q,NEQ,QI)
IF (NC .GE. NUPD) THEN DO
CALL STIFF(AREA,BP12,BP22,BP13,BP23,C1PI,C2PI,C3PI,
*      ELENG,EMOD,GMOD,G,MAXA,MCODE,TR12,TR22,
*      TR13,TR23,SKT,SP12,SP22,SP13,SP23,ST12,
*      ST22,ST13,ST23,ZI2,ZI3,ZJ,IELS,NE,NKT,

```

```

*           RM32,RM33,QL,TU,INDTAN,KT,NPRINT,RL1PI,
*           RL2PI,RL3PI,RL1,RL2,RL3,RL,C1,C2,C3,
*           OLAMD,ELENGO)
      NC=0
      END IF
      CALL STORE(RL1PI,RL2PI,RL3PI,ELENG,RL1,RL2,RL3,RL,
*           C1PI,C2PI,C3PI,C1,C2,C3,NE)
      CALL SOLVE(DD1,Z,MAXA,Q,SKT,NC,NEQ,NEGPIV,NKT,NPRINT)
      CALL SOLVE(DD2,F,MAXA,QT,SKT,1,NEQ,NEGPIV,NKT,NPRINT)
      DQI=-(DOTPRD(DD0,DD2,NEQ))/(DOTPRD(DD0,DD1,NEQ)+DQI1)
      DO 20 I=1,NEQ
          DD(I)=DQI*DD1(I)+DD2(I)
          D(I)=D(I)+DD(I)
20      CONTINUE
C
      NN=2
      CALL FORCES(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,DD,DE,
*           ELENG,EMOD,GMOD,F,FG,FL,MCODE,QT,R12,R22,
*           R13,R23,TORS,TR12,TR22,TR13,TR23,TTORS,SP12,
*           SP22,SP13,SP23,ST12,ST22,ST13,ST23,ZI2,
*           ZI3,ZJ,IMAX,NE,NEQ,Q,QI,C1PI,C2PI,C3PI,RM32,
*           RM33,QL,TU,INDTAN,NPRINT,RL1PI,RL2PI,RL3PI,
*           RL1,RL2,RL3,RL,TUL,TCB2,TCB3,DDJ,JCODE,AP,
*           APNEW,RJ,MINC,OL,CO1,CO2,CO3,DL,R,OLAMD,
*           OLANEW,DQQR,FQR,ELENGO,NJ,DD0,ITIND,NN)
      QI=QI+DQI
      CALL TEST(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,D,DD,DD1,DE,
*           ELENG,EMOD,GMOD,F,FG,FL,FP,FPI,JCODE,MCODE,
*           QT,R12,R22,R13,R23,TORS,SP12,SP22,SP13,SP23,
*           ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,CPDB,CPDC,
*           CPF,ICI,IMAX,NE,NEQ,NJ,Q,QI,C11,C12,C13,
*           ITIND,NPRINT)
      NC=NC+1
      IT=IT+1
      IF (NPRINT.EQ.3) THEN DO
          PRINT 600,IT,QI,D(1),ICI
600      *   FORMAT(// ' ITERATION ',I3,' QI=',F13.7,' D= ',
*           F13.7,' ICI=',I5)
      END IF
      END WHILE
      ITCT=ITCT+1
      DO 50 I=1,NEQ
          FP(I)=F(I)
50      CONTINUE
      IF (NPRINT.EQ.3) THEN DO
          PRINT 700,IT
700      *   FORMAT(// ' ',I5,' ITERATIONS' )
      END IF
      CALL RESULT(C1,C2,C3,D,DJ,ELENG,EMOD,GMOD,FG,FL,JCODE,
*           MCODE,MINC,P,Q,QT,ZI2,ZI3,ZJ,NE,NEQ,NJ,
*           QI,IMP,NPRINT)
      IF (ICI.NE.0) THEN DO
          PRINT 800
800      *   FORMAT(' *** LAST SOLUTION IS NOT CONVERGED ***' )
          STOP
      END IF
      DS=DS*DSORT(DFLOAT(ITDES)/DFLOAT(IT))
      IF (NPRINT.EQ.3) THEN DO
          PRINT 350,DS
350      *   FORMAT(// ' DS= ',F14.8,' ***' )
      END IF
      END WHILE
      RETURN
      END

```

```

C
C
C *****
*                               STIFF                               *
C *****
SUBROUTINE STIFF(AREA,BP12,BP22,BP13,BP23,C1PI,C2PI,C3PI,
*          ELENG,EMOD,GMOD,G,MAXA,MCODE,TR12,TR22,TR13,
*          TR23,SKT,SP12,SP22,SP13,SP23,ST12,ST22,ST13,
*          ST23,ZI2,ZI3,ZJ,IELS,NE,NKT,RM32,RM33,QL,
*          TU,INDTAN,KT,NPRINT,RL1PI,RL2PI,RL3PI,
*          RL1,RL2,RL3,RL,C1,C2,C3,OLAMD,ELENGO)
C  INITIALIZE THE SYSTEM TANGENT STIFFNESS MATRIX,SKT,TO ZERO;
C  FOR EACH ELEMENT CALL ELEMS1 (BEAM-COLUMN MODEL) OR ELEMS2
C  (FINITE ELEMENT MODEL) AND ASSEMS.
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION AREA(1),BP12(1),BP22(1),BP13(1),BP23(1),C1PI(1),
*          C2PI(1),C3PI(1),ELENG(1),EMOD(1),GMOD(1),G(1),
*          MAXA(1),MCODE(12,1),TR12(1),TR22(1),TR13(1),TR23(1),
*          SKT(1),SP12(1),SP22(1),SP13(1),SP23(1),ST12(1),
*          ST22(1),ST13(1),ST23(1),ZI2(1),ZI3(1),ZJ(1),
*          RM32(1),RM33(1),QL(1),TU(1),KT(1),RL1PI(1),RL2PI(1),
*          RL3PI(1),RL1(1),RL2(1),RL3(1),RL(1),C1(1),C2(1),
*          C3(1),OLAMD(3,3,60),ELENGO(1)
  IF (NPRINT .EQ. 3) THEN DO
    PRINT 100
100  FORMAT(// ' STIFF CALLED' )
  END IF
  DO 10 I=1,NKT
    SKT(I)=0. D00
10  CONTINUE
  I=1
C
  WHILE (I .LE. NE) DO
    IF (INDTAN .EQ. 1) THEN DO
      CALL ELEMS1(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,ELENG,
*          EMOD,GMOD,G,TR12,TR22,TR13,TR23,SP12,SP22,
*          SP13,SP23,ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,I,
*          IELS,NE,RM32,RM33,QL,NPRINT,KT,OLAMD,ELENGO)
    END IF
    IF (INDTAN .EQ. 2) THEN DO
      CALL ELEMS2(AREA,C1,C2,C3,ELENGO,EMOD,GMOD,G,TR12,
*          TR22,TR13,TR23,TU,I,ZI2,ZI3,ZJ,KT,NPRINT,QL,
*          OLAMD)
    END IF
    CALL ASSEMS(G,MAXA,MCODE,SKT,I,NKT,NPRINT)
    I=I+1
  END WHILE
C
  IF (NPRINT .EQ. 3) THEN DO
    PRINT 200
200  FORMAT(// ' 0' ,7X, ' STIFFNESS MATRIX' )
    PRINT 300,(SKT(K),K=1,NKT)
300  FORMAT(' ',6(5X,D15.7))
  END IF
  RETURN
  END
C
C
C *****
*                               ELEMS1                               *
C *****
SUBROUTINE ELEMS1(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,ELENG,
*          EMOD,GMOD,G,TR12,TR22,TR13,TR23,SP12,SP22,
*          SP13,SP23,ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,
*          I,IELS,NE,RM32,RM33,QL,NPRINT,KT,OLAMD,
*          ELENGO)

```

C COMPUTE THE COMPONENTS OF THE ELEMENT TANGENT STIFFNESS MATRIX  
C FOR THE BEAM-COLUMN MODEL.

```

IMPLICIT REAL*8(A-H,O-Z)
DIMENSION AREA(1),BP12(1),BP22(1),BP13(1),BP23(1),C1(1),
*      C2(1),C3(1),ELENG(1),EMOD(1),GMOD(1),G(1),
*      TR12(1),TR22(1),TR13(1),TR23(1),SP12(1),SP22(1),
*      SP13(1),SP23(1),ST12(1),ST22(1),ST13(1),ST23(1),
*      ZI2(1),ZI3(1),ZJ(1),RM32(1),RM33(1),QL(1),KT(1),
*      OLAMD(3,3,60),ELENGO(1)
PI=3.141592653589793
IF (IELS .EQ. 0) THEN DO
  ST12(I)=4. D00
  ST22(I)=2. D00
  ST13(I)=4. D00
  ST23(I)=2. D00
  G12=0. D00
  G22=0. D00
  G13=0. D00
  G23=0. D00
  H=(PI**2*ZI2(I))/(AREA(I)*ELENGO(I)**2)
  IF (I .EQ. NE) THEN DO
    IELS=1
  END IF
ELSE DO
  G12=SP12(I)*TR12(I)+SP22(I)*TR22(I)
  G22=SP22(I)*TR12(I)+SP12(I)*TR22(I)
  G13=SP13(I)*TR13(I)+SP23(I)*TR23(I)
  G23=SP23(I)*TR13(I)+SP13(I)*TR23(I)
  H=(PI**2*ZI2(I))/(AREA(I)*ELENGO(I)**2)+BP12(I)*(TR12(I)
*      +TR22(I))**2+BP22(I)*(TR12(I)-TR22(I))**2+ZI2(I)/ZI3(I)*
*      (BP13(I)*(TR13(I)+TR23(I))**2
*      +BP23(I)*(TR13(I)-TR23(I))**2)
END IF

```

C  
C  
C

```

COMPUTE THE COMPONENTS OF THE CONVENTIONAL LINEAR STIFFNESS
MATRIX IN THE LOCAL COORDINATE.
ETA=GMOD(1)*ZJ(1)/(EMOD(1)*ZI2(1))
GP1=PI**2/(H*ELENGO(I)**2)
GP2=(G13+G23)/(H*ELENGO(I)**2)
GP3=(-G12-G22)/(H*ELENGO(I)**2)
GP4=G12/(H*ELENGO(I))
GP5=G13/(H*ELENGO(I))
GP6=G22/(H*ELENGO(I))
GP7=G23/(H*ELENGO(I))
GP8=2. D0*ZI3(I)/ZI2(I)*(ST13(I)+ST23(I))/ELENGO(I)**2
*      +(G13+G23)**2/(ELENGO(I)**2*PI**2*H)
GP9=ZI3(I)/ZI2(I)*(ST13(I)+ST23(I))/ELENGO(I)+(G13**2
*      +G13*G23)/(ELENGO(I)*PI**2*H)
GP10=ZI3(I)/ZI2(I)*(ST13(I)+ST23(I))/ELENGO(I)+(G23**2
*      +G13*G23)/(ELENGO(I)*PI**2*H)
GP11=2. D0*(ST12(I)+ST22(I))/ELENGO(I)**2+(G12+G22)**2/
*      (ELENGO(I)**2*PI**2*H)
GP12=- (ST12(I)+ST22(I))/ELENGO(I)-(G12**2+G12*G22)/
*      (ELENGO(I)*PI**2*H)
GP13=- (ST12(I)+ST22(I))/ELENGO(I)-(G22**2+G12*G22)/
*      (ELENGO(I)*PI**2*H)
GP14=ETA
GP15=ST12(I)+G12**2/(PI**2*H)
GP16=ST22(I)+G12*G22/(PI**2*H)
GP17=ZI3(I)/ZI2(I)*ST13(I)+G13**2/(PI**2*H)
GP18=ZI3(I)/ZI2(I)*ST23(I)+G13*G23/(PI**2*H)
GP19=ST12(I)+G22**2/(PI**2*H)
GP20=ZI3(I)/ZI2(I)*ST13(I)+G23**2/(PI**2*H)

```

C  
C  
C

COMPUTE THE COMPONENTS OF THE CONVENTIONAL LINEAR STIFFNESS  
MATRIX IN THE GLOBAL COORDINATE.



```

ALPHA=EMOD(I)*ZI2(I)/ELENGO(I)
G(1)=(C1(I)**2*GP1+2. DO*C1(I)*OLAMD(2,1,I)*GP2
* +2. DO*C1(I)*OLAMD(3,1,I)*GP3+OLAMD(2,1,I)**2*GP8
* +OLAMD(3,1,I)**2*GP11)*ALPHA
G(2)=(C1(I)*C2(I)*GP1+(C1(I)*OLAMD(2,2,I)+C2(I)*
* OLAMD(2,1,I))*GP2+(C1(I)*OLAMD(3,2,I)+C2(I)*
* OLAMD(3,1,I))*GP3+OLAMD(2,1,I)*OLAMD(2,2,I)*GP8
* +OLAMD(3,1,I)*OLAMD(3,2,I)*GP11)*ALPHA
G(3)=(C1(I)*C3(I)*GP1+(C1(I)*OLAMD(2,3,I)+C3(I)*
* OLAMD(2,1,I))*GP2+(C1(I)*OLAMD(3,3,I)+C3(I)*
* OLAMD(3,1,I))*GP3+OLAMD(2,1,I)*OLAMD(2,3,I)*GP8
* +OLAMD(3,1,I)*OLAMD(3,3,I)*GP11)*ALPHA
G(4)=(C1(I)*OLAMD(2,1,I)*GP4+C1(I)*OLAMD(3,1,I)*GP5
* +OLAMD(2,1,I)*OLAMD(3,1,I)*GP9+OLAMD(3,1,I)*OLAMD(2,1,I)*
* GP12)*ALPHA
G(5)=(C1(I)*OLAMD(2,2,I)*GP4+C1(I)*OLAMD(3,2,I)*GP5
* +OLAMD(2,1,I)*OLAMD(3,2,I)*GP9+OLAMD(3,1,I)*OLAMD(2,2,I)*
* GP12)*ALPHA
G(6)=(C1(I)*OLAMD(2,3,I)*GP4+C1(I)*OLAMD(3,3,I)*GP5
* +OLAMD(2,1,I)*OLAMD(3,3,I)*GP9+OLAMD(3,1,I)*OLAMD(2,3,I)*
* GP12)*ALPHA
G(7)=(C1(I)*OLAMD(2,1,I)*GP6+C1(I)*OLAMD(3,1,I)*GP7
* +OLAMD(2,1,I)*OLAMD(3,1,I)*GP10+OLAMD(3,1,I)*
* OLAMD(2,1,I)*GP13)*ALPHA
G(8)=(C1(I)*OLAMD(2,2,I)*GP6+C1(I)*OLAMD(3,2,I)*GP7
* +OLAMD(2,1,I)*OLAMD(3,2,I)*GP10+OLAMD(3,1,I)*
* OLAMD(2,2,I)*GP13)*ALPHA
G(9)=(C1(I)*OLAMD(2,3,I)*GP6+C1(I)*OLAMD(3,3,I)*GP7
* +OLAMD(2,1,I)*OLAMD(3,3,I)*GP10+OLAMD(3,1,I)*
* OLAMD(2,3,I)*GP13)*ALPHA
G(10)=(C2(I)**2*GP1+2. DO*C2(I)*OLAMD(2,2,I)*GP2+2. DO*C2(I)*
* OLAMD(3,2,I)*GP3+OLAMD(2,2,I)**2*GP8+OLAMD(3,2,I)**2*
* GP11)*ALPHA
G(11)=(C2(I)*C3(I)*GP1+(C2(I)*OLAMD(2,3,I)+C3(I)*
* OLAMD(2,2,I))*GP2+(C2(I)*OLAMD(3,3,I)+C3(I)*
* OLAMD(3,2,I))*GP3+OLAMD(2,2,I)*OLAMD(2,3,I)*GP8
* +OLAMD(3,2,I)*OLAMD(3,3,I)*GP11)*ALPHA
G(12)=(C2(I)*OLAMD(2,1,I)*GP4+C2(I)*OLAMD(3,1,I)*GP5
* +OLAMD(2,2,I)*OLAMD(3,1,I)*GP9+OLAMD(3,2,I)*
* OLAMD(2,1,I)*GP12)*ALPHA
G(13)=(C2(I)*OLAMD(2,2,I)*GP4+C2(I)*OLAMD(3,2,I)*GP5
* +OLAMD(2,2,I)*OLAMD(3,2,I)*GP9+OLAMD(3,2,I)*
* OLAMD(2,2,I)*GP12)*ALPHA
G(14)=(C2(I)*OLAMD(2,3,I)*GP4+C2(I)*OLAMD(3,3,I)*GP5
* +OLAMD(2,2,I)*OLAMD(3,3,I)*GP9+OLAMD(3,2,I)*
* OLAMD(2,3,I)*GP12)*ALPHA
G(15)=(C2(I)*OLAMD(2,1,I)*GP6+C2(I)*OLAMD(3,1,I)*GP7
* +OLAMD(2,2,I)*OLAMD(3,1,I)*GP10+OLAMD(3,2,I)*
* OLAMD(2,1,I)*GP13)*ALPHA
G(16)=(C2(I)*OLAMD(2,2,I)*GP6+C2(I)*OLAMD(3,2,I)*GP7
* +OLAMD(2,2,I)*OLAMD(3,2,I)*GP10+OLAMD(3,2,I)*
* OLAMD(2,2,I)*GP13)*ALPHA
G(17)=(C2(I)*OLAMD(2,3,I)*GP6+C2(I)*OLAMD(3,3,I)*GP7
* +OLAMD(2,2,I)*OLAMD(3,3,I)*GP10+OLAMD(3,2,I)*
* OLAMD(2,3,I)*GP13)*ALPHA
G(18)=(C3(I)**2*GP1+2. DO*C3(I)*OLAMD(2,3,I)*GP2+2. DO*
* C3(I)*OLAMD(3,3,I)*GP3+OLAMD(2,3,I)**2*GP8
* +OLAMD(3,3,I)**2*GP11)*ALPHA
G(19)=(C3(I)*OLAMD(2,1,I)*GP4+C3(I)*OLAMD(3,1,I)*GP5
* +OLAMD(2,3,I)*OLAMD(3,1,I)*GP9+OLAMD(3,3,I)*
* OLAMD(2,1,I)*GP12)*ALPHA
G(20)=(C3(I)*OLAMD(2,2,I)*GP4+C3(I)*OLAMD(3,2,I)*GP5
* +OLAMD(2,3,I)*OLAMD(3,2,I)*GP9+OLAMD(3,3,I)*
* OLAMD(2,2,I)*GP12)*ALPHA
G(21)=(C3(I)*OLAMD(2,3,I)*GP4+C3(I)*OLAMD(3,3,I)*GP5
* +OLAMD(2,3,I)*OLAMD(3,3,I)*GP9+OLAMD(3,3,I)*

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*      OLAMD(2,3,I)*GP12)*ALPHA
G(22)=(C3(I)*OLAMD(2,1,I)*GP6+C3(I)*OLAMD(3,1,I)*GP7
*      +OLAMD(2,3,I)*OLAMD(3,1,I)*GP10+OLAMD(3,3,I)*
*      OLAMD(2,1,I)*GP13)*ALPHA
G(23)=(C3(I)*OLAMD(2,2,I)*GP6+C3(I)*OLAMD(3,2,I)*GP7
*      +OLAMD(2,3,I)*OLAMD(3,2,I)*GP10+OLAMD(3,3,I)*
*      OLAMD(2,2,I)*GP13)*ALPHA
G(24)=(C3(I)*OLAMD(2,3,I)*GP6+C3(I)*OLAMD(3,3,I)*GP7
*      +OLAMD(2,3,I)*OLAMD(3,3,I)*GP10+OLAMD(3,3,I)*
*      OLAMD(2,3,I)*GP13)*ALPHA
G(25)=(C1(I)**2*GP14+OLAMD(2,1,I)**2*GP15+OLAMD(3,1,I)**2*
*      GP17)*ALPHA
G(26)=(C1(I)*C2(I)*GP14+OLAMD(2,1,I)*OLAMD(2,2,I)*GP15
*      +OLAMD(3,1,I)*OLAMD(3,2,I)*GP17)*ALPHA
G(27)=(C1(I)*C3(I)*GP14+OLAMD(2,1,I)*OLAMD(2,3,I)*GP15
*      +OLAMD(3,1,I)*OLAMD(3,3,I)*GP17)*ALPHA
G(28)=(-C1(I)**2*GP14+OLAMD(2,1,I)**2*GP16+OLAMD(3,1,I)**2*
*      GP18)*ALPHA
G(29)=(-C1(I)*C2(I)*GP14+OLAMD(2,1,I)*OLAMD(2,2,I)*GP16
*      +OLAMD(3,1,I)*OLAMD(3,2,I)*GP18)*ALPHA
G(30)=(-C1(I)*C3(I)*GP14+OLAMD(2,1,I)*OLAMD(2,3,I)*GP16
*      +OLAMD(3,1,I)*OLAMD(3,3,I)*GP18)*ALPHA
G(31)=(C2(I)**2*GP14+OLAMD(2,2,I)**2*GP15+OLAMD(3,2,I)**2*
*      GP17)*ALPHA
G(32)=(C2(I)*C3(I)*GP14+OLAMD(2,2,I)*OLAMD(2,3,I)*GP15
*      +OLAMD(3,2,I)*OLAMD(3,3,I)*GP17)*ALPHA
G(33)=(-C2(I)**2*GP14+OLAMD(2,2,I)**2*GP16+OLAMD(3,2,I)**2*
*      GP18)*ALPHA
G(34)=(-C2(I)*C3(I)*GP14+OLAMD(2,2,I)*OLAMD(2,3,I)*GP16
*      +OLAMD(3,2,I)*OLAMD(3,3,I)*GP18)*ALPHA
G(35)=(C3(I)**2*GP14+OLAMD(2,3,I)**2*GP15+OLAMD(3,3,I)**2*
*      GP17)*ALPHA
G(36)=(-C3(I)**2*GP14+OLAMD(2,3,I)**2*GP16+OLAMD(3,3,I)**2*
*      GP18)*ALPHA
G(37)=(C1(I)**2*GP14+OLAMD(2,1,I)**2*GP19+OLAMD(3,1,I)**2*
*      GP20)*ALPHA
G(38)=(C1(I)*C2(I)*GP14+OLAMD(2,1,I)*OLAMD(2,2,I)*GP19
*      +OLAMD(3,1,I)*OLAMD(3,2,I)*GP20)*ALPHA
G(39)=(C1(I)*C3(I)*GP14+OLAMD(2,1,I)*OLAMD(2,3,I)*GP19
*      +OLAMD(3,1,I)*OLAMD(3,3,I)*GP20)*ALPHA
G(40)=(C2(I)**2*GP14+OLAMD(2,2,I)**2*GP19+OLAMD(3,2,I)**2*
*      GP20)*ALPHA
G(41)=(C2(I)*C3(I)*GP14+OLAMD(2,2,I)*OLAMD(2,3,I)*GP19
*      +OLAMD(3,2,I)*OLAMD(3,3,I)*GP20)*ALPHA
G(42)=(C3(I)**2*GP14+OLAMD(2,3,I)**2*GP19+OLAMD(3,3,I)**2*
*      GP20)*ALPHA

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C

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IF (KT(2).GT. 0) THEN DO
G(1)=G(1)+RM33(I)*(2.DO*C1(I)*OLAMD(2,1,I))
*      +RM32(I)*(-2.DO*C1(I)*OLAMD(3,1,I))
*      +QL(I)*(-OLAMD(2,1,I)**2-OLAMD(3,1,I)**2)
G(2)=G(2)+RM33(I)*(C2(I)*OLAMD(2,1,I)+C1(I)*OLAMD(2,2,I))
*      +RM32(I)*(-C2(I)*OLAMD(3,1,I)-C1(I)*OLAMD(3,2,I))
*      +QL(I)*(-OLAMD(2,1,I)*OLAMD(2,2,I)
*      -OLAMD(3,1,I)*OLAMD(3,2,I))
G(3)=G(3)+RM33(I)*(C3(I)*OLAMD(2,1,I)+C1(I)*OLAMD(2,3,I))
*      +RM32(I)*(-C3(I)*OLAMD(3,1,I)-C1(I)*OLAMD(3,3,I))
*      +QL(I)*(-OLAMD(2,1,I)*OLAMD(2,3,I)
*      -OLAMD(3,1,I)*OLAMD(3,3,I))
G(10)=G(10)+RM33(I)*(2.DO*C2(I)*OLAMD(2,2,I))
*      +RM32(I)*(-2.DO*C2(I)*OLAMD(3,2,I))
*      +QL(I)*(-OLAMD(2,2,I)**2-OLAMD(3,2,I)**2)
G(11)=G(11)+RM33(I)*(C3(I)*OLAMD(2,2,I)
*      +C2(I)*OLAMD(2,3,I))
*      +RM32(I)*(-C3(I)*OLAMD(3,2,I)
*      -C2(I)*OLAMD(3,3,I))

```

```

*          +QL(I)*(-OLAMD(2,2,I)*OLAMD(2,3,I)
*          -OLAMD(3,2,I)*OLAMD(3,3,I))
G(18)=G(18)+RM33(I)*(2.DO*C3(I)*OLAMD(2,3,I)
*          +RM32(I)*(-2.DO*C3(I)*OLAMD(3,3,I))
*          +QL(I)*(-OLAMD(2,3,I)**2-OLAMD(3,3,I)**2)
END IF

```

C

```

IF (NPRINT .EQ. 3) THEN DO
PRINT 100,I,(G(KI),KI=1,42)
100 FORMAT(//'-',2X,'G(42) FOR ELE.',I3/5(5X,D15.7)/
*       5(5X,D15.7)/5(5X,D15.7)/5(5X,D15.7)/5(5X,D15.7)/
*       5(5X,D15.7)/5(5X,D15.7)/5(5X,D15.7)/2(5X,D15.7)/)
END IF
RETURN
END

```

C

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C

C

C

C

C

```

*****
*          ELEMS2          *
*****
SUBROUTINE ELEMS2(AREA,C1,C2,C3,ELENGO,EMOD,GMOD,G,TR12,
*          TR22,TR13,TR23,TU,I,ZI2,ZI3,ZJ,KT,NPRINT,
*          QL,OLAMD)
C COMPUTE THE COMPONENTS OF THE ELEMENT TANGENT STIFFNESS MATRIX
C FOR THE FINITE ELEMENT MODEL.
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION AREA(1),C1(1),C2(1),C3(1),ELENGO(1),EMOD(1),GMOD(1),
*          G(1),TR12(1),TR22(1),TR13(1),TR23(1),TU(1),ZI2(1),
*          ZI3(1),ZJ(1),KT(1),QL(1),OLAMD(3,3,60)

```

C

C

```

A=AREA(I)
E=EMOD(I)
GGMOD=GMOD(I)
EL=ELENGO(I)
Z2=ZI2(I)
Z3=ZI3(I)
ZZJ=ZJ(I)
CC1=C1(I)
CC2=C2(I)
CC3=C3(I)
T1=T2=T3=T4=T5=T6=T7=T8=T9=T10=0.DO

```

C

C

C

```

COMPUTE THE CONTRIBUTION OF THE CONVENTIONAL LINEAR STIFFNESS
MATRIX IN THE LOCAL COORDINATE.
IF (KT(1) .GT. 0) THEN DO
T1=T1+E*A/EL
T2=T2+12.DO*E*Z3/EL**3
T3=T3+6.DO*E*Z3/EL**2
T4=T4+12.DO*E*Z2/EL**3
T5=T5-6.DO*E*Z2/EL**2
T6=T6+GGMOD*ZZJ/EL
T7=T7+4.DO*E*Z2/EL
T8=T8+2.DO*E*Z2/EL
T9=T9+4.DO*E*Z3/EL
T10=T10+2.DO*E*Z3/EL
END IF

```

C

C

C

```

COMPUTE THE CONTRIBUTION OF THE INITIAL STRESS (GEOMETRIC)
STIFFNESS MATRIX IN THE LOCAL COORDINATE.
IF (KT(2) .GT. 0) THEN DO
T2=T2+1.2DO*QL(I)/EL
T3=T3+QL(I)/10.DO
T4=T4+1.2DO*QL(I)/EL
T5=T5-QL(I)/10.DO
T7=T7+2.DO*QL(I)*EL/15.DO
T8=T8-QL(I)*EL/30.DO

```

T9=T9+2. D0\*QL(I)\*EL/15. D0  
 T10=T10-QL(I)\*EL/30. D0  
 END IF

C  
 C

COMPUTE THE COMPONENTS OF THE GLOBAL ELEMENT STIFFNESS MATRIX.  
 G(1)=T1\*CC1\*\*2+T2\*OLAMD(2,1,I)\*\*2+T4\*OLAMD(3,1,I)\*\*2  
 G(2)=T1\*CC1\*CC2+T2\*OLAMD(2,1,I)\*OLAMD(2,2,I)  
 \* +T4\*OLAMD(3,1,I)\*OLAMD(3,2,I)  
 G(3)=T1\*CC1\*CC3+T2\*OLAMD(2,1,I)\*OLAMD(2,3,I)  
 \* +T4\*OLAMD(3,1,I)\*OLAMD(3,3,I)  
 G(4)=T3\*OLAMD(2,1,I)\*OLAMD(3,1,I)+T5\*OLAMD(3,1,I)\*OLAMD(2,1,I)  
 G(5)=T3\*OLAMD(2,1,I)\*OLAMD(3,2,I)+T5\*OLAMD(3,1,I)\*OLAMD(2,2,I)  
 G(6)=T3\*OLAMD(2,1,I)\*OLAMD(3,3,I)+T5\*OLAMD(3,1,I)\*OLAMD(2,3,I)  
 G(7)=T3\*OLAMD(2,1,I)\*OLAMD(3,1,I)+T5\*OLAMD(3,1,I)\*OLAMD(2,1,I)  
 G(8)=T3\*OLAMD(2,1,I)\*OLAMD(3,2,I)+T5\*OLAMD(3,1,I)\*OLAMD(2,2,I)  
 G(9)=T3\*OLAMD(2,1,I)\*OLAMD(3,3,I)+T5\*OLAMD(3,1,I)\*OLAMD(2,3,I)  
 G(10)=T1\*CC2\*\*2+T2\*OLAMD(2,2,I)\*\*2+T4\*OLAMD(3,2,I)\*\*2  
 G(11)=T1\*CC2\*CC3+T2\*OLAMD(2,2,I)\*OLAMD(2,3,I)  
 \* +T4\*OLAMD(3,2,I)\*OLAMD(3,3,I)  
 G(12)=T3\*OLAMD(2,2,I)\*OLAMD(3,1,I)+T5\*OLAMD(3,2,I)\*  
 \* OLAMD(2,1,I)  
 G(13)=T3\*OLAMD(2,2,I)\*OLAMD(3,2,I)+T5\*OLAMD(3,2,I)\*  
 \* OLAMD(2,2,I)  
 G(14)=T3\*OLAMD(2,2,I)\*OLAMD(3,3,I)+T5\*OLAMD(3,2,I)\*  
 \* OLAMD(2,3,I)  
 G(15)=T3\*OLAMD(2,2,I)\*OLAMD(3,1,I)+T5\*OLAMD(3,2,I)\*  
 \* OLAMD(2,1,I)  
 G(16)=T3\*OLAMD(2,2,I)\*OLAMD(3,2,I)+T5\*OLAMD(3,2,I)\*  
 \* OLAMD(2,2,I)  
 G(17)=T3\*OLAMD(2,2,I)\*OLAMD(3,3,I)+T5\*OLAMD(3,2,I)\*  
 \* OLAMD(2,3,I)  
 G(18)=T1\*CC3\*\*2+T2\*OLAMD(2,3,I)\*\*2+T4\*OLAMD(3,3,I)\*\*2  
 G(19)=T3\*OLAMD(2,3,I)\*OLAMD(3,1,I)+T5\*OLAMD(3,3,I)\*  
 \* OLAMD(2,1,I)  
 G(20)=T3\*OLAMD(2,3,I)\*OLAMD(3,2,I)+T5\*OLAMD(3,3,I)\*  
 \* OLAMD(2,2,I)  
 G(21)=T3\*OLAMD(2,3,I)\*OLAMD(3,3,I)+T5\*OLAMD(3,3,I)\*  
 \* OLAMD(2,3,I)  
 G(22)=T3\*OLAMD(2,3,I)\*OLAMD(3,1,I)+T5\*OLAMD(3,3,I)\*  
 \* OLAMD(2,1,I)  
 G(23)=T3\*OLAMD(2,3,I)\*OLAMD(3,2,I)+T5\*OLAMD(3,3,I)\*  
 \* OLAMD(2,2,I)  
 G(24)=T3\*OLAMD(2,3,I)\*OLAMD(3,3,I)+T5\*OLAMD(3,3,I)\*  
 \* OLAMD(2,3,I)  
 G(25)=T6\*CC1\*\*2+T9\*OLAMD(3,1,I)\*\*2+T7\*OLAMD(2,1,I)\*\*2  
 G(26)=T6\*CC1\*CC2+T9\*OLAMD(3,1,I)\*OLAMD(3,2,I)  
 \* +T7\*OLAMD(2,1,I)\*OLAMD(2,2,I)  
 G(27)=T6\*CC1\*CC3+T9\*OLAMD(3,1,I)\*OLAMD(3,3,I)  
 \* +T7\*OLAMD(2,1,I)\*OLAMD(2,3,I)  
 G(28)=-T6\*CC1\*\*2+T10\*OLAMD(3,1,I)\*\*2+T8\*OLAMD(2,1,I)\*\*2  
 G(29)=-T6\*CC1\*CC2+T10\*OLAMD(3,1,I)\*OLAMD(3,2,I)  
 \* +T8\*OLAMD(2,1,I)\*OLAMD(2,2,I)  
 G(30)=-T6\*CC1\*CC3+T10\*OLAMD(3,1,I)\*OLAMD(3,3,I)  
 \* +T8\*OLAMD(2,1,I)\*OLAMD(2,3,I)  
 G(31)=T6\*CC2\*\*2+T9\*OLAMD(3,2,I)\*\*2+T7\*OLAMD(2,2,I)\*\*2  
 G(32)=T6\*CC2\*CC3+T9\*OLAMD(3,2,I)\*OLAMD(3,3,I)  
 \* +T7\*OLAMD(2,2,I)\*OLAMD(2,3,I)  
 G(33)=-T6\*CC2\*\*2+T10\*OLAMD(3,2,I)\*\*2+T8\*OLAMD(2,2,I)\*\*2  
 G(34)=-T6\*CC2\*CC3+T10\*OLAMD(3,2,I)\*OLAMD(3,3,I)  
 \* +T8\*OLAMD(2,2,I)\*OLAMD(2,3,I)  
 G(35)=T6\*CC3\*\*2+T9\*OLAMD(3,3,I)\*\*2+T7\*OLAMD(2,3,I)\*\*2  
 G(36)=-T6\*CC3\*\*2+T10\*OLAMD(3,3,I)\*\*2+T8\*OLAMD(2,3,I)\*\*2  
 G(37)=T6\*CC1\*\*2+T9\*OLAMD(3,1,I)\*\*2+T7\*OLAMD(2,1,I)\*\*2  
 G(38)=T6\*CC1\*CC2+T9\*OLAMD(3,1,I)\*OLAMD(3,2,I)  
 \* +T7\*OLAMD(2,1,I)\*OLAMD(2,2,I)  
 G(39)=T6\*CC1\*CC3+T9\*OLAMD(3,1,I)\*OLAMD(3,3,I)

```

*      +T7*OLAMD(2,1,I)*OLAMD(2,3,I)
G(40)=T6*CC2**2+T9*OLAMD(3,2,I)**2+T7*OLAMD(2,2,I)**2
G(41)=T6*CC2*CC3+T9*OLAMD(3,2,I)*OLAMD(3,3,I)
*      +T7*OLAMD(2,2,I)*OLAMD(2,3,I)
G(42)=T6*CC3**2+T9*OLAMD(3,3,I)**2+T7*OLAMD(2,3,I)**2
RETURN
END

```

C  
C  
C  
C  
C  
C  
C

```

*****
*                                     ASSEMS                                     *
*****

```

```

SUBROUTINE ASSEMS(G,MAXA,MCODE,SKT,I,NKT,NPRINT)
ASSEMBLE THE CONTRIBUTIONS FROM EACH ELEMENT INTO THE GLOBAL
SYSTEM STIFFNESS MATRIX.
IMPLICIT REAL*8(A-H,O-Z)

```

```

DIMENSION G(1),MAXA(1),MCODE(12,1),SKT(1)

```

```

INTEGER INDEX(12,12)/1,2,3,4,5,6,-1,-2,-3,7,8,9,
*      2,10,11,12,13,14,-2,-10,-11,15,16,17,
*      3,11,18,19,20,21,-3,-11,-18,22,23,24,
*      4,12,19,25,26,27,-4,-12,-19,28,29,30,
*      5,13,20,26,31,32,-5,-13,-20,29,33,34,
*      6,14,21,27,32,35,-6,-14,-21,30,34,36,
*      -1,-2,-3,-4,-5,-6,1,2,3,-7,-8,-9,
*      -2,-10,-11,-12,-13,-14,2,10,11,-15,-16,-17,
*      -3,-11,-18,-19,-20,-21,3,11,18,-22,-23,-24,
*      7,15,22,28,29,30,-7,-15,-22,37,38,39,
*      8,16,23,29,33,34,-8,-16,-23,38,40,41,
*      9,17,24,30,34,36,-9,-17,-24,39,41,42/

```

```

DO 20 JE=1,12
  J=MCODE(JE,I)
  IF (J.NE.0) THEN DO
    DO 10 NE=1,JE
      N=MCODE(NE,I)
      IF (N.NE.0) THEN DO
        K=MAXA(J)+J-N
        L=INDEX(NE,JE)
        IF (L.GT.0) THEN DO
          SKT(K)=SKT(K)+G(L)
        ELSE DO
          SKT(K)=SKT(K)-G(-L)
        END IF
      END IF
    END IF
  CONTINUE
END IF
10  CONTINUE
END IF
20  CONTINUE
IF (NPRINT.EQ.3) THEN DO
  PRINT 100,I,(SKT(KI),KI=1,NKT)
100  FORMAT(//',7X,'STIFFNESS FROM ELEMENT ',I3,/6(5X,D15.7))
END IF
RETURN
END

```

C  
C  
C  
C  
C  
C  
C

```

*****
*                                     STORE                                     *
*****

```

```

SUBROUTINE STORE(RL1PI,RL2PI,RL3PI,ELENG,RL1,RL2,RL3,RL,C1PI,
*      C2PI,C3PI,C1,C2,C3,NE)

```

```

STORE THE PREVIOUS ITERATION LENGTH AND DIRECTION COSINES OF
ELEMENT I

```

```

IMPLICIT REAL*8(A-H,O-Z)

```

```

DIMENSION RL1PI(1),RL2PI(1),RL3PI(1),ELENG(1),RL1(1),RL2(1),
*      RL3(1),RL(1),C1PI(1),C2PI(1),C3PI(1),C1(1),C2(1),

```

```

*          C3(1)
C          DO 30 I=1,NE
            RL1PI(I)=RL1(I)
            RL2PI(I)=RL2(I)
            RL3PI(I)=RL3(I)
            ELENG(I)=RL(I)
            C1PI(I)=C1(I)
            C2PI(I)=C2(I)
            C3PI(I)=C3(I)
30 CONTINUE
RETURN
END

C
C
C
C
*****
*          FORCES          *
*****
SUBROUTINE FORCES(AREA, BP12, BP22, BP13, BP23, C1, C2, C3, DD, DE,
*          ELENG, EMOD, GMOD, F, FG, FL, MCODE, QT, R12, R22,
*          R13, R23, TORS, TR12, TR22, TR13, TR23, TTORS, SP12,
*          SP22, SP13, SP23, ST12, ST22, ST13, ST23, ZI2, ZI3,
*          ZJ, IMAX, NE, NEQ, O, QI, C1PI, C2PI, C3PI, RM32,
*          RM33, QL, TU, INDTAN, NPRINT, RL1PI, RL2PI, RL3PI,
*          RL1, RL2, RL3, RL, TUL, TCB2, TCB3, DDJ, JCODE, AP,
*          APNEW, RJ, MINC, OL, C01, C02, C03, DL, R, OLAMD,
*          OLANEW, DQQR, FQR, ELENGO, NJ, DDO, ITIND, NN)
C  FOR EACH ELEMENT CALL ELEMF.
  IMPLICIT REAL*8(A-H, O-Z)
  DIMENSION AREA(1), BP12(1), BP22(1), BP13(1), BP23(1), C1(1), C2(1),
*          C3(1), DD(1), DE(1), ELENG(1), EMOD(1), GMOD(1), F(1),
*          FG(12,1), FL(12,1), MCODE(12,1), QT(1), R12(1), R22(1),
*          R13(1), R23(1), TORS(1), TR12(1), TR22(1), TR13(1),
*          TR23(1), TTORS(1), SP12(1), SP22(1), SP13(1), SP23(1),
*          ST12(1), ST22(1), ST13(1), ST23(1), ZI2(1), ZI3(1), ZJ(1),
*          Q(1), C1PI(1), C2PI(1), C3PI(1), RM32(1), RM33(1), QL(1),
*          U(60), RL1PI(1), RL2PI(1), RL3PI(1), RL1(1), RL2(1),
*          RL3(1), RL(1), TUL(1), TCB2(1), TCB3(1), TU(1), DDJ(6,1),
*          JCODE(6,1), AP(3,3,60), APNEW(3,3,60), RJ(3,1),
*          MINC(2,1), OL(1), C01(1), C02(1), C03(1), DL(1),
*          R(3,3,60), OLAMD(3,3,60), OLANEW(3,3,60), ELENGO(1),
*          DDO(1)
C
C
C
C
GENERATE THE INCREMENTAL NODAL DISPLACEMENT MATRIX
DDJ(IDIR, J).
DO 28 J=1, NJ
  DO 12 IDIR=1, 6
    DDJ(IDIR, J)=0. DO
    N=JCODE(IDIR, J)
    IF (N .NE. 0) THEN DO
      IF (NN .EQ. 1 .AND. ITIND .EQ. 1) THEN DO
        DDJ(IDIR, J)=DDO(N)
      ELSE DO
        DDJ(IDIR, J)=DD(N)
      END IF
    END IF
  END IF
12 CONTINUE
C
C
CALCULATE THE INCREMENTAL NODAL ROTATIONS AT NODE J.
*  OME1=AP(1, 1, J)*DDJ(4, J)+AP(1, 2, J)*DDJ(5, J)
*      +AP(1, 3, J)*DDJ(6, J)
*  OME2=AP(2, 1, J)*DDJ(4, J)+AP(2, 2, J)*DDJ(5, J)
*      +AP(2, 3, J)*DDJ(6, J)
*  OME3=AP(3, 1, J)*DDJ(4, J)+AP(3, 2, J)*DDJ(5, J)

```

```

*          +AP(3,3,J)*DDJ(6,J)
C
C   YIELD THE COMPONENTS OF THE ROTATION MATRIX OF NODE J.
RJ(1,1)=1. DO
RJ(1,2)=OME3
RJ(1,3)=-OME2
RJ(2,1)=-OME3
RJ(2,2)=1. DO
RJ(2,3)=OME1
RJ(3,1)=OME2
RJ(3,2)=-OME1
RJ(3,3)=1. DO

C
C   FOR EACH ITERATION EMPLOYING THE LAST NODAL ORIENTATION
C   AND ROTATION MATRICES OF NODE J TO GENERATE THE NEW
C   NODAL ORIENTATION MATRIX.
DO 25 L=1,3
  DO 20 M=1,3
    APNEW(L,M,J)=0. DO
    DO 18 K=1,3
      APNEW(L,M,J)=RJ(L,K)*AP(K,M,J)+APNEW(L,M,J)
18    CONTINUE
20    CONTINUE
25    CONTINUE
    DO 27 L=1,3
      DO 26 K=1,3
        AP(L,K,J)=APNEW(L,K,J)
26    CONTINUE
27    CONTINUE
28  CONTINUE
DO 10 I=1,NEQ
  F(I)=0. DO
10 CONTINUE

C
  I=1
  WHILE (I .LE. NE) DO
    IF (NPRINT .EQ. 3) THEN DO
      PRINT 100,I
100   FORMAT(/' ELEMENT ',I3)
    END IF
    CALL ELEMFB(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,DD,DE,ELENG,
*          EMOD,GMOD,F,FG,FL,MCODE,R12,R22,R13,R23,TORS,
*          TR12,TR22,TR13,TR23,TTORS,SP12,SP22,SP13,SP23,
*          ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,I,IMAX,C1PI,C2PI,
*          C3PI,RM32,RM33,QL,TU,INDTAN,NPRINT,RL1PI,RL2PI,
*          RL3PI,RL1,RL2,RL3,RL,TUL,TCB2,TCB3,DDJ,JCODE,AP,
*          APNEW,RJ,MINC,OL,C01,C02,C03,DL,R,OLAMD,
*          OLANEW,DQQR,FQR,ELENGO,DDO,ITIND,NN)
    I=I+1
  END WHILE
  RETURN
END

C
C
C
C
C *****
*          ELEMFB
*****
SUBROUTINE ELEMFB(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,DD,DE,
*          ELENG,EMOD,GMOD,F,FG,FL,MCODE,R12,R22,
*          R13,R23,TORS,TR12,TR22,TR13,TR23,TTORS,
*          SP12,SP22,SP13,SP23,ST12,ST22,ST13,ST23,
*          ZI2,ZI3,ZJ,I,IMAX,C1PI,C2PI,C3PI,RM32,RM33,
*          QL,TU,INDTAN,NPRINT,RL1PI,RL2PI,RL3PI,RL1,
*          RL2,RL3,RL,TUL,TCB2,TCB3,DDJ,JCODE,AP,APNEW,
*          RJ,MINC,OL,C01,C02,C03,DL,R,OLAMD,OLANEW,

```

```

*          DOOR, FOR, ELENGO, DDO, ITIND, NN)
C COMPUTE THE LOCAL INTERNAL ELEMENT FORCES, FL, DUE TO THE
C DEFORMATIONS; TRANSFORM THEM TO THE GLOBAL ELEMENT FORCES,
C FG; AND COMPUTE THE NODAL FORCES, F, CONTRIBUTED BY ELEMENT I.
C SUBROUTINE ELEMFI IS EXECUTED FOR EACH ELEMENT, I.
  IMPLICIT REAL*8(A-H, O-Z)
  DIMENSION AREA(1), BP12(1), BP22(1), BP13(1), BP23(1), C1(1), C2(1),
*          C3(1), DD(1), DE(1), ELENG(1), EMOD(1), GMOD(1), F(1),
*          FG(12,1), FL(12,1), MCODE(12,1), R12(1), R22(1), R13(1),
*          R23(1), TORS(1), TR12(1), TR22(1), TR13(1), TR23(1),
*          TTORS(1), SP12(1), SP22(1), SP13(1), SP23(1), ST12(1),
*          ST22(1), ST13(1), ST23(1), Z12(1), Z13(1), ZJ(1), C1PI(1),
*          C2PI(1), C3PI(1), RM32(1), RM33(1), QL(1), U(60),
*          RL1PI(1), RL2PI(1), RL3PI(1), RL1(1), RL2(1), RL3(1),
*          RL(1), TUL(1), TCB2(1), TCB3(1), TU(1), DDJ(6,1),
*          JCODE(6,1), AP(3,3,60), APNEW(3,3,60), RJ(3,1),
*          MINC(2,1), OL(1), C01(1), C02(1), C03(1), DL(1),
*          R(3,3,60), OLAMD(3,3,60), OLANEW(3,3,60),
*          ELENGO(1), DDO(1)
  IF (NPRINT .EQ. 3) THEN DO
    PRINT 100, I
100  FORMAT(' DE FOR ELEMENT', I3)
  END IF
  DO 10 J=1, 12
    DE(J)=0. DDO
    K=MCODE(J, I)
    IF (K .NE. 0) THEN DO
      IF (NN .EQ. 1 .AND. ITIND .EQ. 1) THEN DO
        DE(J)=DDO(K)
      ELSE DO
        DE(J)=DD(K)
      END IF
    END IF
    IF (NPRINT .EQ. 3) THEN DO
      PRINT 200, DE(J)
200  FORMAT(5X, D15.7)
    END IF
  10 CONTINUE
C
  RL1PI(I)=ELENG(I)*C1PI(I)
  RL2PI(I)=ELENG(I)*C2PI(I)
  RL3PI(I)=ELENG(I)*C3PI(I)
  RL1(I)=RL1PI(I)+DE(7)-DE(1)
  RL2(I)=RL2PI(I)+DE(8)-DE(2)
  RL3(I)=RL3PI(I)+DE(9)-DE(3)
C  RL(I)=DSQRT(RL1(I)**2+RL2(I)**2+RL3(I)**2)
C
  C1(I)=RL1(I)/RL(I)
  C2(I)=RL2(I)/RL(I)
  C3(I)=RL3(I)/RL(I)
C
C  GENERATE THE INCREMENTS OF THE LOCAL DISPLACEMENTS OF
C  ELEMENT I.
  DL(1)=OLAMD(1,1,I)*DE(1)+OLAMD(1,2,I)*DE(2)
*  +OLAMD(1,3,I)*DE(3)
  DL(2)=OLAMD(2,1,I)*DE(1)+OLAMD(2,2,I)*DE(2)
*  +OLAMD(2,3,I)*DE(3)
  DL(3)=OLAMD(3,1,I)*DE(1)+OLAMD(3,2,I)*DE(2)
*  +OLAMD(3,3,I)*DE(3)
  DL(4)=OLAMD(1,1,I)*DE(4)+OLAMD(1,2,I)*DE(5)
*  +OLAMD(1,3,I)*DE(6)
  DL(5)=OLAMD(2,1,I)*DE(4)+OLAMD(2,2,I)*DE(5)
*  +OLAMD(2,3,I)*DE(6)
  DL(6)=OLAMD(3,1,I)*DE(4)+OLAMD(3,2,I)*DE(5)
*  +OLAMD(3,3,I)*DE(6)
  DL(7)=OLAMD(1,1,I)*DE(7)+OLAMD(1,2,I)*DE(8)

```



```

*      +OLAMD(1,3,I)*DE(9)
DL(8)=OLAMD(2,1,1)*DE(7)+OLAMD(2,2,I)*DE(8)
*      +OLAMD(2,3,1)*DE(9)
DL(9)=OLAMD(3,1,1)*DE(7)+OLAMD(3,2,I)*DE(8)
*      +OLAMD(3,3,1)*DE(9)
DL(10)=OLAMD(1,1,I)*DE(10)+OLAMD(1,2,I)*DE(11)
*      +OLAMD(1,3,1)*DE(12)
DL(11)=OLAMD(2,1,1)*DE(10)+OLAMD(2,2,I)*DE(11)
*      +OLAMD(2,3,1)*DE(12)
DL(12)=OLAMD(3,1,1)*DE(10)+OLAMD(3,2,I)*DE(11)
*      +OLAMD(3,3,1)*DE(12)

```

C  
C  
C

GENERATE THE COMPONENTS OF THE ROTATION MATRIX OF  
ELEMENT I, R.

```

R(1,1,I)=1. DO
R(1,2,I)=(DL(8)-DL(2))/ELENG(I)
R(1,3,I)=(DL(9)-DL(3))/ELENG(I)
R(2,1,I)=-R(1,2,I)
R(2,2,I)=1. DO
R(2,3,I)=(DL(10)-DL(4))/2. DO
R(3,1,I)=-R(1,3,I)
R(3,2,I)=-R(2,3,I)
R(3,3,I)=1. DO

```

C  
C  
C  
C

FOR EACH ITERATION EMPLOYING THE LAST ELEMENT ORIENTATION  
AND ROTATION MATRICES TO GENERATE THE NEW ELEMENT  
ORIENTATION MATRIX.

```

DO 34 L=1,3
  DO 32 J=1,3
    OLANEW(L,J,I)=0. DO
    DO 30 K=1,3
      OLANEW(L,J,I)=R(L,K,I)*OLAMD(K,J,I)+OLANEW(L,J,I)
30    CONTINUE
32  CONTINUE
34  CONTINUE
  DO 37 L=1,3
    DO 36 J=1,3
      OLAMD(L,J,I)=OLANEW(L,J,I)
36  CONTINUE
37  CONTINUE

```

C

UL=(ELENGO(I)-RL(I))/ELENGO(I)

C  
C  
C

EMPLOYING THE MINC CODE TO OBTAIN THE ELEMENT DEFORMATIONS  
AT BOTH ENDS, A AND B.

```

L=MINC(1,I)
K=MINC(2,I)
R12(I)=C1(I)*(-C03(I)/OL(I)*AP(1,1,L)+C01(I)/OL(I)*AP(3,1,L))
*      +C2(I)*(-C03(I)/OL(I)*AP(1,2,L)+C01(I)/OL(I)*AP(3,2,L))
*      +C3(I)*(-C03(I)/OL(I)*AP(1,3,L)+C01(I)/OL(I)*AP(3,3,L))
R13(I)=-C1(I)*(-C01(I)*C02(I)/OL(I)*AP(1,1,L)+OL(I)*AP(2,1,L))
*      +(-C02(I)*C03(I)/OL(I))*AP(3,1,L)
*      +C2(I)*(-C01(I)*C02(I)/OL(I)*AP(1,2,L)+OL(I)
*      *AP(2,2,L)+(-C02(I)*C03(I)/OL(I))*AP(3,2,L))
*      +C3(I)*(-C01(I)*C02(I)/OL(I)*AP(1,3,L)+OL(I)
*      *AP(2,3,L)+(-C02(I)*C03(I)/OL(I))*AP(3,3,L))
R22(I)=C1(I)*(-C03(I)/OL(I)*AP(1,1,K)+C01(I)/OL(I)*AP(3,1,K))
*      +C2(I)*(-C03(I)/OL(I)*AP(1,2,K)+C01(I)/OL(I)*AP(3,2,K))
*      +C3(I)*(-C03(I)/OL(I)*AP(1,3,K)+C01(I)/OL(I)*AP(3,3,K))
R23(I)=-C1(I)*(-C01(I)*C02(I)/OL(I)*AP(1,1,K)+OL(I)*AP(2,1,K))
*      +(-C02(I)*C03(I)/OL(I))*AP(3,1,K)
*      +C2(I)*(-C01(I)*C02(I)/OL(I)*AP(1,2,K)+OL(I)
*      *AP(2,2,K)+(-C02(I)*C03(I)/OL(I))*AP(3,2,K))
*      +C3(I)*(-C01(I)*C02(I)/OL(I)*AP(1,3,K)+OL(I)
*      *AP(2,3,K)+(-C02(I)*C03(I)/OL(I))*AP(3,3,K))
TORS(I)=2. DO*(OLAMD(3,1,I)*(-C01(I)*C02(I)/OL(I)*AP(1,1,K)

```

```

*          +OL(I)*AP(2,1,K)+(-C02(I)*C03(I)/OL(I))*
*          AP(3,1,K))+OLAMD(3,2,I)*(-C01(I)*C02(I)/OL(I)*
*          AP(1,2,K)+OL(I)*AP(2,2,K)+(-C02(I)*C03(I)/
*          OL(I))*AP(3,2,K))+OLAMD(3,3,I)*(-C01(I)*C02(I)/
*          OL(I))*AP(1,3,K)+OL(I)*AP(2,3,K)
*          +(-C02(I)*C03(I)/OL(I))*AP(3,3,K)))

```

C  
C

COMPUTE THE TOTAL RELATIVE DEFORMATIONS OF ELEMENT I.

```

TUL(I)=UL
TR12(I)=R12(I)
TR22(I)=R22(I)
TR13(I)=R13(I)
TR23(I)=R23(I)
TTORS(I)=TORS(I)

```

C  
C  
C

COMPUTE THE LOCAL ELEMENT FORCES DUE TO THE RELATIVE DEFORMATIONS BY BEAM-COLUMN MODEL.

IF (INDTAN .EQ. 1) THEN DO

```

*      CALL BOWCOR(AREA, BP12, BP22, BP13, BP23, ELENG, R12, R22, R13, R23,
*      SP12, SP22, SP13, SP23, ST12, ST22, ST13, ST23, ZI2,
*      ZI3, I, IMA, TUL, EMOD, NPRINT, TCB2, TCB3, ELENGO,
*      DQOR, FQR)

```

```

RM12=EMOD(I)*ZI2(I)*((ST12(I)*TR12(I)+ST22(I)*TR22(I))/RL(I)
RM22=EMOD(I)*ZI2(I)*((ST22(I)*TR12(I)+ST12(I)*TR22(I))/RL(I)
RM32(I)=(RM12+RM22)/(ELENGO(I)**2)

```

C

```

RM13=EMOD(I)*ZI3(I)*((ST13(I)*TR13(I)+ST23(I)*TR23(I))/RL(I)
RM23=EMOD(I)*ZI3(I)*((ST23(I)*TR13(I)+ST13(I)*TR23(I))/RL(I)
RM33(I)=(RM13+RM23)/(ELENGO(I)**2)

```

C

```

FL(1,I)=EMOD(I)*AREA(I)*(TUL(I)-TCB2(I)-TCB3(I))
QL(I)=FL(1,I)/ELENGO(I)

```

C  
C  
C

OR COMPUTE THE LOCAL ELEMENT FORCES DUE TO THE RELATIVE DEFORMATIONS BY FINITE ELEMENT MODEL.

ELSE DO

U(I)=RL(I)-ELENG(I)

TU(I)=TU(I)+U(I) ← TOTAL Axial Defor.

```

E1=TU(I)
E12=TR12(I)
E22=TR22(I)
E13=TR13(I)
E23=TR23(I)
A=AREA(I)
E=EMOD(I)
EL=RL(I)
Z2=ZI2(I)
Z3=ZI3(I)

```

C  
C  
C

38

PRINT 38, I, E1, E12, E22

FORMAT(2X, I3, 3X, 3(2X, D15.7))

```

*      RM12=E*A*(4. DO*E12-E22)*E1/60. DO+2. DO*E*Z2*(2. DO*E12+E22)/
*      EL+E*A*E1*E12/15. DO+E*A*EL*((12. DO*E12**2-3. DO*E12*E22
*      +E22**2)*E12/420. DO+(-3. DO*E12**2+4. DO*E12*E22-3. DO*
*      E22**2)*E22/840. DO)-E*A*E1*E22/60. DO
*      RM22=E*A*(4. DO*E22-E12)*E1/60. DO+2. DO*E*Z2*(2. DO*E22+E12)/
*      EL-E*A*E1*E12/60. DO+E*A*EL*((12. DO*E22**2-3. DO*E12*E22
*      +E12**2)*E22/420. DO+(-3. DO*E12**2+4. DO*E12*E22-3. DO*
*      E22**2)*E12/840. DO)+E*A*E1*E22/15. DO
*      RM13=E*A*(4. DO*E13-E23)*E1/60. DO+2. DO*E*Z3*(2. DO*E13+E23)/
*      EL+E*A*E1*E13/15. DO+E*A*EL*((12. DO*E13**2-3. DO*E13*E23
*      +E13**2)*E13/420. DO+(-3. DO*E13**2+4. DO*E13*E23-3. DO*
*      E23**2)*E23/840. DO)-E*A*E1*E23/60. DO
*      RM23=E*A*(4. DO*E23-E13)*E1/60. DO+2. DO*E*Z3*(2. DO*E23+E13)/
*      EL-E*A*E1*E13/60. DO+E*A*EL*((12. DO*E23**2-3. DO*E13*E23
*      +E13**2)*E23/420. DO+(-3. DO*E13**2+4. DO*E13*E23-3. DO*

```

(see 14 15)

```

*      E23**2)*E13/840. DO)+E*A*E1*E23/15. DO
FL(1,I)=(-E*A)*(E1/EL+(4. DO*E13-E23)*E13/60. DO
*      +(4. DO*E23-E13)*E23/60. DO+(4. DO*E12-E22)*E12/60. DO
*      +(4. DO*E22-E12)*E22/60. DO)
QL(I)=-FL(1,I)
END IF

```

C  
C

```

COMPUTE THE REMAINING LOCAL ELEMENT FORCES.

```

```

FL(2,I)=(RM13+RM23)/RL(I)
FL(3,I)=- (RM12+RM22)/RL(I)
FL(4,I)=-GMOD(I)*ZJ(I)*TTORS(I)/RL(I)
FL(5,I)=RM12
FL(6,I)=RM13
FL(7,I)=-FL(1,I)
FL(8,I)=-FL(2,I)
FL(9,I)=-FL(3,I)
FL(10,I)=-FL(4,I)
FL(11,I)=RM22
FL(12,I)=RM23

```

C

```

IF (NPRINT .EQ. 3) THEN DO

```

```

PRINT 40,I,RL(I),RM12,RM22,RM13,RM23,ST12(I),ST22(I),
*      ST13(I),ST23(I),TCB2(I),TCB3(I),TTORS(I),TR12(I),
*      TR22(I),TR13(I),TR23(I),RL1PI(I),RL2PI(I),
*      RL3PI(I),RL1(I),RL2(I),RL3(I),TU(I),ELENG(I)
40  FORMAT(/' ELEMENT',I3,' RL=',D15.7/
*      ' RM12=',D15.7,' RM22=',D15.7/
*      ' RM13=',D15.7,' RM23=',D15.7/
*      ' ST12=',D15.7,' ST22=',D15.7/
*      ' ST13=',D15.7,' ST23=',D15.7/
*      ' TCB2=',D15.7,' TCB3=',D15.7,' TTORS=',D15.7/
*      ' TR12=',D15.7,' TR22=',D15.7/
*      ' TR13=',D15.7,' TR23=',D15.7/
*      ' RL1PI=',D15.7,' RL2PI=',D15.7,' RL3PI=',D15.7/
*      ' RL1=',D15.7,' RL2=',D15.7,' RL3=',D15.7/
*      ' TU=',D15.7,' ELENG=',D15.7/)
PRINT 50
50  FORMAT(' LOCAL ELEMENT FORCES')
PRINT 60,FL(1,I),FL(2,I),FL(3,I),FL(4,I),FL(5,I),FL(6,I),
*      FL(7,I),FL(8,I),FL(9,I),FL(10,I),FL(11,I),FL(12,I)
60  FORMAT(5X,D15.7,5X,D15.7,5X,D15.7/5X,D15.7,5X,D15.7,5X,
*      D15.7/5X,D15.7,5X,D15.7,5X,D15.7/5X,D15.7,5X,D15.7,
*      5X,D15.7/)
END IF

```

C  
C

```

TRANSFORM TO THE GLOBAL ELEMENT FORCES.

```

```

FG(1,I)=C1(I)*FL(1,I)+OLAMD(2,1,I)*FL(2,I)
*      +OLAMD(3,1,I)*FL(3,I)
FG(2,I)=C2(I)*FL(1,I)+OLAMD(2,2,I)*FL(2,I)
*      +OLAMD(3,2,I)*FL(3,I)
FG(3,I)=C3(I)*FL(1,I)+OLAMD(2,3,I)*FL(2,I)
*      +OLAMD(3,3,I)*FL(3,I)
FG(4,I)=C1(I)*FL(4,I)+OLAMD(2,1,I)*FL(5,I)
*      +OLAMD(3,1,I)*FL(6,I)
FG(5,I)=C2(I)*FL(4,I)+OLAMD(2,2,I)*FL(5,I)
*      +OLAMD(3,2,I)*FL(6,I)
FG(6,I)=C3(I)*FL(4,I)+OLAMD(2,3,I)*FL(5,I)
*      +OLAMD(3,3,I)*FL(6,I)
FG(7,I)=C1(I)*FL(7,I)+OLAMD(2,1,I)*FL(8,I)
*      +OLAMD(3,1,I)*FL(9,I)
FG(8,I)=C2(I)*FL(7,I)+OLAMD(2,2,I)*FL(8,I)
*      +OLAMD(3,2,I)*FL(9,I)
FG(9,I)=C3(I)*FL(7,I)+OLAMD(2,3,I)*FL(8,I)
*      +OLAMD(3,3,I)*FL(9,I)
FG(10,I)=C1(I)*FL(10,I)+OLAMD(2,1,I)*FL(11,I)
*      +OLAMD(3,1,I)*FL(12,I)

```

```

      FG(11,I)=C2(I)*FL(10,I)+OLAMD(2,2,I)*FL(11,I)
*          +OLAMD(3,2,I)*FL(12,I)
      FG(12,I)=C3(I)*FL(10,I)+OLAMD(2,3,I)*FL(11,I)
*          +OLAMD(3,3,I)*FL(12,I)
C
      DO 80 J=1,12
        K=MCODE(J,I)
        IF (K.NE. 0) THEN DO
          F(K)=F(K)+FG(J,I)
          IF (NPRINT.EQ. 3) THEN DO
            PRINT 70,F(K)
70          FORMAT('  F(K)= ',D15.7)
          END IF
        END IF
      80 CONTINUE
      RETURN
      END
C
C
C
C *****
*          BOWCOR          *
C *****
SUBROUTINE BOWCOR(AREA,BP12,BP22,BP13,BP23,ELENG,R12,R22,
*          R13,R23,SP12,SP22,SP13,SP23,ST12,ST22,
*          ST13,ST23,ZI2,ZI3,I,IMAX,TUL,EMOD,NPRINT,
*          TCB2,TCB3,ELENGO,DQQR,FQR)
C  CALCULATE THE STABILITY FUNCTIONS,ST12,ST22,ST13,ST23; THE
C  BOWING FUNCTIONS,BW12,BW22,BW13,BW23; THE LENGTH CORRECTION
C  FACTORS FOR THE BOWING ACTIONS,TCB2,TCB3; AND THE AXIAL
C  FORCE, QPRINT.
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION AREA(1),BP12(1),BP22(1),BP13(1),BP23(1),ELENG(1),
*          R12(1),R22(1),R13(1),R23(1),SP12(1),SP22(1),SP13(1),
*          SP23(1),ST12(1),ST22(1),ST13(1),ST23(1),ZI2(1),
*          ZI3(1),EMOD(1),TUL(1),TCB2(1),TCB3(1),ELENGO(1)
      PI=3.141592653589793
      QR=TUL(I)*AREA(I)*ELENGO(I)**2/(ZI2(I)*PI**2)
C      QR=0. D0
C      IC=0
C
      WHILE (IC.LT.IMAX) DO
        IF (QR.GT. 0. D00) THEN DO
          W2=QR*PI**2
          W2=DSQRT(W2)
          ST12(I)=W2*(DSIN(W2)-W2*DCOS(W2))/(2. D00-2. D00*DCOS(W2)
*          -W2*DSIN(W2))
          ST22(I)=W2*(W2-DSIN(W2))/(2. D00-2. D00*DCOS(W2)-
*          W2*DSIN(W2))
C
          W3=QR*PI**2*ZI2(I)/ZI3(I)
          W3=DSQRT(W3)
          ST13(I)=W3*(DSIN(W3)-W3*DCOS(W3))/(2. D00-2. D00*DCOS(W3)
*          -W3*DSIN(W3))
          ST23(I)=W3*(W3-DSIN(W3))/(2. D00-2. D00*DCOS(W3)-
*          W3*DSIN(W3))
C
          END IF
C
          IF (QR.EQ. 0. D00) THEN DO
            ST12(I)=4. D00
            ST22(I)=2. D00
            SP12(I)=0. D00
            SP22(I)=0. D00
            BP12(I)=0. D00
            BP22(I)=0. D00
            TCB2(I)=TUL(I)/2. D0

```

```

C      ST13(I)=4. D00
      ST23(I)=2. D00
      SP13(I)=0. D00
      SP23(I)=0. D00
      BP13(I)=0. D00
      BP23(I)=0. D00
      TCB3(I)=TUL(I)/2. D0
      RETURN
END IF
C      IF (QR .LT. 0. D00) THEN DO
      W2=-QR*PI**2
      W2=DSQRT(W2)
      ST12(I)=W2*(W2*DCOSH(W2)-DSINH(W2))/
*      (2. D0-2. D0*DCOSH(W2)+W2*DSINH(W2))
      ST22(I)=W2*(DSINH(W2)-W2)/(2. D00-2. D00*DCOSH(W2)
*      +W2*DSINH(W2))
C      W3=-QR*PI**2*ZI2(I)/ZI3(I)
      W3=DSQRT(W3)
      ST13(I)=W3*(W3*DCOSH(W3)-DSINH(W3))/
*      (2. D0-2. D0*DCOSH(W3)+W3*DSINH(W3))
      ST23(I)=W3*(DSINH(W3)-W3)/(2. D00-2. D00*DCOSH(W3)
*      +W3*DSINH(W3))
END IF
C      BW12=(ST12(I)+ST22(I))*(ST22(I)-2. D00)/(8. D00*QR*PI**2)
      BW22=ST22(I)/(8. D00*(ST12(I)+ST22(I)))
      SP12(I)=-2. D00*PI**2*(BW12+BW22)
      SP22(I)=-2. D00*PI**2*(BW12-BW22)
      BP12(I)=(QR*(ST12(I)*SP22(I)+ST22(I)*SP12(I)+2. D0*ST22(I)*
*      SP22(I)-2. D0*SP12(I)-2. D0*SP22(I))-(ST12(I)
*      +ST22(I))*(ST22(I)-2. D0)/(8. D0*PI**2*QR**2)
      BP22(I)=(SP22(I)*(ST12(I)+ST22(I))-ST22(I)*(SP12(I)
*      +SP22(I)))/(8. D0*(ST12(I)+ST22(I))**2)
      TCB2(I)=BW12*(R12(I)+R22(I))**2+BW22*(R12(I)-R22(I))**2
      TCB2P=BP12(I)*(R12(I)+R22(I))**2+BP22(I)*(R12(I)-R22(I))**2
C      BW13=(ST13(I)+ST23(I))*(ST23(I)-2. D00)*ZI3(I)/
*      (8. D0*QR*PI**2*ZI2(I))
      BW23=ST23(I)/(8. D00*(ST13(I)+ST23(I)))
      SP13(I)=-2. D00*PI**2*(BW13+BW23)
      SP23(I)=-2. D00*PI**2*(BW13-BW23)
      BP13(I)=(QR*ZI2(I)/ZI3(I)*(ST13(I)*SP23(I)+ST23(I)*SP13(I)
*      +2. D0*ST23(I)*SP23(I)-2. D0*SP13(I)-2. D0*SP23(I))-
*      (ST13(I)+ST23(I))*(ST23(I)-2. D0)/
*      (8. D0*PI**2*(QR*ZI2(I)/ZI3(I))**2)
      BP23(I)=(SP23(I)*(ST13(I)+ST23(I))-ST23(I)*(SP13(I)
*      +SP23(I)))/(8. D0*(ST13(I)+ST23(I))**2)
      TCB3(I)=BW13*(R13(I)+R23(I))**2+BW23*(R13(I)-R23(I))**2
      TCB3P=BP13(I)*(R13(I)+R23(I))**2+BP23(I)*(R13(I)-R23(I))**2
C      RLA=ELENGO(I)/(DSQRT(ZI2(I)/AREA(I)))
      FQRC=DABS(QR-RLA**2*(TUL(I)-TCB2(I)-TCB3(I)))/PI**2)
      DQ=-((QR-RLA**2*(TUL(I)-TCB2(I)-TCB3(I)))/PI**2)/
*      (1. D0+RLA**2*(TCB2P+TCB3P*ZI2(I)/ZI3(I)))/PI**2)
C      IF (NPRINT .EQ. 3) THEN DO
      PRINT 100, ST12(I), ST22(I), BW12, BW22, SP12(I), SP22(I),
*      BP12(I), BP22(I), TCB2(I), ST13(I), ST23(I),
*      BW13, BW23, SP13(I), SP23(I), BP13(I), BP23(I),
*      TCB3(I)
100     FORMAT(' ST12=', D15. 7, ' ST22=', D15. 7/
*      ' BW12=', D15. 7, ' BW22=', D15. 7/
*      ' SP12=', D15. 7, ' SP22=', D15. 7/

```

```

*          :   BP12=' ,D15.7,'   BP22=' ,D15.7,'   TCB2=' ,D15.7//
*          :   ST13=' ,D15.7,'   ST23=' ,D15.7//
*          :   BW13=' ,D15.7,'   BW23=' ,D15.7//
*          :   SP13=' ,D15.7,'   SP23=' ,D15.7//
*          :   BP13=' ,D15.7,'   BP23=' ,D15.7,'   TCB3=' ,D15.7//
END IF
C
QR=QR+DQ
DQQR=DABS(DQ/QR)
IC=IC+1
QPRINT=EMOD(I)*AREA(I)*(TUL(I)-TCB2(I)-TCB3(I))
IF (NPRINT .EQ. 3) THEN DO
  PRINT 200,IC,QPRINT
200  FORMAT(' ITERATION ',I4,' Q= ',D15.7)
END IF
IF (DQQR .LE. DQQR .AND. FQRC .LE. FQR) THEN DO
  IF (NPRINT .EQ. 3) THEN DO
    PRINT 300,IC,I,FQRC,DQQR
300  FORMAT(' BOWCOR CONVERGED IN ',I5,' ITERATIONS FOR
*      ELEMENT',I5,' FQRC=',D15.7,' DQQR=',D15.7)
    END IF
    IC=IMAX+1
  END IF
END WHILE
C
IF (DQQR .GT. DQQR .OR. FQRC .GT. FQR) THEN DO
C
C 400  PRINT 400,I
C      FORMAT(' NO CONVERGENCE IN BOWCOR FOR ELEMENT ',I3)
C      STOP
C
END IF
RETURN
END

C
C
C
C
*****
*          SOLVE          *
*****
SUBROUTINE SOLVE(DD,F,MAXA,QT,SKT,NC,NEQ,NEGPIV,NKT,NPRINT)
CALL FACTOR IF THE STIFFNESS MATRIX HAS BEEN UPDATED; CALL
REDUCE AND BACSUB TO SOLVE THE SIMULTANEOUS LINEAR SYSTEM
EQUATIONS FOR THE NODAL DISPLACEMENTS BY GAUSS ELIMINATION.
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION DD(1),F(1),MAXA(1),QT(1),SKT(1)
IF (NPRINT .EQ. 3) THEN DO
100  PRINT 100
      FORMAT(////' SOLVE CALLED')
END IF
C
IF (NC .EQ. 0) THEN DO
CALL FACTOR(MAXA,SKT,NEQ,NEGPIV,NPRINT)
IF (NPRINT .EQ. 3) THEN DO
200  PRINT 200
      FORMAT('/ FACTORED STIFFNESS MATRIX')
300  PRINT 300,(SKT(I),I=1,NKT)
      FORMAT(' ',6(5X,D15.7))
END IF
END IF
C
DO 10 I=1,NEQ
10  DD(I)=QT(I)-F(I)
CALL REDUCE(DD,MAXA,SKT,NEQ)
CALL BACSUB(DD,MAXA,SKT,NEQ)
RETURN
END
C

```

```

C
C
C *****
C *                                  FACTOR                                  *
C *****
C SUBROUTINE FACTOR(MAXA,SKT,NEQ,NEGPIV,NPRINT)
C PERFORM L*D*L(T) FACTORIZATION (DECOMPOSITION) OF STIFFNESS
C MATRIX SKT.
C IMPLICIT REAL*8(A-H,O-Z)
C DIMENSION MAXA(1),SKT(1)
C NEGPIV=0
C IF (NPRINT .EQ. 3) THEN DO
100  PRINT 100
    FORMAT(/'   FACTOR CALLED   DIAGONALS OF FACTORIZATION:')
END IF
DO 80 N=1,NEQ
  KN=MAXA(N)
  KL=KN+1
  KU=MAXA(N+1)-1
  KH=KU-KL
  IF (KH) 70,50,10
10  K=N-KH
  IC=0
  KLT=KU
  DO 40 J=1,KH
    IC=IC+1
    KLT=KLT-1
    KI=MAXA(K)
    ND=MAXA(K+1)-KI-1
    IF (ND) 40,40,20
20  KK=MINO(IC,ND)
    C=0. D00
    DO 30 L=1,KK
30  C=C+SKT(KI+L)*SKT(KLT+L)
    SKT(KLT)=SKT(KLT)-C
40  K=K+1
50  K=N
    B=0. D00
    DO 60 KK=KL,KU
      K=K-1
      KI=MAXA(K)
      C=SKT(KK)/SKT(KI)
      B=B+C*SKT(KK)
60  SKT(KK)=C
    SKT(KN)=SKT(KN)-B
70  IF (SKT(KN) .EQ. 0. D00) THEN DO
    PRINT 200,N,SKT(KN)
200  FORMAT('  STIFFNESS MATRIX IS NOT POSITIVE DEFINITE' /
    *    PIVOT IS ZERO FOR D.O.F. ' ,I4/' OPIVOT = ',
    *    D15. 8)
    STOP
    END IF
    IF (SKT(KN) .LT. 0. D00) THEN DO
      NEGPIV=NEGPIV+1
      IF (NPRINT .EQ. 3) THEN DO
300  PRINT 300,NEGPIV
        FORMAT(/'   NEGATIVE PIVOT ENCOUNTERED IN FACTOR',I5)
      END IF
    END IF
    IF (NPRINT .EQ. 3) THEN DO
400  PRINT 400,SKT(KN)
        FORMAT(/'   SKT(KN)= ' ,D15. 7)
    END IF
80  CONTINUE
  RETURN
END

```

C  
C  
C  
C  
C  
C

```

*****
*                                REDUCE                                *
*****
SUBROUTINE REDUCE(DD,MAXA,SKT,NEQ)
REDUCE THE RIGHT HAND SIDE LOAD VECTOR.
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION DD(1),MAXA(1),SKT(1)
DO 20 N=1,NEQ
  KL=MAXA(N)+1
  KU=MAXA(N+1)-1
  KH=KU-KL
  IF (KH .GE. 0) THEN DO
    K=N
    C=0. D00
    DO 10 KK=KL,KU
      K=K-1
      C=C+SKT(KK)*DD(K)
10    DD(N)=DD(N)-C
  END IF
20 CONTINUE
RETURN
END

```

C  
C  
C  
C  
C  
C

```

*****
*                                BACSUB                                *
*****
SUBROUTINE BACSUB(DD,MAXA,SKT,NEQ)
BACK SUBSTITUTE TO OBTAIN THE SOLUTION.
IMPLICIT REAL*8(A-H,O-Z)
DIMENSION DD(1),MAXA(1),SKT(1)
DO 10 N=1,NEQ
  K=MAXA(N)
  DD(N)=DD(N)/SKT(K)
10 CONTINUE
IF (NEQ .EQ. 1) RETURN
N=NEQ
DO 30 L=2,NEQ
  KL=MAXA(N)+1
  KU=MAXA(N+1)-1
  KH=KU-KL
  IF (KH .GE. 0) THEN DO
    K=N
    DO 20 KK=KL,KU
      K=K-1
      DD(K)=DD(K)-SKT(KK)*DD(N)
20    CONTINUE
  END IF
  N=N-1
30 CONTINUE
RETURN
END

```

C  
C  
C  
C  
C  
C

```

*****
*                                TEST                                *
*****
SUBROUTINE TEST(AREA,BP12,BP22,BP13,BP23,C1,C2,C3,D,DD,DD1,DE,
*              ELENG,EMOD,GMOD,F,FG,FL,FP,FPI,JCÓDE,MCODE,
*              QT,R12,R22,R13,R23,TORS,SP12,SP22,SP13,SP23,
*              ST12,ST22,ST13,ST23,ZI2,ZI3,ZJ,CPDB,CPDC,

```



```

*          CPF,ICI,IMAX,NE,NEQ,NJ,Q,QI,CI1,CI2,CI3,
*          ITIND,NPRINT)
C      AT THE END OF EACH ITERATION, CHECK TO SEE IF HAS CONVERGED
C      TO AN EQUILIBRIUM POINT WITHIN PRESET TOLERANCES OR WHETHER
C      THE ITERATION IS DIVERGING.
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION AREA(1),BP12(1),BP22(1),BP13(1),BP23(1),C1(1),C2(1),
*          C3(1),D(1),DD(1),DD1(1),DE(1),ELENG(1),EMOD(1),
*          GMOD(1),F(1),FG(12,1),FL(12,1),FP(1),FPI(1),
*          JCODE(6,1),MCODE(12,1),QT(1),R12(1),R22(1),R13(1),
*          R23(1),TORS(1),SP12(1),SP22(1),SP13(1),SP23(1),
*          ST12(1),ST22(1),ST13(1),ST23(1),ZI2(1),ZI3(1),ZJ(1),
*          Q(1),CI1(1),CI2(1),CI3(1)
      ICI=0
C      IF (CPDC .LT. 1.D00) THEN DO
      CALL DISPLC(D,DD,JCODE,CPDC,ICI,NJ)
      END IF
C      IF (CPDB .LT. 1.D00) THEN DO
      CALL DISPLB(D,DD,CPDB,ICI,NEQ)
      END IF
C      IF (CPF.LT. 1.D00) THEN DO
      CALL UNBALF(F,FP,QT,CPF,ICI,NEQ)
      END IF
C      IF (NPRINT .EQ. 3) THEN DO
      PRINT 100,ICI
100    FORMAT(// ' ICI= ',I5)
      END IF
      RETURN
      END
C
C
C
C
*****
*          DISPLC          *
*****
C      SUBROUTINE DISPLC(D,DD,JCODE,CPDC,ICI,NJ)
C      CHECK TO SEE IF THE RATIO: THE MAX CHANGE IN TRANSLATION (OR
C      ROTATION) TO THE LARGEST TOTAL NODAL DISPLACEMENT (OR
C      ROTATION) IN STRUCTURE, IS WITHIN A DISPLACEMENT
C      CONVERGENCE TOLERANCE, CPDC.
      IMPLICIT REAL*8(A-H,O-Z)
      DIMENSION D(1),DD(1),JCODE(6,1)
      RMT=0.D00
      RMR=0.D00
      CR=0.D00
      CT=0.D00
      I=1
C      WHILE (I .LE. 6) DO
      J=1
      WHILE (J .LE. NJ) DO
      K=JCODE(I,J)
      IF (K .NE. 0) THEN DO
      IF (I .GE. 4) THEN DO
      RMR=DMAX1(DABS(D(K)),RMR)
      CR=DMAX1(DABS(DD(K)),CR)
      ELSE DO
      RMT=DMAX1(DABS(D(K)),RMT)
      CT=DMAX1(DABS(DD(K)),CT)
      END IF
      END IF
      J=J+1

```

```

      END WHILE
      I=I+1
END WHILE
C
IF (RMR .NE. 0.D00) THEN DO
  CPR=CR/RMR
  IF (CPR .GT. CPDC) THEN DO
    ICI=ICI+1
  END IF
END IF
IF (RMT .NE. 0.D00) THEN DO
  CPT=CT/RMT
  IF (CPT .GT. CPDC) THEN DO
    ICI=ICI+1
  END IF
END IF
C
IF (RMR .EQ. 0.D00 .AND. RMT .EQ. 0.D00) THEN DO
  ICI=ICI+1
100 PRINT 100,CR,RMR,CT,CMT
*   FORMAT(/' MAX D IN DISPLC = 0 CR=' ,D15.7,' RMR=' ,D15.7,
      CT=' ,D15.7,' CMT=' ,D15.7)
      STOP
      END IF
      RETURN
      END
C
C
C
C *****
C *                               DISPLB                               *
C *****
C SUBROUTINE DISPLB(D,DD,CPDB,ICI,NEQ)
C CHECK TO SEE IF THE RATIO: THE DISPLACEMENT INCREMENT DURING
C EACH ITERATION TO THE DISPLACEMENT AT THE END OF EACH
C ITERATION, IS WITHIN A DISPLACEMENT CONVERGENCE
C TOLERANCE, CPDB.
C IMPLICIT REAL*8(A-H,O-Z)
C DIMENSION D(1),DD(1)
C CN=0.D00
C CD=0.D00
C DO 10 I=1,NEQ
C   CN=CN+(DD(I))**2
C   CD=CD+(D(I))**2
10 CONTINUE
IF (CD .EQ. 0.D00) THEN DO
  ICI=ICI+10
  RETURN
END IF
C=(DSQRT(CN))/(DSQRT(CD))
IF (C .GT. CPDB) THEN DO
  ICI=ICI+10
END IF
RETURN
END
C
C
C
C *****
C *                               UNBALF                               *
C *****
C SUBROUTINE UNBALF(F,FP,QT,CPF,ICI,NEQ)
C CHECK TO SEE IF THE RATIO: THE UNBALANCED FORCE DURING EACH
C ITERATION TO THE ORIGINAL LOAD INCREMENT, IS WITHIN A PRESET
C TOLERANCE, CPF.
C IMPLICIT REAL*8(A-H,O-Z)

```

```

DIMENSION F(1),FP(1),QT(1)
CN=0. D00
CD=0. D00
DO 10 I=1,NEQ
  CN=CN+(QT(I)-F(I))**2
  CD=CD+(QT(I)-FP(I))**2
10 CONTINUE
IF (CD .EQ. 0. D00) THEN DO
  ICI=ICI+100
  RETURN
END IF
C=(DSQRT(CN))/(DSQRT(CD))
IF (C .GT. CPF) THEN DO
  ICI=ICI+100
END IF
RETURN
END

```

C  
C  
C  
C  
C  
C

```

*****
*                               RESULT                               *
*****
SUBROUTINE RESULT(C1,C2,C3,D,DJ,ELENG,EMOD,GMOD,FG,FL,
*                   JCODE,MCODE,MINC,P,Q,QT,ZI2,ZI3,
*                   ZJ,NE,NEQ,NJ,QI,IMP,NPRINT)
C   INITIALIZE THE NODAL FORCE MATRIX ,P, TO ZERO; CALL JOINTF
C   AND OUTPUT.
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION C1(1),C2(1),C3(1),D(1),DJ(6,1),ELENG(1),EMOD(1),
*           GMOD(1),FG(12,1),FL(12,1),JCODE(6,1),MCODE(12,1),
*           MINC(2,1),P(6,1),Q(1),QT(1),ZI2(1),ZI3(1),ZJ(1)
  DO 20 J=1,NJ
    DO 10 I=1,6
      P(I,J)=0. D00
10  CONTINUE
20  CONTINUE
  CALL JOINTF(FG,MINC,P,NE)
  CALL OUTPUT(D,DJ,FL,JCODE,P,QT,NE,NEQ,NJ,QI,IMP,NPRINT)
  RETURN
END

```

C  
C  
C  
C  
C  
C

```

*****
*                               JOINTF                               *
*****
SUBROUTINE JOINTF(FG,MINC,P,NE)
C   ASSIGN THE GLOBAL ELEMENT FORCES TO THE NODAL FORCE MATRIX,P.
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION FG(12,1),MINC(2,1),P(6,1)
  DO 20 I=1,NE
    J=MINC(1,I)
    K=MINC(2,I)
    DO 10 L=1,6
      P(L,J)=P(L,J)+FG(L,I)
      P(L,K)=P(L,K)+FG(L+6,I)
10  CONTINUE
20  CONTINUE
  RETURN
END

```

C  
C  
C  
C  
C

```

*****
*                               OUTPUT                               *
*****

```

```

C *****
C SUBROUTINE OUTPUT(D,DJ,FL,JCODE,P,QT,NE,NEQ,NJ,QI,IMP,NPRINT)
C PRINT THE NODAL DISPLACEMENTS (INCLUDING NODAL CONSTRAINTS),
C DJ; PRINT THE LOCAL ELEMENT FORCES,FL; PRINT THE NODAL FORCES,
C P; PRINT THE LOAD VERSUS DISPLACEMENT VECTOR,QT AND D.
  IMPLICIT REAL*8(A-H,O-Z)
  DIMENSION D(1),DJ(6,1),FL(12,1),JCODE(6,1),P(6,1),QT(1)
  DO 20 J=1,NJ
    DO 10 I=1,6
      DJ(I,J)=0.D00
      K=JCODE(I,J)
      IF (K.NE.0) THEN DO
        DJ(I,J)=D(K)
      END IF
    10 CONTINUE
  20 CONTINUE
C
  IF (NPRINT.LE.1) THEN DO
    PRINT 100,QI,D(IMP)
  100 FORMAT(' ',F17.9/ ' ',F17.9)
  END IF
C
  IF (NPRINT.EQ.2.OR.NPRINT.EQ.3) THEN DO
    PRINT 200,QI,D(IMP)
  200 FORMAT(' QI EQUILIB =',F17.9,' D EQUILIB =',F17.9)
  END IF
C
  IF (NPRINT.EQ.4) THEN DO
    PRINT 300,QI
  300 FORMAT(' ',80(' '))//31X,'QI=',F15.7//34X,'LOAD VECTOR'
    * /34X,11('*')
    DO 30 I=1,NEQ
      PRINT 400,QT(I)
    400 FORMAT(' ',31X,F16.5)
    30 CONTINUE
    PRINT 500
    500 FORMAT(' ',26X,'GLOBAL NODAL DISPLACEMENTS'/27X,26('*')
      * //10X,'NODE',7X,'1-DIRECTION',8X,'2-DIRECTION',
      * 8X,'3-DIRECTION',8X,'4-DIRECTION',8X,'5-DIRECTION',
      * 8X,'6-DIRECTION')
    DO 40 J=1,NJ
      PRINT 600,J,(DJ(I,J),I=1,6)
    600 FORMAT(' ',8X,I4,6(2X,F17.9))
    40 CONTINUE
    PRINT 700
    700 FORMAT(' ',29X,'LOCAL ELEMENT FORCES'/30X,20('*')/
      * ÉLE.,30X,'A-END',60X,'B-END'/ NUM.',3X,'1',
      * 9X,'2',9X,'3',9X,'4',9X,'5',9X,'6',9X,'7',9X,'8',9X,
      * '9',9X,'10',9X,'11',9X,'12')
    DO 50 J=1,NE
      PRINT 800,J,(FL(I,J),I=1,12)
    800 FORMAT(' ',12,12(1X,F9.4))
    50 CONTINUE
    PRINT 900
    900 FORMAT(' ',33X,'NODAL FORCES'/34X,12('*')/10X,'NODE',
      * 7X,'1-DIRECTION',8X,'2-DIRECTION',8X,'3-DIRECTION',
      * 8X,'4-DIRECTION',8X,'5-DIRECTION',8X,'6-DIRECTION')
    DO 60 J=1,NJ
      PRINT 600,J,(P(I,J),I=1,6)
    60 CONTINUE
    PRINT 1000
  1000 FORMAT(' ',80(' '))
  END IF
  RETURN
  END
C

```

```

C
C
C *****
C *                               DOTPRD                               *
C *****
C FUNCTION DOTPRD(DOT1,DOT2,N)
C COMPUTE THE DOT PRODUCT OF DOT1 AND DOT2.
C IMPLICIT REAL*8(A-H,O-Z)
C DIMENSION DOT1(1),DOT2(1)
C DOTPRD=0.D00
C DO 10 I=1,N
C   DOTPRD=DOTPRD+DOT1(I)*DOT2(I)
10 CONTINUE
C RETURN
C END

C
C
C *****
C *                               UPDATE                               *
C *****
C SUBROUTINE UPDATE(QT,Q,NEQ,QI)
C IMPLICIT REAL*8(A-H,O-Z)
C DIMENSION QT(1),Q(1)
C DO 10 I=1,NEQ
C   QT(I)=Q(I)*QI
10 CONTINUE
C RETURN
C END

C
C //DATA

```

**The vita has been removed from  
the scanned document**