

THE REACTION OF ORTHO-POSITRONIUM
WITH NITROAROMATICS
VIA COMPLEX FORMATION

by

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INTRODUCTION

The chemical consequences of a nuclear transformation such as electronic excitation, ionization, and bond cleavage allow one to study the effect that such an event has on a particular system. This area is generally referred to as nuclear or radiation chemistry. However, there are certain nuclear transformations that are directly influenced by the physical and chemical nature of their environment. The two most notable of these are the Mössbauer effect and positron annihilation. The study of these types of phenomena cannot be classified as the chemical effect of a nuclear transformation, but rather as the chemical effect on it. The physical and chemical properties of a system alter the rate of positron annihilation and allow the positron to be used as a probe of its environment. For example, recent work in this area has used the positron to measure the electron momentum distributions in certain aliphatic liquids. It is also applicable to the solid state where positron lifetimes have been useful in estimating the concentration of lattice defects in semi-conductors.

The positron itself is a stable particle, but life in one's antimatter world is quite short, typically a fraction of a nanosecond in the condensed phase. By forming a bound state with an electron whose spin is triplet relative to

the positrons' (ortho-positronium), the positron may experience a tenfold enhancement in its lifetime.

A rapid reaction in solution has previously been observed between ortho-positronium and certain diamagnetic organic molecules such as nitroaromatics and quinones. They react rapidly with second-order rate constants typically three orders of magnitude greater than most diamagnetic compounds.

The focus of this research is to determine the mechanism by which ortho-positronium reacts with a certain class of compounds, the nitroaromatics, and to evaluate any of the pertinent conditions necessary for this process.

Determining the nature of this process becomes possible by studying the variation in the positron lifetime distribution on dissolution of small amounts of these compounds in suitable solvents. The role that substituents play will be examined from both a steric and reactivity viewpoint. Investigating the temperature dependence of these reactions will also be helpful. Quantum mechanical calculations will be performed to study the stability of positron and positronium complexes with diamagnetic organic molecules.

REVIEW OF BASIC PRINCIPLES

A. The Positron

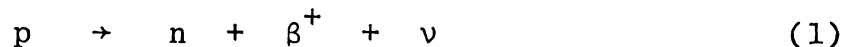
One of the most interesting consequences of the theory of relativity is the division of the universe into matter and anti-matter. Dirac¹ has shown by purely theoretical arguments that the relativistic wave equation for a free electron has both positive and negative energy solutions. The magnitude of these solutions is at least $m_e c^2$, where m_e is the mass of the electron and c is the speed of light. Instead of disregarding the negative solution, he proposed that it corresponded to a particle with a mass equal to that of an electron but carrying the opposite charge -- a positron. Two years later, Anderson² discovered the positron in cloud chamber tracks of cosmic radiation.

The positron states can be intuitively understood in terms of the "theory of holes." Since electrons are observed only in positive energy states, the theory requires that the negative energy states be completely filled with electrons. A photon with energy greater than $2 m_e c^2$ could raise an electron from its negative energy state to a positive one. The "hole" created by this transition in the "sea" of electrons would behave as a positron in that it is defined to have negative energy, negative momentum, and spin opposite to the electron.

There is a second interpretation of the negative energy solution which is equally valid and does not require an unobservable "sea" of electrons. The positron, instead of being a negative energy "hole," can be described as an electron moving backwards in time.³ At the present level of our understanding, the latter interpretation has no physical meaning.

B. Positron Decay

The emission of a positron from a radioactive nucleus is referred to as positron decay. It occurs in nuclei that are proton rich in relation to their most stable isobars. The creation of a positron arises from the transformation of a proton, p , within the nucleus into a neutron, n , a positron, β^+ , and a neutrino, ν , with an accompanying decrease in the atomic number of the parent nucleus by one unit.



The positron cannot exist within a nucleus, because of difficulties arising from the considerations of angular momentum conservation, and is emitted. Once the positron has left the nucleus, it is repelled by the Coulombic interaction between itself and the positively charged nucleus. The release of binding energy during positron

decay along with the Coulomb effect cause the positron to possess a considerable amount of kinetic energy. This may range from a few hundred keV up to several MeV.

The parent and daughter nuclei are in definite energy states for positron decay. Energetically, in order for positron decay to take place, the atomic mass of the parent nucleus, M_p , must be greater than the sum of the rest masses of two beta particles and the atomic mass of the daughter nuclei, M_d .

$$M_p > M_d + 2m_e \quad (2)$$

A neutrino is emitted simultaneously with the positron. Therefore, the energy spectrum of emitted positrons is continuous rather than having a discrete value. The positron energy spectrum⁴ for sodium-22 is shown in Figure 1. The maximum energy, also referred to as the end-point energy, is 0.544 MeV. In general, the maximum number of positrons emitted in the energy distribution occurs at one-third of the end-point energy. The remaining energy is divided between the recoil energy of the daughter nucleus and the neutrino energy.

The elementary theory of beta decay, according to the basic transformations involved, was first given by Fermi.⁵ His descriptive ideas gave satisfactory explanations for the shapes of beta particle energy spectrums, half-lives,

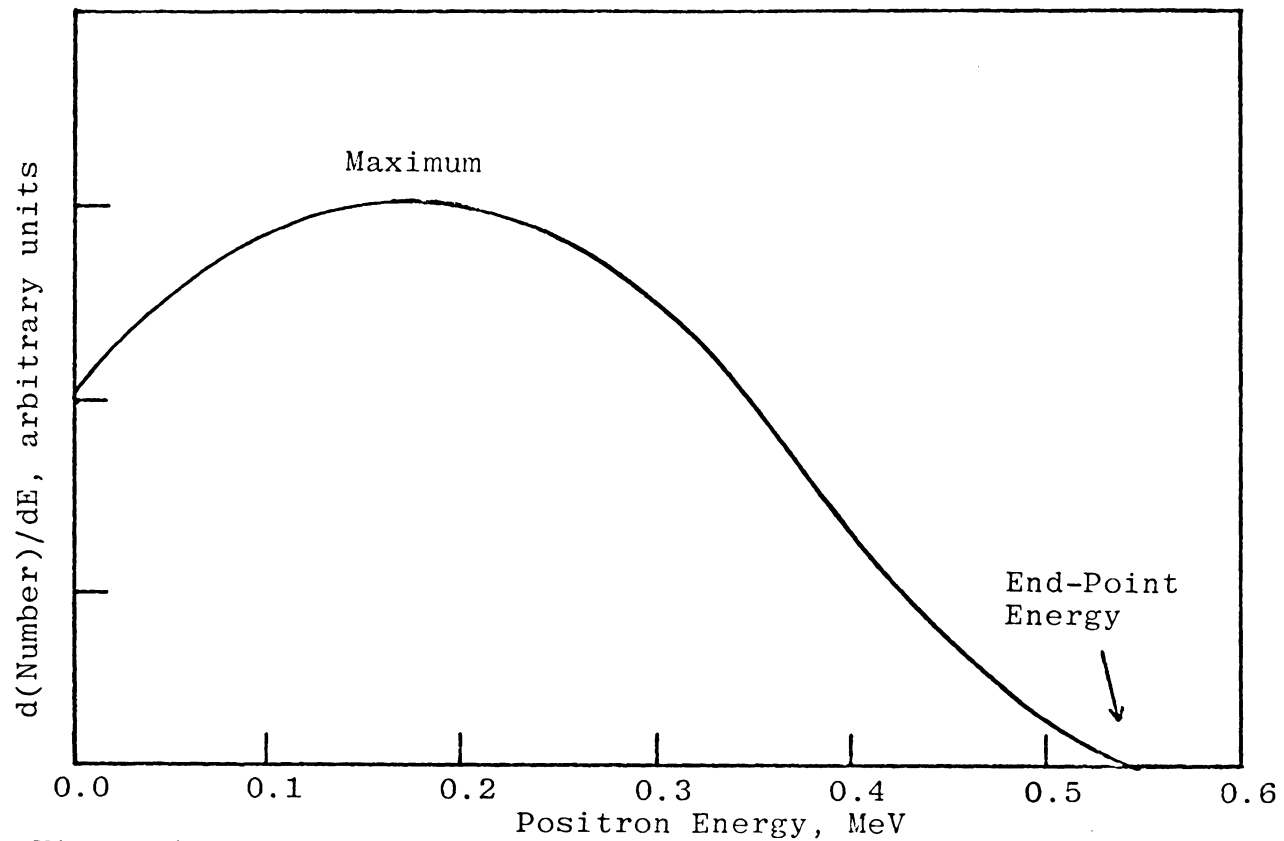


Figure 1. Energy Spectrum of Positrons Emitted from Sodium-22.
(from Ref. 4)

etc. A more modern theory, which takes into account the nonconservation of parity, has been developed by using relativistic quantum mechanics and was introduced by Lee and Yang.⁶

Positrons, like all charged particles, lose their energy while passing through matter. Their interactions can induce ionizations and excitations in the stopping medium. Energy losses due to bremsstrahlung dominate only at higher energies.⁷ Tao and Green have estimated that a high energy positron interacting with a dense material such as water will be slowed down to the first ionization energy of water in 7 psec. The total time needed to thermalize the positron was roughly estimated to be 100 psec. The interactions of positrons with matter are difficult to study at very low energies because of the competing process of mass annihilation whose cross-section is greatest for thermalized positrons.

C. Positron Annihilation

Annihilation is an event unique to matter-anti-matter pairs such as the electron-positron. However, it has been shown by Goldanskii⁸ that fewer than five percent of the positrons emitted into a condensed medium will annihilate during the slowing down process. According to the theory of positron annihilation,⁹ one, two, or three photons may be emitted as the result of a collision between a positron

and an electron. If the two particles meet with their spins antiparallel, two photon emission results. Two photon emission is forbidden for collisions where the two particles meet with their spin parallel and three photons are emitted. Emission of a single photon requires that a third body, M, be present to absorb the recoil momentum of the annihilation photon.



In order to be complete, it should be noted that if two such bodies, M, are present, annihilation may occur without the creation of electromagnetic energy.



The probability of such an occurrence is very small and it has never been observed.

Goldanskii has characterized the three basic mechanisms by which slow-positrons are annihilated.⁸

First, annihilation as "free" positrons in collisions with electrons of their surroundings.

The cross section for the two photon annihilation of a positron and an electron, $\sigma_{2\gamma}$, was first calculated by Dirac.¹ He assumed in the calculation, that the electron of the annihilating pair was at rest. The velocity, v, of thermal positrons is small (6.6×10^6 cm/sec) compared to the speed of light and the equation reduces to

$$\sigma_{2\gamma} = \sigma_S = \pi r_0^2 c/v \quad (4)$$

or

$$\sigma_{2\gamma} v = \pi r_0^2 c = \text{constant} \quad (4a)$$

where $r_0 = 2.8 \times 10^{-13}$ cm is the classical radius of a beta particle.

The two gamma annihilation rate in a multi-electron system is

$$\lambda_{2\gamma} = \lambda_S = \sigma_{2\gamma} v \rho N_0 Z_{\text{EFF}}/M \quad (5)$$

where ρ is the density, N_0 is Avogadro's number, M is the molecular weight, and Z_{EFF} is the effective number of electrons seen by the positron (usually this is taken to be the number of valence electrons). The annihilation rate in water, assuming $Z_{\text{EFF}} = 8$, is calculated at 2.2 nsec^{-1} and corresponds to a free positron lifetime of 0.46 nsec . This value agrees very well with experimental results.¹⁰ Annihilation rates in the gas phase are usually two to three orders of magnitude less, due to lower densities.

The triplet interaction usually results in three photon emission. Ore and Powell¹¹ in 1949 were the first to correctly calculate the ratio of two photon to three photon annihilation rates as 1:1:5. The much lower triplet rate, λ_T , is caused by the multiple character of the three

quantum process. The ratio of the cross sections for these two modes of decay is one-fourth the ratio of the annihilation rates (see p. 13). Summarizing,

$$\lambda_S/\lambda_T = 1115 \quad ; \quad \sigma_S/\sigma_T = 372 \quad (6)$$

A second channel of annihilation involves the binding of a positron to a multi-electron system. This should cause polarization of the electron shells and subsequent formation of positron atomic or molecular orbitals. Much theoretical work in this area is being performed and Chapter 4 of this work deals with this subject.

D. Positronium

The most widely studied interaction of slow positrons with matter is the formation of positronium, Ps, a hydrogen-like bound state of a positron and an electron. This species was proposed by Mohorovicic in 1934¹² to account for the unusual emission spectra of some nebulae. The ionization potential of ortho-positronium was calculated to be 6.77 eV by Ruark,¹³ and he also proposed the name positronium for the e^+e^- atom.

In order to treat positronium as the structural analogue of hydrogen, the reduced mass of hydrogen, $\mu_H \approx m_e$, must be replaced by the reduced mass of the positronium atom, $\mu_{PS} = m_e/2$. The Bohr radius, a_{PS} , is 1.06 \AA , twice

that of hydrogen, and the first ionization potential, binding energy, is half as great in positronium as in hydrogen.

The two ground, $n = 1$, spin states of positronium are the singlet, para-positronium (p-Ps) and the triplet, ortho-positronium (o-Ps). Figure 2 illustrates both forms and lists some of their characteristics. The average lifetime of these two atoms in free space can be calculated from the rates of singlet and triplet annihilation given in Eq. (5) and (6). To do this, the factor, $\rho N_{\text{O}} Z_{\text{EFF}}/M$, must be replaced by the density of the electron wave function at the positron, $|\psi(0)|^2$. Eq. (5) now becomes

$$\lambda_{\text{S}}^{\text{O}} = 4\pi r_{\text{O}}^2 c |\psi(0)|^2 = 1/\tau_{\text{S}}^{\text{O}} \quad (7)$$

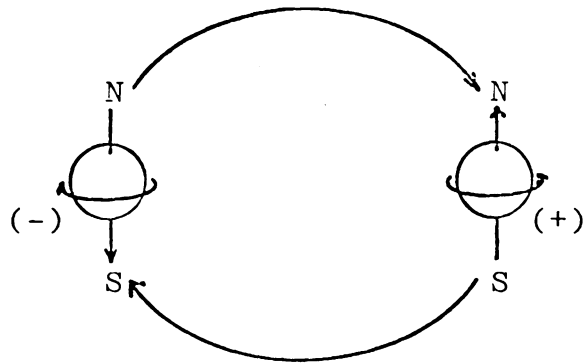
In a free ground-state atom of positronium

$$|\psi(0)|^2 = (1/\pi a_{\text{PS}}^3) = (1/\pi)(\mu_{\text{PS}} e^2/\hbar^2)^3 = 2.68 \times 10^{23} \text{ cm}^{-3}$$

and therefore

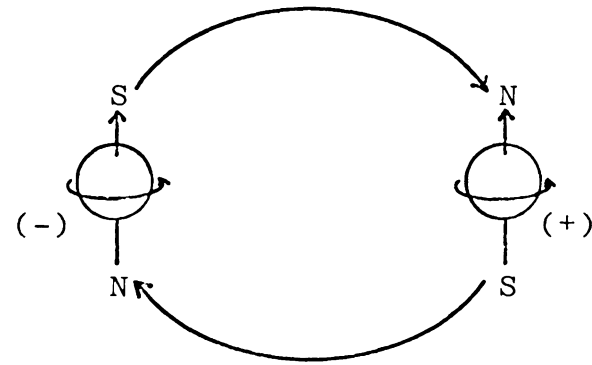
$$\tau_{\text{S}}^{\text{O}} = 0.125 \text{ nsec} \quad (8)$$

for a para-positronium atom in free space. Para-positronium has a shorter lifetime than that of free positrons in most



Para-Positronium

Singlet state 1S
 Spins are antiparallel
 Magnetic poles are parallel
 Disintegrates into 2 photons
 Mean-life = 0.125 nsec



Ortho-Positronium

Triplet State 3S
 Spins are parallel
 Magnetic poles are antiparallel
 Disintegrates into 3 photons
 Mean-life = 140 nsec

Figure 2. Some Physical Properties of Para- and Ortho-Positronium.

substances, the latter has $\tau_F \approx 0.45$ nsec in condensed matter. An ortho-positronium atom in free space has a calculated mean lifetime¹¹

$$\tau_T^O = 140 \text{ nsec} \quad (9)$$

This value is 1115 times longer than the mean lifetime of the singlet atom.

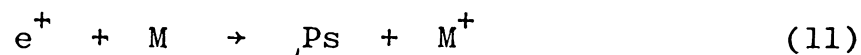
The four ground state wave functions of positronium, $\psi_{J,m}$, are composed of both positron and electron moieties. They are given in Eqs. (10). Para-positronium has a total momentum, $J_S = 0$, and a magnetic quantum number, $m = 0$. There are three substates of ortho-positronium, $J_T = 1$, with $m = +1, 0, -1$.

$$\begin{aligned} \psi_{0,0} &= (1/\sqrt{2})\{\chi_-(1)\chi_+(2) - \chi_-(2)\chi_+(1)\} \quad {}^1S_0 \\ \psi_{1,1} &= \chi_-(1)\chi_+(1) \\ \psi_{1,0} &= (1/\sqrt{2})\{\chi_-(1)\chi_+(2) + \chi_-(2)\chi_+(1)\} \quad {}^3S_1 \\ \psi_{1,-1} &= \chi_-(2)\chi_+(2) \end{aligned} \quad (10)$$

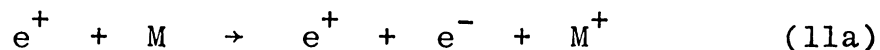
Statistically, ortho-positronium should be formed in 75% of the atoms with para-positronium formation occurring 25% of the time. The factor, $|\psi(0)|^2$, in Eq. (7) can now be more correctly defined to be the density of electrons whose spin is antiparallel to that of the positron. In a multi-electron system, this factor would only be 1/4 of that for

para-positronium, hence the need for the additional factor of 4 in Eq. (7).

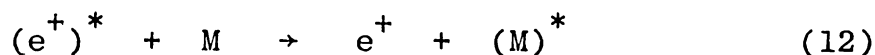
The theory of the formation of positronium in the gas phase, as developed by Ore¹⁴ places certain energy restrictions on this process. The narrow energy region where positronium formation is allowed is called the Ore gap. Formation of positronium can be represented as follows:



A positron collides with a molecule, M, removes an electron from it, and forms positronium. The probability of ionization as described in Eq. (11a) predominates over positronium formation if the positron has kinetic energy greater than the ionization energy of the molecule, I.



This energy level is considered to be the top of the Ore gap, the upper kinetic energy limit at which positronium formation can occur. If the kinetic energy of the positron is greater than the first excitation potential, E^* , of M, electronic excitation of the molecule could occur and would compete with positronium formation.



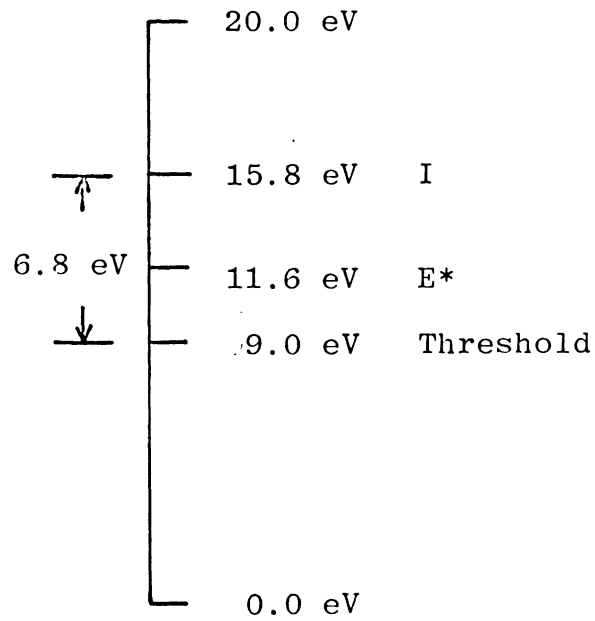
If the ionization potential of M is greater than the binding energy of positronium, 6.8 eV, the reaction is endothermic. Only positrons possessing enough kinetic energy, $E_{\text{Kin}} > I - 6.8$, can form positronium. This makes, $I - 6.8$, the lower positron energy threshold, T, for positronium formations.

The Ore theory is based on three basic assumptions:

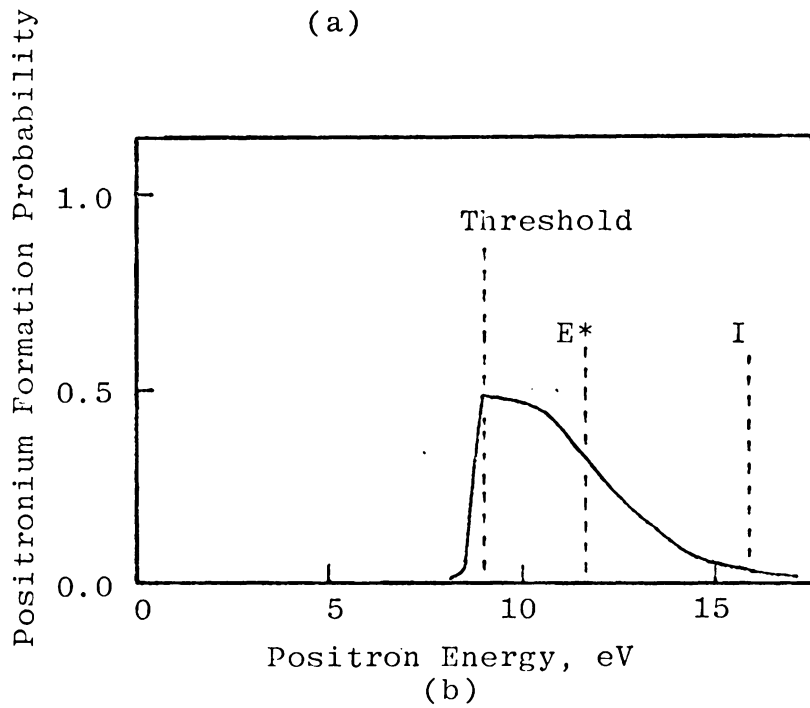
- (1) All positrons reach the top of the Ore gap, I, without being annihilated.
- (2) There is a statistical distribution of positron energies in the range between I and $E = \text{thermal}$.
- (3) All positrons whose kinetic energy lies within the Ore gap will form positronium.

Figure 3a is an energy diagram showing the Ore gap in gaseous argon and Figure 3b is a plot of the hypothetical dependence of the positronium formation probability in argon gas as a function of positron energy.¹⁵ Green and Lee note that the curve is meant to illustrate general features only. The actual formation probability has been measured by Gittleman¹⁶ in a variety of gases and was found to average about 30% of positrons emitted.

Since the binding energy of positronium is less than the ionization potential of most molecules, a considerable number of positrons exist with energies below the Ore gap.



(a)



(b)

Figure 3. (a) Energy Diagram of the Ore Gap and (b) Hypothetical Positronium Formation Probability in Gaseous Argon. (from Ref. 15)

The Ore theory does not place any energy restriction on positron complex formation and positrons may bind to molecules at any energy. Positron complex formation above or within the Ore gap causes inhibition of positronium formation. However, positron complexes are generally thought to form in the energy region beneath the Ore gap due to the lower positron energy and lack of competition with positronium formation.

E. Quenching

Quenching is any means by which the positronium lifetime is shortened from its self-annihilation lifetime as described in the previous section. Quenching occurs in all systems and is most noticeable for the long-lived ortho-positronium. The ortho-positronium lifetime in free space, 140 nsec, may be shortened to a few tenths of a nanosecond in condensed phases. There are three main types of quenching: conversion, pickoff, and chemical reaction.²⁰

Conversion may occur when a positronium atom collides with a molecule containing one or more unpaired electrons. In ortho-positronium, the electron undergoes a spin flip changing it into the shorter-lived para-positronium. Conversion of the singlet atom is also possible but it usually undergoes self-annihilation before it can be converted. The most common and bothersome quencher in liquids is

dissolved oxygen. Care must be taken to remove it from the liquid in order to measure the positronium annihilation rates. The mechanism by which oxygen quenches and the procedure for degassing a liquid are discussed in detail in the chapter on experimental methods. Other molecules that are known to cause conversion quenching are NO, NO₂, and many of the paramagnetic transition metal cations and their complexes.

Pickoff is the annihilation of the bound positron with an electron of a colliding molecule. Again, pickoff is most noticeable for ortho-positronium and occurs when the positron part of the positronium wave function sufficiently overlaps with a singlet electron of the surrounding medium. Hence the positron "pickoff" an electron from the molecule. The ortho-positronium pickoff annihilation rate of 193 organic liquids has been measured by Grey, Sturm, and Cook.²¹ They have found the average annihilation cross section, $\langle\sigma v\rangle$, to correlate directly with the electron polarizability of normal alkanes as calculated from their indices of refraction. Partial quenching cross sections were calculated for 27 structural functional groups (See Table I) and can be used to predict molecular quenching cross sections. The average error of the calculated cross sections for 151 compounds was 1.79%. The cross sections for the remaining 42 compounds could not be calculated because they contained

Table I. Average Quenching Cross Sections.

(From Ref. 21)

	Group	$\langle\sigma v\rangle$, 10^{-14} cm ³ /sec
CH ₃		0.855
CH ₂		0.971
CH		0.924
C		0.705
C ₆ H ₁₁	Monosubstituted Cyclohexanes	5.515
C ₆ H ₁₀	Disubstituted Cyclohexanes	5.373
C ₅ H ₉	Monosubstituted Cyclopentanes	4.469
C ₆ H ₅	Monosubstituted Benzenes	4.679
C ₆ H ₄	Disubstituted Benzenes	4.844
-CH=CH-		1.906
CH ₂ =CH-		1.667
CH ₂ =C(CH ₃)-		2.635
-CH=C(CH ₃)-		3.076
H		0.129
OH		0.749
-O-	Ethers	0.836
CHO	Aldehydes	1.673
COOH	Acids	1.839
C=O	Ketones	1.786
COO-	Esters	1.939
F		0.084
Cl		1.122
Br		1.872
NH ₂		1.624
PO ₃ [≡]	Phosphites	3.382
PO ₄ [≡]	Phosphates	3.754
SH	Thiols	2.330

unusual structures such as keto-enol forms or because the ortho-positronium intensity, I_2 , was so small that accurate lifetime measurements could not be made.

The pickoff process in the condensed phase has been characterized by Brandt¹⁷ using a "free volume model" in which the positronium atom is thought to occupy the interstitial volume between the molecules. The pertinent assumptions are:

- (1) Mutual polarization of the lattice and the positronium is neglected.
- (2) The positronium is considered to be thermalized.
- (3) The lattice potential is rectangular in shape (Figure 4 gives a qualitative description of this arrangement). The potential well has height, U_0 , radius, r_0 corresponding to an excluded volume V_0 and electron density, ρ_0 . A unit cell has radius, $r_1 > r_0$ and cell volume V_1 . The free volume, V^* , is equal to $V_1 - V_0$.

The amount of free volume present in a liquid depends upon the molecular properties of the liquid, the temperature, and the pressure. Henderson and Millett¹⁸ studied the pressure dependence of the ortho-positronium lifetime, τ_2 ,

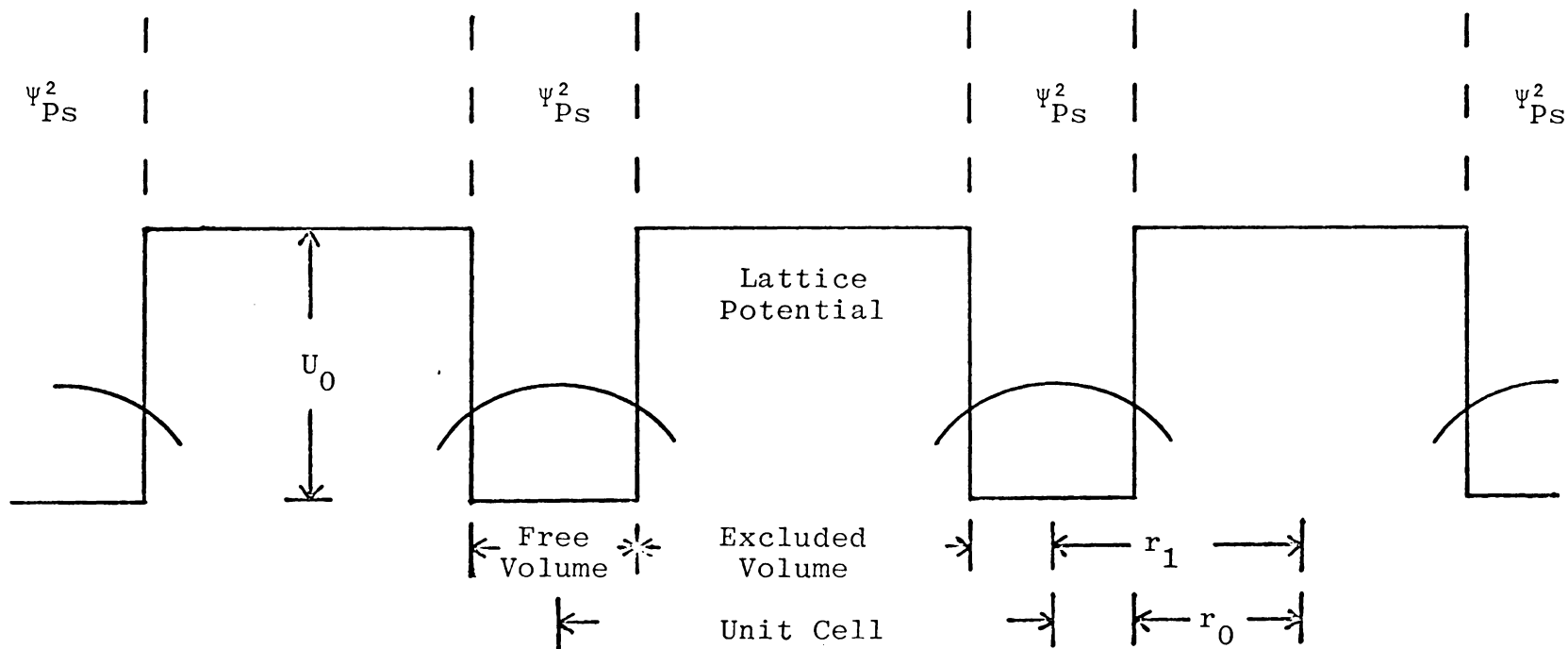


Figure 4. Lattice Potential as Described by the "Free Volume Model".
(from Ref. 17)

in 22 organic liquids and established that it decreased linearly with decreasing specific volume, $V(P)/V_0$, where V_0 is the volume at STP. Wilson, Johnson, and Stump¹⁹ investigated both the temperature and pressure dependence of τ_2 in a series of liquids. They found the lifetimes to follow

$$\tau_2 = \tau_T^0 \exp(-\mu V/V_0) \quad (13)$$

where μ is a characteristic constant for each liquid. The lifetimes of ortho-positronium in liquids that cannot be described by these relationships will be discussed in Chapter 3.

The validity of the "free volume model" is further substantiated by two photon angular correlation experiments. When singlet annihilation occurs, the two photons are emitted in exactly opposite directions, $\theta = \pi$, if the center of mass of the electron-positron pair is at rest. However, the center of mass is in motion for all practical situations and the angular distribution, $C(\theta)$ versus θ , of the annihilating photons appears bell-shaped about $\theta = \pi$ radians. The angular distribution may be used to obtain the momentum (p) distribution, $C(p)$ versus p .²² An example of both the angular and momentum distributions is shown in Figure 5. Momentum distributions generally consist of two components. It was pointed out by Ferrell,²³ that para-

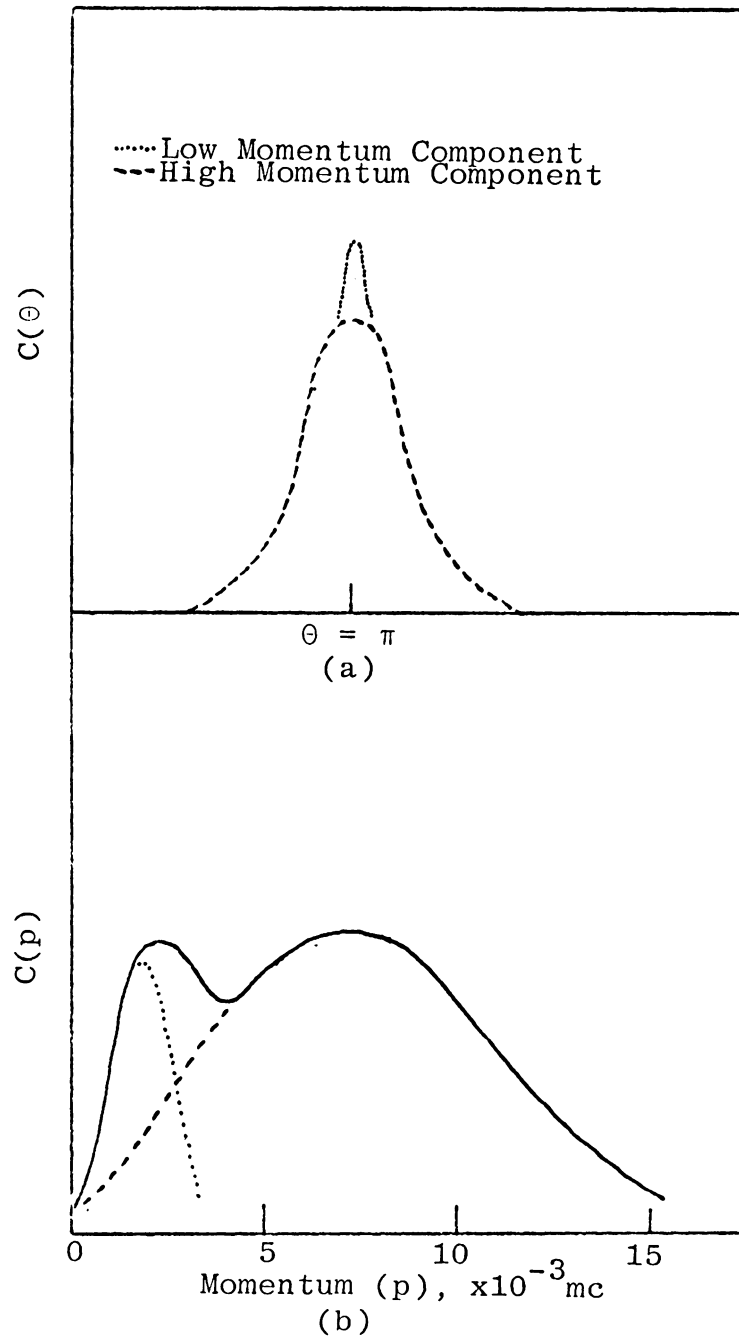


Figure 5. (a) Angular and (b) Momentum Distributions of the Annihilating Photons.

positronium should give rise to a relatively low momentum component since it usually decays by self-annihilation of the electron-positron pair. Direct positron and ortho-positronium pickoff annihilation occur with electrons bound to molecules of the system. These photons should contribute to a relatively higher momentum component because if the positron or positronium is thermalized, as is required by the bubble model, the primary contribution to the momentum of the center of mass of the annihilating pair will come from the high momentum bonding electrons. If this is true, the intensity of the low momentum component should correlate with the intensity of para-positronium as measured from lifetime measurements. Kerr et al²⁴ have compared the two intensities for ten organic liquids. The two measurements, taken from two completely different experiments, agreed within experimental error for all 10 compounds.

The higher momentum component of the momentum distribution can be used as an experimental verification of calculated electron momentum distributions if pickoff annihilation occurs as described by the "free volume model." The momentum distributions for electrons in C-C and C-H bonds have been calculated using analytic SCF functions for atomic carbon orbitals and Heitler-London-type functions for the two paired electrons in the C-C and C-H bonds.²⁵ The authors then computed a theoretical electron momentum distribution

curve for hexane and decane. The results for decane are shown in Figure 6. Component A is approximately gaussian and corresponded to the low momentum part due to para-positronium. Component B is the calculated C-H contribution and Component C is the calculated C-C contribution. When fit to the experimental curve, the relative areas under Components B and C were found to be in the exact ratio of the electrons in the C-H bonds compared to the electrons in the C-C bonds. For hexane and decane these ratios were 28/10 and 44/18 respectively. They concluded that during pickoff annihilation, the positron of ortho-positronium annihilates primarily with the bonding electrons of a molecule.

The third kind of positronium quenching is by chemical reaction. Tao and Green²⁶ have divided this category into four types: (1) oxidation by electron transfer; (2) reduction; (3) compound formation; and (4) double decomposition. This thesis deals primarily with the chemical reactions by compound formation and for that reason, it will be discussed in detail in Chapter 3.

Figure 7 summarizes the many interactions of positrons and positronium in water.²⁰ High energy positrons lose energy rapidly while slowing down. If a collision takes place while the positron is in the Ore gap, positronium atoms are formed with excess kinetic energy. About one-

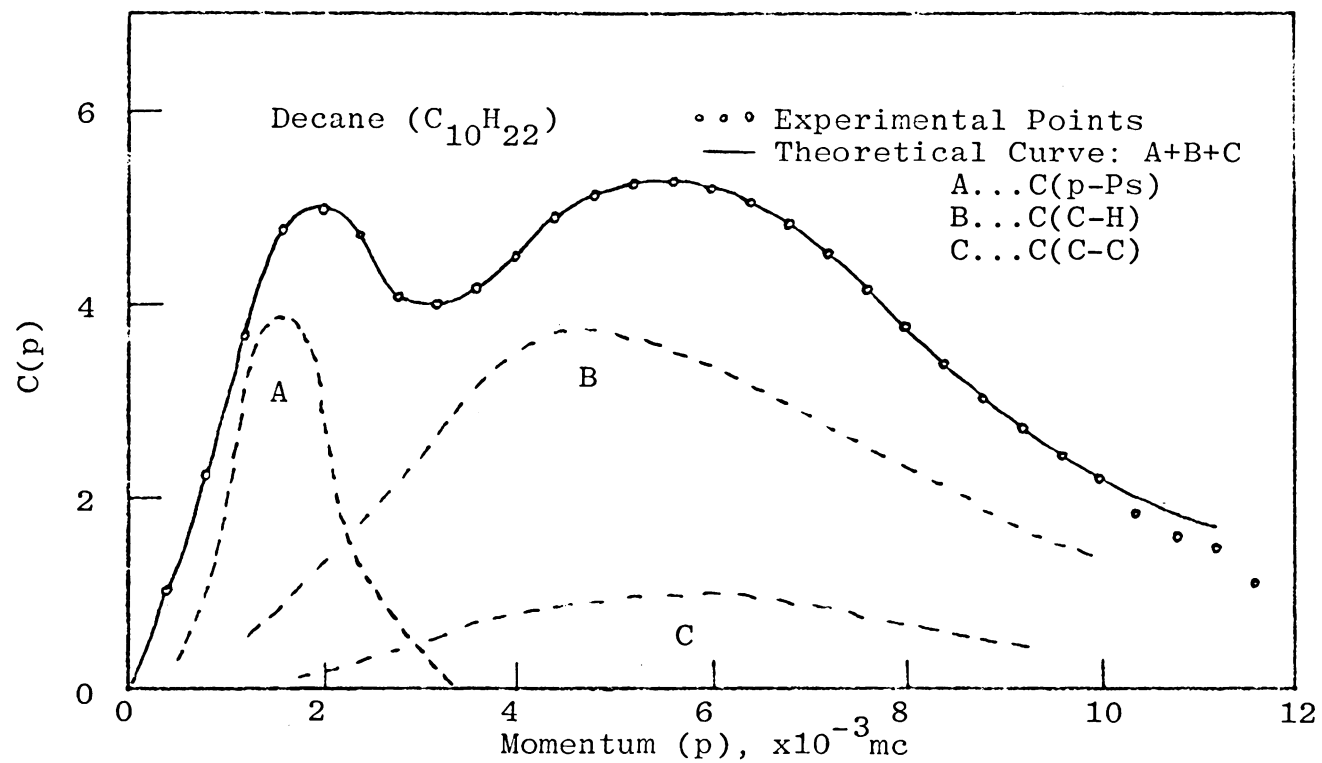


Figure 6. Computed and Experimental Momentum Distributions for Decane. (from Ref. 25)

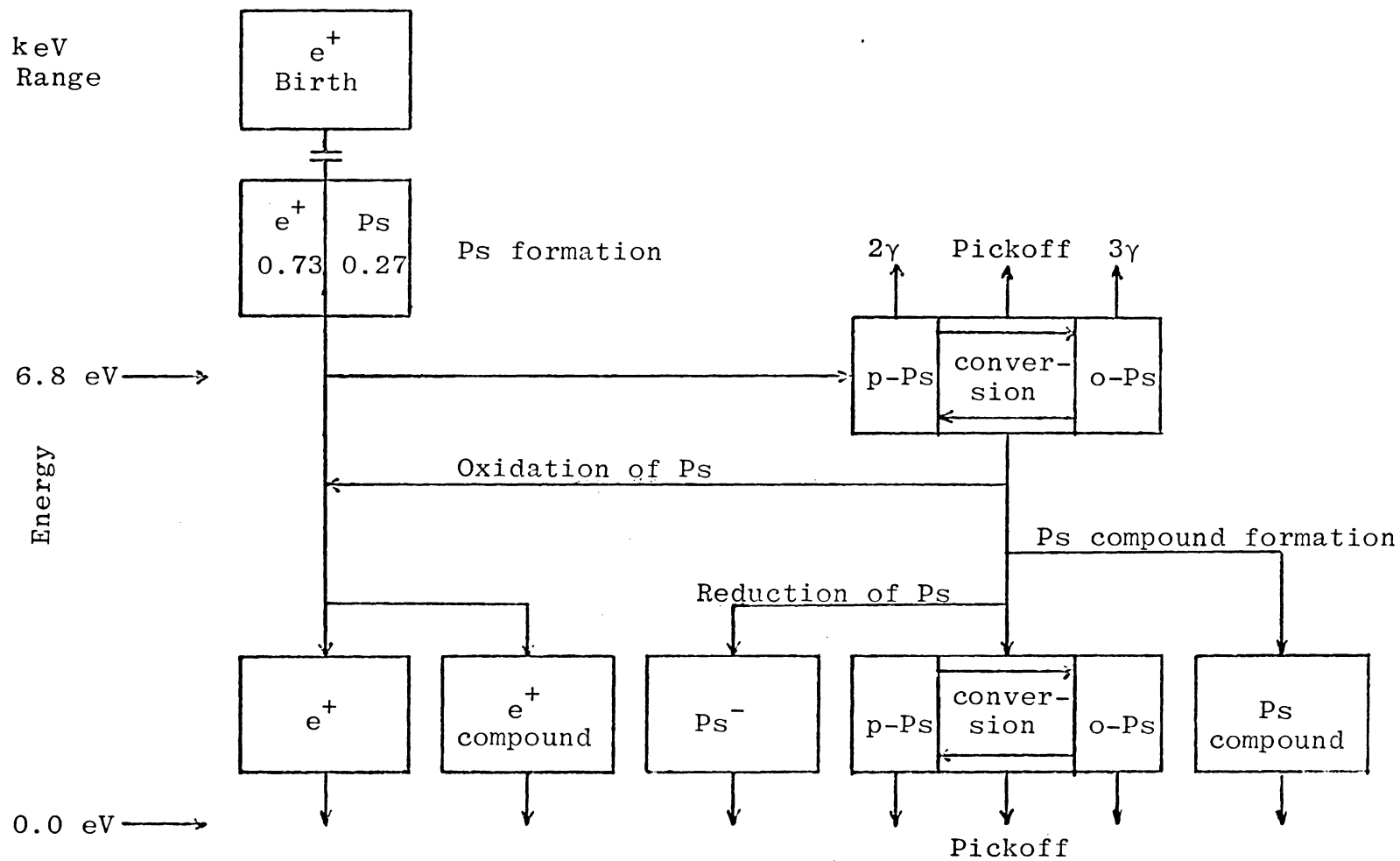


Figure 7. Positron and Positronium Decay in Water. (from Ref. 20)

third of the positrons will do this, three-fourths of them forming ortho-positronium and one-fourth forming para-positronium. If paramagnetic species are present, conversion quenching may occur. Pickoff annihilation is always possible. If strong oxidants are present, positronium may be oxidized back to a free positron. Reduction of positronium would form Ps^- , which has a binding energy of 0.3 eV, however, this species has not been observed.⁸ Positronium compounds may form at any time during or after thermalization. Free positrons may annihilate or form positron complexes and then annihilate. The possibilities for interaction are numerous, some are possible only for positrons or positronium, others are analogous to the reactions of the hydrated electron or hydrogen atom, two species to which positrons and positronium are often compared.

EXPERIMENTAL METHODS

A. Introduction

Positron lifetime measurements are usually carried out by using a delayed coincidence technique.²⁷ This method is the most applicable one for measuring short (less than 10^{-4} s) nuclear mean lives. It requires that a start pulse be present which originates from a zero-time physical event. Also, an end-time event pulse is necessary so that both pulses can be fed into a coincidence circuit. A graph of coincidence rate plotted as a function of time is the lifetime spectrum and can be analyzed to give the lifetime of the positrons in the system. Improvements on the basic experimental arrangement were made by Bell and coworkers when they introduced the fast-slow principle²⁸ and the time-to-pulse-height-converter.²⁹

A large majority of the data presented in this work was obtained by measuring a positron lifetime spectrum under varying conditions. These spectra were then used to measure the lifetime of free positrons as well as the lifetimes of both electron-positron bound states in solution. Positron annihilation rates are defined as the reciprocal of these positron lifetimes. These rates were used to evaluate annihilation rate constants. The measurements of certain rates as a function of temperature may be used to obtain various thermodynamic quantities. Hence, kinetic as well

as thermodynamic information was obtained in order to determine the mechanism by which one of the positron species, ortho-positronium, reacted; this was a major objective of the thesis.

B. Positron Sources

In the study of positronium chemistry, the most widely used source of positrons is sodium-22 which has a half-life of 2.60 years. It may be purchased from some of the major radionuclide distributors either as carrier free $^{22}\text{NaCl}$, $^{22}\text{NaHCO}_3$, or as carrier solutions of these compounds. The decay scheme³⁰ of ^{22}Na is shown in Fig. 8. A majority of the decays (90%) cause positron emission and leave the daughter nuclei, neon-22, in a very short lived ($\tau = 3\text{psec}$) excited state. Electron capture (EC) decay, in which no positron is formed, occurs ten percent of the time. The ground state of sodium-22 lies 2.840 MeV above the ground state of neon-22 and 1.564 MeV above the first excited state of neon-22. However, the maximum positron energy is only 0.544 MeV due to the energy requirements placed upon positron decay as explained in Chapter 1. The transition from the first excited state of neon-22 to its ground state takes place in a time two to three orders of magnitude faster than the average lifetime of a positron in condensed matter. Consequently, the 1.276 MeV gamma-ray can be used to initiate the zero-time start pulse required for coinci-

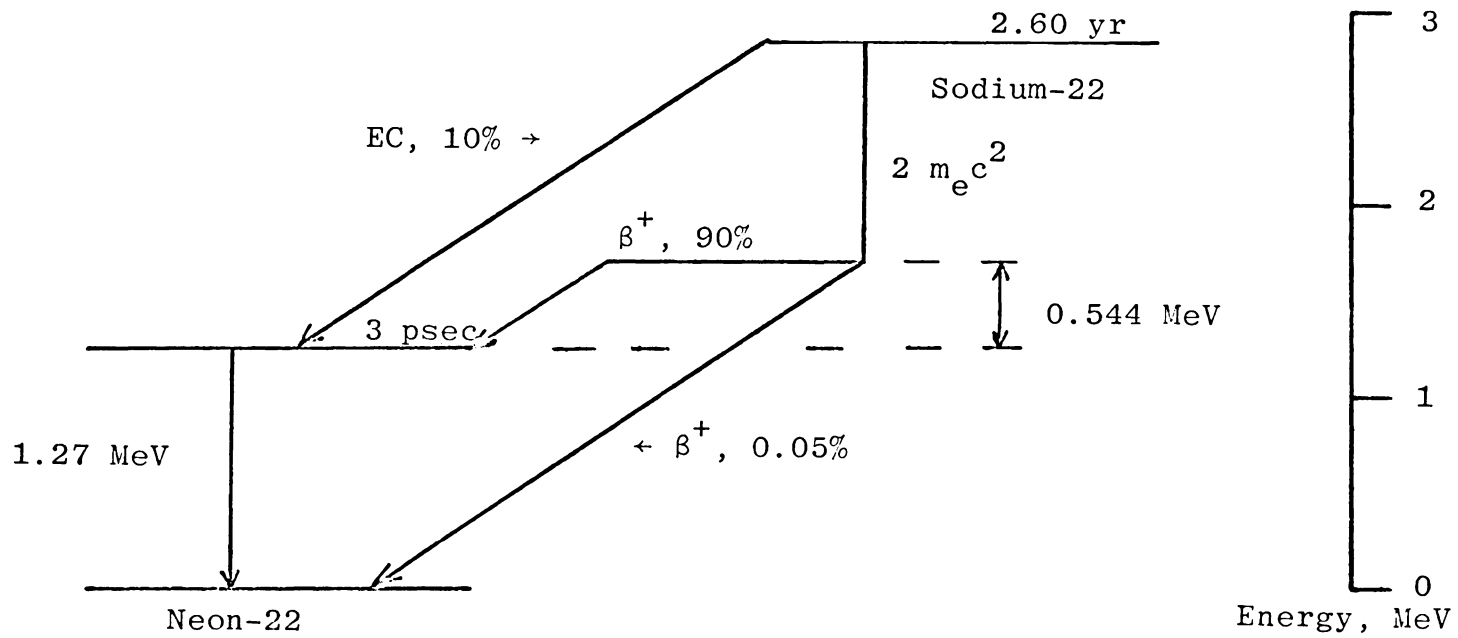


Figure 8. Decay Scheme of Sodium-22. (from Ref. 30)

dence studies, and the annihilation gamma rays used for the end-time signal.

All of the data presented in this study were taken on positrons annihilating in the liquid phase. The solid support for the positron source was a 1 mil thick aluminum foil, 5 mm by 10 mm in size. A water solution containing the sodium-22 was added dropwise to the foil and evaporated at atmospheric pressure under an incandescent lamp. This procedure was repeated until the foil contained the desired activity. This rather arbitrary amount of activity depended upon the specific activity of the sodium-22 solution, the number of drops added and the counting efficiency of the detector system. Ideally, one tried to place enough activity on the foil so that a spectrum containing at least 15,000 counts in the peak channel could be obtained every six to eight hours. In working with most organic liquids, the solubility of the sodium-22 compounds was small and very little activity was lost from the foil during the course of the experiment.

Activity losses due to the degassing procedure will be discussed in the sample preparation and degassing section of this chapter. Positron annihilation in the aluminum foil and Pyrex glass sample tubes has been studied³¹ and found not to be a problem when using the abovementioned procedures.

C. Fast-Slow Delayed Coincidence Counting

A typical arrangement for a fast-slow coincidence timing system is shown in Figure 9. A fraction of the gamma photons emitted during the positron decay and annihilation processes in the source are stopped by the two plastic scintillators (phosphors). Each unit consists of a phosphor optically coupled to a photo-multiplier (PHOTO-MULT) tube. The phosphors used were Nation 156 plastic scintillators, 1" by 1". They have a fast decay constant in comparison to the positron lifetimes and also high light output. The photomultiplier tubes were RCA 8575 mounted on Ortec 265 photomultiplier bases (PM BASE). They were used for their high gain, fast rise time, and are low in the emission of thermal electrons. The output signals from these detectors are simultaneously fed into two coupled electronic circuits. The circuits consist of a fast inner loop and a slow outer loop. Each circuit has two sides, a start side and a stop side.

The inner circuit was used to accurately determine the time interval between the formation of the positron and its subsequent decay. The output from the anode of the detector is sent to an Ortec 417 fast discriminator (FAST DISC). The units are operated at a biased voltage in order to reduce the noise level in the circuit by setting a lower threshold limit for accepting pulses from the photomultiplier tube. Both fast discriminators send out fast timing pulses to an Ortec

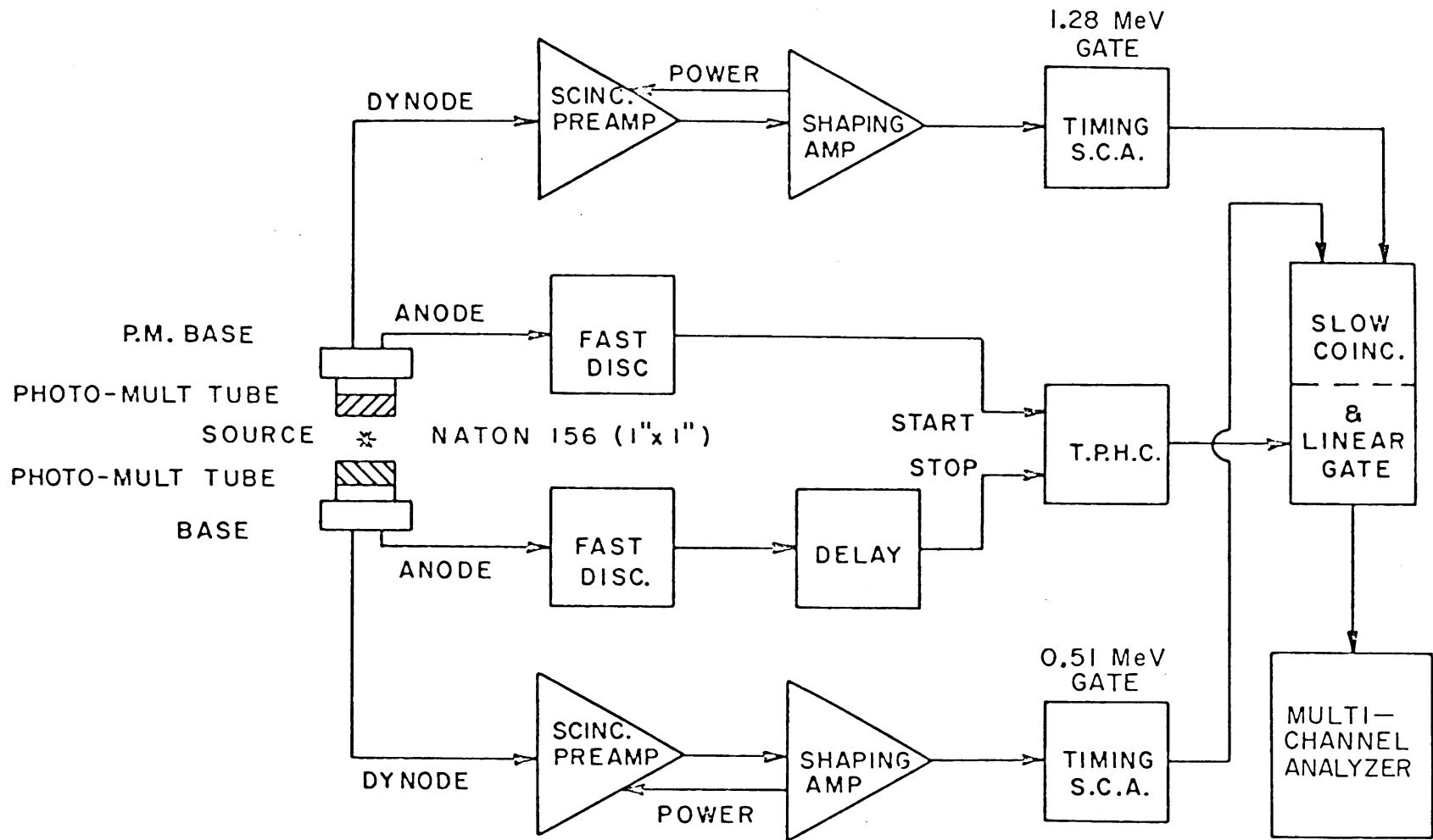


Figure 9. Simple Fast-Slow Timing System.

437 time-to-pulse-height-converter (TPHC). This unit converts the time difference between the start and stop pulses into a proportional pulse height signal. Before this pulse is allowed to pass into the multichannel analyzer (MCA), certain energy and coincidence conditions must be met. An Ortec 425 variable delay unit is inserted into the fast circuit to calibrate the channel width, time increment per channel, of the multichannel analyzer. The calibration procedure and the importance of this quantity will be discussed later.

The outer or slow loop takes its signals from the dynodes of the photomultiplier tubes. These signals have not been fully amplified in the detectors and therefore must first be amplified and shaped by means of an Ortec 113 scintillation preamplifier (SCINT PREAMP) and then an Ortec 440 shaping amplifier (AMP). Both sides of the slow circuit contain an Ortec 420 timing single channel analyzer (TIMING SCA). On the start side, this unit is biased so that an energy window is created that will only accept pulses that correspond in energy to very nearly 1.276 MeV. This is the energy of the de-excitation photon emitted during positron decay. Pulses substantially larger or smaller than this are rejected by the unit. The timing single channel analyzer in the stop side of this circuit performs the very same function for 0.511 MeV pulses, those from the annihilation process. If both the start and stop pulses reach the Ortec 409 slow

coincidence (COINC) and linear gate unit (this time is called the resolving time and is usually less than a μs) within a short time interval, the slow coincidence unit opens the linear gate and the signal from the time-to-pulse-height converter is allowed to pass into the multichannel analyzer. Overall, three different multichannel analyzers were used in this study. They were a Nuclear Data 2200, a Northern 600, and a Packard. They were all operated in a pulse height analysis (PHA) mode. These units will accept pulses of varying heights and sort them accordingly. These data representing the lifetime distribution are stored in the memory and can be read-out of the units in hard copy form onto paper, paper tape, or magnetic tape. Once the lifetime distribution data are taken out of the memory, it can then be computer analyzed into its various lifetime components.

The positron lifetime measurement begins when the starting detector accepts the zero-time pulse corresponding to the emission of a positron and continues until the end-time pulse from the decay of that positron is received. The lifetime of that positron is stored in memory and the coincidence circuits are cleared in order to measure the lifetime of another positron, etc. This procedure is analogous to having a large number of positrons at time $t=0$ and observing their lifetime distribution.

D. Positron Lifetime Spectra

The end result of the experiment described in the previous section is a positron lifetime spectrum. A typical spectrum is shown in Fig. 10. The ordinate of Fig. 10 gives the logarithm of the number of two-gamma decays while the abscissa is the time axis in nanosecond units. A point on the curve, $(\log N(t), t)$, represents the number of positrons that have lived for a certain time t and then decayed by two-gamma emission. This two component exponential decay spectrum is typical of positrons decaying in a condensed phase. The steeper short-lived component, λ_1 , is attributed to the decay of free positrons, those which have not formed a bound state with an electron, and para-positronium, the singlet bound state. The more shallow long-lived component, λ_2 , is assigned to ortho-positronium. Even though this species lives in solution as a triplet, its positron annihilates with a singlet electron by two-gamma emission before it reaches its expected lifetime.

In general, a multicomponent exponential decay spectrum has the following form.

$$N(t) = \sum_i N_i \exp(-\lambda_i t) + \text{background}$$

where the N_i 's are the pre-exponential constants that weight the individual components, and the λ_i 's are the annihilation rates for the various modes of positron decay. A mean life,

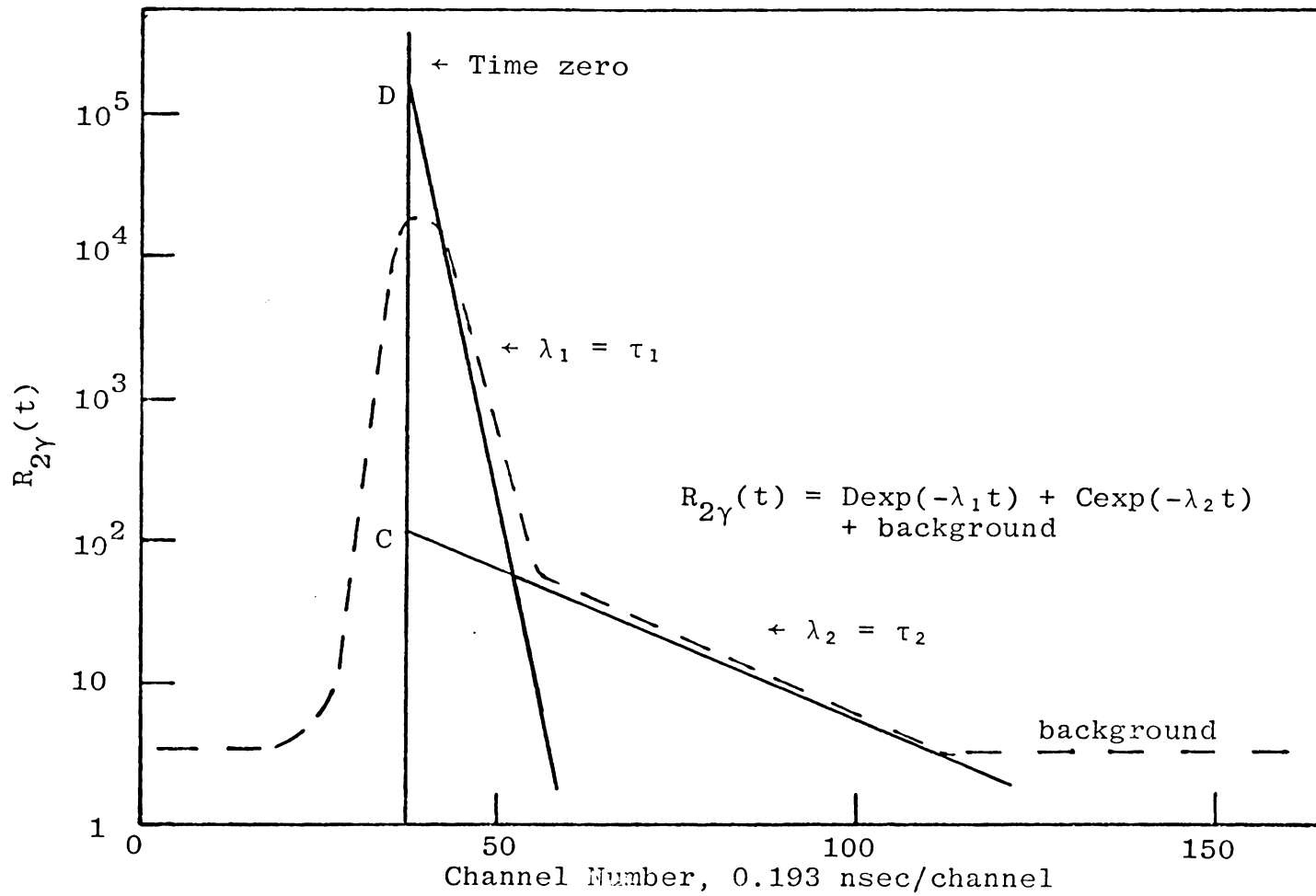


Figure 10. Typical Positron Lifetime Spectrum.

τ_i , is shown as the reciprocal of the *i*th annihilation rate. These spectra can be computer resolved into their various components using programs originally written by Cumming³² or Tao³³ and then modified in this laboratory. The programs were run on the IBM 360/155 computer located at the Virginia Tech Computing Center. Briefly, the programs perform a linear least squares analysis on each component present in the spectrum starting with the longest lived one and then constructs a calculated spectrum. The lifetime, annihilation rate, and intensity of each component is output along with a calculated fit parameter. This fit is a rather complicated function of the standard deviation and values less than 2.00 are generally acceptable. In addition to the printer output is a Calcomp plot of the experimental data and the calculated spectrum which can be obtained by using the PAL program.³⁴ An example of the Calcomp plot is given by Fig. 11.

This is a plot of the base ten logarithm of the coincidence rate versus time. The calculated fit is 1.05. The calculated lifetime and intensity of the two components are $\tau_1 = 0.398$ nsec, $I_1 = 80.4\%$ and $\tau_2 = 2.162$ nsec, $I_2 = 19.6\%$. The PAL program requires that the calculation be started beyond the curved top of the peak at a point where the first component appears to decay linearly. This distortion at the peak is caused by the finite resolving power of the instrument and will be discussed later. The total area under the

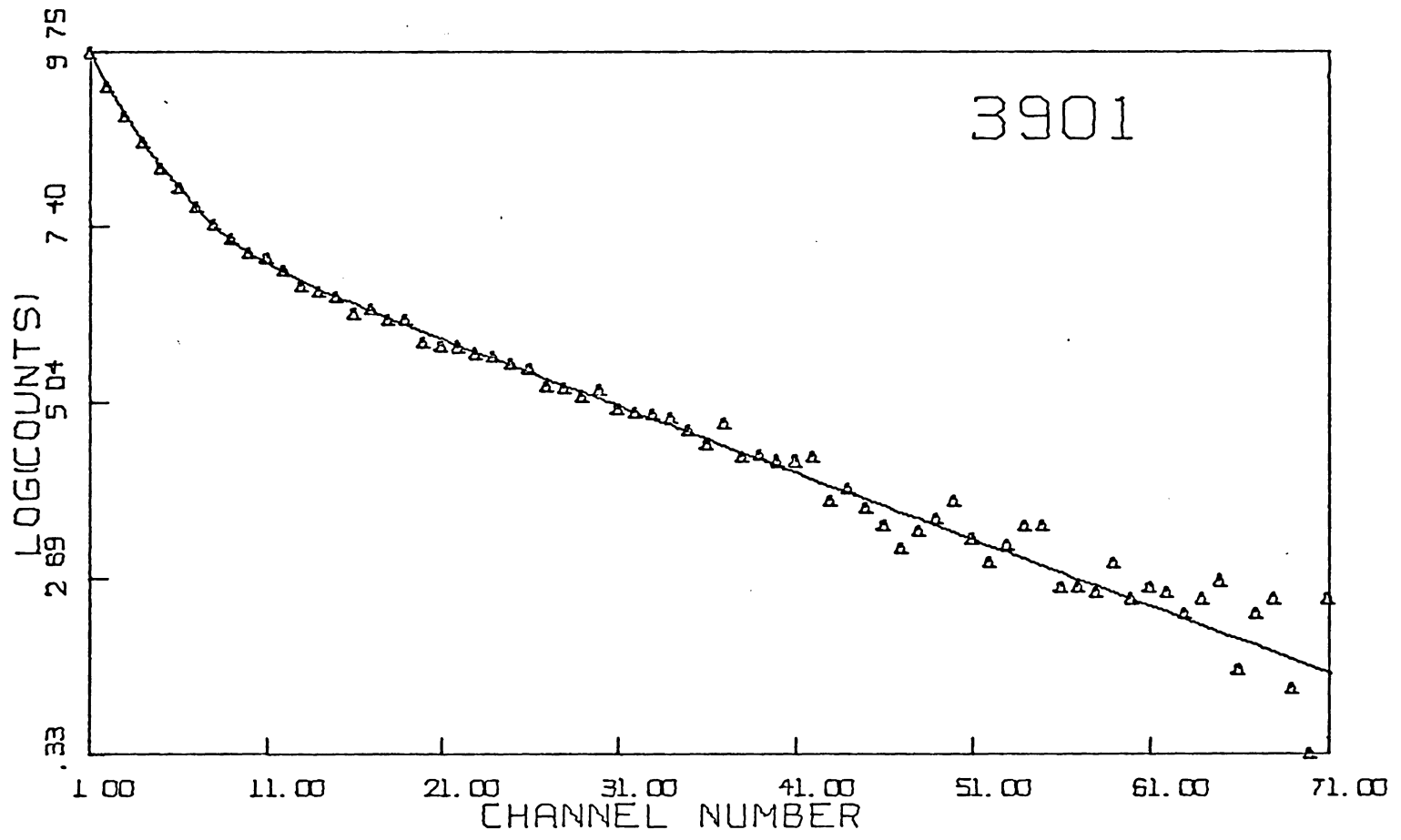


Figure 11. Typical Calcomp Plot of the Lifetime Data.

two components is normalized by the program to 100% so that by reporting the intensity of either component, one automatically sets the intensity of the other. No such relationship can be applied to the annihilation rates, however λ_1 remains fairly constant at 2.5 ns^{-1} in the liquid phase and, therefore, will not be reported here. It is λ_2 , the annihilation rate of ortho-positronium and its intensity, I_2 , that have been shown to be of interest chemically.³⁵ A quantity that is used to evaluate the ability of any fast-slow coincidence unit to resolve a multicomponent spectrum is termed the resolution of such a unit. This parameter has time units and is defined as the full-width-at-half-maximum (FWHM) of a prompt spectrum. A prompt spectrum is one in which the zero-time and end-time pulses are received by the detectors simultaneously. Ideally, a single spike should be seen in the lifetime spectrum corresponding to time $t=0$. Experimentally, this can easily be done by placing a cobalt-60 source between the two detectors. Cobalt-60 decays by negatron emission and emits a 1.173 and a 1.333 MeV gamma ray coincidentally. If the two timing single channel analyzers are set so that these two photons are used to start and stop the timing measurement, a prompt spectrum can be obtained. Due to the inadequate time resolution of the instrument, the prompt spectrum appears broadened about $t=0$ as shown in Fig. 12. The resolution, FWHM, was typically 300-400 psec for the experiments reported here.

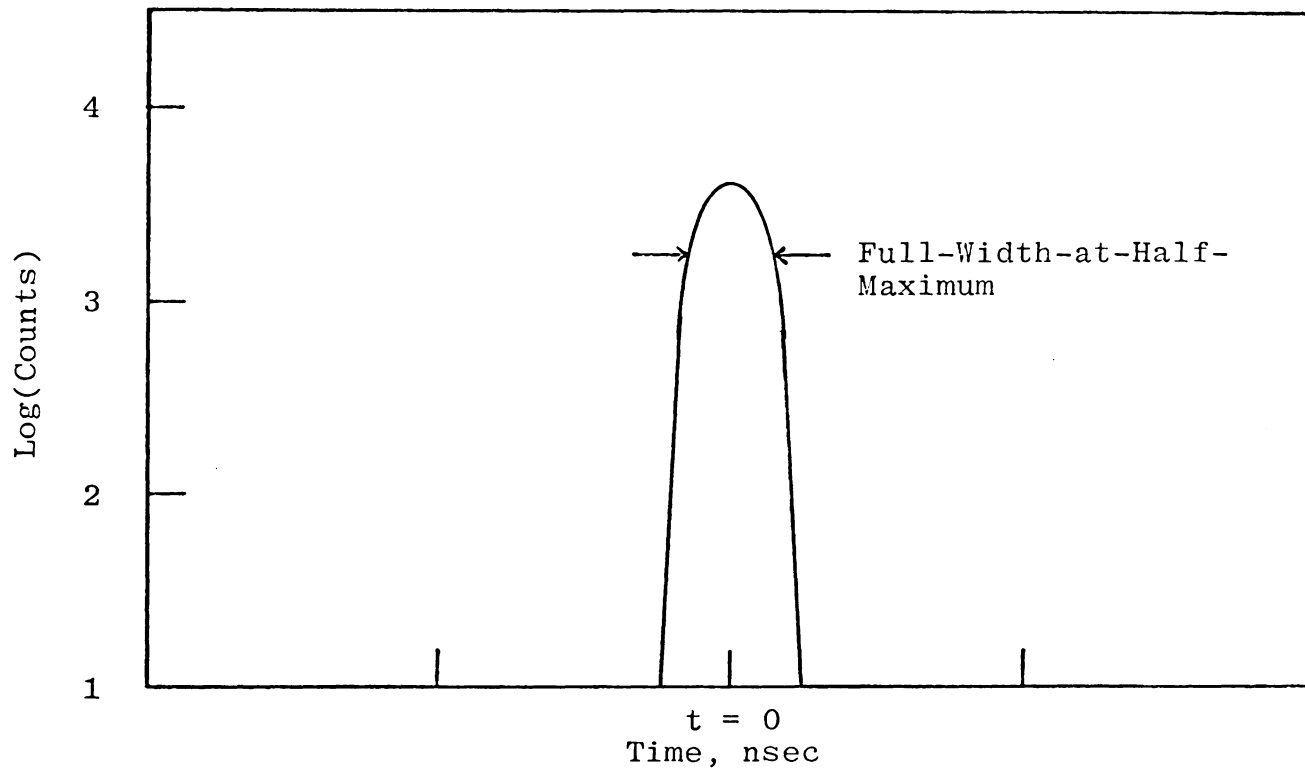


Figure 12. Prompt Spectrum of Cobalt-60.

By inserting a known delay, i.e. 16 nsec, into the timing circuit of Fig. 9, the prompt spectrum can be moved 16 nsec down the time axis. The number of channels between the maximum of the prompt spectrum with no delay and the maximum of the prompt spectrum with a 16 nsec delay can then be used to calibrate the channel width, time increment per channel. For this work, a 16 nsec delay caused the prompt maximum to move 83 channels which represents a channel width of 0.193 nsec/channel. This parameter is essential for accurately evaluating the individual components (slopes) of the lifetime spectrum and has been checked regularly.

E. Sample Preparation and Degassing

Measuring the positron lifetime distribution in the liquid phase requires a sample vial capable of holding the aluminum foil positron source and a few milliliters of solution. The vial has been constructed so that its contents could be easily degassed.

The rate constants for the reactions of ortho-positronium were calculated by measuring the ortho-positronium annihilation rates at various concentrations of the reactive compounds. The sample tube used in this type of experiment is shown in Fig. 13. It was constructed of cylindrical Pyrex glass tubing, approximately 200 mm long, 10 mm i.d. at the narrow counting end, and 25 mm i.d. at the top.

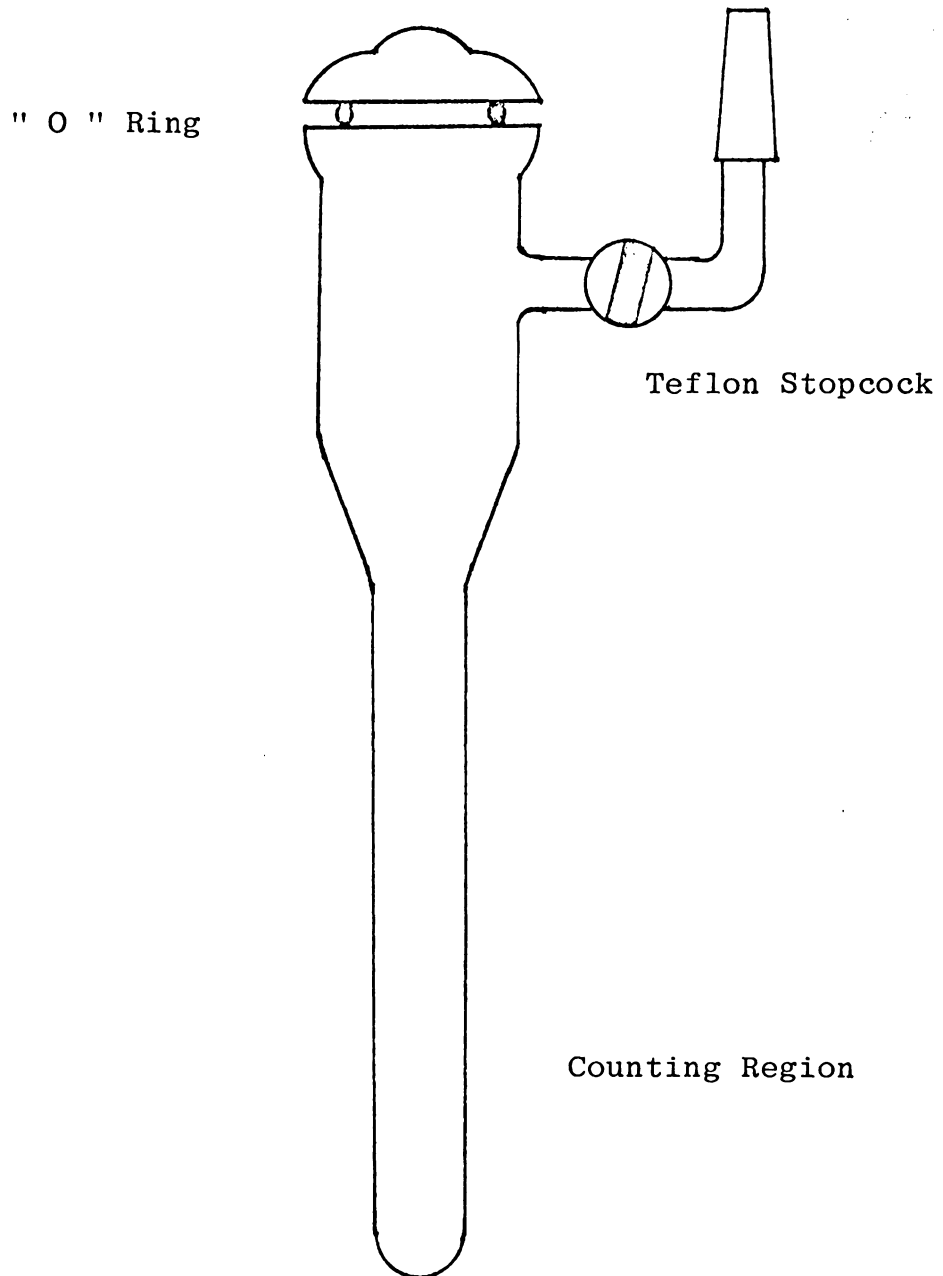


Figure 13. Sample Tube used for Variable Concentration Experiments.

Perkal³¹ has shown that a positron in a liquid such as water will travel approximately 0.5 mm before annihilating. A right angle side arm with Teflon stopcock was fitted to the wider portion and a mole 10/40 ground glass joint was attached so that the sample tube could be connected to the vacuum manifold. The lid was sealed to the main body of the tube by a Teflon "O" ring and held by a clamp (not shown). This arrangement has only been used for lifetime measurements taken at room temperature.

For measuring annihilation rates at temperatures above and below 295^oK, a special sample tube was used. The degassing bulb and sample tube for this type of measurement are shown in Fig. 14. The sample tube was a Pyrex glass tube about 100 mm long and 10 mm i.d. at the bottom. It was tapered to a capillary and had a male 10/40 ground glass joint on top. The wide bottom portion was broken so that the aluminum foil positron source could be inserted and then it was resealed. The solution to be counted, about 2 ml, was introduced via a capillary syringe through the top. The tube containing the source and solution was then fitted into the female 10/20 joint of the degassing bulb. The male joint of the sample tube protruded past the end of the female joint of the bulb by approximately 8 mm (see dotted line in Fig. 14). This prevented the solution that splashed up into the bulb during degassing from becoming

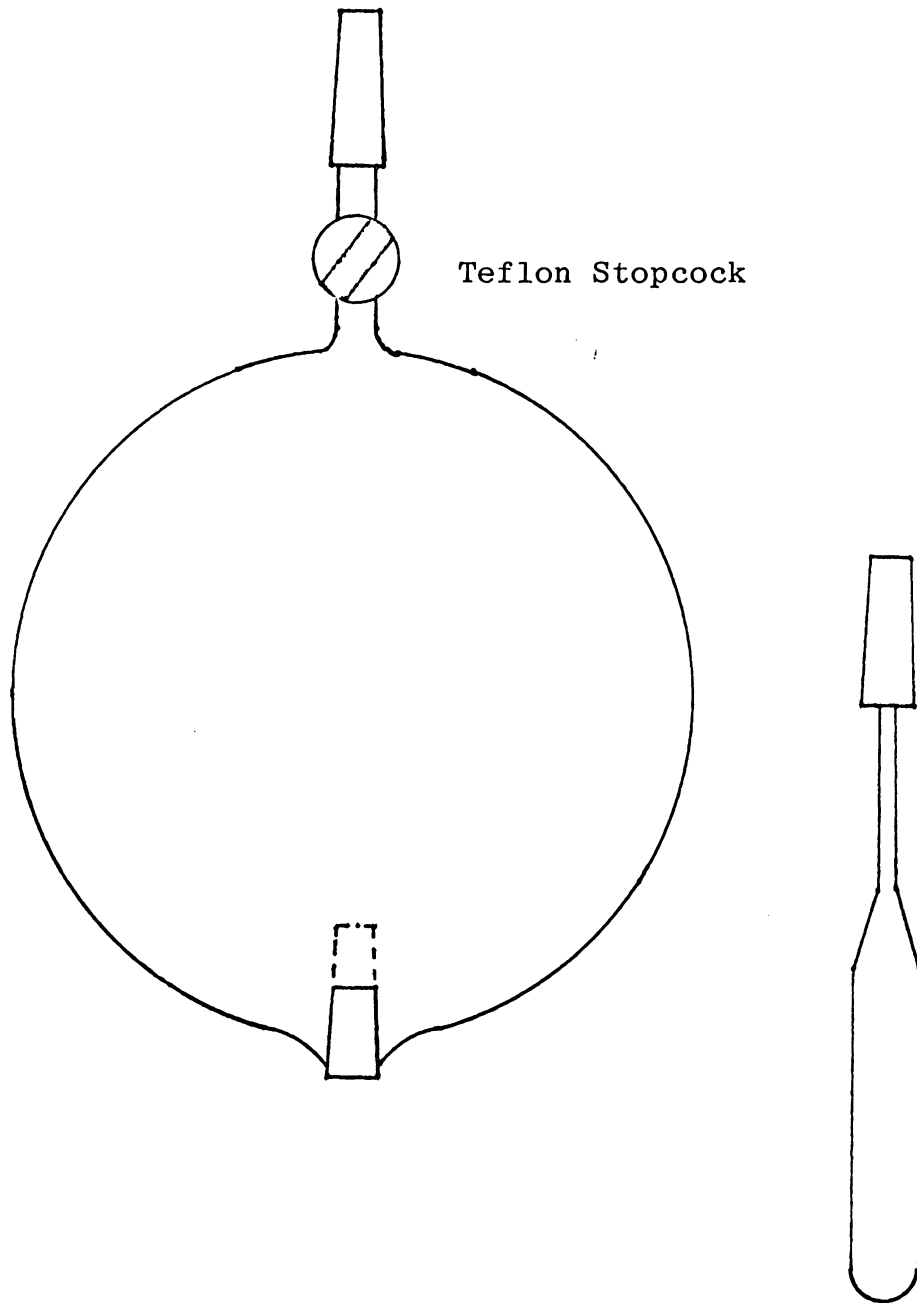


Figure 14. Degassing Bulb and Sample Tube used for Variable Temperature Experiments.

contaminated by the stopcock grease and running back into the sample tube. The degassing bulb was 75 mm in diameter, contained a Teflon stopcock to isolate the sample from the vacuum, and was fitted with male 10/40 ground glass joint for attaching it to the vacuum manifold. This relatively easy design proved quite adequate.

In dealing with organic liquids, the presence of dissolved oxygen drastically reduces τ_2 . This was first shown by Lee and Celitans³⁶ in 1966 and later by Cooper et al³⁷ in 1967. Even today, data that was taken in the presence of dissolved oxygen is still appearing in the literature and must be discounted. The process responsible for the shortening of the ortho-positronium lifetime has been found to be conversion quenching.³⁸ Arguments are made indicating that the triplet ortho-positronium reacts with the unpaired electrons of the oxygen molecules in such a way that the electron of the positronium atom is caused to spin-flip creating the much shorter lived para-positronium. In order to remove this quencher, each sample must be thoroughly degassed by a simple freeze-thaw technique. The sample tube is attached to the vacuum manifold and the section containing the solution is frozen by liquid nitrogen to avoid any evaporation losses. The Teflon stopcock is opened and the vial is evacuated to 10^{-4} torr using an oil diffusion pump. The stopcock is then closed and the frozen

sample is allowed to come to room temperature. This allows the dissolved oxygen to equilibrate between the solution and the partial vacuum above it. The sample is then refrozen, the stopcock opened and the oxygen above the frozen sample is pumped off. This procedure should be repeated 5-7 times or until the solution no longer is seen to give off bubbles upon thawing. For the variable concentration studies, the sample tube requires no further preparation and can be counted immediately after degassing. For the variable temperature studies, the sample must still be sealed at the capillary. So, after the last degassing, the sample was refrozen, the stopcock opened, and the capillary heated with a propane torch so that the capillary collapsed under the reduced pressure and the sealed sample tube removed from its ground glass joint.

Repeated degassing causes the aluminum foil backing to become brittle. When this occurs, the activity can be removed by dissolving the sodium-22 compound in water and then discarding the foil. The recovered activity can be reused. The degassing process also causes some activity to be worn off of the source. This will appear as a slower counting rate each time the source is used. It cannot be avoided and more activity must be added to the source periodically.

The counting room temperature was maintained at 22°C (295°K) and no special temperature regulation was used for the variable concentration studies. Measurements below room temperature were carried out in a Dewar flask fitted with a cold finger as shown in Fig. 15. The sealed cylindrical sample tube is fastened to a glass rod by a piece of rubber tubing and the temperature is measured by a Fisher low temperature thermometer. A variety of low temperature liquid nitrogen³⁹ slush baths were used along with isopropanol/dry ice (-75°C) and water/ice (4°C) baths. The Dewar was filled with the coolant and the cold finger placed between the two detectors for counting.

Temperatures above room temperature were maintained by the thermostated oil bath shown in Fig. 16. The oil, Fisher oil bath oil, was stirred by a magnetic stirring bar and insulated by the Dewar flask. The sample tube, thermometer, and platinum heating coil were submerged in the oil. The heating coil current was controlled by two thermistors and regulated by an external potentiometer. The apparatus was capable of reaching $220 (\pm 2)^{\circ}\text{C}$.

F. Solvents and Solutes

The solvents used in this study were benzene, toluene, n-heptane, and n-pentanol. Benzene and toluene were both Fisher Spectroanalyzed, the n-pentanol was Fisher Research Grade, and the n-heptane was Phillips Research Grade. All

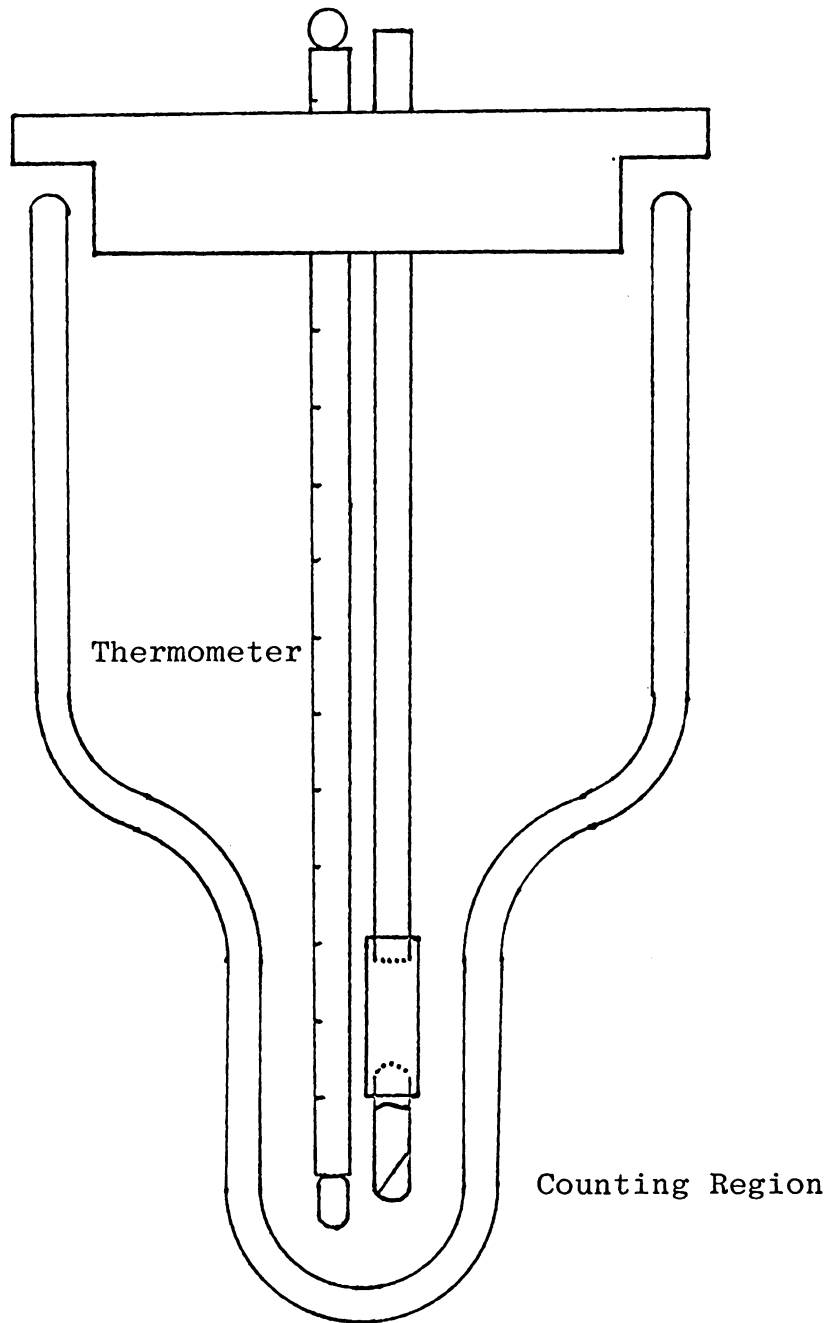


Figure 15. Experimental Arrangement used for Low Temperature Lifetime Measurements.

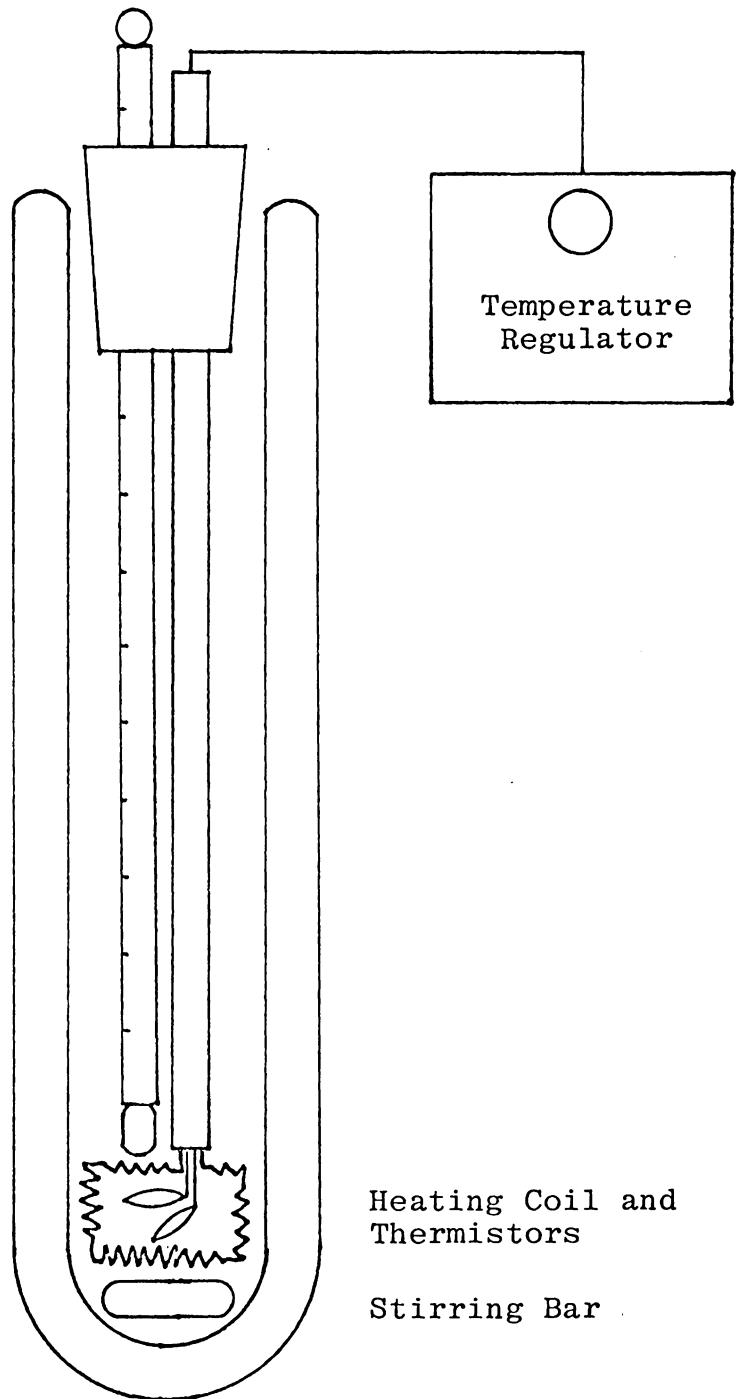


Figure 16. Experimental Arrangement used for High Temperature Lifetime Measurements.

were used without further purification; however, molecular sieve material was added to remove dissolved water.

The organic nitro compounds are all obtainable from Aldrich Chemical Company and were of the highest purity available. Most solids were recrystallized and most liquids were distilled until their melting points or refractive indices agreed with handbook values.

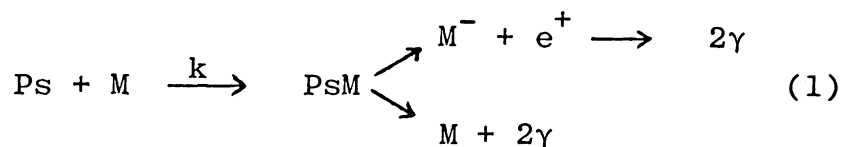
CHEMICAL REACTIONS OF ORTHO-POSITRONIUM

A. Diamagnetic Quenchers

There are certain diamagnetic compounds that are able to quench the lifetime of ortho-positronium to such an extent that the long-lived component of the positron lifetime spectrum is absent. Among them are the nitroaromatics, quinones, certain compounds that resemble maleic anhydride in structure and composition and some nitriles. Goldanskii^{40,41} was the first to observe this behavior in dilute solutions of nitrobenzene in benzene.

If the ortho-positronium lifetime is shortened so that it becomes indistinguishable with the lifetimes of free positrons and para-positronium, then the positron lifetime spectrum appears as a rapidly decaying one-component exponential decay curve. This is thought to be caused by rapid chemical reaction of the ortho-positronium. The ortho-positronium annihilation rate, the one that gives rise to the shallow component in the spectrum, increased to such an extent that the long-lived component was "pushed" into the shorter-lived steep one. However, inhibition of positronium formation will also cause the long-lived component to be absent from the spectrum. If ortho-positronium is not formed in any significant amount, $I_2 \approx 0$ and the long-lived component "drop out" of the spectrum.

It is rather simple to determine whether inhibition of positronium formation or rapid chemical reaction is taking place. The lifetime spectrum is observed as a function of increasing solute concentration in a suitable solvent. One that will give both a long-lived second component, $\tau_2 > 2$ nsec, and a fairly high second component intensity, $I_2 > 20\%$. (These values should serve only as general guidelines). Positronium atoms that react while still possessing a considerable amount of kinetic energy must do so very early in their lifetimes and annihilations by this mode appear in the short-lived component. If on the other hand, the annihilation rate remains constant but I_2 decreases from the value in the pure solvent, I_2^0 , to zero, inhibition of ortho-positronium formation is occurring. Goldanskii⁴¹ proposed the following mechanism for the reaction of ortho-positronium with nitrobenzene and other rapidly reacting compounds.



The positronium atom is thought to form a complex with the reacting molecule, M, and either annihilate while in the complex (lower path) or, since all diamagnetic quenchers are known to be strong electron acceptors, annihilate as a free-positron following a charge transfer mechanism (upper

path), $\tau_F \approx 0.45$ nsec in condensed matter. Two arguments can be made against the charge transfer mechanism.

Although all quenchers are known to be strong electron acceptors, not all strong acceptors quench ortho-positronium, i.e. benzonitrile, benzaldehyde, and naphthalene. Charge transfer is energetically possible in solution only if the following inequality is satisfied.

$$-I + A + Q + S > 0 \quad (2)$$

where I is the ionization potential of the donor, A is the electron affinity of the acceptor, Q is the Coulomb interaction energy of the ions, and S is the solvation energy of the ions.⁴¹ The ionization potential of positronium is constant and all compounds with large electron affinities should at least react to some extent with ortho-positronium. It is also difficult to envision the electron transfer step of the mechanism. The positronium atom does not have a nucleus in the usual chemical sense. Its positron and electron are light particles and should be delocalized in the PsM complex. Charge transfer in a complex of this nature has no significance in terms of molecular orbital theory. This does not exclude electron-transfer reactions of positronium via some other mechanism and indeed this has been observed.⁴²

The second order rate constant for the reaction of ortho-positronium with M in Eq. (1) can be obtained as follows:⁴⁹

$$-d[\text{Ps}]/dt = k[\text{Ps}][\text{M}] \quad (3)$$

Dividing by [Ps] yields

$$-d\ln[\text{Ps}]/dt = k[\text{M}] \quad (4)$$

Since $M \gg \text{Ps}$, the reaction will appear to be pseudo-first-order. The left-hand-side of Eq. (4) is defined to be the rate of disappearance of positronium, λ_M . Therefore,

$$\lambda_M = k[\text{M}] \quad (5)$$

The ortho-positronium annihilation rate in a pure liquid such as benzene is $\lambda_2 (= \lambda_{\text{Benzene}})$ and

$$k_{\text{Benzene}} = \lambda_2 / [\text{Benzene}] \quad (6)$$

The annihilation rate in benzene was measured to be 0.33 nsec^{-1} , this corresponds to a rate constant for the reaction of $2.96 \times 10^7 \text{ M}^{-1}\text{sec}^{-1}$.

The annihilation rate for a solution is actually the sum of two rates, λ_S , the rate in the solvent, and λ_M , the rate due to the solute, M. If the solute is present in small concentration, i.e. millimolar, the following approximations are valid:⁴³

$$[\text{Solvent}] \approx \text{constant} \quad (7)$$

$$\lambda_S \approx \text{constant}$$

and

$$\lambda_2 = \lambda_S + k_M[M] \quad (8)$$

Any rate increase, $\lambda_2 - \lambda_S$ is due entirely to the solute. A plot of λ_2 versus $[M]$ has a slope that is equal to the rate constant for the reaction between the ortho-positronium and M. An example of one such plot is given in Fig. 17 for nitrobenzene in benzene. The rate constant was found to be $27 \pm 2 \text{ M}^{-1} \text{ nsec}^{-1}$, almost three orders of magnitude higher than rate constant for the reaction with benzene.

The general statements made in this section apply to all rapid reactions. For reasons of availability and numerous possible substituents, the nitroaromatics were chosen to be examined in this study. The following sections will deal with the ortho-positronium reaction with them.

B. Conjugation Effects

Conjugation, or lack of it as the case may be, is a major factor in determining the ability of a molecule to chemically quench ortho-positronium. The microwave spectrum of nitrobenzene shows that the molecule is planar⁴⁴ with an internal barrier to rotation of $1000 \pm 500 \text{ cm}^{-1}$. When then the nitro group is co-planar with the aromatic ring, it has

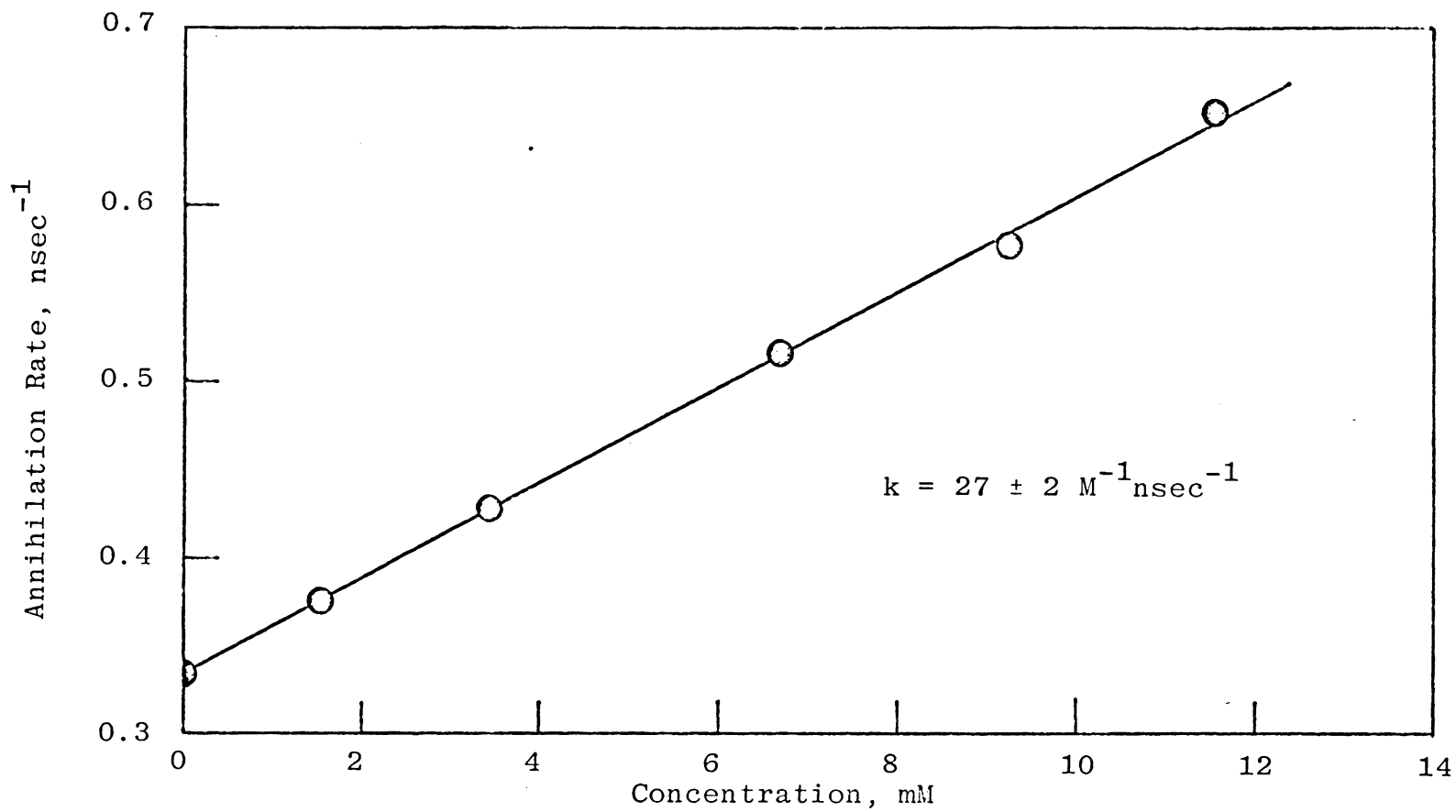


Figure 17. λ_2 versus Nitrobenzene Concentration in Benzene.

its maximum electron withdrawing effect and is best able to remove electron density from the ring.⁴⁵

The steric interaction between the nitro group and the methyl group(s) in the ortho position was used to measure the effect that conjugation losses had on the values of the measured rate constants. Methyl groups were convenient because their inductive and mesomeric effects are small. Toluene, like benzene, reacts very slowly with positronium, $k_{\text{Tol.}} \approx 10^7 \text{ M}^{-1} \text{ sec}^{-1}$. The observed rate constants for some model compounds are given in Fig. 18. Para- and meta-nitrotoluene react slightly slower than nitrobenzene itself, but the conjugation losses in the ortho isomer cause it to react only one-third as fast. If a second methyl group is placed adjacent to the nitro group, i.e. 2,6-dimethylnitrobenzene, the nitro group is twisted to an even greater extent out of the plane of the ring. It is therefore less able to conjugate with the ring and its corresponding rate constant drops further. This behavior was observed every time a large group was placed next to the nitro group (Appendix I lists all of the rate data taken in benzene solution at 295°K).

Isolating the nitro group from the π system was also found to have a drastic effect on the measured rate constants. In Figure 19, the rate constant for α -nitrotoluene is compared to that of nitrobenzene and found to be drastically

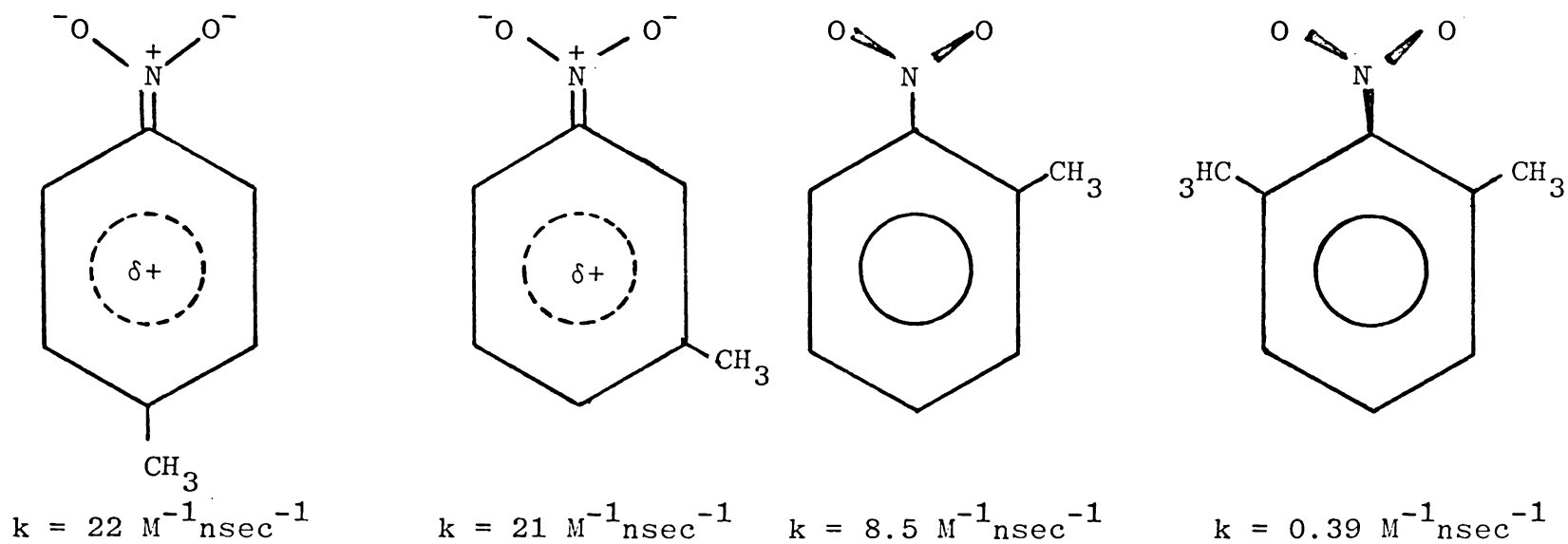


Figure 18. Steric Interactions that Force the Nitro Group out of the Plane of the Ring Cause a Reduction in the Rate Constant.

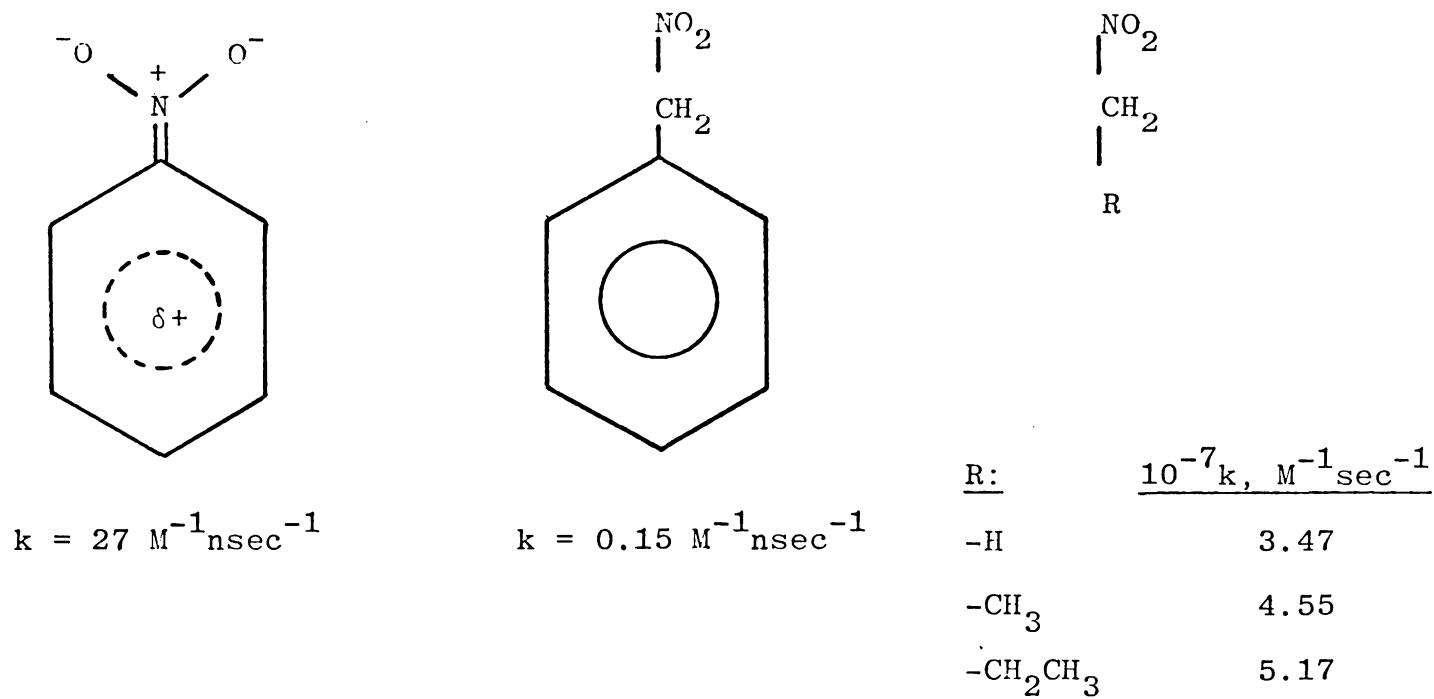


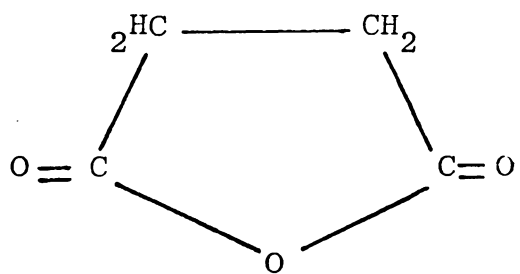
Figure 19. Nitro Compounds in Which the Nitro Group is Not Attached to a Conjugated System React Slowly with Ortho-Positronium.

reduced. Similar small reaction rate constants are observed for the nitroalkanes where there is no conjugation between the nitro group and the alkyl group.

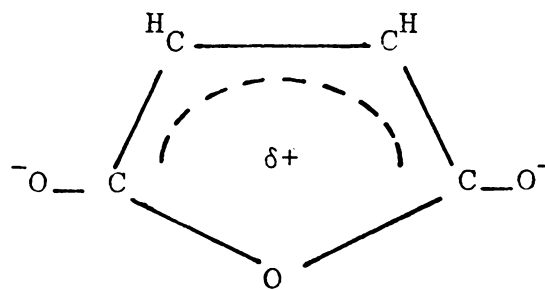
A third way of examining the effect of conjugation is by measuring the rate constant for the reaction with certain anhydrides that contain varying amounts of conjugation. Fig. 20 shows a comparison of the observed positronium rate constants for the reaction with succinic and maleic anhydride. Succinic anhydride is unconjugated and exhibits a rate constant indicative of pickoff quenching. Conversely, maleic anhydride is highly conjugated and has resonance structures that place a delocalized positive charge on all four carbon atoms. The addition of the double bond caused the positronium rate constant to increase by a factor greater than 200, compared to succinic anhydride. These two are compared to the highly conjugated chemical quencher, parabenzoquinone which shows a much higher reaction rate constant.

C. Temperature Effects

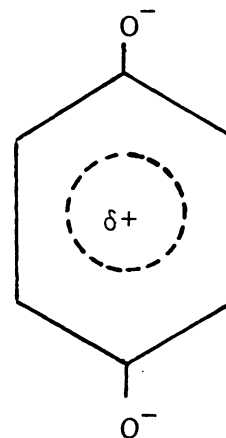
If the mechanism given in Eq. (1) is correct, a study of the temperature dependence of the rate constants should provide information about the activation energy required to form the various P_sM complexes. The solvent was changed from benzene to toluene for these experiments because it could be used over a much wider temperature range (200-500°K).



$$k = 0.03 \text{ M}^{-1} \text{ nsec}^{-1}$$



$$k = 7.1 \text{ M}^{-1} \text{ nsec}^{-1}$$



$$k = 50 \text{ M}^{-1} \text{ nsec}^{-1}$$

Figure 20. Effect of Conjugation on Positronium Rate Constants.

First the temperature dependence of the solvent annihilation rate had to be determined in order to use Eq. (8) for calculating the rate constants. Fig. 21 shows a plot of λ_2 ($=\lambda_S$) versus $T(^{\circ}\text{C})$ for toluene. A linear least squares program was used to find the best fit to the data. It was found that

$$\lambda_S(T) = 0.3368 - (6.60 \times 10^{-4}) T \text{ nsec}^{-1} \quad (9)$$

The second-order rate constant for toluene changed regularly from $3.45 \times 10^7 \text{ M}^{-1}\text{sec}^{-1}$ at -75°C to $2.97 \times 10^7 \text{ M}^{-1}\text{sec}^{-1}$ at 170°C .

When the annihilation rate of a 9.77 mM solution of nitrobenzene in toluene was measured over this temperature range, a rather interesting behavior was observed. The observed rate constant, k_{obs} , when plotted against the reciprocal Kelvin temperature, increased linearly from -75°C up to approximately room temperature, $10^3/T^{\circ}\text{K} = 3.4$, and then decreased linearly as the temperature was raised further.⁴⁶ This is shown in Fig. 22. A comparison is made in this figure between the temperature dependence of the rate constants in toluene and the nitrobenzene solution. The difference between the two rate constants is greatest around room temperature. But at the high and low temperature ends they begin to approach each other. Deviations from typical Arrhenius behavior was found for other nitroaromatics. Fig. 23 shows the $\log k$ versus $1/T$ curves for p-

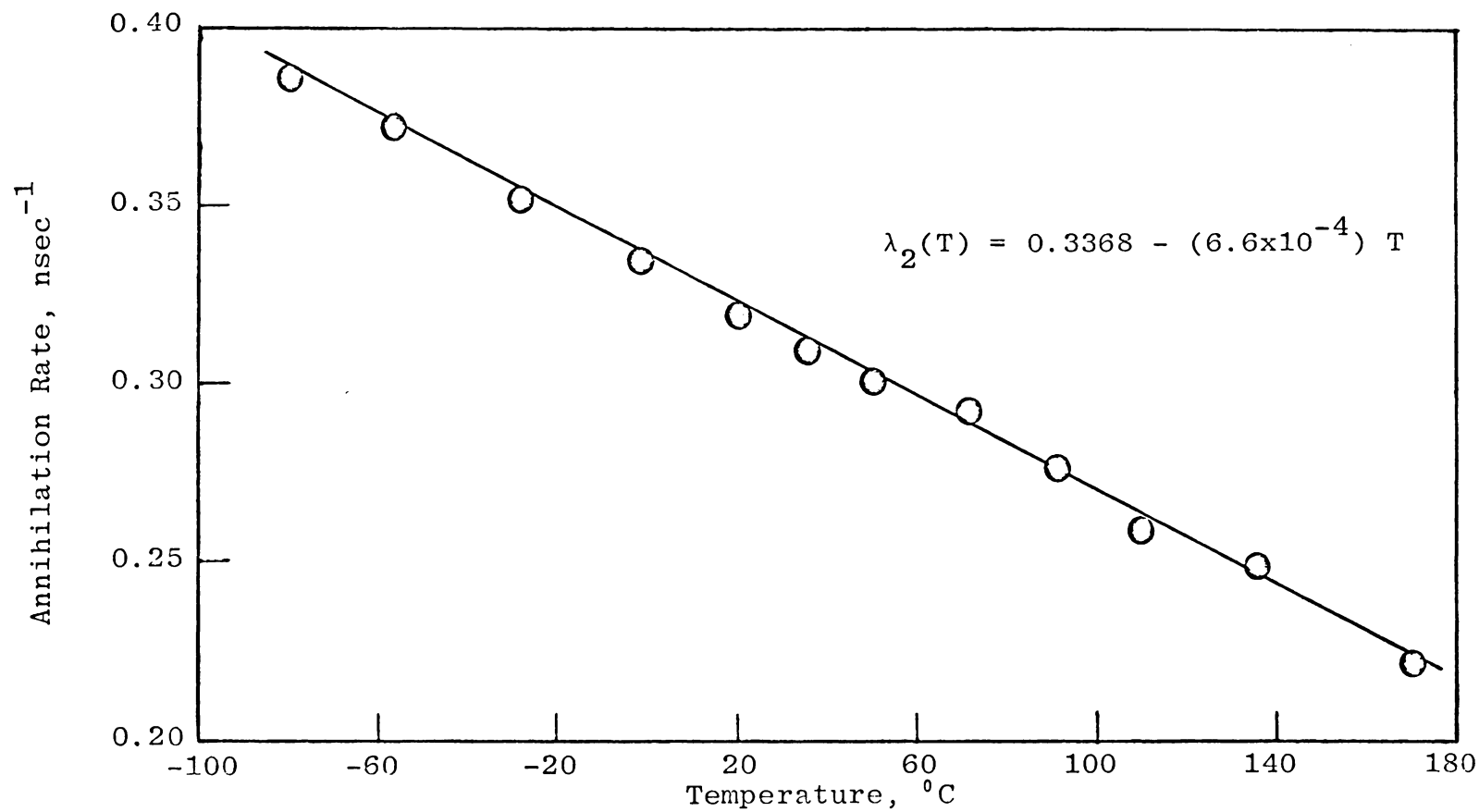


Figure 21. λ_2 versus T (°C) in Toluene.

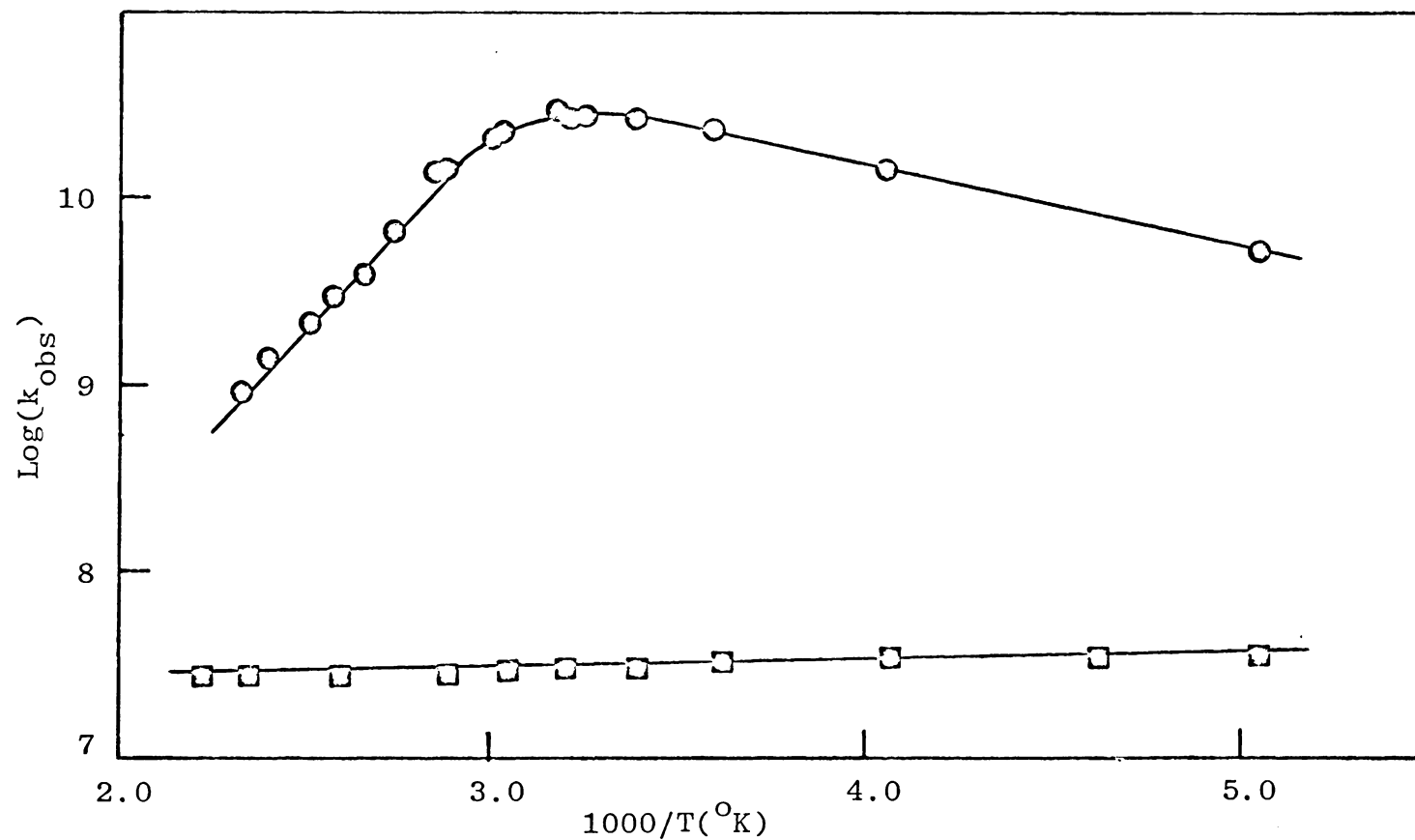


Figure 22. Arrhenius Plots for the Reaction of Ortho-Positronium With 9.77 mM Nitrobenzene in Toluene (○) and Pure Toluene (□).

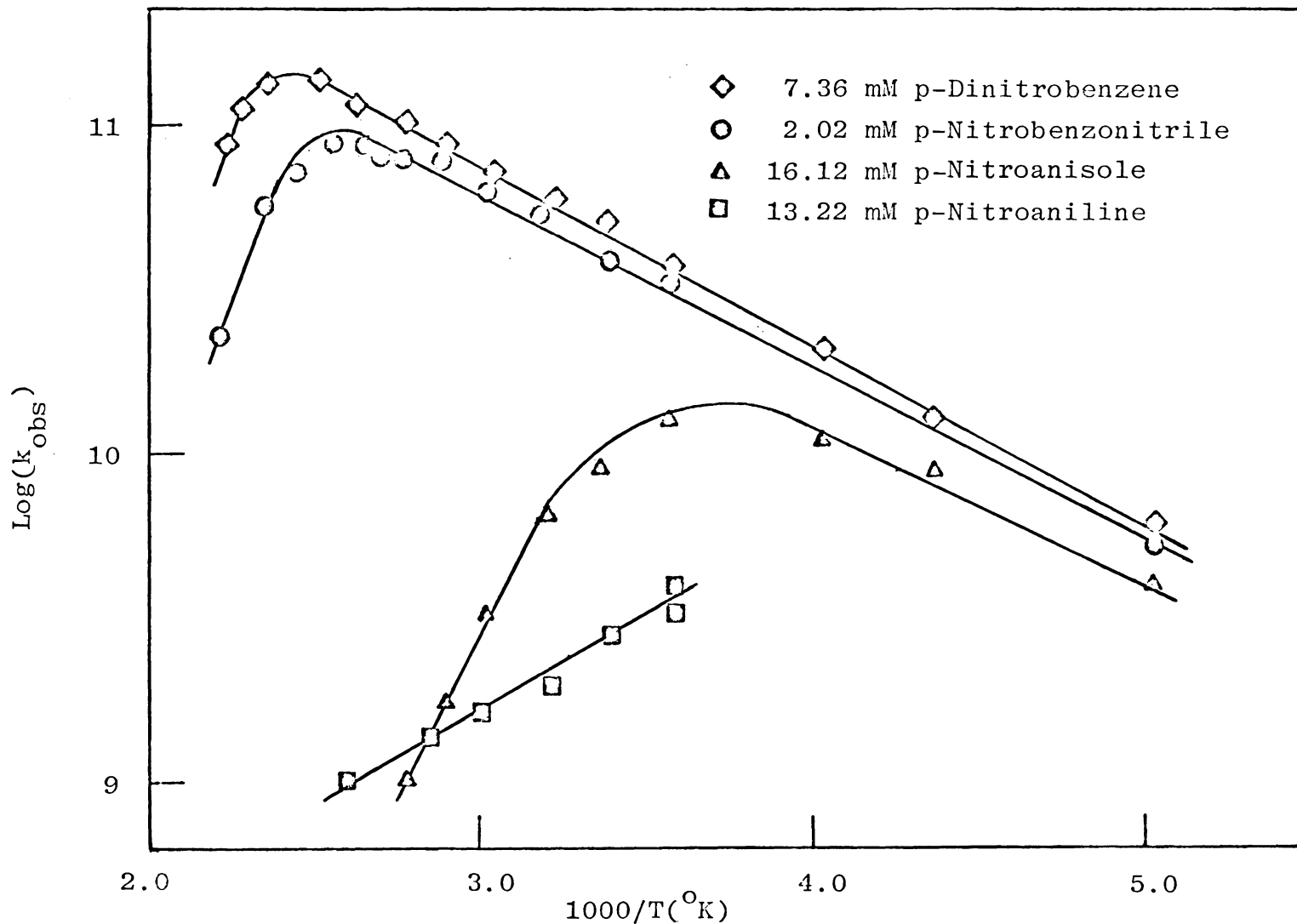


Figure 23. Arrhenius Plots for the Reaction of Ortho-Positronium with Various Nitroaromatics in Toluene.

dinitrobenzene, p-benzonitrile, p-nitroanisole, and p-nitroaniline. The Arrhenius portions of these five curves (this refers to the section of the curve with negative slope) all have approximately the same slope, while the non-Arrhenius sections (positive slope) have different slopes. p-Nitroaniline could not be run below 4°C due to its insufficient solubility in toluene below this temperature.

There are a variety of reactions that can cause deviation from the Arrhenius law. In principle, even the simplest kinetic processes should show a temperature dependence of the activation energy,⁴⁷ although it is rarely observed due to limited experimental temperature ranges or accuracy. Deviation from the Arrhenius law of a different nature can be found for reactions involving, for example equilibria, parallel, or consecutive processes.⁴⁸ This may be caused by a changing reaction pathway, rate determining step, or any process that causes the observed activation energy to change with temperature. For the type of deviation shown in Figures 22 and 23, where the measured activation energy actually becomes negative at high temperatures, only a few of these reactions are possible.

The inversion in the log k versus 1/T plot eliminates the possibility of a reaction involving a consecutive or parallel pathway. Deviation in the Arrhenius plot could

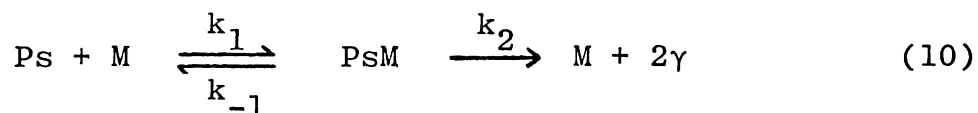
occur by this method but both activation energies would be positive.

Formation of an unreactive substance in the sample might explain the observed behavior. The nitroaromatics could possibly decompose at higher temperature giving rise to a product that reacts more slowly. This possibility was ruled out experimentally. The data was obtained by first measuring the rate constant at room temperature followed by kinetic runs at increasingly higher temperatures. After the measurement was taken at the highest temperature, the sample was rerun at room temperature to check against decomposition. In all cases, the two room temperature values agreed within experimental error. Reversible formation of an unreactive species has not been ruled out by this procedure. For example, the positronium could react with a solvent-solute complex. Although nitroaromatics do form these complexes in toluene, the results in other solvents discounted this explanation. See the section on solvent effects in this chapter.

The positronium atom is expected to tunnel readily due to its low reduced mass.⁴⁹ However, the tunnel effect, when observed, usually causes a $\log k$ versus $1/T$ plot to curve rather than invert as in the present case. A tunneling calculation has been made for both a positron and positronium atom incident on a potential barrier of varying height and geometry.⁵⁰ The results indicate that the prob-

ability of tunneling is very high. To apply these results to the present situation is impractical at the present time. The shape of the barrier and the nature of the tunneling particle are not known. Figure 24 is a plot of log of the transmission coefficient versus particle energy for a positronium atom tunneling through a parabolic, square wave, and Eckart barrier, 1 Å wide and 20 kcal/mole high.⁴⁹ The initial results are encouraging and further work seems warranted.

A mechanism that satisfactorily explains the observed behavior is the following:



ortho-Positronium combines with the nitroaromatic molecule to form a PsM complex in a reversible step. This complex is thought to be of a nature in which the positron and electrons are completely delocalized. The electron and positron enter into different molecular orbitals of the complex. The complex can either go on to annihilate the positron, k_2 , or decompose into an ortho-positronium atom and molecule, k_{-1} . Since the concentration of M remains essentially constant throughout the experiment, the mechanism can be simplified to

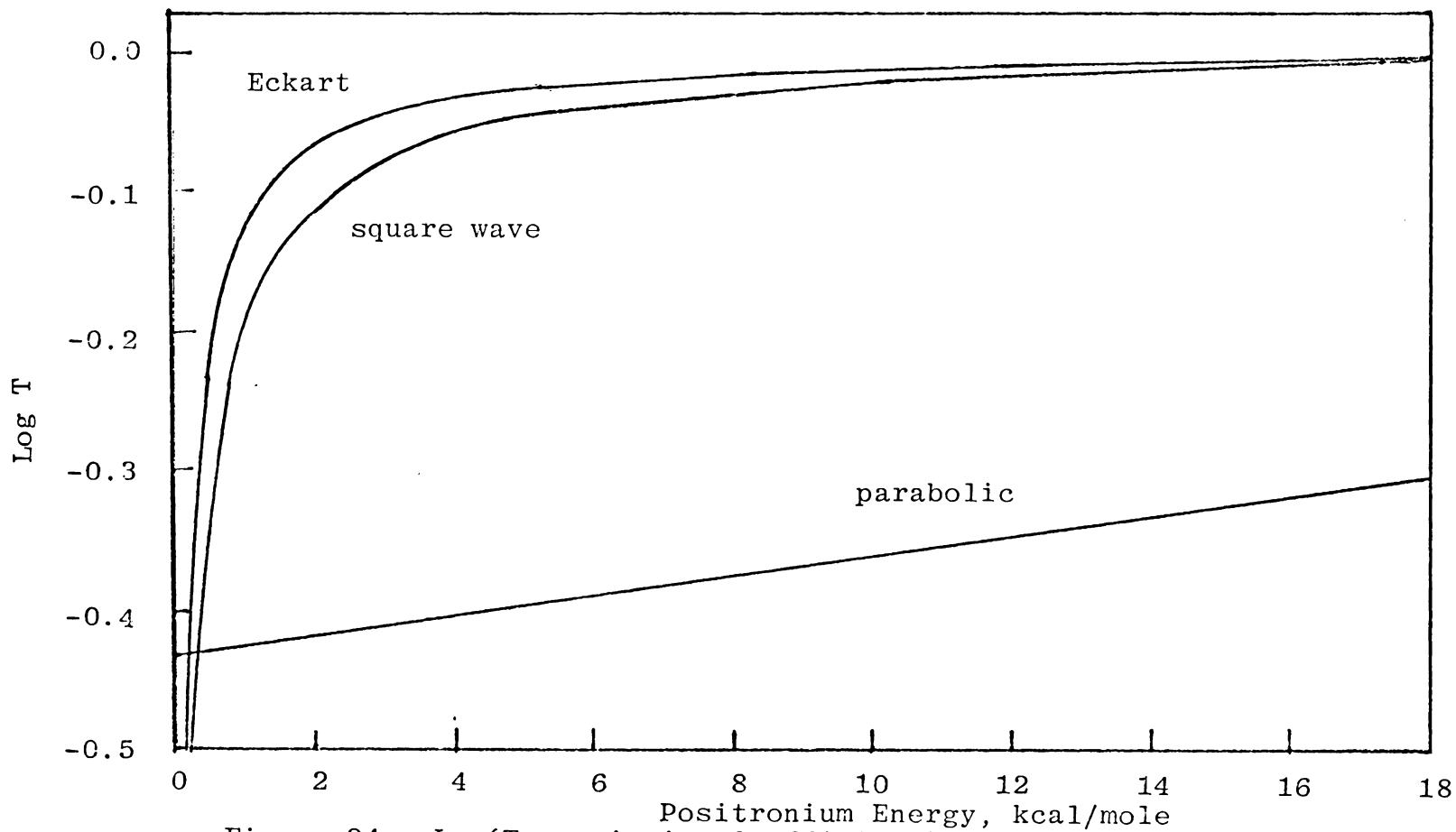
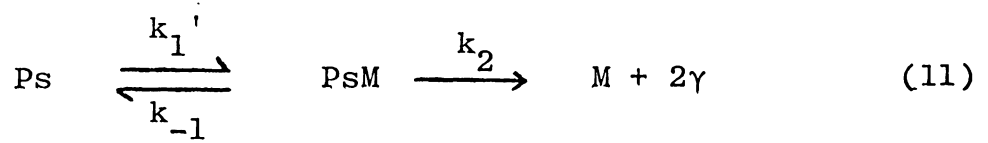


Figure 24. Log(Transmission Coefficient) versus Positronium Energy for Three Different Potential Barriers.



for $k_1' = k_1 [M]$. For this mechanism, one obtains the following equations.

$$\begin{aligned} -d[\text{Ps}]/dt &= k_1'[\text{Ps}] - k_{-1}[\text{PsM}] \\ d[\text{PsM}]/dt &= k_1'[\text{Ps}] - (k_{-1} + k_2)[\text{PsM}] \\ d[2\gamma]/dt &= k_2[\text{PsM}] \end{aligned} \quad (12)$$

These differential equations can be solved exactly^{51,52} for the time dependence of the concentration of Ps, PsM and M (or 2γ). The exact solutions are found to be functions of three parameters, S_1 , S_2 and S_3 .

$$\begin{aligned} S_1 &= (p-q)/2 \\ S_2 &= (p+q)/2 \\ S_3 &= 0 \end{aligned} \quad (13)$$

where $p = k_1' + k_{-1} + k_2$ and $q = (p^2 - 4k_1'k_2)^{1/2}$. S_3 is related to the total equilibrium solution, the time independent terms. If $S_1 \ll S_2$, there is a simple time dependence of $[\text{Ps}]$, $[\text{PsM}]$ and $[\text{M}]$ which involves only $\exp(-S_1 t)$ after a short induction period. This inequality can be met if $k_{-1}, k_2 \gg k_1'$ which is equivalent to invoking the steady-state assumption for the PsM complex. Letting $d[\text{PsM}]/dt = 0$,

$$S_1 \approx k_1'k_2/(k_{-1}+k_2) \approx [k_1k_2/(k_{-1}+k_2)] [M]$$

and

$$S_2 \approx k_1+k_2 \tag{14}$$

Therefore, if the steady-state assumption is valid, the observed rate constant for the disappearance of positronium (or the formation of 2γ) is $k_1k_2/(k_{-1}+k_2)$. In relation to Figures 22 and 23, the low temperature limiting line corresponds to a situation in which $k_2 \gg k_{-1}$ so that the reaction rate, r , becomes

$$r = k_1 [M] [Ps] \tag{15}$$

The high temperature limit corresponds to the situation in which $k_{-1} \gg k_2$ and the reaction rate changes to

$$r = k_{EQ}k_2 [M] [Ps] \tag{16}$$

where $k_{EQ} = k_1/k_{-1}$. The correct mechanism must predict that both the low and high temperature rate expressions be first order in M since this is observed experimentally.

The activation parameter obtained from the low temperature line represents the activation energy, $E_a(1)$, of the associative step, k_1 . The high temperature line can be

$$k_1 = A_1 \exp(-E_a(1)/RT) \tag{17}$$

used to measure the sum of two activation parameters, ΔH_{EQ} for the equilibrium step and the activation energy, $E_a(2)$ for the annihilation step.

$$k_2 k_{\text{EQ}} = A_2 A_1 / A_{-1} \exp -\Delta H_{\text{EQ}} - E_a(2) / RT \quad (18)$$

These are shown schematically in Figures 25 and 26.

A non-Arrhenius section is present in all the plots. The activation energy, $E_a(2)$, is always positive, therefore the results indicate that the PSM complex is more stable than the separated moieties. Chapter 4 deals primarily with this topic. Representing the barrier to annihilation as a thermodynamic barrier, as was done in Fig. 26, presented a problem that was never fully clarified.

The activation parameters for the reactions of the five nitroaromatics with positronium are listed in Table II. (The activation parameters for all reactions appear in Section E of Appendix I).

Two points can be made concerning the data taken in toluene. The activation energies for the first step were all small, approximately 1.0 kcal/mole, and fairly constant throughout. This is indicative of a fast reaction. It is interesting to note that if the second substituent on the nitrobenzene molecule was electron withdrawing, the ΔS_a (-26°C) was more unfavorable. The order being $\text{OCH}_3 > \text{H} > \text{CN} > \text{NO}_2$. The change in entropy was calculated from the average value of the activation energy, $E_a(1)$, in toluene by using the following equation:

$$k_1 = \kappa kT/h \exp(\Delta S_a/R) \exp(-E_a(1)/RT) \quad (19)$$

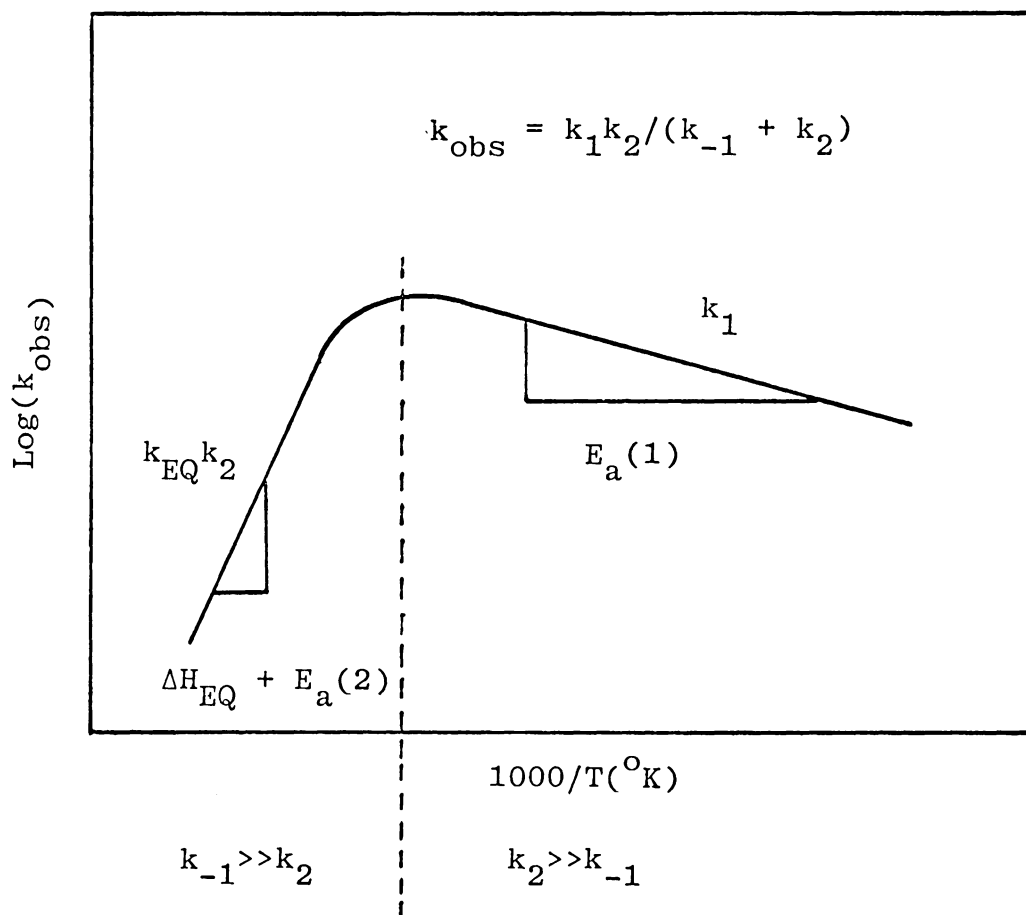


Figure 25. Typical Arrhenius Plot for the Reaction of Ortho-Positronium with Nitroaromatics Showing the Thermodynamic Parameters.

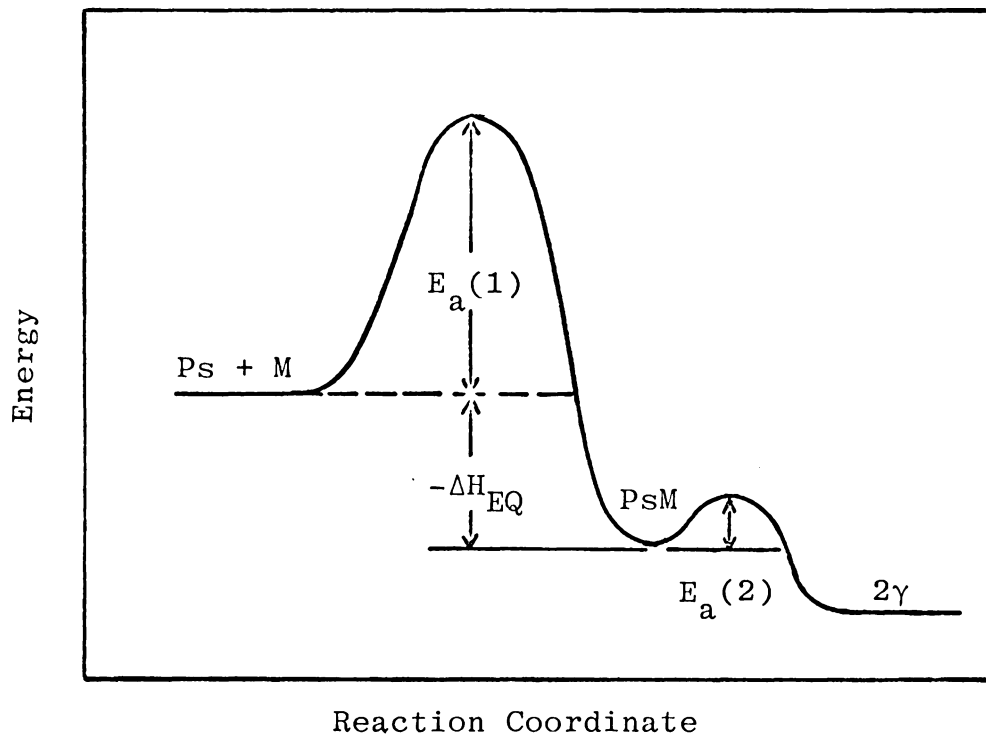


Figure 26. Potential Energy Surface for Ortho-Positronium Reacting with a Nitroaromatic Molecule.

Table II. Thermodynamic Data for the Reactions of
o-Ps in Toluene

<u>Compound</u>	Low Temperature Limit		High Temperature Limit
	$E_a(1)$, <u>kcal/mole</u>	$\Delta S_a(-26^\circ\text{C})$, <u>eu*</u>	$\Delta H_{\text{EQ}} + E_a(2)$, <u>kcal/mole</u>
p-Nitroaniline	---	---	-1.1 ± 0.1
2,6-Dimethylnitrobenzene	---	---	-2.7 ± 0.3
p-Nitroanisole	0.9 ± 0.1	-10.2	-3.6 ± 0.3
Nitrobenzene	0.8 ± 0.1	-9.7	-4.5 ± 0.3
p-Nitrobenzotrile	1.0 ± 0.1	-9.2	-6.0 ± 0.5
p-Dinitrobenzene	1.1 ± 0.1	-8.9	<-4.0
p-Benzoquinone	1.0 ± 0.1	-10.5	-7.5 ± 0.5

* Calculated using the average $E_a(1)$ value.

The concept of entropy changes as applied to positronium complexes has not been reported previously. In terms of reorganizing the positronium atom and molecule into a complex, the results would tend to indicate that the complex would more readily proceed to the transition state with molecules containing electron deficient substituents.

A qualitative analysis of the stability of the complexes can be obtained from the high temperature region of the graphs. $\Delta H_{EQ} + E_a(2)$ is always negative. This indicates that $E_a(2)$ is relatively small because even when complex formation is least likely, ΔH_{EQ} is negative but small, the sum of the two energies is still negative.

The ability of nitrobenzene to complex with positronium is probably the reason that it is the only mono-substituted benzene known to rapidly react with ortho-positronium. In order to react, the main requirements appear to be: (1) conjugation, (2) strong electron withdrawing substituents, and (3) an ability of the molecule to form strong molecular complexes.⁵³ It has been shown⁴¹, that compounds that form very strong molecular complexes, such as 7,7,8,8-tetracyanoquinodimethane (TCNQ) and tetracyanoethylene (TCNE), react at least an order of magnitude faster than nitrobenzene, which forms rather weak complexes.

It was initially thought that the inversion in the log k versus $1/T$ plots was caused by a conjugation effect as

described in the previous section. As the temperature was increased, the nitro group would begin to rotate out of the plane of the ring, and the observed rate constant would decrease due to a loss of conjugation. For this reason, the annihilation rate for a p-benzoquinone was measured as a function of temperature. The Arrhenius plot is given in Figure 27. p-Benzoquinone is known to react faster than nitrobenzene with positronium and it also readily forms molecular complexes. If the inversion was caused by a loss of conjugation, the effect should not be seen for p-benzoquinone which does not have a functional group that could rotate out of the plane of the ring. However, it behaved in a similar manner to the nitroaromatics, $E_a(1) = 1.0$ kcal/mole and $\Delta H_{EQ} + E_a(2) = -7.5$ kcal/mole, and the mechanism given in Eq. 10 was retained.

The $\log k_{obs}$ versus $1/T$ plot for 2,6-dimethylnitrobenzene is also shown in Figure 27. This compound showed unusual behavior in that no Arrhenius portion of this curve was found. Instead of inverting, the curve remains constant over a 50°C range. Either the complex is so weakly bound due to the large loss of conjugation caused by the steric-interaction between the two methyl groups and the nitro group that it decomposes rapidly or lower temperature data are needed to see the break.

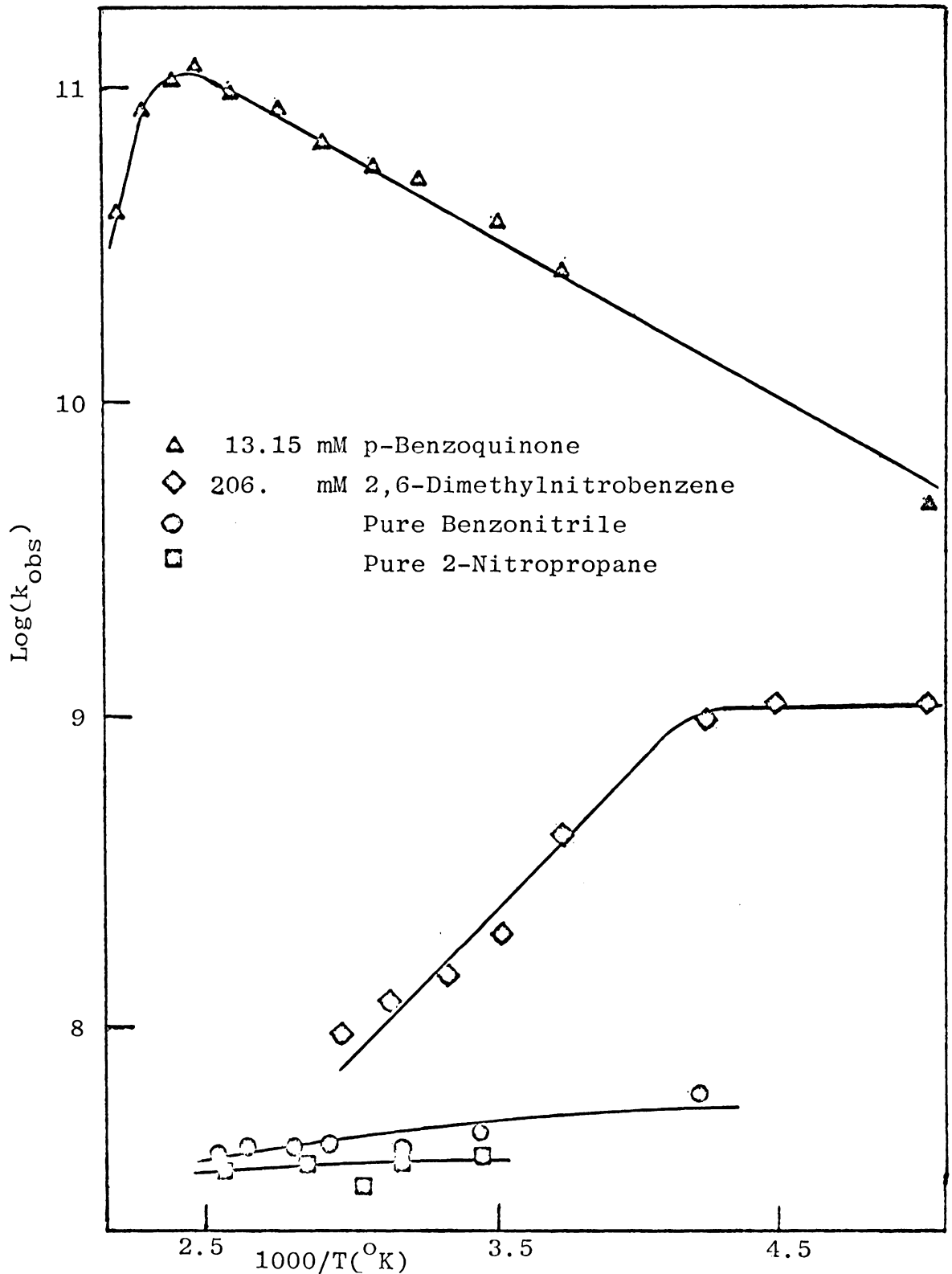


Figure 27. Arrhenius Plots for the Reaction of Ortho-Positronium with p-Benzoquinone, 2,6-Dimethylnitrobenzene, Benzonitrile and 2-Nitropropane in Toluene.

It was mentioned previously that benzonitrile and 2-nitropropane react very slowly with positronium at room temperature. Their complexes might not be stable at room temperature but data taken at other temperatures could possibly show compound formation. Figure 27 clearly shows no evidence for compound formation with these two compounds in the temperature range studied.

D. Solvent Effects

Similar Arrhenius data were taken using n-heptane and n-pentanol as the solvents. The temperature dependence of the solvent annihilation rate was measured and can be found in Section D of Appendix I. Only three nitroaromatic compounds were used to study the effect of the solvent on these reactions; nitrobenzene, p-nitroanisole, and p-dinitrobenzene. The Arrhenius plots and thermodynamic data appear in Figures 28 and 29 and in Table III respectively. Similar trends were found in these solvents. Figure 30 shows a comparison of the Arrhenius plots for the reaction of ortho-positronium with nitrobenzene in all three solvents. Since an inversion was seen in the Arrhenius plot obtained in n-heptane, the possibility of a solvent-solute species being responsible for the rapid reaction was now eliminated. Nitroaromatics are not known to form molecular complexes with alkanes.

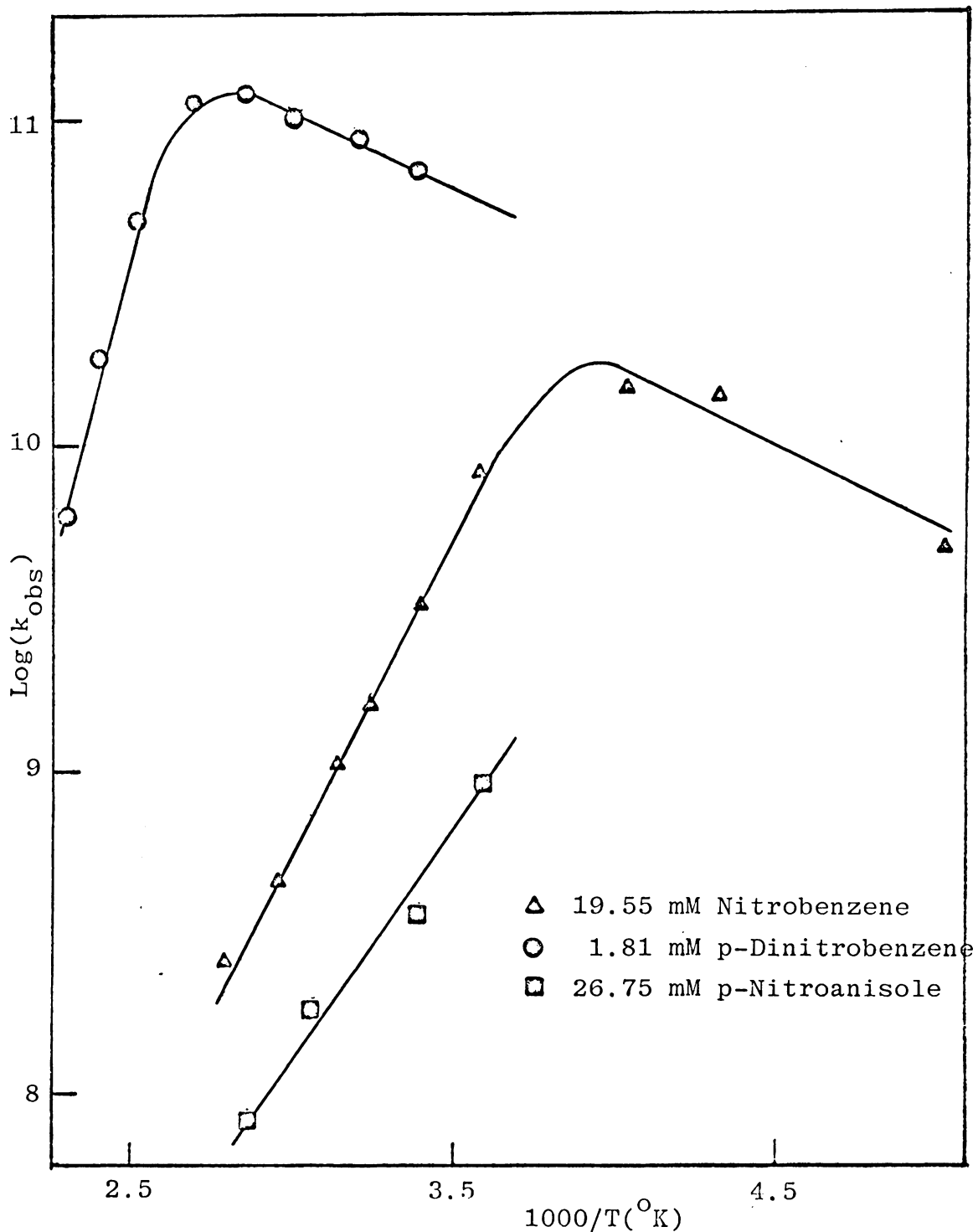


Figure 28. Arrhenius Plots for the Reaction of Ortho-Positronium with Nitrobenzene, p-Dinitrobenzene and p-Nitroanisole in n-Heptane.

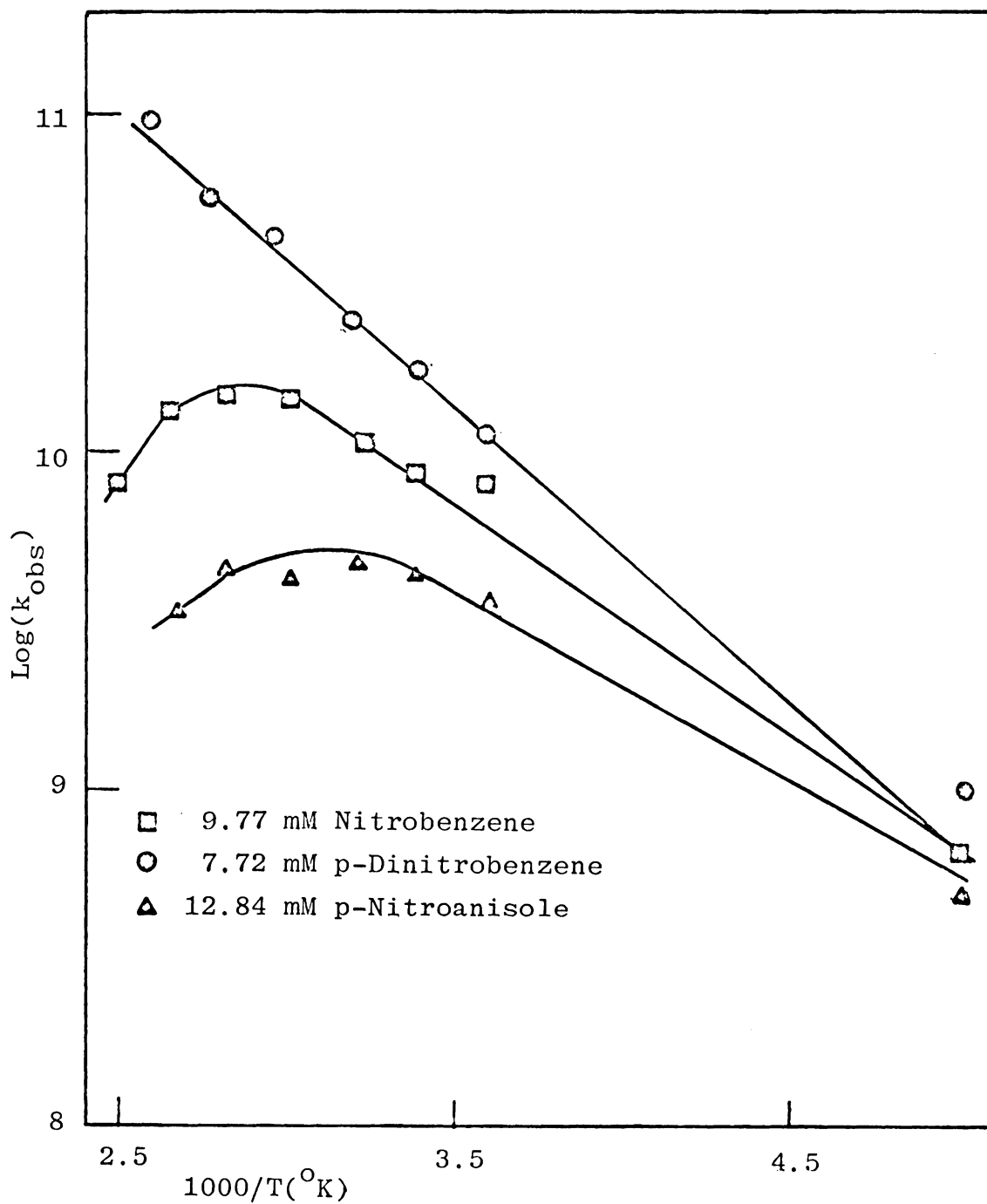


Figure 29. Arrhenius Plots for the Reaction of Ortho-Positronium with Nitrobenzene, p-Dinitrobenzene and p-Nitroanisole in n-Pentanol.

Table III. Thermodynamic Data for the Reactions of
o-Ps in n-Heptane and n-Pentanol

<u>Compound</u>	Low Temperature Limit		High Temperature Limit
	<u>$E_a(1)$,</u> <u>kcal/mole</u>	<u>$\Delta S_a(-26^\circ\text{C})$,</u> <u>eu*</u>	<u>$\Delta H_{\text{EQ}} + E_a(2)$,</u> <u>kcal/mole</u>
Solvent: n-Heptane			
p-Nitroanisole	---	---	-2.9 ± 0.4
Nitrobenzene	1.3 ± 0.2	-8.1	-4.4 ± 0.4
p-Dinitrobenzene	1.0 ± 0.2	-7.1	-8.0 ± 1.0
Solvent: n-Pentanol			
p-Nitroanisole	1.2 ± 0.1	-7.3	-1.3 ± 0.8
Nitrobenzene	1.3 ± 0.2	-6.2	-3.5 ± 0.6
p-Dinitrobenzene	1.8 ± 0.3	-5.5	---

* Calculated using the average $E_a(1)$ value.

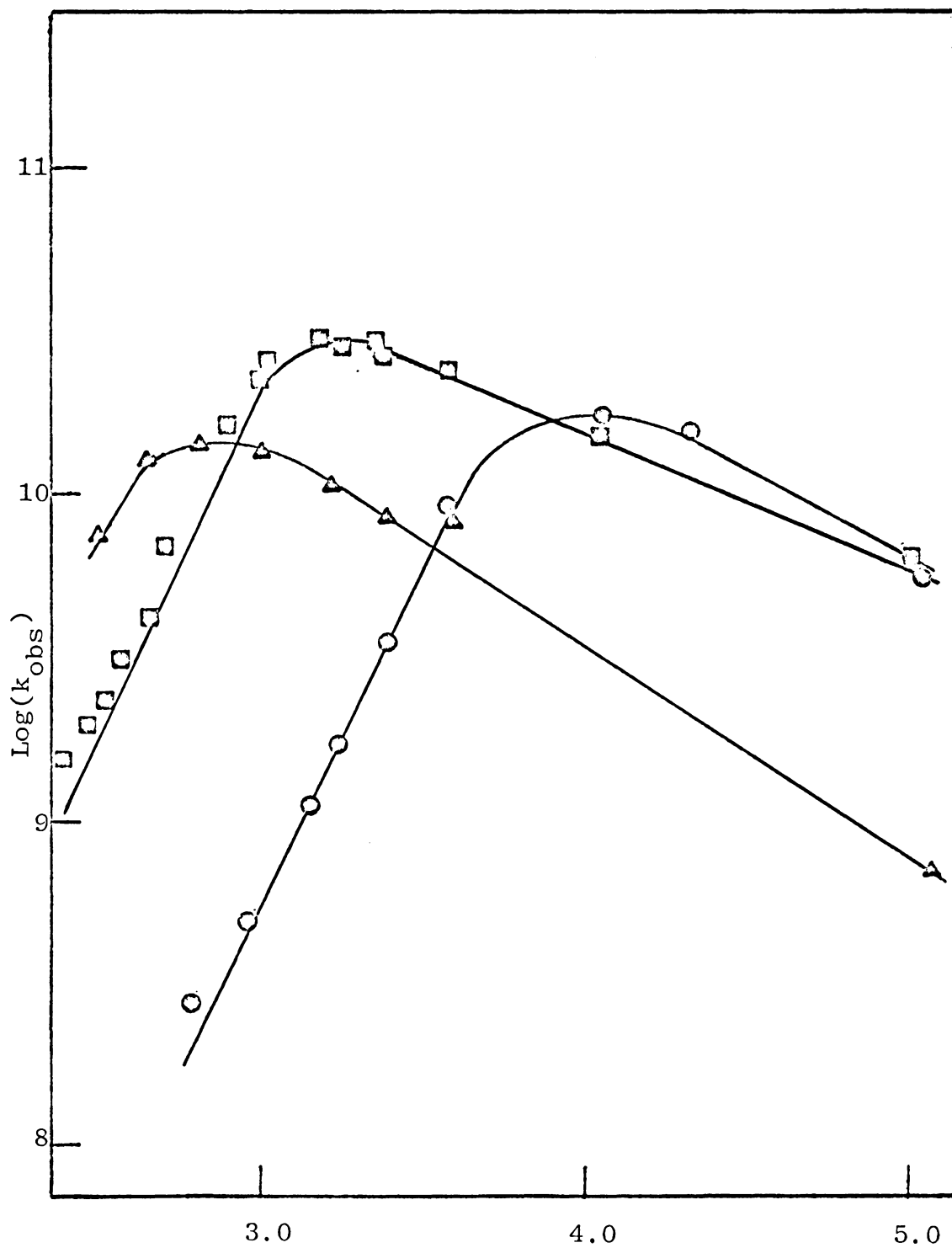
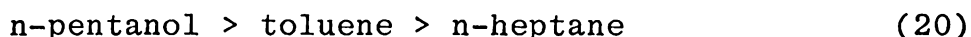


Figure 30. Arrhenius Plots for the Reaction of Ortho-Positronium with Nitrobenzene in (□) Toluene, (○) n-Heptane and (△) n-Pentanol.

The point of inversion in the curves is determined by the slopes of the two limiting lines and the pre-exponential factor, A , for each observed rate constant. These are A_1 for the low temperature limiting line and A_1A_2/A_{-1} for the high temperature limiting line. According to absolute rate theory, these factors contain two terms that could possibly affect the point of inversion. The first is the transmission coefficient, κ , and the second is the entropy change. The inversion point in the nitrobenzene curves of Figure 30 can be used to determine the ability of each solvent to assist in reaching the transition state before forming the PsM complex. For nitrobenzene and the other two compounds, this ability increased following the order



This is the order expected when neutral species combine to form charged or polar complexes.⁵³ The more polar the solvent, the more able it is to neutralize the charge separation.

E. Correlations

Many correlations were attempted before the temperature effect of Section C was discovered.⁵⁴⁻⁵⁶ Numerous rate constants for substituted nitrobenzene compounds were measured in benzene at 295°K. The log of the relative rate constants (relative to nitrobenzene) was correlated to a

variety of molecular parameters that included the energy of the lowest unoccupied molecular orbital (LUMO), the net π charge density of the aromatic ring, free valence, radical localization energy, etc. The results were moderately good considering that the observed rate constants were for different combinations of k_1 , k_{-1} and k_2 .

The reactivity of these compounds toward positronium atoms at room temperature was correlated by use of the Hammett equation for electrophilic and nucleophilic attack.⁵⁷ Again, the observed rate constants could be k_1 , $k_1 k_2 / (k_{-1} + k_2)$, or $k_{EQ} k_2$ depending on where the inversion point occurred in the Arrhenius plot. An example of this can be found in Figure 23. At room temperature, the observed rate constant for p-nitroaniline is $k_{EQ} k_2$, for p-nitroanisole it is $k_1 k_2 / (k_{-1} + k_2)$, and for p-dinitrobenzene it is k_1 . However, it was found that around -25°C , most measured rate constants are k_1 . Figure 31 is a Hammett plot for nucleophilic attack⁵⁸ using only k_1 rate constants as measured in toluene at -25°C .⁵⁹ These preliminary results are encouraging and further work at this low temperature seems warranted.

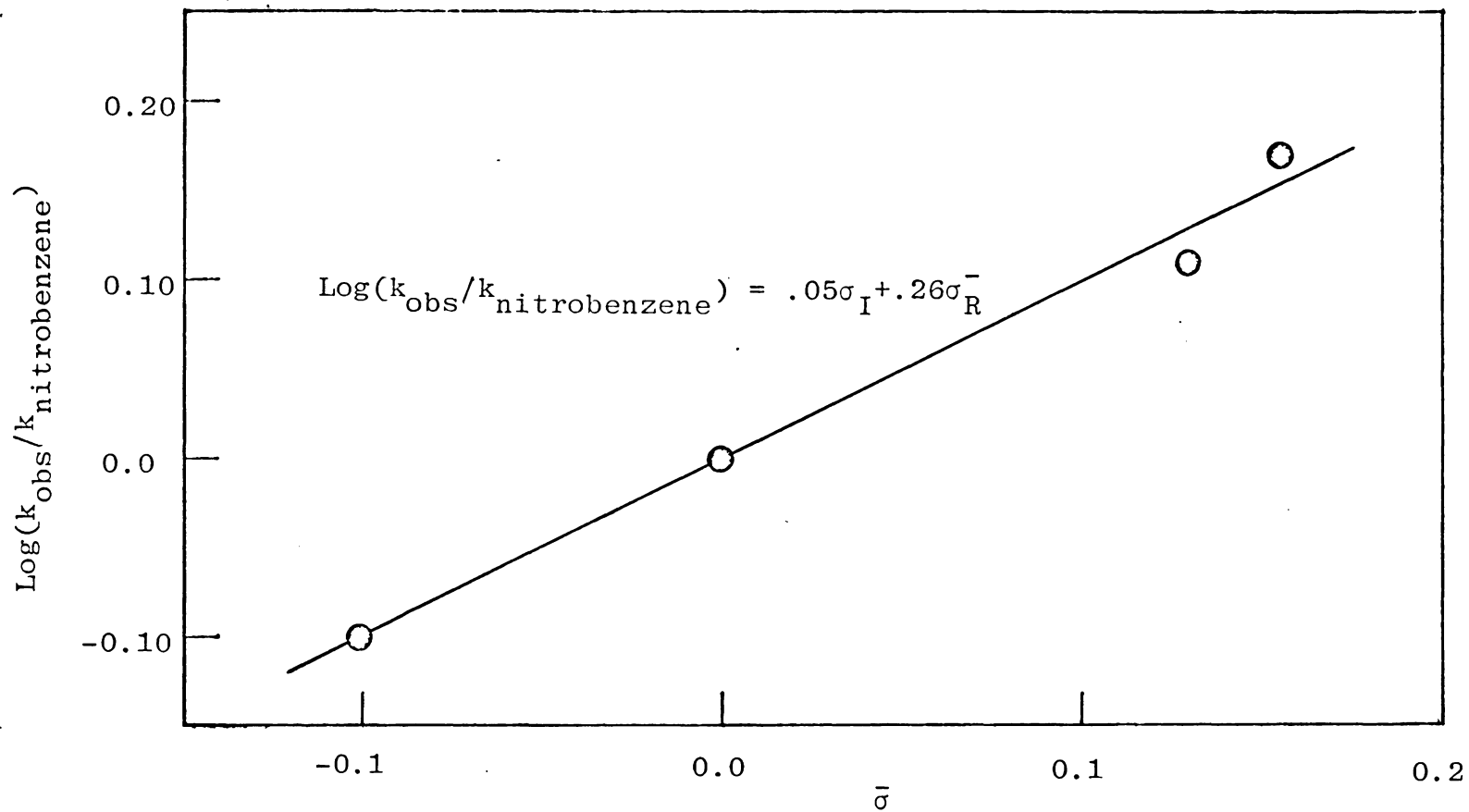


Figure 31. Hammett Plot for the Reaction of Ortho-Positronium with Substituted Nitroaromatics in Toluene at -26°C .

POSITRON AND POSITRONIUM COMPLEXES⁶⁰

A. Introduction

There is now a considerable amount of evidence available both from this work as well as lifetime and angular correlation experiments to indicate the existence of stable positron and positronium complexes with atoms, ions, and molecules.^{8,26,61-65} Theoretical calculations of the stability of positron complexes with atoms and atomic ions, and of the positron lifetimes in these complexes have generally been rewarding.^{61,66-69} The most extensive theoretical work in this area has been done by Schrader and co-workers: a 1970 paper⁷⁰ provides a general non-relativistic self-consistent-field theory of one-positron, many electron systems and explains the approximations involved in most of the earlier work; a more recent paper describes a variational treatment of e^+H^- which incorporates interparticle coordinates in the wavefunction.⁶⁹ The most accurate result of Navin, Schrader, and Lebeda⁶⁹ indicates that e^+H^- is 0.794 eV more stable than separated positronium and hydrogen atoms. This means that the positron binding energy (positron affinity) of the hydride ion is on the order of 6.8 eV. Earlier calculations indicated positron affinities on the order of 2 eV for Cl^- and alkali atoms.^{61,67}

This chapter is concerned with the existence of positron complexes with relatively large organic compounds. For large

compounds, containing many electrons, a variational treatment similar to that for e^+H^- is presently intractable, and it is necessary to retreat to some form of the self-consistent-field theory.⁷⁰ We take the natural approach of employing the Born-Oppenheimer separation and the Hartree-Fock approximation. The former implies that we treat the nuclei as fixed point charges, and the latter, that we treat the positron on an equal basis with the electrons, calculating only average Coulombic interactions between all light particles. In general this approach should give reasonable results. The necessary corrections for more accurate calculations will arise from short-range interparticle potentials, and should be amenable to treatment by an extension of Sinanoglu's theory of purely electronic correlation effects.⁷¹ The positron-electron correlation effects are expected to be extremely important in the prediction of positron annihilation rates, so we have not tried to calculate these; we are concerned here only with the stability of the positron and positronium complexes.

We make use of the approximate all-valence-electron self-consistent-field theory developed by Pople and co-workers,⁷² the so-called CNDO/2 approximation. We seek positronic and electronic molecular orbitals and express both as linear combinations of the same set of Slater-type atomic orbitals.^{72,73} This has the advantage of providing a basis for estimating the necessary integrals over positron

coordinates without a great deal of computation. For all electronic matrix elements, the CNDO/2 approximations and the parameterization suggested by Pople and Beveridge were retained. The positronic matrix elements of the core Hamiltonian were estimated by comparison with the corresponding electronic integrals. The latter were approximately decomposed into kinetic and potential energy parts using the virial theorem. A sign change of the potential part provided the approximation for the one-center positron matrix elements. For the two-center core matrix elements, these considerations were combined with the Wolfsberg-Helmholtz approximation⁷⁴ and the formulation given by Cusachs and Cusachs.⁷⁵ All Coulomb integrals were evaluated using the CNDO/2 approximations, but accounting of course for the positive charge of the positron. It is recognized that these procedures only provide first approximations. In particular, the virial theorem is not strictly applicable to the electronic integrals involving Slater-type orbitals, and the Slater-type orbitals provide improper cusp behavior for the positron wavefunctions.

The calculations in this chapter deal primarily with positron complexes with closed shell molecules. By making use also of the lowest unoccupied electronic molecular orbital, we obtain a first approximation for the binding energies of positron complexes with molecular anions.

(These are alternatively referred to as positronium-molecule complexes).

B. Hartree-Fock Theory

The total Hamiltonian, in atomic units, for a one-positron 2n-electron system is given by

$$H = \sum_{\mu} h_{\mu}^{\text{core}(-)} + \sum_{\mu < \nu} g_{\mu\nu} + h_p^{\text{core}(+)} + \sum_{\mu} g_{\mu p}, \quad (1)$$

where (using atomic units)

$$h_{\mu}^{\text{core}(-)} = -\nabla_{\mu}^2/2 - \sum_{\alpha} Z_{\alpha}/r_{\mu\alpha}$$

$$h_p^{\text{core}(+)} = -\nabla_p^2/2 + \sum_{\alpha} Z_{\alpha}/r_{p\alpha}$$

$$g_{\mu\nu} = r_{\mu\nu}^{-1}$$

$$g_{\mu p} = -r_{\mu p}^{-1}$$

The subscripts μ and ν are used to indicate electron terms, p indicates positron terms, and α refers to nuclei. The superscripts $(-)$ and $(+)$ will be used with electron and positron terms respectively for clarity. The core Hamiltonian terms include the kinetic energies and Coulombic interactions with all nuclei of the molecule. For a system containing 2n electrons, and one positron, the wave function is

$$\psi = \psi_p(p) [(2n)!]^{-1/2} \det |\psi_1(1) \psi_2(2) \cdots \psi_{2n}(2n)|, \quad (2)$$

where each function, ψ_i , is a product of a molecular orbital, ϕ_i , and a spin wavefunction, either α or β . The positron spin is arbitrary. Since we are dealing with a closed-shell electronic system, each occupied electronic orbital occurs twice, being associated once with spin function α and once with β . We make the usual assumptions of orthonormality:

$$\langle \phi_p | \phi_p \rangle = 1; \langle \phi_i | \phi_j \rangle = \delta_{ij} . \quad (3)$$

The energy expectation value is

$$\begin{aligned} E = \langle \psi | H | \psi \rangle = & \langle \psi | \sum_{\mu} h_{\mu}^{\text{core}(-)} | \psi \rangle + \langle \psi | \sum_{\mu < \nu} g_{\mu\nu} | \psi \rangle \\ & + \langle \psi | h_p^{\text{core}(+)} | \psi \rangle + \langle \psi | \sum_{\mu} g_{\mu p} | \psi \rangle . \quad (4) \end{aligned}$$

Using the fact that the electronic part of ψ consists of n doubly occupied orbitals, this can be written in terms of summations over those molecular orbitals,

$$\begin{aligned} E = & 2 \sum_i^n H_{ii}^{(-)} + 2 \sum_{i < j}^n (2J_{ij}^{(-)} - K_{ij}^{(-)}) + H_{pp}^{(+)} \\ & + 2 \sum_i^n J_{pi}^{(+)} \quad (5) \end{aligned}$$

Here, $H_{pp}^{(+)}$ and $J_{pi}^{(+-)}$ are respectively the positron-core and positron-electron contributions,

$$H_{pp}^{(+)} = \langle \phi_p | h_p^{\text{core}(+)} | \phi_p \rangle, \quad (6)$$

$$J_{pi}^{(+)} = - \langle \phi_p(p) \phi_i(i) | r_{ip}^{-1} | \phi_p(p) \phi_i(i) \rangle, \quad (7)$$

and $H_{ii}^{(-)}$, $J_{ij}^{(-)}$, and $K_{ij}^{(-)}$ are the purely electronic core interactions, Coulomb and exchange integrals involving the molecular orbitals indicated by the subscripts. The molecular orbitals are expanded as linear combinations of the atomic valence-shell orbitals $\{\chi_\mu\}$,

$$\phi_i = \sum_{\mu}^m \chi_{\mu} C_{\mu i}, \quad \phi_p = \sum_{\mu}^m \chi_{\mu} C_{\mu p}. \quad (8)$$

If we introduce charge-density and bond-order matrices for both electrons and the positron,

$$P_{\mu\nu}^{(-)} = 2 \sum_i^n C_{\mu i}^* C_{\nu i}, \quad P_{\mu\nu}^{(+)} = C_{\mu p}^* C_{\nu p}, \quad (9)$$

the energy can be written in terms of integrals over atomic orbitals:

$$\begin{aligned} E = & \sum_{\mu, \nu}^m P_{\mu\nu}^{(-)} H_{\mu\nu}^{(-)} \\ & + (1/2) \sum_{\mu, \nu, \lambda, \sigma}^m P_{\mu\nu}^{(-)} P_{\lambda\sigma}^{(-)} [\langle \chi_{\mu}(1) \chi_{\lambda}(2) | r_{12}^{-1} | \chi_{\nu}(1) \chi_{\sigma}(2) \rangle \\ & - (1/2) \langle \chi_{\mu}(1) \chi_{\lambda}(2) | r_{12}^{-1} | \chi_{\sigma}(1) \chi_{\nu}(2) \rangle] + \sum_{\mu, \nu}^m P_{\mu\nu}^{(+)} H_{\mu\nu}^{(+)} \\ & - \sum_{\mu, \nu, \lambda, \sigma}^m P_{\nu\sigma}^{(-)} P_{\mu\lambda}^{(+)} \langle \chi_{\mu}(p) \chi_{\nu}(1) | r_{1p}^{-1} | \chi_{\lambda}(p) \chi_{\sigma}(1) \rangle. \quad (10) \end{aligned}$$

Again, the first two terms are identical with those arising in the simple closed shell electronic case,⁷² and the latter two are caused by the positron,

$$H_{\mu\nu}^{(+)} = \langle \chi_{\mu}(p) | h_p^{\text{core}(+)} | \chi_{\nu}(p) \rangle \quad (11)$$

Minimizing the energy with respect to variations in the electronic and positronic orbitals leads to two sets of Hartree-Fock equations,

$$\begin{aligned} F^{(-)} \phi_i &= \epsilon_i \phi_i \\ F^{(+)} \phi_p &= \epsilon_p \phi_p \quad , \end{aligned} \quad (12)$$

where the electron and positron Fock operators are defined as

$$F^{(-)} = h^{\text{core}(-)} + \sum_j (2 J_j^{(-)} - K_j^{(-)}) + J_p^{(+)} \quad (13)$$

and

$$F^{(+)} = h^{\text{core}(+)} + 2 \sum_j J_j^{(+)} .$$

J_j and K_j are the usual electronic Coulomb and exchange operators associated with the j th molecular orbital; $J_p^{(+-)}$ is the positronic Coulomb operator,

$$J_p^{(+)}(\underline{r}_i) = - \int \phi_p^*(\underline{r}_p) [1/r_{ip}] \phi_p(\underline{r}_p) d\tau_p ;$$

and $J_j^{(+-)}$ is the operator providing the Coulomb potential seen by a positron due to an electron in the j th orbital,

$$J_j^{(+)}(\underline{r}_p) = - \int \phi_j^*(\underline{r}_i) [1/r_{ip}] \phi_j(\underline{r}_i) d\tau_i .$$

For computational purposes, it is convenient to express the energy as

$$E = (1/2) \sum_{\mu, \nu} P_{\mu\nu}^{(-)} (H_{\mu\nu}^{(-)} + F_{\mu\nu}^{(-)}) + (1/2) \sum_{\mu, \nu} P_{\mu\nu}^{(+)} (H_{\mu\nu}^{(+)} + F_{\mu\nu}^{(+)}) , \quad (14)$$

where all matrix elements are in the basis set of atomic orbitals. It should also be noted that in forming the wavefunction (2) from the self-consistent solutions to Eqs. (12) we consistently employ the Aufbau Principle. That is, the $2n$ electrons are put into the n lowest lying electronic orbitals, and the positron is assumed to occupy the positron orbital of lowest orbital energy, ϵ_p .

C. Approximations and Parameterization

By using the zero-differential overlap approximation, the multi-centered interaction integrals are simplified so that

$$\langle \chi_\mu \chi_\nu | r_{12}^{-1} | \chi_\lambda \chi_\alpha \rangle = \langle \chi_\mu \chi_\nu | r_{12}^{-1} | \chi_\mu \chi_\nu \rangle \delta_{\mu\lambda} \delta_{\nu\alpha} = \gamma_{\mu\nu} \delta_{\mu\lambda} \delta_{\nu\alpha} \quad (15)$$

and the Fock matrix elements become

$$\begin{aligned}
F_{\mu\mu}^{(+)} &= H_{\mu\mu}^{(+)} - \sum_{\lambda} P_{\lambda\lambda}^{(-)} \gamma_{\mu\lambda} \\
F_{\mu\nu}^{(+)} &= H_{\mu\nu}^{(+)} \quad (\mu \neq \nu) \\
F_{\mu\mu}^{(-)} &= H_{\mu\mu}^{(-)} - (1/2) P_{\mu\mu}^{(-)} \gamma_{\mu\mu} + \sum_{\lambda} P_{\lambda\lambda}^{(-)} \gamma_{\mu\lambda} \\
&\quad - \sum_{\lambda} P_{\lambda\lambda}^{(+)} \gamma_{\mu\lambda} \\
F_{\mu\nu}^{(-)} &= H_{\mu\nu}^{(-)} - (1/2) P_{\mu\nu}^{(-)} \gamma_{\mu\nu} \quad (\mu \neq \nu) .
\end{aligned} \tag{16}$$

One of the CNDO approximations is to reduce Coulomb integrals to one per atom pair, i.e., $\gamma_{\mu\nu} = \gamma_{AB}$, where χ_{μ} is located on atom A and χ_{ν} on atom B.

It should be noted that the electron Fock matrix contains a contribution of the form, $-P_{\mu\nu}^{(-)} \gamma_{\mu\nu}/2$ and that this term is absent in the corresponding $F_{\mu\nu}^{(+)}$. This exchange term arises from the two-electron repulsions; since our system contains only one positron, this term does not appear in the positron matrix.

Eqs. (12) can now be solved by iterative techniques and the electron and positron wave functions can be obtained in which the electronic orbitals have been altered by the presence of the positron, and the positron orbital has been altered by the electrons.

The diagonal positron core matrix elements are approximated as

$$H_{\mu\mu}^{(+)} = U_{\mu\mu}^{(+)} + \sum_{B \neq A} Z_B \gamma_{AB} ; \quad (\chi_{\mu} \text{ on atom A}) \tag{17}$$

$U_{\mu\mu}^{(+)}$ is an atomic term containing the kinetic energy and potential energy of the positron in the field of the Ath core; the second term describes the positron interactions with the other atomic cores.

The CNDO/2 parameterization as described by Pople and Beveridge⁷³ is used for all electron terms. However, for the positron no data such as ionization potentials and positron affinities are available. In order to parametrize the positron Fock terms we relate them to the corresponding electronic integrals as follows. For kinetic energy (T) and potential energy (V) matrix elements,

$$T_{\mu\nu}^{(+)} = T_{\mu\nu}^{(-)} \quad \text{and} \quad V_{\mu\nu}^{(+)} = - V_{\mu\nu}^{(-)} \quad (18)$$

so that

$$U_{\mu\mu}^{(+)} = T_{\mu\mu}^{(+)} + V_{\mu\mu}^{(+)} = T_{\mu\mu}^{(-)} - V_{\mu\mu}^{(-)} \quad (19)$$

Applying the virial theorem,

$$T_{\mu\mu}^{(-)} = - (1/2) V_{\mu\mu}^{(-)} = - U_{\mu\mu}^{(-)} \quad (20)$$

We see that

$$U_{\mu\mu}^{(+)} = - 3 U_{\mu\mu}^{(-)} \quad (21)$$

Treatment of the off-diagonal core terms is not as direct. Because of the empirical nature of the bonding parameters, $\beta_{AB}^{o(-)}$, used by Pople and Beveridge, no direct conversion to

$\beta_{AB}^{O(+)}$ was found. Therefore, we decided to combine the Wolfsberg-Helmholtz approximation⁷⁴ with the ideas of Cusachs and Cusachs.⁷⁵ For the electronic core Hamiltonian matrix, we have

$$H_{\mu\nu}^{(-)} = T_{\mu\nu}^{(-)} + V_{\mu\nu}^{(-)} \quad (22a)$$

$$\approx k^{(-)} S_{\mu\nu} (H_{\mu\mu}^{(-)} + H_{\nu\nu}^{(-)})/2 \quad (22b)$$

$$\approx S_{\mu\nu} (2 - |S_{\mu\nu}|) (H_{\mu\mu}^{(-)} + H_{\nu\nu}^{(-)})/2 . \quad (22c)$$

The analogous expressions for the positron are

$$H_{\mu\nu}^{(+)} = T_{\mu\nu}^{(+)} + V_{\mu\nu}^{(+)} \quad (23a)$$

$$\approx k^{(+)} S_{\mu\nu} (H_{\mu\mu}^{(+)} + H_{\nu\nu}^{(+)})/2 \quad (23b)$$

The terms in Eq. (23a) are now treated by the method of Cusachs and Cusachs:

$$T_{\mu\nu}^{(+)} \approx S_{\mu\nu} |S_{\mu\nu}| (T_{\mu\mu}^{(+)} + T_{\nu\nu}^{(+)})/2,$$

$$V_{\mu\nu}^{(+)} \approx S_{\mu\nu} (V_{\mu\mu}^{(+)} + V_{\nu\nu}^{(+)})/2,$$

and, substituting Eqs. (18) into the analogue of Eq. (19), yields

$$H_{\mu\mu}^{(+)} = (3/2) V_{\mu\mu}^{(+)} = 3T_{\mu\mu}^{(+)} .$$

This gives directly

$$H_{\mu\nu}^{(+)} \approx (1/3) S_{\mu\nu} (2 + |S_{\mu\nu}|) (H_{\mu\mu}^{(+)} + H_{\nu\nu}^{(+)}) / 2 . \quad (23c)$$

Values typically used⁷⁶ for the Wolfsberg-Helmholtz $k^{(-)}$ in Eq. (22b) range from 1.75 to 2.0, and comparison with (22c) shows that this corresponds to a range of $|S_{\mu\nu}|$ between 0 and 0.25. When these values are employed in Eq. (23c) and the result equated to (23b), it appears that $k^{(+)}$ should be in the range 2/3 to 3/4.

A computer program, PSCF, was developed by modifying the CNDO/2 program listed by Pople and Beveridge.⁷³ A listing of the pertinent subroutines is given in Section E of Appendix II. Along with the other approximations already described, we employed Eq. (23b) for $H_{\mu\nu}^{(+)}$ and varied $k^{(+)}$ over the indicated range.

D. Results

Using this formalism, we have studied the stability of positron and positronium complexes of benzene and a number of substituted benzene derivatives. In the process of carrying out the iterations to self-consistency, two slight problems arose. A number of calculations, especially with the larger values of $k^{(+)}$, oscillated rather than converging directly to an energy minimum. The oscillations were successfully eliminated by taking the modified density matrices and averaging the new calculated matrices with the old ones

on each iteration. Also, for some of the symmetric molecules, the symmetry was lost during the iteration process. We believe that this resulted from round-off errors in the positron density matrix, where the elements are generally quite small. To eliminate the need for symmetry projection, we restricted our solutions to the symmetric cases by symmetrizing the molecular orbitals at each iteration.

The important results are listed in Tables IV and V for the two limiting values of $k^{(+)}$ mentioned previously. All calculations indicated that the positron should be bound by these molecules. Two approximations for the positron binding energies are given: $(-\epsilon_p)$ is the negative of the positron orbital energy, and provides a Koopmans' theorem⁷⁷ estimate; the other entry (B.E.) was a direct calculation of the binding energy as the difference between the total energy of the neutral molecule and that of the positron-molecule complex. The two estimates are fairly comparable to one another, but the former is always of greater magnitude than the latter. The calculated positron binding energy depends critically on the value used for $k^{(+)}$. This is illustrated graphically in Fig. 32, which shows the two estimates of the positron affinity of nitrobenzene as a function of $k^{(+)}$. Corresponding to this, the number of bound positronium orbitals also increased with $k^{(+)}$. For the majority of the complexes, we found one or two bound orbitals for $k^{(+)} = 2/3$, and six to eight for $k^{(+)} = 3/4$.

Table IV. Positron and Positronium Binding Energies,
a.u. (C) $k^{(+)} = 2.3$

Compound	$-\epsilon_p$	B.E. (e^+)	e^+ site ^(a)	B.E. (Ps)
Benzene	0.0082	0.0081	C-H	-0.2197
Toluene	0.1877	0.1427	Alpha C-H	-0.0608
Fluorobenzene	0.2083	0.1030	Para C-H	-0.0809 ^(b)
Aniline	0.3375	0.2736	N-H	+0.1184 ^(b)
Benzonitrile	0.0780	0.0371	Meta C-H	-0.1726
Benzaldehyde	0.2256	0.1234	Carbonyl C-H	-0.0002
Nitrobenzene	0.0608	0.0200	Meta C-H	-0.1512
p-Benzoquinone	0.0213	0.0131	C-H	-0.0232

(a) greater than 90% of the e^+ charge density occurs here

(b) LUMO of the positronium complex is a σ orbital; in all other cases it is a π orbital.

(c) 1 a.u. = 27.21 eV = 627.71 kcal/mole.

Table V, Positron and Positronium Binding Energies,
 a.u. (C) $k^{(+)} = 3/4$

Compound	$-\epsilon_p$	B.E. (e^+)	e^+ site ^(a)	B.E. (Ps)
Benzene	+0.2840	+0.2831	C-H	+0.0597
Toluene	+0.4717	+0.4270	Alpha C-H	+0.2225
Fluorobenzene	+0.4426	+0.3451	Para C-H	+0.1582 ^(b)
Aniline	+0.6528	+0.5890	N-H	+0.4339 ^(b)
Benzonitrile	+0.3671	+0.3194	C≡N	+0.1573
Benzaldehyde	+0.4686	+0.3757	Carbonyl C-H	+0.2533
Nitrobenzene	+0.4268	+0.3250	Para C-H	+0.1740
p-Benzoquinone	+0.3137	+0.3084	C-H	+0.2802

(a) greater than 90% of the e^+ charge density occurs here

(b) LUMO of the positronium complex is a σ orbital; in all cases it is a π orbital.

(c) 1 a.u. = 27.21 eV = 627.71 kcal/mole.

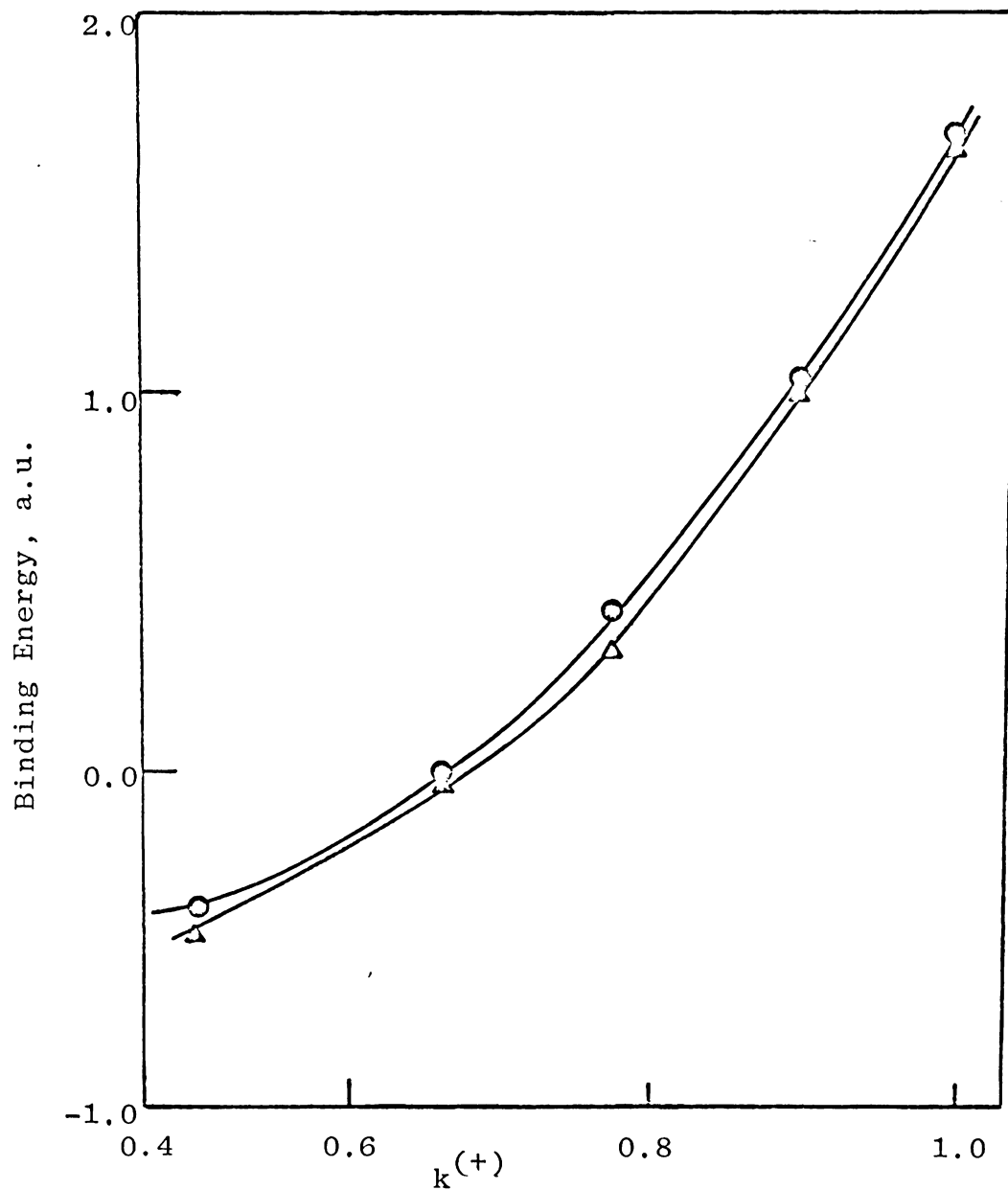


Figure 32. Variation in the Positron Binding Energy in Nitrobenzene with $k^{(+)}$. Koopman's Theorem Estimate (○) and Change in Energy Directly Calculated (△) for the Complex.

Also indicated in Tables IV and V are the locations of maximum positron density found for the several complexes. The positron was always found to be essentially localized in a single bond (or in equivalent bonds when this was required by symmetry). The position at which the positron was found varies widely through the series of molecules, and for some cases the positron's location was seen to change as $k^{(+)}$ was varied. We have not been able so far to reconcile these findings consistently with the results of Chapter 3. For example, if the complex formation is considered to be an example of nucleophilic attack, one might expect the positron to locate in the para position of nitrobenzene. This was observed for $k^{(+)} = 3/4$, but when $k^{(+)}$ was taken as $2/3$, the positron moved to the meta position. Similar shifts were observed for benzonitrile, but for the larger $k^{(+)}$ the positron localized on the cyanide group. A number of additional points could be made, but at the present level of approximation extended discussion is probably not justified.

It is of some interest to note that the positron was always found to be most stable in a sigma rather than a pi molecular orbital. This is in fact expected if one looks at the positron's Fock matrix elements. In our approximations, all diagonal elements, $F_{\mu\mu}^{(+)}$, are positive, but the minimum values always occur for the hydrogen 1s orbitals. This situation could easily be expected to change if different approximations are employed.

Addition of the positron to a molecule is of course expected to cause some shifting of electronic charge density. As a typical case, we list in Table VI the calculated total atomic charge densities for the aniline molecule and its positron complex. The atomic charge density is obtained by summing the diagonal elements of the electronic charge-density and bond-order matrix (see Eq. (9)) over all the valence orbitals for each atom. The same data are shown graphically in slightly different form in Fig. 33. This shows profiles of net pi charge and net total charge before and after positron addition. The net total charge is the difference between the calculated charge density and the atomic core charge. The net pi charge is the pi-orbital density minus one. The figure is intended to show electron build-up at the various atoms; the numbers between the atoms are not meaningful. In all cases, the electron density is enhanced in the vicinity of the positron, as expected from the mutual attraction. The largest changes seen in the electron density upon positron addition were about 0.1 e, and most changes were appreciably less.

The positron densities listed in Table VI illustrate the high degree of localization found for the positron as indicated earlier. In this case, more than 99% of the total positron density is located in the NH_2 group.

Table VI. Calculated Total Charge Densities for
Aniline and a Positron-Aniline Complex

$$k^{(+)} = 3/4$$

Atom	Charge Density in Molecule	Charge Density Electron	Charge Density in Complex Positron
N	5.214	5.342	0.194
H _N	0.939	1.041	0.399
C _{SUBST}	3.770	3.789	0.002
C _{ORTHO}	4.010	3.971	0.001
H _{ORTHO}	1.057	1.032	0.001
C _{META}	3.912	3.897	0.000
H _{META}	1.066	1.014	0.000
C _{PARA}	3.982	3.942	0.000
H _{PARA}	1.064	1.016	0.000

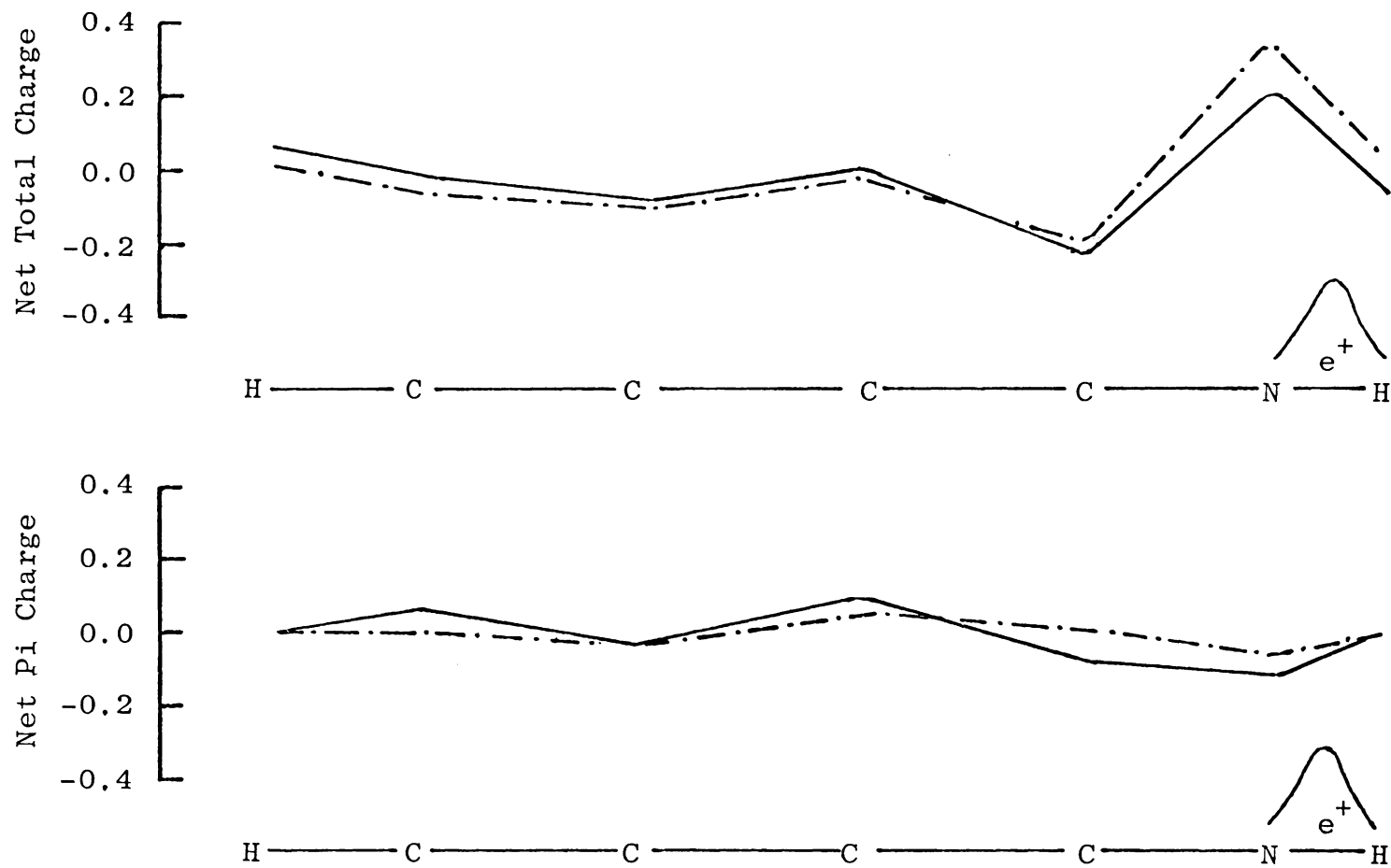
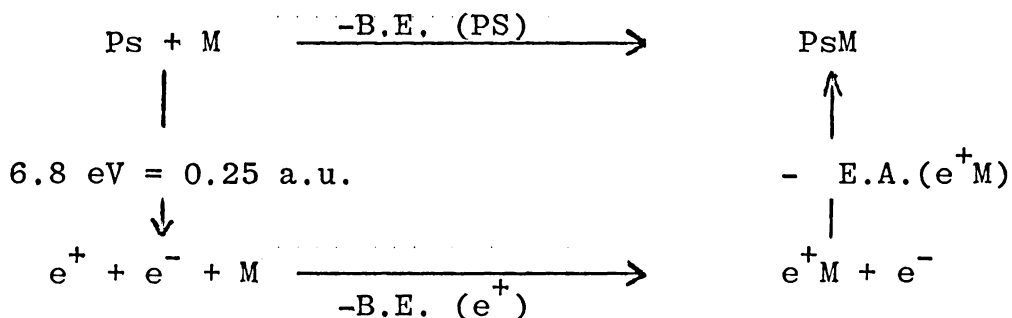


Figure 33. Profiles of Electronic Total Charge and Pi Charge for Aniline, Before (—) and After (---) Positron Addition. The Position of the Highest Positron Density (not to scale) is Indicated in the N-H Bond.

It is possible to estimate positronium binding energies for the molecules considered. This is done by considering the following cycle of processes, in which M denotes a neutral molecule, and Ps indicates the positronium atom (e^+e^-):



Thus, the positronium binding energy is approximately

$$\text{B.E. (Ps)} = \text{B.E.}(e^+) + \text{E.A.}(e^+\text{M}) - 0.25 \text{ a.u.} \quad (24)$$

Here, $\text{B.E.}(e^+)$ is the molecular positron affinity discussed earlier, $\text{E.A.}(e^+\text{M})$ is the electron affinity of the positron-molecule complex, obtained by Koopmans' theorem from the orbital energy of the first unoccupied electronic orbital in the complex, and the last term is the dissociation energy of the ground state of the hydrogen-like positronium atom.⁸⁰ The values obtained from Eq. (24) are listed in Tables IV and V for each molecule.

These calculations do not unambiguously show whether these molecules form stable positronium complexes. For all but aniline, the PsM complex was found to be unstable with

respect to dissociation to Ps + M when $k^{(+)} = 2/3$. But with $k^{(+)} = 3/4$, all positronium-molecule complexes were found to be stable.

Positronium complex formation in liquids is entirely consistent with the "free volume model" that is used to describe pickoff quenching and may be considered as a special case of it. In the "free volume model," the positronium is repelled by the lattice potential of the molecules of the liquid and forced to remain in the free volume. The pickoff annihilation rate is determined by the small amount of overlap between the wave functions of the positronium and the molecules of the liquid. In the case of nitroaromatic liquids, a free volume exists, they have a density, viscosity and surface tension that is similar to that of many non-reactive liquids. However, if complex formation occurs, as indicated by our present experimental and theoretical results, the positronium atom is not repelled by the molecules of the liquid but is incorporated into them. In this reversible complex, the positron occupies a molecular orbital and is no longer associated with the electron of the positronium atom. The observed rapid annihilation rate in nitroaromatics is caused by the enhanced electron density at the positron while in the complex.

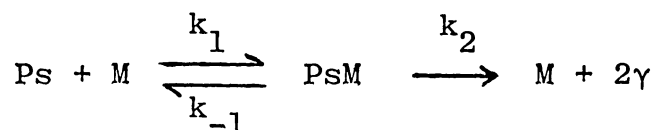
The calculations do not entirely agree with the experimental results of the previous chapter that showed complex

formation for only two of the compounds studied, nitrobenzene and p-benzoquinone. Schrader is presently performing similar calculations for positronium complexes⁸¹ and he also finds them to be bound but not quite as strongly. Other recent experimental studies^{78,79} suggest that the heats of complexation are negative and on the order of a few kcal/mole. This apparent discrepancy could be caused by any number of reasons. Inaccurate parameterization of the PSCF program is one possibility. The calculations were only for a 1:1 complex and in solution this is not necessarily the case especially if one is to use the "free volume model." More results are needed for gas phase reactions in order to properly test our calculations.

CONCLUSIONS

A. Summary

The mechanism for the rapid chemical reaction of ortho-positronium with a nitroaromatic molecule, M, in solution was found to be



The annihilation step is preceded by reversible complex formation. Only diamagnetic compounds that form strong molecular complexes were found to react by this mechanism.

Nitroaromatics are known to be strong electron accepting molecules and this is directly related to their ability to form these complexes. Because of this, substituents on the aromatic ring were found to affect the reactivity of these molecules toward the positronium atom. Nitroaromatics with electron withdrawing groups attached to the aromatic ring were found to form more stable complexes with positronium than those that contained electron donating ones.

The complexing ability of these molecules was found to be heavily dependent upon conjugation of the nitro group with the aromatic ring. Steric interactions that force the nitro group out of the plane of the ring caused the forward rate constant, k_1 , to be reduced.

A solvent effect was observed for this mechanism. Polar solvents were more able to assist the reactants in reaching the transition state associated with the complex formation step as compared to non-polar ones.

The results of CNDO calculations regarding the stability of positron and positronium complexes with large organic molecules showed that these complexes should be stable. However, complex formation was predicted for molecules in which no complex formation was observed.

B. Future Possibilities

The results of angular correlation experiments using the nitroaromatics could be used to determine the type(s) of electrons that annihilate the positron. Experiments similar to those mentioned for hexane and decane could be used to locate the bond(s) that contain the annihilation electrons, i.e., C-C, C-H, C-N, or N-O.

Many improvements can be made concerning the approximations used in the molecular orbital calculation. For example, the basis set used to construct the positron molecular orbitals was by no means complete. Atomic positron wave functions could be calculated and used here. These would also serve to better parameterize the calculations. These complexes were found to be stable in the absence of any external forces. Unlike the liquid phase, gas phase experiments should provide more realistic conditions under which the results of

these calculations could be evaluated. Since positronium is formed above thermal energy, excited positronium complexes are also possible and configuration interaction (CI) calculations should be examined.

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APPENDIX I. Data Tables

A. Rate Constants for the Reaction of Ortho-Positronium with Aromatic Nitro Compounds in Benzene* at 295°K.

<u>Solute</u> <u>k, M⁻¹nsec⁻¹</u>	<u>Concentra-</u> <u>tion, mM</u>	<u>λ₂, nsec⁻¹</u>	<u>I₂, %</u>
Nitrobenzene 27 ± 2	1.62	0.376	37.0
	3.68	0.430	35.1
	6.73	0.511	34.7
	9.33	0.575	33.5
	11.56	0.653	34.3
p-Dinitrobenzene 50 ± 2	1.52	0.417	36.5
	3.46	0.516	35.8
	6.35	0.669	36.8
	7.14	0.666	36.9
	10.90	0.882	36.4
	13.61	1.021	39.7
m-Dinitrobenzene 47 ± 2	1.38	0.391	35.0
	3.14	0.483	34.5
	5.77	0.605	34.7
	6.92	0.660	33.4
	8.56	0.697	33.0
	9.88	0.817	32.9
o-Dinitrobenzene 33 ± 2	1.40	0.400	34.9
	2.70	0.453	32.5
	5.84	0.544	36.4
	8.09	0.650	29.5
	10.01	0.728	29.2
	13.48	0.841	28.6
	15.94	0.770	21.6
20.45	0.997	26.4	
p-Nitrobenzotrile 39 ± 1	2.71	0.450	36.1
	5.42	0.557	34.5
	8.12	0.667	35.4
	10.83	0.817	35.6
	13.54	0.879	35.8
	27.07	1.404	39.3

*Pure benzene has λ₂ = 0.336 nsec⁻¹ and I₂ = 38.0%

Solute $k, M^{-1}nsec^{-1}$	Concentration, mM	$\lambda_2, nsec^{-1}$	$I_2, \%$
m-Nitrobenzonitrile 39 ± 2	2.58	0.434	35.1
	5.16	0.521	34.3
	7.74	0.636	33.5
	10.32	0.704	33.5
	12.91	0.830	35.6
	15.49	0.963	37.2
o-Nitrobenzonitrile 39 ± 2	2.49	0.416	33.2
	4.97	0.476	29.7
	7.46	0.627	31.3
	9.94	0.722	31.8
	12.43	0.794	31.8
	14.92	0.917	33.2
m-Nitrobenzotri- fluoride 38 ± 1	1.88	0.403	36.8
	3.76	0.475	36.4
	7.52	0.605	35.6
	11.27	0.748	36.1
	15.03	0.911	38.7
	16.91	0.982	38.6
o-Nitrobenzotri- fluoride 25 ± 1	3.14	0.409	34.6
	6.27	0.501	32.4
	9.41	0.582	31.5
	12.54	0.644	30.4
	15.68	0.723	29.1
	18.82	0.810	28.3
p-Nitrobenzaldehyde 40 ± 2	2.50	0.434	36.5
	4.99	0.533	34.6
	7.49	0.659	34.4
	9.99	0.743	34.5
	12.49	0.839	34.7
	14.98	0.936	35.1
m-Nitrobenzaldehyde 35 ± 1	2.42	0.424	35.9
	4.84	0.506	36.2
	7.27	0.594	34.5
	9.69	0.679	35.0
	12.11	0.741	32.0
	14.53	0.860	35.0

Solute $k, M^{-1}nsec^{-1}$	Concentration, mM	$\lambda_2, nsec^{-1}$	$I_2, \%$
o-Nitrobenzaldehyde 32 ± 2	3.00	0.428	35.3
	6.00	0.513	32.2
	9.00	0.614	31.7
	12.00	0.709	33.2
	15.00	0.800	33.8
	18.00	0.932	34.4
p-Nitroacetophenone 38 ± 2	2.52	0.411	33.0
	5.05	0.533	32.6
	7.57	0.626	31.2
	10.10	0.758	33.3
	12.62	0.803	28.6
	15.14	0.906	28.7
m-Nitroacetophenone 30 ± 2	2.32	0.401	35.9
	4.64	0.485	36.4
	6.96	0.546	33.9
	9.28	0.613	34.1
	11.61	0.701	34.4
	13.93	0.753	33.6
o-Nitroacetophenone 30 ± 2	2.88	0.412	34.4
	5.76	0.504	32.1
	8.64	0.599	31.3
	11.52	0.649	30.1
	14.40	0.736	27.8
	17.27	0.873	31.6
p-Nitrobenzoic acid 29 ± 5	0.17	0.349	37.9
	0.33	0.356	36.9
	0.50	0.357	37.2
	0.66	0.364	36.8
m-Nitrobenzoic acid 28 ± 2	4.35	0.481	35.9
	8.70	0.576	32.2
	13.05	0.694	32.9
	17.4	0.816	30.8
	26.1	1.092	32.2
o-Nitrobenzoic acid 20 ± 2	1.04	0.366	36.5
	2.60	0.399	34.5
	3.64	0.415	32.6
	5.20	0.428	33.0
	7.79	0.527	33.0
	10.39	0.538	29.8

<u>Solute</u> <u>k, M⁻¹nsec⁻¹</u>	<u>Concentra-</u> <u>tion, mM</u>	<u>λ₂, nsec⁻¹</u>	<u>I₂, %</u>
p-Nitrotoluene 22 ± 2	5.89	0.465	38.1
	11.60	0.580	33.9
	22.60	0.826	37.2
	29.45	0.972	34.4
m-Nitrotoluene 21 ± 2	1.53	0.392	34.3
	5.62	0.487	35.0
	9.20	0.581	35.3
	16.95	0.743	35.1
	22.50	0.917	36.6
	33.75	1.018	29.7
o-Nitrotoluene 8.5 ± 2	2.23	0.358	34.3
	4.24	0.360	32.0
	8.48	0.420	32.9
	12.72	0.455	32.0
	16.35	0.457	31.2
	23.10	0.515	29.9
2,6-Dimethylnitro- benzene 0.39 ± 0.05	36.75	0.346	64.2
	73.50	0.372	60.5
	98.00	0.370	34.0
	110.25	0.370	33.1
	147.00	0.393	32.1
p-Nitroanisole 13 ± 2	2.63	0.383	35.2
	5.25	0.396	35.1
	7.88	0.455	33.1
	10.51	0.477	33.9
	13.41	0.517	32.6
	15.76	0.549	32.2
m-Nitroanisole 27 ± 2	3.33	0.420	36.4
	6.65	0.521	36.8
	9.98	0.580	34.7
	13.31	0.686	35.8
	16.63	0.762	35.6
	19.96	0.883	36.7
o-Nitroanisole 49 ± 8	6.15	0.374	34.7
	12.30	0.411	33.0
	18.44	0.439	31.2
	24.59	0.449	27.5
	30.73	0.481	28.8
	36.88	0.532	28.6

Solute $k, M^{-1}nsec^{-1}$	Concentration, mM	$\lambda_2, nsec^{-1}$	$I_2, \%$
p-Nitrophenol 13 \pm 2	1.47	0.351	37.0
	2.94	0.371	34.2
	5.87	0.406	33.4
	8.81	0.452	31.6
	11.74	0.468	32.3
	14.68	0.528	30.8
m-Nitrophenol 26 \pm 2	2.88	0.409	35.4
	5.76	0.486	34.2
	8.63	0.560	32.8
	11.51	0.623	33.6
	14.39	0.698	32.9
	17.27	0.802	33.2
o-Nitrophenol 33 \pm 2	3.17	0.438	36.7
	6.33	0.532	35.5
	9.62	0.635	36.1
	12.66	0.768	37.2
	15.83	0.876	37.0
	18.99	0.954	36.4
p-Nitro-N,N- dimethylaniline 18 \pm 2	2.41	0.381	35.9
	4.81	0.425	36.2
	7.21	0.463	34.8
	9.62	0.512	35.9
	12.02	0.551	35.2
	14.42	0.592	35.5
p-Nitroaniline 5.1 \pm 0.6	1.78	0.327	35.0
	2.62	0.355	36.3
	3.71	0.349	35.1
	4.58	0.369	35.1
	6.55	0.367	35.6
	6.85	0.376	34.2
	9.17	0.384	33.7
	10.28	0.372	33.5
	11.78	0.395	32.9
	13.10	0.394	33.0
	15.71	0.427	32.5
	18.33	0.421	31.5
m-Nitroaniline 19 \pm 2	1.62	0.362	36.9
	3.36	0.401	35.8
	6.20	0.454	33.9
	8.07	0.492	34.9
	9.31	0.529	33.1
	11.52	0.557	34.1
	13.44	0.586	33.3

<u>Solute</u> <u>k, M⁻¹nsec⁻¹</u>	<u>Concentra-</u> <u>tion, mM</u>	<u>λ₂, nsec⁻¹</u>	<u>I₂, %</u>
o-Nitroaniline 10 ± 2	4.35	0.374	35.6
	8.70	0.415	34.2
	15.53	0.476	33.2
	21.76	0.521	32.0
	29.01	0.630	31.9
	36.26	0.643	31.6
	43.52	0.789	31.1
p-Nitrophenyl Acetate 26 ± 1	2.29	0.394	35.6
	4.59	0.459	33.7
	6.88	0.515	33.6
	9.17	0.569	34.3
	11.47	0.626	32.2
	13.76	0.698	32.4
o-Nitrophenyl Acetate 23 ± 1	2.47	0.388	35.5
	4.94	0.442	33.4
	7.41	0.497	32.2
	9.88	0.558	33.0
	12.35	0.619	32.7
	14.82	0.680	30.8
p-Nitrobenzyl alcohol 25 ± 1	1.37	0.372	37.0
	2.74	0.409	35.2
	5.49	0.469	33.3
	8.23	0.531	33.7
	10.98	0.600	32.3
	13.72	0.687	33.0
m-Nitrobenzyl alcohol 24 ± 2	2.97	0.432	35.6
	5.93	0.497	34.7
	8.90	0.542	33.2
	11.87	0.633	33.3
	14.84	0.681	31.7
	17.80	0.780	33.7
o-Nitrobenzyl alcohol 20 ± 2	2.99	0.401	34.2
	5.99	0.447	34.3
	8.98	0.507	32.0
	11.98	0.576	32.6
	14.97	0.638	32.9
	17.96	0.684	31.9

Solute $k, M^{-1}nsec^{-1}$	Concentration, mM	$\lambda_2, nsec^{-1}$	$I_2, \%$
p-Nitrofluoro- benzene 27 ± 2	2.36	0.397	36.1
	4.72	0.467	35.8
	7.07	0.531	36.2
	9.43	0.586	35.3
	11.79	0.654	35.3
	14.15	0.733	35.3
m-Nitrofluoro- benzene 35 ± 2	2.35	0.415	35.5
	4.70	0.497	35.9
	7.05	0.573	34.7
	9.40	0.663	36.3
	11.75	0.743	35.3
	14.10	0.835	36.6
o-Nitrofluoro- benzene 29 ± 2	2.36	0.407	36.0
	4.71	0.477	34.8
	7.07	0.522	32.7
	9.42	0.606	33.9
	11.78	0.667	33.4
	14.13	0.762	34.1
p-Nitrochloro- benzene 31 ± 2	1.65	0.380	36.4
	3.44	0.433	36.4
	3.85	0.482	35.5
	6.89	0.507	33.5
	9.55	0.658	35.0
	11.82	0.698	35.4
m-Nitrochloro- benzene 35 ± 2	1.62	0.391	36.5
	3.69	0.424	33.2
	6.76	0.546	33.1
	9.36	0.609	34.9
	11.59	0.745	33.7
	o-Nitrochloro- benzene 22 ± 2	1.44	0.369
3.28		0.420	34.5
6.02		0.479	33.4
8.33		0.529	32.4
10.32		0.596	32.9
p-Nitrobromo- benzene 32 ± 2		2.48	0.410
	4.96	0.492	33.7
	7.44	0.588	33.4
	9.92	0.655	32.7
	12.40	0.762	30.5
	14.88	0.782	33.9

<u>Solute</u> <u>k, M⁻¹nsec⁻¹</u>	<u>Concentra-</u> <u>tion, mM</u>	<u>λ₂, nsec</u>	<u>I₂, %</u>
m-Nitrobromo- benzene 36 ± 2	3.33	0.446	33.5
	6.66	0.536	31.9
	9.98	0.712	34.3
	13.31	0.814	33.6
	16.63	0.923	34.9
	19.96	1.048	36.8
o-Nitrobromo- benzene 22 ± 2	1.48	0.374	34.1
	2.95	0.418	33.2
	5.90	0.481	32.3
	8.85	0.560	29.4
	11.80	0.583	30.5
	14.75	0.680	31.3
p-Nitroiodo- benzene 32 ± 2	2.05	0.418	35.8
	4.10	0.476	34.6
	8.19	0.634	34.2
	10.24	0.691	33.2
	12.29	0.729	31.4
	14.34	0.796	32.0
m-Nitroiodo- benzene 35 ± 2	3.06	0.410	33.1
	6.11	0.505	31.4
	9.17	0.587	30.5
	12.22	0.635	28.2
	15.28	0.722	30.9
	18.34	0.753	28.8
o-Nitroiodo- benzene 23 ± 2	2.79	0.443	34.8
	5.58	0.519	33.0
	8.38	0.610	30.9
	11.17	0.756	34.1
	13.96	0.789	30.3
	16.75	0.945	34.1

B. Rate Constants for the Reaction of Ortho-Positronium
with Aliphatic Nitro Compounds in Benzene at 295°K.

Solute $10^{-7} k, M^{-1} \text{sec}^{-1}$	Concentration, M	$\lambda_2, \text{nsec}^{-1}$	$I_2, \%$
Nitromethane 1.4 ± 0.3	1.23	0.374	25.1
	2.47	0.383	19.2
	3.70	0.407	16.5
	4.96	0.426	14.8
	6.27	0.462	13.6
	8.64	0.464	12.1
Nitroethane 2.4 ± 0.3	0.46	0.362	30.9
	0.93	0.364	27.1
	1.87	0.383	21.7
	2.80	0.420	18.8
	3.73	0.447	16.9
	4.67	0.451	14.8
	6.53	0.500	12.7
	8.40	0.510	12.0
1-Nitropropane 2.0 ± 0.3	0.38	0.355	28.2
	1.50	0.363	21.4
	2.25	0.387	20.4
	3.00	0.398	17.2
	4.50	0.435	15.7
	5.25	0.454	14.1
	6.00	0.473	13.4
2-Nitropropane 1.85 ± 0.2	0.14	0.343	36.4
	0.37	0.345	32.5
	1.11	0.358	24.9
	2.22	0.376	18.7
	4.45	0.430	15.1
	6.68	0.464	12.4
	8.91	0.499	10.9
	11.14	0.551	9.5
α -Nitrotoluene 0.11 ± 0.02	0.85	0.432	19.0
	1.70	0.516	15.7
	2.55	0.619	13.6
	3.40	0.741	15.0
	8.46	1.235	10.3

C. Temperature Dependence of the Observed Rate Constants for the Reaction of Ortho-Positronium with Substituted Nitrobenzenes.

Solvent: Toluene

<u>Solute</u>	<u>T, °C</u>	<u>λ_2, nsec⁻¹</u>	<u>I₂, %</u>	<u>k, M⁻¹nsec⁻¹</u>
9.77 mM Nitrobenzene	-75	0.433	27.3	5.1
	-27	0.497	29.9	14.6
	5	0.567	32.7	23.9
	21	0.575	33.3	25.8
	23	0.594	32.2	27.9
	33	0.579	31.5	27.1
	38	0.576	33.8	27.0
	42	0.597	34.6	29.6
	56	0.531	34.2	23.7
	58	0.511	34.5	21.8
	75	0.436	36.2	15.3
	77	0.434	35.3	15.1
	93	0.340	37.2	6.6
	100	0.310	36.9	4.0
	113	0.292	37.3	3.1
122	0.277	37.6	2.2	
129	0.270	38.4	1.9	
142	0.257	38.5	1.5	
154	0.244	39.4	1.0	
7.36 mM p-Dinitrobenzene	-75	0.430	33.3	6.0
	-45	0.459	32.4	12.7
	-26	0.507	32.4	20.8
	5	0.595	32.4	35.6
	21	0.700	34.9	51.3
	35	0.752	35.3	59.5
	54	0.819	35.0	70.4
	70	0.927	36.5	86.5
	85	1.043	40.0	103.6
	105	1.125	40.9	116.5
	123	1.262	46.0	136.7
	149	1.231	46.5	134.8
	163	1.046	46.8	111.0
173	0.856	44.4	86.1	

<u>Solute</u>	<u>T, °C</u>	<u>λ_2, nsec⁻¹</u>	<u>I₂, %</u>	<u>k, M⁻¹nsec⁻¹</u>
2.02 mM p-Nitrobenzo- nitrile	-75	0.396	35.4	5.2
	6	0.400	34.2	33.5
	21	0.402	34.4	39.5
	40	0.418	34.5	53.4
	56	0.428	33.8	63.8
	71	0.447	35.8	77.8
	87	0.441	34.9	80.2
	95	0.434	36.2	79.5
	102	0.441	36.6	85.0
	117	0.437	36.8	88.2
	132	0.397	38.1	72.9
	152	0.350	40.0	56.5
	176	0.267	41.1	23.0
	16.12 mM p-Nitroanisole	-75	0.444	27.1
-45		0.510	28.9	8.9
-26		0.530	30.2	11.0
6		0.538	31.1	12.8
23		0.466	36.6	9.0
38		0.416	33.0	6.5
54		0.353	35.1	3.3
71		0.318	37.2	1.8
86		0.296	36.3	1.0
107		0.282	37.0	1.0
130		0.262	38.6	0.7
146		0.254	38.9	0.8
173		0.235	40.3	0.8
13.22 mM p-Nitroaniline	5	0.385	28.6	3.9
	5	0.376	30.1	3.2
	21	0.359	32.2	2.8
	38	0.337	34.1	2.0
	59	0.319	36.0	1.6
	77	0.304	37.6	1.4
	96	0.295	38.2	1.7
	111	0.277	37.4	1.0
	128	0.268	38.7	1.2
145	0.254	38.8	1.0	
160	0.242	39.2	0.2	

<u>Solute</u>	<u>T, °C</u>	<u>$\lambda_{2'}$ nsec⁻¹</u>	<u>I_{2'}, %</u>	<u>k, M⁻¹nsec⁻¹</u>
206. mM	-75	0.601	11.8	1.04
2,6-Dimethyl- nitrobenzene	-40	0.587	14.3	1.10
	-27	0.550	17.4	0.95
	5	0.419	27.8	0.42
	21	0.364	31.1	0.20
	37	0.342	32.8	0.15
	56	0.324	35.0	0.12
	74	0.307	35.7	0.09
	91	0.284	35.8	0.04
	106	0.277	36.8	0.05
	117	0.273	35.8	0.07
	130	0.261	36.5	0.05
	142	0.261	36.7	0.09
<u>Miscellaneous</u>				
13.15 mM	-75	0.435	33.3	3.7
p-Benzoquinone	5	0.725	33.7	29.8
	23	0.846	37.3	39.9
	41	1.035	43.1	55.2
	58	1.014	38.7	54.4
	78	1.193	48.5	69.0
1.315 mM	-75	0.390	38.9	2.8
p-Benzoquinone	5	0.367	36.4	25.8
	22	0.372	37.1	37.6
	46	0.373	37.7	51.2
	62	0.369	37.3	56.0
	82	0.369	38.3	66.2
	100	0.383	38.8	85.7
	123	0.382	42.4	96.6
	142	0.397	42.6	117.7
	157	0.374	43.1	107.2
	174	0.337	43.3	87.8
	192	0.263	47.1	40.4

Solvent: n-Heptane		λ_2 ,	I_2 , %	k , $M^{-1}nsec^{-1}$
<u>Solute</u>	<u>T, °C</u>	<u>nsec⁻¹</u>	<u>I₂, %</u>	<u>k, M⁻¹nsec⁻¹</u>
19.55 mM	-75	0.425	27.0	4.9
Nitrobenzene	-42	0.589	28.8	14.3
	-26	0.608	32.0	15.8
	6	0.443	33.7	8.3
	21	0.335	35.0	3.3
	35	0.294	35.9	1.6
	44	0.277	36.1	1.0
	65	0.253	36.7	0.5
	86	0.237	36.7	0.3
1.814 mM	22	0.403	33.8	73.0
p-Dinitrobenzene	22	0.411	33.5	77.6
	38	0.422	34.3	88.8
	58	0.432	35.0	100.8
	76	0.456	35.9	120.4
	97	0.430	48.0	112.9
	122	0.297	50.8	49.1
	144	0.231	52.8	18.7
	161	0.198	52.5	6.2
26.75 mM	5	0.305	35.6	0.92
p-Nitroanisole	22	0.280	32.3	0.38
	22	0.279	33.6	0.34
	54	0.256	32.4	0.19
	76	0.240	40.0	0.08

Solvent: n-Pentanol		λ_2	I_2 , %	k , $M^{-1}nsec^{-1}$
<u>Solute</u>	<u>T, °C</u>	<u>nsec⁻¹</u>	<u>I₂, %</u>	<u>k, M⁻¹nsec⁻¹</u>
9.77 mM	-75	0.372	22.0	0.7
Nitrobenzene	5	0.406	21.1	8.3
	22	0.409	20.8	8.6
	37	0.428	21.0	10.6
	59	0.461	23.3	13.9
	82	0.461	23.3	14.6
	105	0.444	24.0	13.2
	127	0.392	23.8	8.0

<u>Solute</u>	<u>T, °C</u>	<u>$\lambda_2,$ nsec⁻¹</u>	<u>I₂, %</u>	<u>k, M⁻¹nsec⁻¹</u>
7.72 mM p-Dinitrobenzene	-75	0.374	22.0	1.0
	5	0.414	19.9	11.6
	22	0.462	19.6	17.8
	40	0.516	20.6	24.9
	65	0.657	24.2	43.1
	88	0.758	27.3	56.9
	113	0.924	34.3	99.2
12.84 mM p-Nitroanisole	-75	0.371	37.2	0.5
	4	0.373	25.4	3.8
	22	0.379	21.0	4.3
	39	0.385	21.0	4.7
	59	0.378	21.5	4.1
	81	0.377	22.3	4.6
	101	0.359	23.4	3.4

D. Temperature Dependence of the Observed Rate Constants for the Reaction Ortho-Positronium with Solvents and Weakly Reactive Compounds.

<u>Compound</u>	<u>T, °C</u>	<u>λ_2, nsec⁻¹</u>	<u>I₂, %</u>	<u>10⁻⁷k, M⁻¹ sec⁻¹</u>
n-Heptane	-75	0.319	31.6	3.90
	-75	0.325	30.4	3.98
	-45	0.323	32.0	4.20
	-23	0.300	32.4	4.06
	6	0.286	34.8	4.09
	21	0.269	35.8	3.94
	21	0.263	36.1	3.86
	49	0.255	36.9	3.91
	65	0.248	37.8	3.91
	78	0.243	39.3	3.90
	92	0.224	39.1	3.68
	110	0.219	41.7	3.70
	131	0.208	41.5	3.61
	148	0.192	42.5	3.42
	162	0.184	42.9	3.34
179	0.173	44.6	3.23	
Toluene	-75	0.384	38.5	3.45
	-75	0.381	37.3	3.42
	-57	0.373	36.4	3.44
	-27	0.355	36.4	3.46
	3	0.333	37.3	3.43
	3	0.337	33.8	3.47
	21	0.319	36.8	3.39
	37	0.312	36.5	3.40
	53	0.303	36.2	3.38
	72	0.294	37.1	3.39
	91	0.277	37.8	3.28
	110	0.260	38.1	3.17
	136	0.250	40.5	3.17
	170	0.223	41.1	2.97
n-Pentanol	-75	0.366	23.6	3.97
	22	0.324	22.2	3.51
	22	0.326	22.8	3.53
	36	0.326	22.7	3.53
	65	0.325	22.9	3.52
	88	0.319	22.8	3.46
	108	0.315	23.1	3.42

<u>Compound</u>	<u>T, °C</u>	<u>$\lambda_2,$ nsec⁻¹</u>	<u>I₂, %</u>	<u>10⁻⁷k, M⁻¹ sec⁻¹</u>
Benzonitrile	5	0.541	12.4	5.52
	23	0.458	10.5	4.67
	40	0.428	9.7	4.36
	67	0.457	12.2	4.66
	86	0.435	12.1	4.44
	112	0.443	14.0	4.52
	135	0.379	12.6	3.87
	157	0.379	13.8	3.86
2-Nitropropane	22	0.452	6.5	4.06
	41	0.481	8.1	4.31
	62	0.429	6.0	3.85
	91	0.447	9.0	4.01
	120	0.467	11.3	4.18

E. Thermodynamic Data for the Reactions of o-Ps

<u>Compound</u>	Low Temperature Limit		High Temperature Limit
	$E_a(1)$, <u>kcal/mole</u>	$\Delta S_a(-26^\circ\text{C})$, <u>eu*</u>	$\Delta H_{\text{EQ}} + E_a(2)$, <u>kcal/mole</u>
Solvent: Toluene			
p-Nitroaniline	---	---	-1.1 ± 0.1
2,6-Dimethylnitrobenzene	---	---	-2.7 ± 0.3
p-Nitroanisole	0.9 ± 0.1	-10.2	-3.6 ± 0.3
Nitrobenzene	0.8 ± 0.1	-9.7	-4.5 ± 0.3
p-Nitrobenzotrile	1.0 ± 0.1	-9.2	-6.0 ± 0.5
p-Dinitrobenzene	1.1 ± 0.1	-8.9	<-4.0
p-Benzoquinone	1.0 ± 0.1	-10.5	-7.5 ± 0.5
Solvent: n-Heptane			
p-Nitroanisole	---	---	-2.9 ± 0.4
Nitrobenzene	1.3 ± 0.3	-8.1	-4.4 ± 0.4
p-Dinitrobenzene	1.0 ± 0.2	-7.1	-8.0 ± 1.0
Solvent: n-Pentanol			
p-Nitroanisole	1.2 ± 0.1	-7.3	-1.3 ± 0.8
Nitrobenzene	1.3 ± 0.2	-6.2	-3.5 ± 0.6
p-Dinitrobenzene	1.8 ± 0.3	-5.5	---

* Calculated using the average $E_a(1)$ value.

APPENDIX II. Computer Programs

A. Introduction

Many trivial but time consuming calculations were performed by the three convenience programs (SLOW RATE, ARRHENIUS, AND HAMMETT) listed in this appendix. They are written in either FORTRAN (IV) or WATFIV for use on an autobatcher.

Section E contains the four subroutines needed to convert the CNDO/2 program listed in reference 68 into the PSCF program used to do the positron complex calculations. By substituting these four subroutines for the three (HUCKCL, SCFCLO, and CPRINT) in CNDO/2, one obtains the PSCF program.

The last section contains the program, PAL, that was used to calculate the lifetime, annihilation rate and intensity data from the positron lifetime spectra. The Colcomp output is generated in the POL subroutine of this program.

C	PROGRAM SLOW RATE	SLOW 10
C		SLOW 20
C	THIS FORTRAN(IV) PROGRAM CALCULATES RATE CONSTANTS FOR THE	SLOW 30
C	REACTIONS OF O-PS WHEN LO, THE SOLVENT ANNIHILATION RATE,	SLOW 40
C	CHANGES DUE TO HIGH CONCENTRATIONS OF A LIQUID SOLUTE	SLOW 50
C	IDEAL MIXING IS ASSUMED	SLOW 60
C	THE EQUATION NECESSARY IS $L2(V/MN+MS)=KS+FN(KN-KS)$	SLOW 70
C	A PLOT OF THE LEFT HAND SIDE VERSUS FN SHOULD PRODUCE A	SLOW 80
C	STRAIGHT LINE WITH SLOPE= $(KN-KS)$ AND Y-INTERCEPT= KS	SLOW 90
C		SLOW 100
	DIMENSION X(20), Y(20), VS(20), VN(20), COM(10), FS(20), FN(20)	SLOW 110
	REAL KS,MDS,MS(20),MN(20),MDN,MWS,MWN	SLOW 120
	READ (5,60) NCASE	SLOW 130
C	NUMBER OF CASES INPUT	SLOW 140
20	READ (5,70) (COM(I),I=1,10)	SLOW 150
	WRITE (6,70) (COM(I),I=1,10)	SLOW 160
	READ (5,60) N	SLOW 170
C	NUMBER OF DATA POINTS	SLOW 180
	READ (5,80) DN,MWN	SLOW 190
C	DENSITY OF SOLUTE, GM/ML, MOLECULAR WEIGHT OF SOLUTE	SLOW 200
	DO 30 I=1,N	SLOW 210
	READ (5,80) X(I),Y(I)	SLOW 220
C	MOLARITY OF SOLUTE, LAMBDA-2(TIMES 10(-9))	SLOW 230
30	CCONTINUE	SLOW 240
	KS=2.963*10**(7)	SLOW 250
C	RATE CONSTANT FOR PURE SOLVENT	SLOW 260
	DS=0.87865	SLOW 270
C	DENSITY OF PURE SOLVENT, GM/ML	SLOW 280
	MWS=78.12	SLOW 290
C	MOLECULAR WEIGHT OF THE SOLVENT	SLOW 300
	MDS=DS/MWS	SLOW 310
C	CONCENTRATION OF THE PURE SOLVENT	SLOW 320

	MDN=DN/MWN	SLOW 330
C	CONCENTRATION OF THE PURE SOLUTE	SLOW 340
	DO 40 I=1,N	SLOW 350
	VN(I)=X(I)/(MDN*1000.)	SLOW 360
C	VOLUME OF SOLUTE IN 1ML OF SOLUTION	SLOW 370
	VS(I)=1.0-VN(I)	SLOW 380
C	VOLUME OF SOLVENT IN 1ML OF SOLUTION	SLOW 390
	MN(I)=X(I)/1000.	SLOW 400
C	MILLIMOLES OF SOLUTE	SLOW 410
	MS(I)=VS(I)*MDS	SLOW 420
C	MILLIMOLES OF SOLVENT	SLOW 430
	FS(I)=MS(I)/(MS(I)+MN(I))	SLOW 440
C	MOLE FRACTION OF SOLVENT	SLOW 450
	FN(I)=1.0-FS(I)	SLOW 460
C	MOLE FRACTION OF SOLUTE	SLOW 470
	Y(I)=Y(I)/((MS(I)+MN(I))*1000.)	SLOW 480
	X(I)=FN(I)	SLOW 490
40	CCONTINUE	SLOW 500
C	A LINEAR LEAST SQUARES PROGRAM FOR N DATA POINTS	SLOW 510
	SUMX=0.0	SLOW 520
	SUMX2=0.0	SLOW 530
	SUMXY=0.0	SLOW 540
	SUMY=0.00	SLOW 550
	DO 50 I=1,N	SLOW 560
	PRINT , 'X=',X(I),'Y=',Y(I)	SLOW 570
	SUMX=SUMX+X(I)	SLOW 580
	SUMX2=SUMX2+X(I)**2	SLOW 590
	SUMXY=SUMXY+X(I)*Y(I)	SLOW 600
50	SUMY=SUMY+Y(I)	SLOW 610
	SLOPE=(N*SUMXY-SUMX*SUMY)/(N*SUMX2-SUMX**2)	SLOW 620
	YINTR=(SUMX2*SUMY-SUMX*SUMXY)/(N*SUMX2-SUMX**2)	SLOW 630
	SLOPE=(SLOPE*10**9)	SLOW 640

	YINTR=(YINTR*10**9)	SLOW 650
C	RATE CONSTANT OF THE SOLVENT, USE AS A CHECK	SLOW 660
	SLOPE=SLOPE+KS	SLOW 670
C	RATE CONSTANT OF THE SOLUTE	SLOW 680
	WRITE (6,90) SLOPE,YINTR	SLOW 690
	NCASE=NCASE-1	SLOW 700
	IF (NCASE.LE.0) STOP	SLOW 710
	GO TO 20	SLOW 720
C		SLOW 730
C		SLOW 740
6C	FORMAT (I2)	SLOW 750
7C	FORMAT (2X,10A4)	SLOW 760
8C	FORMAT (2F9.4)	SLOW 770
9C	FORMAT (20X,13H K(SOLUTE)=,E11.3,13H K(SOLVENT)=,E11.3)	SLOW 780
	END	SLOW 790

C	PROGRAM ARRHENIUS	ARHN	10
C		ARHN	15
C	THIS FORTRAN(IV) PROGRAM CALCULATES POINTS ON AN ARRHENIUS	ARHN	20
C	PLOT FROM O-PS ANNIHILATION RATES	ARHN	30
C		ARHN	40
	REAL L2,LO,K,LOGK	ARHN	50
	DIMENSION L2(50), LO(50), TC(50), TK(50), TKINV(50), IDENT(20)	ARHN	60
	DIMENSION K(50), LOGK(50)	ARHN	70
	DIMENSION CCN(50)	ARHN	80
	READ (5,60) NCASE	ARHN	90
C	NUMBER OF CASES INPUT	ARHN	100
	NCASE1=NCASE	ARHN	110
20	READ (5,90) IDENT	ARHN	120
C	IDENTIFIED, ONE FOR EACH CASE	ARHN	130
	WRITE (6,50)	ARHN	140
	WRITE (6,90) IDENT	ARHN	150
	READ (5,60) N	ARHN	160
C	NUMBER OF DATA POINTS FOR EACH CASE	ARHN	170
	READ (5,70) CONC	ARHN	180
C	CONCENTRATION OF SOLUTE IN MOLAR UNITS	ARHN	190
	WRITE (6,100)	ARHN	200
	DO 30 I=1,N	ARHN	210
	READ (5,80) L2(I),TC(I)	ARHN	220
C	LAMBDA-2, CENTIGRADE TEMPERATURE	ARHN	230
C	L2 TIMES 10(-9)	ARHN	240
	LO(I)=0.28397-(0.0006005*TC(I))	ARHN	250
C	TEMPERATURE DEPENDENCE OF SOLVENT ANNIHILATION RATE	ARHN	260
	TK(I)=TC(I)+273.	ARHN	270
C	KELVIN TEMPERATURE	ARHN	280
	K(I)=(L2(I)-LO(I))/CONC	ARHN	290
C	OBSERVED RATE CONSTANT	ARHN	300
C	K EXPRESSED TIMES 10(-9)	ARHN	310

	TKINV(I)=1000.0/TK(I)	ARHN 320
C	INVERSE KELVIN TEMPERATURE TIMES 1000	ARHN 330
30	CCONTINUE	ARHN 340
	DO 40 I=1,N	ARHN 350
	WRITE (6,110) TC(I),L2(I),TKINV(I),K(I)	ARHN 360
40	CCONTINUE	ARHN 370
	NCASE=NCASE-1	ARHN 380
	IF (NCASE.NE.0) GO TO 20	ARHN 390
	WRITE (6,50)	ARHN 400
	STOP	ARHN 410
C		ARHN 420
C		ARHN 430
C		ARHN 440
50	FORMAT (1H1)	ARHN 450
60	FORMAT (I2)	ARHN 460
70	FORMAT (F12.9)	ARHN 470
80	FCRMT (2F9.4)	ARHN 480
90	FORMAT (20A4)	ARHN 490
100	FORMAT (20X,43H TC L2 1000/TK K*10(-9))	ARHN 500
110	FORMAT (20X,F10.1,F10.4,F10.3,F10.4)	ARHN 510
	END	ARHN 520

C	PROGRAM HAMMETT	HMET	10
C		HMET	15
C	THIS PROGRAM CALCULATES POINTS ON A HAMMETT PLOT WHERE THE	HMET	20
C	HAMMETT EQUATION IS	HMET	30
C		HMET	40
C	$\text{LOG}(K/K_{\text{REF}}) = \rho(\text{IND}) * \sigma(\text{IND}) + \rho(\text{RES}) * \sigma(\text{RES})$	HMET	50
C		HMET	60
	REAL LOGY(100), LOGYC(100)	HMET	70
	REAL LOGYAV	HMET	80
	DIMENSION ID(20)	HMET	90
	DIMENSION SUB(100,9), SBAR(100)	HMET	100
	DIMENSION Y(100), SI(100), SR(100), YCAL(100), DEL(100)	HMET	110
	READ , NCASE	HMET	120
C	NUMBER OF CASES INPUT	HMET	130
20	READ (5,100) ID	HMET	140
C	IDENTIFIER, ONE FOR EACH CASE	HMET	150
	WRITE (6,120)	HMET	160
	WRITE (6,110) ID	HMET	170
	READ , N	HMET	180
C	NUMBER OF DATA POINTS	HMET	190
	DO 30 I=1,N	HMET	200
	READ (5,80) (SUB(I,J),J=1,9),Y(I),SI(I),SR(I)	HMET	210
C	SUBSTITUENT ON NITROBENZENE, K OBS, SIGMA(IND), SIGMA(RES)	HMET	220
C	Y(1) MUST BE FOR REFERENCE COMPOUND, I.E. NITROBENZENE	HMET	230
30	CONTINUE	HMET	240
	DO 40 I=1,N	HMET	250
	LOGY(I)=ALOG10(Y(I)/Y(1))	HMET	260
40	CONTINUE	HMET	270
C	LINEAR LEAST SQUARES PORTION OF THE PROGRAM	HMET	280
	SLYSI=0.0	HMET	290
	SLYSR=0.0	HMET	300
	SSRSI=0.0	HMET	310

	SSR2=0.0	HMET 320
	SSI2=0.0	HMET 330
	DO 50 I=1,N	HMET 340
	SLYSI=SLYSI+LOGY(I)*SI(I)	HMET 350
	SLYSR=SLYSR+LOGY(I)*SR(I)	HMET 360
	SSRSI=SSRSI+SR(I)*SI(I)	HMET 370
	SSI2=SSI2+SI(I)*SI(I)	HMET 380
	SSR2=SSR2+SR(I)*SR(I)	HMET 390
50	CCONTINUE	HMET 400
	SSRSI2=SSRSI**2	HMET 410
	PI=((SSR2*SLYSI)-(SSRSI*SLYSR))/((SSI2*SSR2)-(SSRSI2))	HMET 420
C	RHO INDUCTIVE	HMET 430
	PR=((SSI2*SLYSR)-(SSRSI*SLYSI))/((SSI2*SSR2)-(SSRSI2))	HMET 440
C	RHO RESONANCE	HMET 450
	PRINT , 'P(I)=' ,PI,'P(R)=' ,PR	HMET 460
	PRINT , 'SUBST. RATEK RATEKCAL DEL LOGK/KNB LOGK/KC SBAR'	HMET 470
	DO 60 I=1,N	HMET 480
	SBAR(I)=PI*SI(I)+PR*SR(I)	HMET 490
C	SIGMA BAR	HMET 500
	YCAL(I)=10**((SBAR(I)+ALOG10(Y(1))))	HMET 510
	LOGYC(I)=ALOG10(YCAL(I)/Y(1))	HMET 520
	DEL(I)=LOGYC(I)-LOGY(I)	HMET 530
C	DIFFERENCE BETWEEN LOG(K) AND LOG(KCALC)	HMET 540
	WRITE (6,90) (SUR(I,J),J=1,9),Y(I),YCAL(I),DEL(I),LOGY(I),LOGYC(I)	HMET 550
	1,SBAR(I)	HMET 560
60	CCONTINUE	HMET 570
	LOGYAV=0.0	HMET 580
	DELAV=0.0	HMET 590
	DO 70 I=1,N	HMET 600
	DELAV=DELAV+DEL(I)**2	HMET 610
	LOGYAV=LOGYAV+LOGY(I)**2	HMET 620
70	CCONTINUE	HMET 630

	SD=(DELAV/FLOAT(N))*0.5	HMET 640
C	STANDARD DEVIATION	HMET 650
	RMS=(LOGYAV/FLOAT(N))*0.5	HMET 660
C	ROOT MEAN SQUARE	HMET 670
	F=SD/RMS	HMET 680
C	EHRENSON'S FIT PARAMETER	HMET 690
	PRINT , 'STDDEV=',SD,'RMS=',RMS,'F=',F	HMET 700
	NCASE=NCASE-1	HMET 710
	IF (NCASE.NE.0) GO TO 20	HMET 720
	WRITE (6,120)	HMET 730
	STOP	HMET 740
C		HMET 750
80	FORMAT (9A1,3F10.5)	HMET 760
90	FORMAT (5X,9A1,6F10.4)	HMET 770
100	FORMAT (20A4)	HMET 780
110	FORMAT (10X,20A4)	HMET 790
120	FORMAT (1H1)	HMET 800
	END	HMET 810

C	PROGRAM PSCF	PSCF	10
	SUBROUTINE HUCKCL(ICORE)	HUCK	10
C	CALCULATIONS DONE IN THIS SUBROUTINE USE EXTENDED HUCKEL	HUCK	20
C	THEORY FOR CLOSED SHELLS ONLY	HUCK	30
C	OVERLAPS ARE IN MATRIX A, COULOMB INTEGRALS (GAMMA) ARE IN MATRIX G	HUCK	40
C	A SUBROUTINE CORMOD(ICORE) MUST BE SUPPLIED. THIS IS USED TO	HUCK	50
C	MODIFY THE CORE INTEGRALS IF (ICORE.NE.0)	HUCK	60
C	FOR MINDO, USE ICORE=2	HUCK	70
	IMPLICIT REAL*8(A-H,C-Z)	HUCK	80
	COMMON /ARRAYS/ A(80,80),B(80,80),PS(80,80),PSB(80,80)	HUCK	90
	COMMON /INFO/ NATOMS,CHARGE,MULTIP,AN(35),C(35,3),N	HUCK	100
	COMMON /INFO1/ CZ(35),U(80),ULIM(35),LLIM(35),NELECS,OCCA,OCCB	HUCK	110
	COMMON /GAB/ XXX(400),G(35,35),Q(80),YYY(80),ENERGY,XXY(214)	HUCK	120
	COMMON /OPTICM/ OPTION,OPNCLO,HUCKEL,CNDO,INDO,CLOSED,OPEN	HUCK	130
	COMMON /CORE/ G1(18),F2(18),ENEG(18,3),BETA0(18)	HUCK	140
	COMMON /PS/ PSQ(80),PSK	HUCK	150
	INTEGER CHARGE,OCCA,CCCB,UL,AN,CZ,U,ULIM,ANI	HUCK	160
	INTEGER OPTICN,OPNCLO,HUCKEL,CNDO,INDO,CLOSED,OPEN	HUCK	170
	G1(3)=.092012D0	HUCK	180
	G1(4)=.1407D0	HUCK	190
	G1(5)=.199265D0	HUCK	200
	G1(6)=.267708D0	HUCK	210
	G1(7)=.346029D0	HUCK	220
	G1(8)=.43423D0	HUCK	230
	G1(9)=.532305D0	HUCK	240
	F2(3)=.049865D0	HUCK	250
	F2(4)=.089125D0	HUCK	260
	F2(5)=.13041D0	HUCK	270
	F2(6)=.17327D0	HUCK	280
	F2(7)=.219055D0	HUCK	290

F2(8)=.266415D0
F2(9)=.31580D0
ENEG(1,1)=7.1761D0
ENEG(3,1)=3.1055D0
ENEG(3,2)=1.258D0
ENEG(4,1)=5.94557D0
ENEG(4,2)=2.563D0
ENEG(5,1)=9.59407D0
ENEG(5,2)=4.001D0
ENEG(6,1)=14.051D0
ENEG(6,2)=5.572D0
ENEG(7,1)=19.31637D0
ENEG(7,2)=7.275D0
ENEG(8,1)=25.39017D0
ENEG(8,2)=9.111D0
ENEG(9,1)=32.2724D0
ENEG(9,2)=11.08D0
ENEG(11,1)=2.804D0
ENEG(11,2)=1.302D0
ENEG(11,3)=0.150D0
ENEG(12,1)=5.1254D0
ENEG(12,2)=2.0516D0
ENEG(12,3)=0.16195D0
ENEG(13,1)=7.7706D0
ENEG(13,2)=2.9951D0
ENEG(13,3)=0.22425D0
ENEG(14,1)=10.0327D0
ENEG(14,2)=4.1325D0
ENEG(14,3)=0.337D0
ENEG(15,1)=14.0327D0
ENEG(15,2)=5.4638D0
ENEG(15,3)=0.500D0

HUCK 300
HUCK 310
HUCK 320
HUCK 330
HUCK 340
HUCK 350
HUCK 360
HUCK 370
HUCK 380
HUCK 390
HUCK 400
HUCK 410
HUCK 420
HUCK 430
HUCK 440
HUCK 450
HUCK 460
HUCK 470
HUCK 480
HUCK 490
HUCK 500
HUCK 510
HUCK 520
HUCK 530
HUCK 540
HUCK 550
HUCK 560
HUCK 570
HUCK 580
HUCK 590
HUCK 600
HUCK 610

ENEG(16,1)=17.6496D0	HUCK 620
ENEG(16,2)=6.989D0	HUCK 630
ENEG(16,3)=0.71325D0	HUCK 640
ENEG(17,1)=21.5906D0	HUCK 650
ENEG(17,2)=8.7081D0	HUCK 660
ENEG(17,3)=0.97695D0	HUCK 670
BETA0(1)=-9.D0	HUCK 680
BETA0(3)=-9.D0	HUCK 690
BETA0(4)=-13.D0	HUCK 700
BETA0(5)=-17.D0	HUCK 710
BETA0(6)=-21.D0	HUCK 720
BETA0(7)=-25.D0	HUCK 730
BETA0(8)=-31.D0	HUCK 740
BETA0(9)=-39.D0	HUCK 750
BETA0(11)=-7.7203D0	HUCK 760
BETA0(12)=-9.4471D0	HUCK 770
BETA0(13)=-11.3011D0	HUCK 780
BETA0(14)=-13.065D0	HUCK 790
BETA0(15)=-15.070D0	HUCK 800
BETA0(16)=-18.150D0	HUCK 810
BETA0(17)=-22.330D0	HUCK 820
IF (ICORE.NE.0) CALL CORMOD (ICORE)	HUCK 830
C FIND NELECS AND FILL W CORE(DIAGONAL) WITH (I+A)/2	HUCK 840
NELECS=0	HUCK 850
DO 50 I=1,NATCMS	HUCK 860
NELECS=NELECS+CZ(I)	HUCK 870
LL=LLIM(I)	HUCK 880
UL=ULIM(I)	HUCK 890
ANI=AN(I)	HUCK 900
L=0	HUCK 910
DO 40 J=LL,UL	HUCK 920
L=L+1	HUCK 930

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IF (L.EQ.1) GO TO 30                                HUCK 940
IF (L.LT.5) GO TO 20                                HUCK 950
A(J,J)=-ENEG(ANI,3)/27.21D0                          HUCK 960
PS(J,J)=-3.0D0*(A(J,J)-(DFLOAT(CZ(I))-0.5D0)*G(I,I)) HUCK 970
GO TO 40                                              HUCK 980
20 A(J,J)=-ENEG(ANI,2)/27.21D0                        HUCK 990
PS(J,J)=-3.0D0*(A(J,J)-(DFLOAT(CZ(I))-0.5D0)*G(I,I)) HUCK1000
C MINDG CORRECTION FOR P-ORBITALS                     HUCK1010
IF (ICORE.EQ.2) A(J,J)=A(J,J)+0.5D0*G(I,I)+G1(ANI)*(DFLOAT(CZ(I)-1
1)/6.0D0-1.0D0)/3.0D0+F2(ANI)*(0.28D0/3.0D0*DFLOAT(CZ(I)-1)-0.08D0*
2(DFLOAT(CZ(I))-2.5D0)) HUCK1020
GO TO 40                                              HUCK1030
30 A(J,J)=-ENEG(ANI,1)/27.21D0                        HUCK1040
PS(J,J)=-3.0D0*(A(J,J)-(DFLOAT(CZ(I))-0.5D0)*G(I,I)) HUCK1050
C MINDG CORRECTION FOR S-ORBITALS                     HUCK1060
C H ATOMS                                              HUCK1070
IF (ICORE.EQ.2.AND.ANI.EQ.1) A(J,J)=A(J,J)+0.5D0*G(I,I) HUCK1080
C OTHER ATOMS THAN H                                  HUCK1090
IF (ICORE.EQ.2.AND.ANI.NE.1) A(J,J)=A(J,J)+0.5D0*G(I,I)+G1(ANI)/12
1.0D0 HUCK1100
40 CCONTINUE                                          HUCK1110
50 CONTINUE                                          HUCK1120
NELECS=NELECS-CHARGE                                HUCK1130
OCCA=NELECS/2                                        HUCK1140
C FORM HUCKEL HAMILTONIAN IN A (OFF DIAGONAL TWO CENTER TERMS) HUCK1150
DO 90 I=2,N                                          HUCK1160
K=U(I)                                              HUCK1170
L=AN(K)                                             HUCK1180
UL=I-1                                             HUCK1190
DO 90 J=1,UL                                       HUCK1200
KK=U(J)                                            HUCK1210
LL=AN(KK)                                          HUCK1220

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	IF ((L.GT.9).OR.(LL.GT.9)) GO TO 80	HUCK1260
	IF (ICORE.NE.2) GO TO 60	HUCK1270
C	MINDO BETAS	HUCK1280
	A(I,J)=A(I,J)*A(J,I)	HUCK1290
	GO TO 70	HUCK1300
60	PS(I,J)=0.75D0*A(I,J)*(PS(I,I)+PS(J,J))/2.0D0	HUCK1310
	A(I,J)=A(I,J)*(BETA0(L)+BETA0(LL))/54.42D0	HUCK1320
	PS(J,I)=PS(I,J)	HUCK1330
70	A(J,I)=A(I,J)	HUCK1340
	GO TO 90	HUCK1350
80	PS(I,J)=0.75D0*A(I,J)*(PS(I,I)+PS(J,J))/2.0D0	HUCK1360
	A(I,J)=0.750D0*A(I,J)*(BETA0(L)+BETA0(LL))/54.42D0	HUCK1370
	PS(J,I)=PS(I,J)	HUCK1380
	A(J,I)=A(I,J)	HUCK1390
90	CONTINUE	HUCK1400
C	OFF DIAG TERMS CALCD BY A WOLFSBERG-HELMHOLTZ METHOD	HUCK1410
	DO 100,I=1,N	HUCK1420
100	Q(I)=A(I,I)	HUCK1430
	RHO=1.D-8	HUCK1440
	CALL EIGN (N,RHO)	HUCK1450
C	EIGENVECTORS(IN B) ARE CONVERTED INTO DENSITY MATRIX(IN B)	HUCK1460
	CALL DEN (OCCA,2,N)	HUCK1470
C	ADD V(AB) TO HCORE--CNDD	HUCK1480
	DO 120 I=1,N	HUCK1490
	J=U(I)	HUCK1500
	Q(I)=Q(I)+0.5D0*G(J,J)	HUCK1510
	DO 110 K=1,NATOMS	HUCK1520
	IF (K.EQ.J) GO TO 110	HUCK1530
	PS(I,I)=PS(I,I)+DFLOAT(CZ(K))*G(J,K)	HUCK1540
110	Q(I)=Q(I)-DFLOAT(CZ(K))*G(J,K)	HUCK1550
	PSQ(I)=PS(I,I)	HUCK1560
120	CONTINUE	HUCK1570

C	EXIT SEGMENT IF ONLY CNDO APPROXIMATIONS ARE DESIRED	HUCK1580
	IF (OPTION.EQ.CNDO) GO TO 200	HUCK1590
C	INDO MODIFICATION(CORRECTION TO U(I,I))	HUCK1600
	DO 190 I=1,NATGMS	HUCK1610
	K=AN(I)	HUCK1620
	J=LLIM(I)	HUCK1630
	IF ((K.GT.1).AND.(K.LT.10)) GO TO 130	HUCK1640
	GO TO 190	HUCK1650
130	IF (K.LE.3) GO TO 140	HUCK1660
	Q(J)=Q(J)+(DFLOAT(CZ(I))-1.500)*G1(K)/6.00	HUCK1670
140	IF (K.EQ.3) GO TO 160	HUCK1680
	IF (K.EQ.4) GO TO 150	HUCK1690
	TEMP=G1(K)/3.00+(DFLOAT(CZ(I))-2.500)*2.00*F2(K)/25.00	HUCK1700
	GO TO 170	HUCK1710
150	TEMP=G1(K)/4.00	HUCK1720
	GO TO 170	HUCK1730
160	TEMP=G1(K)/12.00	HUCK1740
170	CCONTINUE	HUCK1750
	DO 180 L=1,3	HUCK1760
180	Q(J+L)=Q(J+L)+TEMP	HUCK1770
190	CCONTINUE	HUCK1780
200	CCONTINUE	HUCK1790
	DO 220 I=1,N	HUCK1800
	DO 210 J=I,N	HUCK1810
	PSB(I,J)=0.000	HUCK1820
	PSB(J,I)=PSB(I,J)	HUCK1830
210	A(J,I)=A(I,J)	HUCK1840
220	A(I,I)=Q(I)	HUCK1850
	WRITE (6,230)	HUCK1860
	CALL SCFCUT (0,1)	HUCK1870
	WRITE (6,240)	HUCK1880
	CALL SCFCUT (0,3)	HUCK1890


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WRITE (6,250) HUCK1900
RETURN HUCK1910
C HUCK1920
230 FORMAT (1X,28H ELECTRONIC CORE HAMILTONIAN/) HUCK1930
240 FDMAT (1X,29H POSITRONIC CORE HAMILTONIAN /) HUCK1940
250 FORMAT (1X,73H INITIAL ELECTRON DENSITY MATRIX AND EIGENVALUES CAL HUCK1950
ICULATED FROM HCORE(-)/) HUCK1960
END HUCK1970
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	SUBROUTINE DEN (IOCC,MOP,N)	DENS 10
	IMPLICIT REAL*8(A-H,O-Z)	DENS 20
	COMMON /ARRAYS/ A(80,80,4)	DENS 30
	DIMENSION X(400)	DENS 40
	L=1	DENS 50
	IF (IOCC.GT.1) L=2	DENS 60
	DO 40 I=1,N	DENS 70
	DO 20 J=I,N	DENS 80
	X(J)=0.000	DENS 90
	DO 10 K=1,IOCC	DENS 100
10	X(J)=X(J)+L*A(I,K,MOP)*A(J,K,MOP)	DENS 110
20	CONTINUE	DENS 120
	DO 30 J=I,N	DENS 130
30	A(I,J,MOP)=X(J)	DENS 140
40	CONTINUE	DENS 150
	DO 50 I=1,N	DENS 160
	DO 50 J=I,N	DENS 170
50	A(J,I,MOP)=A(I,J,MOP)	DENS 180
	RETURN	DENS 190
	END	DENS 200

	SUBROUTINE SCFCLO	SCFC	10
C	CNDO/INDO CLOSED SHELL SCF SEGMENT	SCFC	20
C	GAMMA MATRIX CONTAINED IN G, CORE HAMILTONIAN CONTAINED IN Q AND	SCFC	30
C	UPPER TRIANGLE OF A, AND INITIAL DENSITY MATRIX CONTAINED IN B	SCFC	40
C	CPTICNS CNDO OR INDO	SCFC	50
	IMPLICIT REAL*8(A-H,O-Z)	SCFC	60
	COMMON /ARRAYS/ A(80,80),B(80,80),PS(80,80),PSB(80,80)	SCFC	70
	COMMON /INFO/ NATOMS,CHARGE,MULTIP,AN(35),C(35,3),N	SCFC	80
	COMMON /GAB/ XXX(400),G(35,35),Q(80),YYY(80),ENERGY,XXY(214)	SCFC	90
	COMMON /INFC1/ CZ(35),U(80),ULIM(35),LLIM(35),NELECS,OCCA,OCCB	SCFC	100
	COMMON /CPTICM/ OPTION,OPNCLO,HUCKEL,CNDO,INDO,CLOSED,OPEN	SCFC	110
	COMMON /CORE/ G1(18),F2(18),ENEG(18,3),BETAG(18)	SCFC	120
	COMMON /PS/ PSQ(80),PSK	SCFC	130
	DIMENSION HELP1(80,80), HELP2(80,80), S(80,80), P(80,80)	SCFC	140
	INTEGER OPTION,OPNCLO,HUCKEL,CNDO,INDO,CLOSED,OPEN	SCFC	150
	INTEGER CHARGE,OCCA,OCCB,UL,ULIM,U,AN,CZ,Z	SCFC	160
	Z=0	SCFC	170
	IT=99	SCFC	180
	RHO=1.0-8	SCFC	190
	DO 20 I=1,N	SCFC	200
	II=U(I)	SCFC	210
	DO 20 K=1,N	SCFC	220
	JJ=U(K)	SCFC	230
20	PS(I,I)=PS(I,I)-B(K,K)*G(II,JJ)	SCFC	240
	DO 30 I=1,N	SCFC	250
	DO 30 J=1,N	SCFC	260
	HELP1(I,J)=A(I,J)	SCFC	270
	HELP2(I,J)=B(I,J)	SCFC	280
	A(I,J)=PS(I,J)	SCFC	290
30	B(I,J)=PSB(I,J)	SCFC	300
	CALL EIGN (N,RHO)	SCFC	310
	CALL DEN (1,2,N)	SCFC	320

	WRITE (6,230)	SCFC 330
	CALL SCFOUT (1,2)	SCFC 340
	DO 40 I=1,N	SCFC 350
	DO 40 J=1,N	SCFC 360
	PS(I,J)=A(I,J)	SCFC 370
	PSB(I,J)=B(I,J)	SCFC 380
	A(I,J)=HELP1(I,J)	SCFC 390
40	B(I,J)=HELP2(I,J)	SCFC 400
50	CONTINUE	SCFC 410
	Z=Z+1	SCFC 420
	ENERGY=0.0D0	SCFC 430
	DO 60 I=1,N	SCFC 440
	A(I,I)=Q(I)	SCFC 450
	PS(I,I)=PSQ(I)	SCFC 460
	DO 60 J=I,N	SCFC 470
	S(I,J)=B(I,J)	SCFC 480
	S(J,I)=S(I,J)	SCFC 490
	P(I,J)=PSB(I,J)	SCFC 500
	P(J,I)=P(I,J)	SCFC 510
60	A(J,I)=A(I,J)	SCFC 520
	DO 70 I=1,N	SCFC 530
	II=U(I)	SCFC 540
	A(I,I)=A(I,I)-B(I,I)*G(II,II)*0.5D0	SCFC 550
	DO 70 K=1,N	SCFC 560
	JJ=U(K)	SCFC 570
	PS(I,I)=PS(I,I)-B(K,K)*G(II,JJ)	SCFC 580
70	A(I,I)=A(I,I)+B(K,K)*G(II,JJ)-PSB(K,K)*G(II,JJ)	SCFC 590
	NM=N-1	SCFC 600
	DO 80 I=1,NM	SCFC 610
	II=U(I)	SCFC 620
	LL=I+1	SCFC 630
	DO 80 J=LL,N	SCFC 640

	JJ=U(J)	SCFC 650
80	A(J,I)=A(J,I)-B(J,I)*G(II,JJ)*0.5D0	SCFC 660
C	INDO MODIFICATION	SCFC 670
	IF (OPTICN.EQ.CND0) GO TO 110	SCFC 680
	DO 100 II=1,NATOMS	SCFC 690
	K=AN(II)	SCFC 700
	I=LLIM(II)	SCFC 710
	IF (K.EQ.1) GO TO 100	SCFC 720
	PAA=B(I,I)+B(I+1,I+1)+B(I+2,I+2)+B(I+3,I+3)	SCFC 730
	A(I,I)=A(I,I)-(PAA-B(I,I))*G1(K)/6.D0	SCFC 740
	DO 90 J=1,3	SCFC 750
	A(I+J,I+J)=A(I+J,I+J)-B(I,I)*G1(K)/6.D0-(PAA-B(I,I))*7.D0*F2(K)/50	SCFC 760
	1.D0+B(I+J,I+J)*11.D0*F2(K)/50.D0	SCFC 770
90	A(I+J,I)=A(I+J,I)+B(I,I+J)*G1(K)/2.D0	SCFC 780
	I1=I+1	SCFC 790
	I2=I+2	SCFC 800
	I3=I+3	SCFC 810
	A(I2,I1)=A(I2,I1)+B(I2,I1)*11.D0*F2(K)/50.D0	SCFC 820
	A(I3,I1)=A(I3,I1)+B(I3,I1)*11.D0*F2(K)/50.D0	SCFC 830
	A(I3,I2)=A(I3,I2)+B(I3,I2)*11.D0*F2(K)/50.D0	SCFC 840
100	CCONTINUE	SCFC 850
110	CGCONTINUE	SCFC 860
	DO 120 I=1,N	SCFC 870
	KK=U(I)	SCFC 880
120	ENERGY=ENERGY+.5D0*B(I,I)*(A(I,I)+Q(I))+.5D0*PSB(I,I)*(PS(I,I)+PSQ	SCFC 890
	1(I))	SCFC 900
	DO 130 I=1,NM	SCFC 910
	II=U(I)	SCFC 920
	LL=I+1	SCFC 930
	DO 130 J=LL,N	SCFC 940
	JJ=U(J)	SCFC 950
130	ENERGY=ENERGY+B(I,J)*(A(I,J)+A(J,I))+PSB(I,J)*(PS(I,J)+PS(J,I))	SCFC 960

	WRITE (6,240) ENERGY	SCFC 970
	IF (DABS(ENERGY-OLDENG).GE.1.00D-6) GO TO 140	SCFC 980
	Z=100	SCFC 990
	WRITE (6,250)	SCFC1000
	GO TO 150	SCFC1010
140	CONTINUE	SCFC1020
	OLDENG=ENERGY	SCFC1030
150	CONTINUE	SCFC1040
	IF (Z.LE.IT) GO TO 170	SCFC1050
C	SYMMETRIZE F FOR PRINTING (MATRIX A)	SCFC1060
	DO 160 I=1,N	SCFC1070
	DO 160 J=1,N	SCFC1080
160	A(I,J)=A(J,I)	SCFC1090
	WRITE (6,260)	SCFC1100
	CALL SCFCOUT (0,1)	SCFC1110
	WRITE (6,270)	SCFC1120
	CALL SCFCOUT (0,3)	SCFC1130
170	CONTINUE	SCFC1140
	CALL EIGN (N,RHO)	SCFC1150
	IF (Z.LT.IT) GO TO 180	SCFC1160
	WRITE (6,280)	SCFC1170
	CALL SCFCOUT (1,2)	SCFC1180
180	DO 190 I=1,N	SCFC1190
	DO 190 J=1,N	SCFC1200
	HELP1(I,J)=A(I,J)	SCFC1210
	HELP2(I,J)=B(I,J)	SCFC1220
190	A(I,J)=PS(I,J)	SCFC1230
	CALL EIGN (N,RHO)	SCFC1240
	DO 200 I=1,N	SCFC1250
	DO 200 J=1,N	SCFC1260
	PS(I,J)=A(I,J)	SCFC1270
	PSB(I,J)=B(I,J)	SCFC1280

	A(I,J)=HELP1(I,J)	SCFC1290
200	B(I,J)=HELP2(I,J)	SCFC1300
	IF (Z.LE.IT) GO TO 210	SCFC1310
	WRITE (6,290)	SCFC1320
	CALL SCFOUT (1,4)	SCFC1330
210	CONTINUE	SCFC1340
C	EIGENVECTORS (IN B) ARE CONVERTED INTO DENSITY MATRIX (IN B)	SCFC1350
	CALL DEN (OCCA,2,N)	SCFC1360
	CALL DEN (1,4,N)	SCFC1370
	DO 220 I=1,N	SCFC1380
	DO 220 J=1,N	SCFC1390
	PSB(I,J)=0.5D0*(PSB(I,J)+P(I,J))	SCFC1400
220	B(I,J)=0.5D0*(B(I,J)+S(I,J))	SCFC1410
	IF (Z.LE.IT) GO TO 50	SCFC1420
	RETURN	SCFC1430
C		SCFC1440
230	FORMAT (1X,69H INITIAL POSITRON DENSITY MATRIX AND EIGENVALUES CALCULATED FROM F(+))	SCFC1450
240	FORMAT (//,10X,21H ELECTRONIC ENERGY ,F16.10)	SCFC1460
250	FORMAT (5X,18H ENERGY SATISFIED /)	SCFC1470
260	FORMAT (1X,36H ELECTRON HARTREE-FOCK ENERGY MATRIX/)	SCFC1480
270	FORMAT (1X,36H POSITRON HARTREE-FOCK ENERGY MATRIX/)	SCFC1490
280	FORMAT (1X,33H ELECTRON EIGENVALUES AND VECTORS/)	SCFC1500
290	FORMAT (1X,33H POSITRON EIGENVALUES AND VECTORS/)	SCFC1510
	END	SCFC1520
		SCFC1530

	SUBROUTINE CPRINT	CPRT 10
	IMPLICIT REAL*8(A-H,O-Z)	CPRT 20
C	CNDC-INDO SCF CLOSED SHELL - PRINTOUT	CPRT 30
	CCOMON /ARRAYS/ A(80,80),B(80,80),D(80,80),PSB(80,80)	CPRT 40
	COMMON /GAB/ XXX(400),G(35,35),Q(80),YYY(80),ENERGY,XXY(214)	CPRT 50
	COMMON /INFO/ NATOMS,CHARGE,MULTIP,AN(35),C(35,3),N	CPRT 60
	COMMON /INFO1/ CZ(35),U(80),ULIM(35),LLIM(35),NELECS,OCCA,OCCB	CPRT 70
	CCOMON /PERTEL/ EL(18)	CPRT 80
	COMMON /OPTICM/ OPTION,OPNCLO,HUCKEL,CNDC,INDO,CLOSED,OPEN	CPRT 90
	CCOMON /ORB1/ ORB(9)	CPRT 100
	INTEGER OPTION,OPNCLO,HUCKEL,CNDC,INDO,CLOSED,OPEN	CPRT 110
	INTEGER CHARGE,OCCA,OCCB,UL,ULIM,U,AN,CZ,Z,EL,ANI,ORB	CPRT 120
	INTEGER ULI,ULJ	CPRT 130
	DIMENSION DPM(3), DM(3), DMSP(3), DMPD(3)	CPRT 140
	DIMENSION ATENG(18)	CPRT 150
	IF (OPTION.EQ.CNDC) GO TO 10	CPRT 160
	ATENG(1)=-0.6387302462D0	CPRT 170
	ATENG(3)=-.2321972405D0	CPRT 180
	ATENG(4)=-1.1219620354D0	CPRT 190
	ATENG(5)=-2.8725750048D0	CPRT 200
	ATENG(6)=-5.9349548261D0	CPRT 210
	ATENG(7)=-10.6731741251D0	CPRT 220
	ATENG(8)=-17.2920850650D0	CPRT 230
	ATENG(9)=-26.2574377875D0	CPRT 240
	GO TO 20	CPRT 250
10	CONTINUE	CPRT 260
	ATENG(1)=-0.6387302462D0	CPRT 270
	ATENG(3)=-.2321972405D0	CPRT 280
	ATENG(4)=-1.1454120355D0	CPRT 290
	ATENG(5)=-2.9774239048D0	CPRT 300
	ATENG(6)=-6.1649936261D0	CPRT 310
	ATENG(7)=-11.0768746252D0	CPRT 320

	ATENG(8)=-18.081965865100	CPRT 330
	ATENG(9)=-27.549130288000	CPRT 340
	ATENG(11)=-.197700956800	CPRT 350
	ATENG(12)=-.867191383300	CPRT 360
	ATENG(13)=-2.036455774400	CPRT 370
	ATENG(14)=-3.897903468600	CPRT 380
	ATENG(15)=-6.796600916300	CPRT 390
	ATENG(16)=-10.765817434100	CPRT 400
	ATENG(17)=-16.046701794000	CPRT 410
20	CONTINUE	CPRT 420
	K=NATOMS-1	CPRT 430
	WRITE (6,160)	CPRT 440
	CALL SCFOUT (0,2)	CPRT 450
	WRITE (6,170)	CPRT 460
	CALL SCFCUT (0,4)	CPRT 470
	IF (NATOMS.EQ.1) GO TO 40	CPRT 480
	DO 30 I=1,K	CPRT 490
	L=I+1	CPRT 500
	DO 30 J=L,NATOMS	CPRT 510
	RAD=DSQRT((C(I,1)-C(J,1))**2+(C(I,2)-C(J,2))**2+(C(I,3)-C(J,3))**2)	CPRT 520
	1)	CPRT 530
30	ENERGY=ENERGY+(DFLOAT(CZ(I))*DFLOAT(CZ(J)))/RAD	CPRT 540
40	WRITE (6,180) ENERGY	CPRT 550
	DO 50 I=1,NATOMS	CPRT 560
	ANI=AN(I)	CPRT 570
50	ENERGY=ENERGY-ATENG(ANI)	CPRT 580
	WRITE (6,190) ENERGY	CPRT 590
	WRITE (6,200)	CPRT 600
	DO 70 I=1,NATOMS	CPRT 610
	TCHG=0.00	CPRT 620
	PSCH=0.000	CPRT 630
	LL=LLIM(I)	CPRT 640

	UL=ULIM(I)	CPRT 650
	DO 60 J=LL,UL	CPRT 660
	PSCH=PSCH+PSB(J,J)	CPRT 670
60	TCHG=TCHG+B(J,J)	CPRT 680
	ANI=AN(I)	CPRT 690
	WRITE (6,210) I,EL(ANI),TCHG,PSCH	CPRT 700
	XXX(I)=TCHG	CPRT 710
70	CONTINUE	CPRT 720
	DO 80 I=1,3	CPRT 730
	DM(I)=0.CDO	CPRT 740
	DMSP(I)=0.DO	CPRT 750
80	DMPD(I)=C.DO	CPRT 760
	DO 140 J=1,NATCMS	CPRT 770
	IF (AN(J).LT.3) GO TO 120	CPRT 780
	IF (AN(J).LT.11) GO TO 100	CPRT 790
	SLTR1=(.65DO*DFLOAT(AN(J))-4.95DO)/3.DO	CPRT 800
	FACTOR=2.5416DO*7.DO/(DSQRT(5.DO)*SLTR1)	CPRT 810
	INDEX=LLIM(J)	CPRT 820
	DO 90 K=1,3	CPRT 830
90	DMSP(K)=DMSP(K)-B(INDEX,INDEX+K)*10.27175DO/SLTR1	CPRT 840
	DMPD(1)=DMPD(1)-FACTOR*(B(INDEX+2,INDEX+8)+B(INDEX+3,INDEX+5)+B(INDEX+4,INDEX+7)-1.DO/DSQRT(3.DO)*B(INDEX+1,INDEX+4))	CPRT 850
	DMPD(2)=DMPD(2)-FACTOR*(B(INDEX+1,INDEX+8)+B(INDEX+3,INDEX+6)+B(INDEX+4,INDEX+7)-1.DO/DSQRT(3.DO)*B(INDEX+2,INDEX+4))	CPRT 860
	DMPD(3)=DMPD(3)-FACTOR*(B(INDEX+1,INDEX+5)+B(INDEX+2,INDEX+6)+2.DO/DSQRT(3.DO)*B(INDEX+3,INDEX+4))	CPRT 870
	GO TO 120	CPRT 880
100	INDEX=LLIM(J)	CPRT 890
	DO 110 K=1,3	CPRT 900
110	DMSP(K)=DMSP(K)-B(INDEX,INDEX+K)*7.33697DO/(.325DO*DFLOAT(AN(J)-1))	CPRT 910
	1)	CPRT 920
120	DO 130 I=1,3	CPRT 930
		CPRT 940
		CPRT 950
		CPRT 960

130	DM(I)=DM(I)+(DFLOAT(CZ(J))-XXX(J))*C(J,I)*2.5416D0	CPRT 970
140	CCONTINUE	CPRT 980
	DO 150 I=1,3	CPRT 990
150	DPM(I)=DM(I)+DMSP(I)+DMPD(I)	CPRT1000
	WRITE (6,220)	CPRT1010
	WRITE (6,230)	CPRT1020
	WRITE (6,240) DM(1),DM(2),DM(3)	CPRT1030
	WRITE (6,250) DMSP(1),DMSP(2),DMSP(3)	CPRT1040
	WRITE (6,260) DMPD(1),DMPD(2),DMPD(3)	CPRT1050
	WRITE (6,270) DPM(1),DPM(2),DPM(3)	CPRT1060
	DP=DSQRT(DPM(1)**2+DPM(2)**2+DPM(3)**2)	CPRT1070
	WRITE (6,280) DP	CPRT1080
	RETURN	CPRT1090
C		CPRT1100
160	FORMAT (1X,24H ELECTRON DENSITY MATRIX//)	CPRT1110
170	FORMAT (1X,24H POSITRON DENSITY MATRIX//)	CPRT1120
180	FCRMT (//,10X,16H TOTAL ENERGY = ,F16.10)	CPRT1130
190	FORMAT (//,10X,16H BINDING ENERGY= ,F16.10,5H A.U./)	CPRT1140
200	FORMAT (10X,32H(-)CHG DENSITY (+)CHG DENSITY)	CPRT1150
210	FCRMT (13,A4,8X,F7.4,11X,F7.4)	CPRT1160
220	FORMAT (//,20X,16H DIPOLE MOMENTS,/)	CPRT1170
230	FCRMT (5X,11H COMPONENTS,3X,2H X,8X,2H Y,8X,2H Z)	CPRT1180
240	FCRMT (5X,10H DENSITIES,3(1X,F9.5))	CPRT1190
250	FORMAT (5X,4H S,P,6X,3(1X,F9.5))	CPRT1200
260	FORMAT (5X,4H P,D,6X,3(1X,F9.5))	CPRT1210
270	FORMAT (5X,6H TOTAL,4X,3(1X,F9.5),/)	CPRT1220
280	FORMAT (3X,15H DIPOLE MOMENT=,F9.5,7H DEBYES,//)	CPRT1230
	END	CPRT1240

C	PROGRAM PAL	PAL	10
C	A PROGRAM FOR THE ANALYSIS OF DECAY CURVES	PAL	20
C	VERSION FOR IBM 360	PAL	30
C	INCLUDING AN ITERATIVE PROCEDURE FOR DETERMINATION OF HALF LIVES	PAL	40
C	WITH REVISED RJT ROUTINE. - CAUSES OLD ROUTINE TO BE USED.	PAL	50
C	USING DOUBLE PRECISION INVERSION ROUTINE BNSINV	PAL	60
	DOUBLE PRECISION B(10,10)	PAL	70
	DOUBLE PRECISION U(10,10),BR(10,10)	PAL	80
	DOUBLE PRECISION T(512),F(512),SFSQ(512),H(10),D(10),A(512,10),APFPAL	PAL	90
	1(10),FCALC(512),SF(200),SIGMAH(10),X(10),SIGMAD(10),DELTAD(10,9),DPAL	PAL	100
	2DELTA(10,9),XPDTO(10),XEOB(10),SXEOB(10),UNITK(10),TAU(10),ANN(10)PAL	PAL	110
	DOUBLE PRECISION BADT(512),BADFS(512),BINVRS(10,10),FKCS(5),SFKCS(PAL	PAL	120
	15),RATIO(512),V(512)	PAL	130
	DOUBLE PRECISION DT,DABS,DSQRT,DEXP	PAL	140
	EQUIVALENCE (BADT,SF), (BADFS,V), (BINVRS,B)	PAL	150
	DIMENSION AA(5), SLOP(5), YCAL(512)	PAL	160
	DIMENSION CCMNT(8), NAME(11), JUNIT(10)	PAL	170
	INTEGER HR,DAY,YR	PAL	180
	INTEGER CMIT	PAL	190
	DATA HR,DAY,YR/1HH,1HD,1HY/	PAL	200
	READ (5,890) COMNT	PAL	210
C	ENTRANCE TO PROGRAM	PAL	220
C	INPUT	PAL	230
	AIM=0.0	PAL	240
20	LT1=0	PAL	250
	READ (5,890) IDENT,NAME,IENDS	PAL	260
	IF (IENDS.EQ.0) GO TO 40	PAL	270
	IF (AIM.EQ.0.0) GO TO 30	PAL	280
	CALL PLOT (0.0,0.0,-4)	PAL	290
30	CALL EXIT	PAL	300
C	NC IS NO. OF COMPONENTS	PAL	310
C	NV IS NO. OF UNKNOWN HALF LIVES	PAL	320

C	NCNV ALWAYS EQUALS ONE	PAL	330
C	CNV GOVERNS HOW FAR ITERATIONS WILL PROCEED; IF CNV IS ZERO A	PAL	340
C	MAX. OF NINE ITERATIONS WILL BE PERFORMED	PAL	350
C	BGD IS THE COUNTER BACKGROUND FOR SUBTRACTION	PAL	360
C	SBGD IS THE STANDARD DEVIATION FOR ROOT-MEAN-SQUARE ADDITION TO	PAL	370
C	THE STANDARD DEVIATION OF EACH POINT (USUALLY ZERO)	PAL	380
C	IN... INPUT CONTROL (ZERO OR BLANK CAUSES USE OF PROGRAM INPUT	PAL	390
C	FORMAT; POSITIVE INTEGER TRANSFERS TO SUBROUTINE INPUT)	PAL	400
C	IT...NUMBER HERE CAUSES INTERMEDIATE MATRICES TO BE PRINTED	PAL	410
C	(USUALLY LEFT BLANK)	PAL	420
C	BLOCK IS THE COUNTER DEAD TIME (USUALLY LEFT BLANK)	PAL	430
C	SCOFF IS CUTOFF (IN PERCENT) ON THE SMALLEST VALUE AT WHICH THE	PAL	440
C	PROGRAM WILL USE THE STANDARD DEVN. FROM STATS. ALONE (0.5)	PAL	450
C	RJT...IF NOT BLANK OR ZERO PROGRAM EXAMINES ITS OUTPUT AND	PAL	460
C	REJECTS POINTS WHICH FALL FURTHER THAN RJT TIMES THE STANDARD	PAL	470
C	DEVN. FROM THE CURVE. THE FIT IS THEN REPEATED	PAL	480
C	KCS...CAUSES THE PROGRAM TO ENTER A KNOWN COMPT. SUBTRACTION	PAL	490
C	ROUTINE (USUALLY SET AT ZERO)	PAL	500
C	ICARD IF UNEQUAL TO ZERO CAUSES PROGRAM TO ENTER INTENSITY	PAL	510
C	ROUTINE (USUALLY EQUALS ZERO)	PAL	520
C	INP...IF EQUAL TO ZERO OR BLANK ALLOWS HALF-LIFE GUESSES TO BE	PAL	530
C	ENTERED AS HALF-LIVES AFTER TIME-PER-CHANNEL CORRECTION; IF	PAL	540
C	EQUAL TO ONE ALLOWS TAU GUESSES TO BE USED	PAL	550
C	OMIT...DATA OUTPUT OPTIONS	PAL	560
C	OMIT=1:OUTPUT INCLUDES ALL DATA, INCLUDES GRAPH	PAL	570
C	OMIT=2:OMITS SUPERFLUOUS DATA, INCLUDES GRAPH	PAL	580
C	OMIT=3:INCLUDES ALL DATA, OMITS GRAPH	PAL	590
C	OMIT=4: OMITS SUPERFLUOUS DATA, OMITS GRAPH	PAL	600
40	READ (5,900) NC,NV,NCNV,CNV,BGD,SBGD,IN,IT,BLOCK,SCOFF,RJT,KCS,ICARD	PAL	610
	1RD,INP,OMIT	PAL	620
	IF (OMIT) 50,50,60	PAL	630
50	WRITE (6,910)	PAL	640

	GO TO 30	PAL 650
60	IF (OMIT-5) 70,50,50	PAL 660
70	IF (INP.EQ.0) GO TO 90	PAL 670
	READ (5,920) CALIB	PAL 680
	DO 80 I=1,NC	PAL 690
	READ (5,920) TAU(I),JUNIT(I)	PAL 700
	H(I)=(TAU(I)*0.69314718)/CALIB	PAL 710
80	D(I)=0.69314718/H(I)	PAL 720
	GO TO 110	PAL 730
90	DO 100 I=1,NC	PAL 740
	READ (5,920) H(I),JUNIT(I)	PAL 750
C	HALF LIFE INPUT CONVERSION. H,D, OR Y ARE RECOGNIZED AS SUCH.	PAL 760
C	OTHER CHARACTERS ARE TAKEN TO MEAN M.	PAL 770
	UNITK(I)=0.69314718D0	PAL 780
	IF (JUNIT(I).EQ.HR) UNITK(I)=0.11552453D-1	PAL 790
	IF (JUNIT(I).EQ.DAY) UNITK(I)=0.48135221D-3	PAL 800
	IF (JUNIT(I).EQ.YR) UNITK(I)=0.13178987D-5	PAL 810
	D(I)=UNITK(I)/H(I)	PAL 820
100	CONTINUE	PAL 830
110	NP=0	PAL 840
	NSCHAN=0	PAL 850
C	INPUT OPTION	PAL 860
	IF (IN.EQ.0) GO TO 120	PAL 870
	CALL INPUT (T,F,SFSQ,NP,BGD,NSCHAN)	PAL 880
	GO TO 130	PAL 890
120	NP=NP+1	PAL 900
	READ (5,930) T(NP),F(NP),SFSQ(NP)	PAL 910
	IF (SFSQ(NP).NE.0.) GO TO 120	PAL 920
C	TIME NORMALIZATION	PAL 930
130	DO 140 I=2,NP	PAL 940
	T(I)=T(I)-T(1)	PAL 950
140	CGCONTINUE	PAL 960

	T(1)=0.000	PAL 970
	TO=T(NP)	PAL 980
	NP=NP-1	PAL 990
	NP1=NP	PAL 1000
	NC1=NC	PAL 1010
	IF (BLOCK.EQ.0.) GO TO 160	PAL 1020
C	DEAD TIME CORRECTION	PAL 1030
	DEADT=BLOCK*(1.67E-8)	PAL 1040
	DO 150 I=1,NP	PAL 1050
150	F(I)=F(I)/(1.-DEADT*F(I))	PAL 1060
160	IF (BGD.LE.0.) GO TO 180	PAL 1070
	SBGDSQ=SBGD**2	PAL 1080
	DO 170 I=1,NP	PAL 1090
	F(I)=F(I)-BGD	PAL 1100
170	SFSQ(I)=SFSQ(I)+SBGDSQ	PAL 1110
C	SUBROUTINE FOR SIGMA CUTOFF	PAL 1120
180	IF (SCOFF.EQ.0.) GO TO 200	PAL 1130
	SCOFSQ=SCOFF**2/10000.	PAL 1140
	DO 190 I=1,NP	PAL 1150
	TSFSQ=SCOFSQ*F(I)**2	PAL 1160
	IF (SFSQ(I).LT.TSFSQ) SFSQ(I)=TSFSQ	PAL 1170
190	CONTINUE	PAL 1180
C	KNOWN COMPONENT SUBTRACTION (KCS) ROUTINE	PAL 1190
200	IF (KCS.EQ.0) GO TO 240	PAL 1200
	READ (5,940) (FKCS(I),SFKCS(I),I=1,KCS)	PAL 1210
	DO 230 J=1,KCS	PAL 1220
	ND=NC-KCS+J	PAL 1230
	DO 210 I=1,NP	PAL 1240
	FACTOR=DEXP((TO-T(I))*D(ND))	PAL 1250
	IF ((F(I)-(FKCS(J)*FACTOR)).LE.0.000) GO TO 220	PAL 1260
	F(I)=F(I)-FKCS(J)*FACTOR	PAL 1270
	SFSQ(I)=SFSQ(I)+(SFKCS(J)*FACTOR)**2	PAL 1280

210	CONTINUE	PAL 1290
220	NP=NP-1	PAL 1300
230	CONTINUE	PAL 1310
	NC=NC-KCS	PAL 1320
C	INITIALIZATION AND RETURNS	PAL 1330
240	NITER=-1	PAL 1340
	IF (NV.EQ.0) GO TO 250	PAL 1350
	LT1=1	PAL 1360
	NVTEMP=NV	PAL 1370
	NV=0	PAL 1380
C	RETURN AFTER FIXED HALF LIFE PASS	PAL 1390
250	NT=NC+NV	PAL 1400
	LV=NC+1	PAL 1410
C	CALCULATION OF A MATRIX	PAL 1420
C	ALSO RETURN FOR NEXT ITERATION	PAL 1430
260	NITER=NITER+1	PAL 1440
	DO 300 I=1,NP	PAL 1450
	DO 280 J=1,NC	PAL 1460
	DT=-D(J)*T(I)	PAL 1470
	IF ((DT).GT.(-34.5D0)) GO TO 270	PAL 1480
	A(I,J)=0.0D0	PAL 1490
	GO TO 280	PAL 1500
270	A(I,J)=DEXP(DT)	PAL 1510
280	CONTINUE	PAL 1520
	IF (NV.LE.0) GO TO 300	PAL 1530
	DO 290 J=LV,NT	PAL 1540
	K=J-NC	PAL 1550
290	A(I,J)=T(I)*A(I,K)	PAL 1560
300	CONTINUE	PAL 1570
	IF (IT) 310,320,310	PAL 1580
310	WRITE (6,950) ((A(I,J),J=1,NT),I=1,2)	PAL 1590
C	CONSTRUCT MATRIX APF	PAL 1600

C	ALSO RETURN AFTER DATA REJECTION	PAL 1610
320	DO 330 I=1,NT	PAL 1620
	APF(I)=0.000	PAL 1630
	DO 330 K=1,NP	PAL 1640
	APF(I)=APF(I)+A(K,I)*F(K)/SFSQ(K)	PAL 1650
330	CONTINUE	PAL 1660
	IF (IT) 340,350,340	PAL 1670
340	WRITE (6,950) (APF(I),I=1,NT)	PAL 1680
C	CONSTRUCT MATRIX B	PAL 1690
350	DO 360 I=1,NT	PAL 1700
	DO 360 J=1,NT	PAL 1710
	B(I,J)=0.000	PAL 1720
	DO 360 K=1,NP	PAL 1730
	B(I,J)=B(I,J)+A(K,I)*A(K,J)/SFSQ(K)	PAL 1740
360	CONTINUE	PAL 1750
	IF (IT) 370,390,370	PAL 1760
370	WRITE (6,950) ((B(I,J),J=1,NT),I=1,NT)	PAL 1770
	DO 380 KI=1,NT	PAL 1780
	DO 380 KJ=1,NT	PAL 1790
380	BR(KI,KJ)=B(KI,KJ)	PAL 1800
C	B MATRIX INVERSION	PAL 1810
390	EPSIL=0.000	PAL 1820
	CALL BNSINV (B,NT,DETERM,EPSIL,ITEST)	PAL 1830
C	B MATRIX IS NOW REPLACED BY ITS INVERSE, BINVRS.	PAL 1840
	IF (ITEST.NE.0) GO TO 580	PAL 1850
	IF (IT) 400,420,400	PAL 1860
400	WRITE (6,950) ((BINVRS(I,J),J=1,NT),I=1,NT)	PAL 1870
	DO 410 I=1,NT	PAL 1880
	DO 410 J=1,NT	PAL 1890
	U(I,J)=0.000	PAL 1900
	DO 410 K=1,NT	PAL 1910
410	U(I,J)=U(I,J)+BINVRS(I,K)*BR(K,J)	PAL 1920

	WRITE (6,950) ((U(I,J),J=1,NT),I=1,NT)	PAL 1930
C	SOLUTION OF NORMAL EQUATIONS	PAL 1940
420	DO 430 I=1,NT	PAL 1950
	X(I)=0.000	PAL 1960
	DO 430 J=1,NT	PAL 1970
	X(I)=X(I)+BINVRS(I,J)*APF(J)	PAL 1980
430	CONTINUE	PAL 1990
C	ENTRANCE TO CALCULATIONS FOR OUTPUT(127) OR ITERATION TESTS(301)	PAL 2000
	IF (NV.NE.0) GO TO 500	PAL 2010
C	CALCULATIONS FOR OUTPUT DATA	PAL 2020
440	DO 450 I=1,NC	PAL 2030
	SIGMAH(I)=0.000	PAL 2040
450	XPDTO(I)=DEXP(-TO*D(I))	PAL 2050
C	HALF LIFE OUTPUT CONVERSION	PAL 2060
	IF (NV.EQ.0) GO TO 580	PAL 2070
	IF (INP.EQ.0) GO TO 470	PAL 2080
	DO 460 I=1,NV	PAL 2090
	H(I)=0.69314718/D(I)	PAL 2100
460	SIGMAH(I)=H(I)*SIGMAD(I)/D(I)	PAL 2110
	GO TO 490	PAL 2120
470	DO 480 I=1,NV	PAL 2130
	H(I)=UNITK(I)/D(I)	PAL 2140
480	SIGMAH(I)=H(I)*SIGMAD(I)/D(I)	PAL 2150
490	GO TO 580	PAL 2160
C	TESTS FOR COMPLETION OF ITERATION	PAL 2170
500	NON=1	PAL 2180
	DO 520 I=1,NV	PAL 2190
	L=I+NC	PAL 2200
	DELTAD(I,NITER)=-X(L)/X(I)	PAL 2210
	DDELTA(I,NITER)=DABS(DELTAD(I,NITER))	PAL 2220
510	IF (DABS(DELTAD(I,NITER)).LT.0.5DO*DABS(D(I))) GO TO 520	PAL 2230
C	CHANGE IN THE DECAY CONSTANT CANNOT EXCEED HALF ITS VALUE	PAL 2240

	DELTA(I,NITER)=0.75D0*DELTA(I,NITER)	PAL 2250
	GO TO 510	PAL 2260
520	D(I)=D(I)+DELTA(I,NITER)	PAL 2270
	IF (NITER.LE.2) GO TO 260	PAL 2280
	IF (NCNV.GT.0) GO TO 540	PAL 2290
	DO 530 I=1,NV	PAL 2300
	IF (DDELTA(I,NITER).LE.DDELTA(I,NITER-1)) GO TO 530	PAL 2310
	IF (DDELTA(I,NITER-1).GT.DDELTA(I,NITER-2)) GO TO 550	PAL 2320
530	CONTINUE	PAL 2330
540	NON=0	PAL 2340
550	DO 560 I=1,NV	PAL 2350
	L=I+NC	PAL 2360
560	SIGMAD(I)=DSQRT(BINVRS(L,L))/DABS(X(I))	PAL 2370
	IF (NON.EQ.1) GO TO 440	PAL 2380
	IF (NITER.GE.9) GO TO 440	PAL 2390
	DO 570 I=1,NV	PAL 2400
	IF (DABS(DELTA(I,NITER)).GE.CNV*SIGMAD(I)) GO TO 260	PAL 2410
570	CONTINUE	PAL 2420
	GO TO 440	PAL 2430
C	PRINT OUT RESULTS	PAL 2440
580	WRITE (6,960) COMNT	PAL 2450
	WRITE (6,970) IDENT,NAME	PAL 2460
	IF (ITEST.EQ.0) GO TO 590	PAL 2470
	WRITE (6,980)	PAL 2480
	GO TO 20	PAL 2490
590	WRITE (6,990) NP,NC,NV,CNV,BGD,SBGD,BLOCK,SCOFF,RJT,KCS	PAL 2500
	IF (NV.EQ.0) GO TO 630	PAL 2510
	IF (NON.EQ.1) GO TO 600	PAL 2520
	WRITE (6,1000) NITER	PAL 2530
	GO TO 610	PAL 2540
600	WRITE (6,1010) NITER	PAL 2550
610	WRITE (6,1020)	PAL 2560

	WRITE (6,1030) (D(I),I=1,NV)	PAL 2570
	DO 620 J=1,NITER	PAL 2580
	WRITE (6,1040) J,(DELTAD(I,J),I=1,NV)	PAL 2590
620	CCONTINUE	PAL 2600
	WRITE (6,1050) (SIGMAD(I),I=1,NV)	PAL 2610
C	CONVERSICN TO EOB	PAL 2620
630	DO 640 I=1,NC	PAL 2630
	XEOB(I)=X(I)*XPDTO(I)	PAL 2640
	SXEOB(I)=BINVRS(I,I)	PAL 2650
	IF (I.GT.NV) GO TO 640	PAL 2660
	L=I+NC	PAL 2670
	SXEOB(I)=SXEOB(I)+2.*TO*BINVRS(I,L)+TO**2*BINVRS(L,L)	PAL 2680
640	SXEOB(I)=XPDTO(I)*DSQRT(DABS(SXEOB(I)))	PAL 2690
	DO 650 I=1,NC	PAL 2700
650	TAU(I)=(H(I)*CALIB)/0.69314718	PAL 2710
	WRITE (6,1060)	PAL 2720
	WRITE (6,1070) (I,H(I),JUNIT(I),SIGMAH(I),JUNIT(I),XEOB(I),SXEOB(I)	PAL 2730
	1),XPDTO(I),TAU(I),I=1,NC)	PAL 2740
	IF (KCS.EQ.0) GO TO 670	PAL 2750
	I=NC	PAL 2760
	J=0	PAL 2770
660	I=I+1	PAL 2780
	J=J+1	PAL 2790
	WRITE (6,1080) I,H(I),JUNIT(I),FKCS(J),SFKCS(J)	PAL 2800
	IF (KCS.GT.J) GO TO 660	PAL 2810
670	VPV=0.0	PAL 2820
	DO 690 I=1,NP	PAL 2830
	FCALC(I)=0.0DO	PAL 2840
	DO 680 J=1,NC	PAL 2850
	FCALC(I)=FCALC(I)+A(I,J)*X(J)	PAL 2860
680	CONTINUE	PAL 2870
	V(I)=F(I)-FCALC(I)	PAL 2880

	SF(I)=DSQRT(SFSQ(I))	PAL 2890
	RATIO(I)=V(I)/SF(I)	PAL 2900
	VPV=VPV+V(I)**2/SFSQ(I)	PAL 2910
690	CONTINUE	PAL 2920
	DF=NP-NT	PAL 2930
	FIT=SQRT(VPV/DF)	PAL 2940
	WRITE (6,1090) FIT	PAL 2950
	GO TO (700,710,700,710), OMIT	PAL 2960
700	WRITE (6,1100)	PAL 2970
	WRITE (6,1110) (T(I),F(I),FCALC(I),V(I),SF(I),RATIO(I),I=1,NP)	PAL 2980
710	IF (RJT.EQ.0.) GO TO 800	PAL 2990
	IF (NV.NE.0) GO TO 800	PAL 3000
C	DATA REJECTION SUBROUTINE	PAL 3010
	I=0	PAL 3020
	NRJT=0	PAL 3030
720	I=I+1	PAL 3040
730	IF (I.GT.NP) GO TO 780	PAL 3050
	IF (DABS(RATIO(I)).LE.ABS(RJT)) GO TO 720	PAL 3060
	IF (RJT.LT.0.) GO TO 760	PAL 3070
	IF (I.EQ.NP) GO TO 750	PAL 3080
	IF (RATIO(I)*RATIO(I+1)) 740,740,720	PAL 3090
740	IF (I.EQ.1) GO TO 760	PAL 3100
750	IF (RATIO(I)*RATIO(I-1)) 760,760,720	PAL 3110
760	NRJT=NRJT+1	PAL 3120
	BADT(NRJT)=T(I)	PAL 3130
	BADFS(NRJT)=F(I)	PAL 3140
	NP=NP-1	PAL 3150
C	MOVE UP REMAINING POINTS	PAL 3160
	DO 770 J=I,NP	PAL 3170
	T(J)=T(J+1)	PAL 3180
	F(J)=F(J+1)	PAL 3190
	SFSQ(J)=SFSQ(J+1)	PAL 3200

	RATIO(J)=RATIO(J+1)	PAL 3210
	DO 770 K=1,NT	PAL 3220
	A(J,K)=A(J+1,K)	PAL 3230
770	CONTINUE	PAL 3240
	GO TO 730	PAL 3250
780	IF (NRJT.NE.0) GO TO 790	PAL 3260
	WRITE (6,1120)	PAL 3270
	GO TO 800	PAL 3280
790	WRITE (6,1130) NRJT	PAL 3290
	WRITE (6,1140) (BADT(I),BADFS(I),I=1,NRJT)	PAL 3300
	WRITE (6,1150)	PAL 3310
	GO TO 320	PAL 3320
800	IF (LT1.NE.1) GO TO 810	PAL 3330
	LT1=0	PAL 3340
	NV=NVTEMP	PAL 3350
	GO TO 250	PAL 3360
810	GO TO (820,820,830,830), OMIT	PAL 3370
820	CALL POL (F,FCALC,NP,IDENT)	PAL 3380
	AIM=1.0	PAL 3390
830	IF (ICARD.EQ.0) GO TO 20	PAL 3400
	T(1)=-T0	PAL 3410
	STX=T(1)	PAL 3420
	DO 840 L=2,NP1	PAL 3430
	T(L)=T(L-1)+1.0DO	PAL 3440
840	CONTINUE	PAL 3450
	DO 850 K=1,NP1	PAL 3460
	FCALC(K)=0.0DO	PAL 3470
	DO 850 J=1,NC1	PAL 3480
	FCALC(K)=FCALC(K)+XEOB(J)*DEXP(-D(J)*T(K))	PAL 3490
850	CONTINUE	PAL 3500
	DO 860 I=1,NP1	PAL 3510
	YCAL(I)=FCALC(I)	PAL 3520

860	CONTINUE	PAL 3530
	DO 870 I=1,NC1	PAL 3540
	AA(I)=XEOB(I)	PAL 3550
	SLOP(I)=D(I)	PAL 3560
870	CONTINUE	PAL 3570
	CALL INTENS (IDENT,NAME,BGD,STX,NC,NP,AA,SLOP,YCAL,OMIT)	PAL 3580
C	CAL. OF ANNIHILATION RATES	PAL 3590
	DO 880 I=1,NC	PAL 3600
	ANN(I)=1.0/TAU(I)	PAL 3610
880	CONTINUE	PAL 3620
	WRITE (6,1160) (TAU(I),I=1,NC)	PAL 3630
	WRITE (6,1170) (ANN(I),I=1,NC)	PAL 3640
	WRITE (6,1180) CALIB	PAL 3650
	WRITE (6,1190) FIT	PAL 3660
	GO TO 20	PAL 3670
C		PAL 3680
890	FORMAT (A4,11A4,I2)	PAL 3690
900	FORMAT (2I3,I2,F6.4,2F8.2,2I2,3F5.2,4I2)	PAL 3700
910	FORMAT (1H1,49H PROGRAM FAILURE,OMIT ERROR: DATA CANCELLED)	PAL 3710
920	FORMAT (F10.5,A1)	PAL 3720
930	FORMAT (3E15.8)	PAL 3730
940	FORMAT (F16.8,F14.8)	PAL 3740
950	FORMAT (/1H ,9E14.7)	PAL 3750
960	FORMAT (1H1,12A4)	PAL 3760
970	FORMAT (1H0,A4,11A4)	PAL 3770
980	FORMAT (50HOMATRIX INVERSION HAS FAILED. MORE DATA IS NEEDED.)	PAL 3780
990	FORMAT (//4H NP=,I3,4H NC=,I2,4H NV=,I1,5H CNV=,F6.4,5H BGD=,F8.2,	PAL 3790
	16H SBGD=,F8.2,7H BLOCK=,F5.2,7H SCOFF=,F5.2,5H RJT=,F6.4,5H KCS=,I	PAL 3800
	22)	PAL 3810
1000	FORMAT (//25H ITERATIONS PERFORMED=,I2,15H CONVERGENT)	PAL 3820
1010	FORMAT (//25H ITERATIONS PERFORMED=,I2,15H NONCONVERGENT)	PAL 3830
1020	FORMAT (//86H 1ST COMP 2ND COMP 3RD COMP	PAL 3840

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1P          4TH COMP          5TH COMP)          PAL 3850
1030  FORMAT (10H  D          ,5E16.8)          PAL 3860
1040  FORMAT (7H DELTA(,I2,1H),5E16.8)          PAL 3870
1050  FORMAT (10H SIGMA          ,5E16.8)          PAL 3880
1060  FORMAT (//99H          HALF LIFE  SIGMA H          CPM AT EOBPAL 3890
1          SIGMA          DECAY FACTOR          TAU)          PAL 3900
1070  FORMAT (6H COMP(,I2,2H) ,F13.8,A1,F13.8,A1,3E17.8,F15.8)          PAL 3910
1080  FORMAT (6H COMP(,I2,2H) ,F10.5,A1,9X,2E16.8,17H KNOWN COMPONENT) PAL 3920
1090  FORMAT (//8H  FIT=,F10.5//)          PAL 3930
1100  FORMAT (98H          T(I)          F(I)          FCALC(I)          PAL 3940
1          V(I)          SIGMAF(I)          RATIO(I))          PAL 3950
1110  FORMAT (1H ,5E17.8,2X,F10.5)          PAL 3960
1120  FORMAT (20H OALL DATA POINTS OK.)          PAL 3970
1130  FORMAT (22H ODATA POINTS REJECTED(,I3,18H) ARE GIVEN BELOW./28H          PAL 3980
1  BAD T          BAD F)          PAL 3990
1140  FORMAT (1H ,2E16.8)          PAL 4000
1150  FORMAT (41H OREPEAT CALCULATION WITHOUT THESE POINTS.)          PAL 4010
1160  FORMAT (//5X,24H LIFETIME (TAU) X 10-9SEC,4X,5F16.8)          PAL 4020
1170  FORMAT (//5X,28H ANNIHILATION RATE X 10+9SEC,5F16.8)          PAL 4030
1180  FORMAT (//37H  CALIBRATION (10-9SEC PER CHAN.) =,F10.5)          PAL 4040
1190  FORMAT (//16H          FIT=,F10.5)          PAL 4050
END          PAL 4060
SUBROUTINE INPUT (T,F,SFSQ,NP,BGD,NSCHAN)          INPT 10
C  INPUT REVISED FOR POSITRON LIFETIME DATA 4/30/70          INPT 20
C  NZCHAN EQUALS ZERO CHANNEL NUMBER          INPT 30
C  NSCHAN EQUALS STARTING CHANNEL NO. FOR DATA USED IN CALC LIFETIME INPT 40
C  NP EQUALS NO. OF DATA POINTS USED IN CALC TAU          INPT 50
DOUBLE PRECISION T(512),F(512),SFSQ(512)          INPT 60
READ (5,50) NZCHAN,NSCHAN,NP,CONTR          INPT 70
IF (CONTR.NE.0) GO TO 10          INPT 80
READ (5,60) (F(I),I=1,NP)          INPT 90
GO TO 20          INPT 100

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10	READ (5,70) (F(I),I=1,NP)	INPT 110
20	DO 30 I=1,NP	INPT 120
30	SFSQ(I)=F(I)	INPT 130
	NP=NP+1	INPT 140
	DO 40 I=1,NP	INPT 150
40	T(I)=NSCHAN-1+I	INPT 160
	T(NP)=NZCHAN	INPT 170
	RETURN	INPT 180
C		INPT 190
50	FORMAT (4I5)	INPT 200
60	FORMAT (F6.0,9F7.0)	INPT 210
70	FORMAT (F9.0,7F10.0)	INPT 220
	END	INPT 230
	SUBROUTINE BNSINV (A,N,DETERM,EPSIL,ITEST)	BNSI 10
C	MATRIX INVERSION SUBROUTINE	BNSI 20
C	CALLING SEQUENCE CALLBNSINV(A,N,DETERM,EPSIL,ITEST)	BNSI 30
C	A IS ORIGINAL MATRIX (REPLACED BY A-INVERSE AT END OF PROGRAM)	BNSI 40
C	EL IS A 10 BY 10 ARRAY USED ONLY FOR WORKING STORAGE	BNSI 50
C	N IS DIMENSION OF MATRIX	BNSI 60
C	DETERM IS LOCATION OF DETERMINANT	BNSI 70
C	EPSIL IS USED TO TEST WHETHER MATRIX IS SINGULAR.	BNSI 80
C	IF,AFTER TRIANGULARIZING THE MATRIX, THE RATIO OF THE	BNSI 90
C	SMALLEST TO LARGEST DIAGONAL ELEMENT IS LESS THEN EPSIL	BNSI 100
C	THE MATRIX IS CONSIDERED SINGULAR. IF ONE SETS EPSIL=0.	BNSI 110
C	THE ROUTINE WILL USE EPSIL=1.D-10	BNSI 120
C	ITEST IS AN ERROR INDICATOR.	BNSI 130
C	IF ITEST = 0 ALL WENT WELL	BNSI 140
C	IF ITEST = 1 MATRIX WAS FOUND TO BE SINGULAR	BNSI 150
C	IF ITEST =-1 DIVIDE CHECK OCCURRED	BNSI 160
	DOUBLE PRECISION A(10,10),EL(10,10),DABS,DSQRT	BNSI 170
	DOUBLE PRECISION D,S,C,DETERM,DETERP,EPSIL,EPSILP	BNSI 180
	DOUBLE PRECISION E	BNSI 190

	IF (N.GT.1) GO TO 10	BNSI 200
C	SPECIAL CASE FOR N=1	BNSI 210
	A(1,1)=1.000/A(1,1)	BNSI 220
	GO TO 170	BNSI 230
C	SET EL = IDENTITY MATRIX	BNSI 240
10	DO 20 J=1,N	BNSI 250
	DO 20 I=1,N	BNSI 260
	IF (I.EQ.J) EL(I,J)=1.000	BNSI 270
	IF (I.NE.J) EL(I,J)=0.000	BNSI 280
20	CONTINUE	BNSI 290
C	TRIANGULARIZE A, FORM EL	BNSI 300
	N1=N-1	BNSI 310
	M=2	BNSI 320
	DO 60 J=1,N1	BNSI 330
	DO 50 I=M,N	BNSI 340
	IF (A(I,J).EQ.0.) GO TO 50	BNSI 350
	D=DSQRT(A(J,J)**2+A(I,J)**2)	BNSI 360
	C=A(J,J)/D	BNSI 370
	S=A(I,J)/D	BNSI 380
	DO 30 K=J,N	BNSI 390
	D=C*A(J,K)+S*A(I,K)	BNSI 400
	A(I,K)=C*A(I,K)-S*A(J,K)	BNSI 410
	A(J,K)=D	BNSI 420
30	CONTINUE	BNSI 430
	DO 40 K=1,N	BNSI 440
	D=C*EL(J,K)+S*EL(I,K)	BNSI 450
	EL(I,K)=C*EL(I,K)-S*EL(J,K)	BNSI 460
	EL(J,K)=D	BNSI 470
40	CONTINUE	BNSI 480
50	CONTINUE	BNSI 490
60	M=M+1	BNSI 500
C	CALCULATE DETERMINANT	BNSI 510

	DETERP=A(1,1)	BNSI 520
	DO 70 I=2,N	BNSI 530
70	DETERP=DETERP*A(I,I)	BNSI 540
	DETERM=DETERP	BNSI 550
C	IS MATRIX SINGULAR	BNSI 560
	D=A(1,1)	BNSI 570
	E=A(1,1)	BNSI 580
	DO 80 I=2,N	BNSI 590
	IF (DABS(D).LT.DABS(A(I,I))) D=A(I,I)	BNSI 600
	IF (DABS(E).GT.DABS(A(I,I))) E=A(I,I)	BNSI 610
80	CONTINUE	BNSI 620
	EPSILP=EPSIL	BNSI 630
	IF (EPSILP.EQ.0.) EPSILP=1.D-10	BNSI 640
	IF (D.EQ.0.) GO TO 190	BNSI 650
	IF (DABS(E/D).LT.DABS(EPSILP)) GO TO 190	BNSI 660
C	INVERT TRIANGULAR MATRIX	BNSI 670
	J=N	BNSI 680
	DO 120 J1=1,N	BNSI 690
	LGT=1	BNSI 700
	I=J	BNSI 710
	DO 110 I1=1,J	BNSI 720
	IF (LGT.EQ.0) GO TO 90	BNSI 730
	LGT=0	BNSI 740
	IF (A(I,I).EQ.0.0) GO TO 200	BNSI 750
	A(I,J)=1.0D0/A(I,I)	BNSI 760
	GO TO 110	BNSI 770
90	KS=I+1	BNSI 780
	D=0.0D0	BNSI 790
	DO 100 K=KS,J	BNSI 800
100	D=D+A(I,K)*A(K,J)	BNSI 810
	IF (A(I,I).EQ.0.0) GO TO 200	BNSI 820
	A(I,J)=-D/A(I,I)	BNSI 830

110	I=I-1	BNSI 840
120	J=J-1	BNSI 850
C	PREMULTIPLY EL BY INVERTED TRIANGULAR MATRIX	BNSI 860
	M=1	BNSI 870
	DO 150 I=1,N	BNSI 880
	DO 140 J=1,N	BNSI 890
	D=0.000	BNSI 900
	DO 130 K=M,N	BNSI 910
130	D=D+A(I,K)*EL(K,J)	BNSI 920
	EL(I,J)=D	BNSI 930
140	CONTINUE	BNSI 940
150	M=M+1	BNSI 950
C	RECOPY EL TO A	BNSI 960
	DO 160 I=1,N	BNSI 970
	DO 160 J=1,N	BNSI 980
160	A(I,J)=EL(I,J)	BNSI 990
170	ITEST=0	BNSI1000
180	RETURN	BNSI1010
190	ITEST=1	BNSI1020
	GO TO 180	BNSI1030
200	ITEST=-1	BNSI1040
	GO TO 180	BNSI1050
	END	BNSI1060
	SUBROUTINE POL (F,FCALC,NP,IDENT)	POL 10
C	PLOT PS ANN. TIME SPECTRUM:THEORETICAL AND EXP. RESULTS	POL 20
	DIMENSION X(512), Y(1024)	POL 30
	DOUBLE PRECISION F(512),FCALC(512)	POL 40
	IA=NP+1	POL 50
	IB=2*NP	POL 60
	IC=1	POL 70
	DO 10 I=1,NP	POL 80
	X(I)=I	POL 90

10	CONTINUE	POL	100
	DO 20 I=1,NP	POL	110
	Y(I)=F(I)	POL	120
20	CONTINUE	POL	130
	DO 30 J=IA,IB	POL	140
	Y(J)=FCALC(J-NP)	POL	150
30	CONTINUE	POL	160
	DO 40 J=1,IB	POL	170
	IF (Y(J).LT.1.0) Y(J)=1.0	POL	180
	Y(J)=ALOG(Y(J))	POL	190
40	CONTINUE	POL	200
	CALL SCALE (X(IC),NP,7.0,XMIN,DX,1)	POL	210
	CALL SCALE (Y(IC),IB,4.0,YMIN,DY,1)	POL	220
	CALL PLOT (2.0,2.0,-3)	POL	230
	CALL AXIS (0.0,0.0,14HCHANNEL NUMBER,14,-7.0,0.0,XMIN,DX)	POL	240
	CALL AXIS (0.0,0.0,14H LOG(COUNTS),-14,-4.0,90.0,YMIN,DY)	POL	250
	CALL PLOT (0.0,4.0,3)	POL	260
	CALL PLOT (7.0,4.0,2)	POL	270
	CALL PLOT (7.0,0.0,2)	POL	280
	CALL SYMBOL (5.0,3.5,.25,IDENT,0.0,4)	POL	290
	CALL PLOT (0.0,0.0,3)	POL	300
	DO 50 I=1,NP	POL	310
50	CALL SYMBOL (X(I),Y(I),.1,2,0.,-1)	POL	320
	CALL PLOT (0.0,0.0,3)	POL	330
C	PLOT OF F(CALC) WILL COME HERE	POL	340
	I=IA	POL	350
	CALL LINE (X(I-NP),Y(I),NP,1)	POL	360
	CALL PLOT (120.0,-1.0,-3)	POL	370
	RETURN	POL	380
	END	POL	390
	SUBROUTINE INTENS (IDENT,NAME,BGD,STX,NC,NP,AA,SLOP,YCAL,OMIT)	INTS	10
C	INTENSITY CALC	INTS	20

C	BGD IS THE COUNTER BACKGROUND FOR SUBTRACTION	INTS 30
C	STX IS THE NUMBER OF CHANNELS BETWEEN ZERO CHANNEL AND	INTS 40
C	FIRST DATA CHANNEL FOR CALC TAU	INTS 50
C	NC IS THE NUMBER OF COMPONENTS	INTS 60
C	NP IS THE NUMBER OF DATA PTS USED IN CALC TAU FOR COMP	INTS 70
C	AA IS THE COMP INTERCEPT VALUES AT ZERO TIME	INTS 80
C	SLOP IS THE COMPONENT SLOPES	INTS 90
C	YCAL IS THE THEORETICAL VALUES FOR NP DATA POINTS	INTS 100
C	YEXP IS THE EXPERIMENTAL VALUES FOR NEP DATA PTS ON	INTS 110
C	LEFT SIDE OF SPECTRUM FOR DETERMING TOTAL AREA UNDER CURVE	INTS 120
	DIMENSION NAME(11), AA(5), SLOP(5), YCAL(512)	INTS 130
	DIMENSION YEXP(200), XX(512), AREXP(5), TEN(5), AT(5)	INTS 140
	INTEGER OMIT	INTS 150
	READ (5,280) NEP,CONTR	INTS 160
C	NEP IS THE NO. OF ADDITIONAL EXPERIMENTAL GRAPH PTS NEEDED	INTS 170
C	TO CALC COMPONENT INTENSITIES	INTS 180
	IF (CONTR.NE.0) GO TO 10	INTS 190
	READ (5,290) (YEXP(I),I=1,NEP)	INTS 200
	GO TO 20	INTS 210
10	READ (5,300) (YEXP(I),I=1,NEP)	INTS 220
20	DO 30 I=1,NEP	INTS 230
	YEXP(I)=YEXP(I)-BGD	INTS 240
30	CONTINUE	INTS 250
	IE=I	INTS 260
C	PTST TEST	INTS 270
	XX(1)=STX	INTS 280
	DO 40 I=2,NP	INTS 290
	XX(I)=XX(I-1)+1.0	INTS 300
40	CONTINUE	INTS 310
	L=0	INTS 320
50	L=L+1	INTS 330
	DT=(-SLOP(NC)*XX(L))	INTS 340

	P=AA(NC)*EXP(DT)	INTS 350
	PTST=YCAL(L)-P	INTS 360
	IF (PTST.GT.5.0) GO TO 50	INTS 370
	PYS=YCAL(L)	INTS 380
	PXS=XX(L)	INTS 390
	J=IE+L	INTS 400
C	CHANGE OF SUBSCRIPTS TO UNIFY THE DATA ORDER	INTS 410
	DO 60 I=1,L	INTS 420
	IE=IE+1	INTS 430
	TEMP=YCAL(I)	INTS 440
	YEXP(IE)=TEMP	INTS 450
60	CONTINUE	INTS 460
C	OUTPUT 1	INTS 470
	WRITE (6,310) IDENT,NAME	INTS 480
	WRITE (6,320) NP,NC,BGD,STX	INTS 490
	WRITE (6,330) (I,AA(I),I,SLOP(I),I=1,NC)	INTS 500
70	GO TO (80,90,80,90), OMIT	INTS 510
80	WRITE (6,340)	INTS 520
	WRITE (6,350) (I,YEXP(I),I=1,J)	INTS 530
C	TOTAL AREA UNDER CURVE(ARSM)	INTS 540
C	ADJUSTMENT OF THE INDEX J	INTS 550
90	ODD=0.0	INTS 560
	EVEN=0.0	INTS 570
	KJ=0	INTS 580
	KI=0	INTS 590
	NEV=J-1	INTS 600
	DO 100 I=2,NEV,2	INTS 610
	KJ=KJ+1	INTS 620
	EVEN=EVEN+YEXP(I)	INTS 630
100	CONTINUE	INTS 640
	NOD=J-2	INTS 650
	DO 110 I=3,NOD,2	INTS 660

	KI=KI+1	INTS 670
	ODD=ODD+YEXP(I)	INTS 680
110	CONTINUE	INTS 690
C	CALCULATION OF AREA	INTS 700
	JK=KJ+KI+2	INTS 710
	IF (JK-J) 120,140,120	INTS 720
120	J=J-1	INTS 730
	DO 130 JK=1,J	INTS 740
	YEXP(JK)=YEXP(JK+1)	INTS 750
130	CONTINUE	INTS 760
	GO TO 70	INTS 770
140	ARSM=0.33333*1.0*(YEXP(1)+2.0*EVEN+4.0*ODD+YEXP(J))	INTS 780
	TA=(-SLOP(NC)*PXS)	INTS 790
	ARTA=(AA(NC)*EXP(TA))/(SLOP(NC))	INTS 800
	ARSM=ARSM+ARTA	INTS 810
	WRITE (6,360) ARSM,ARTA,J,PTST,PXS,PYS	INTS 820
C	SUM OF COMPONENT AREAS(SUM)	INTS 830
	SUM=0.0	INTS 840
	DO 150 I=1,NC	INTS 850
	AREXP(I)=AA(I)/(SLOP(I))	INTS 860
	SUM=SUM+AREXP(I)	INTS 870
150	CONTINUE	INTS 880
	DO 160 I=1,NC	INTS 890
	TEN(I)=AREXP(I)/SUM	INTS 900
160	CONTINUE	INTS 910
	WRITE (6,340)	INTS 920
	WRITE (6,370) SUM	INTS 930
	WRITE (6,380) (I,TEN(I),I,AREXP(I),I,AA(I),I=1,NC)	INTS 940
C	OPTIMIZATION OF ZERO LINE	INTS 950
	T=0.0	INTS 960
	ADD=0.1	INTS 970
170	SUM=0.0	INTS 980

	T=T+ADD	INTS 990
	IF (T.GT.STX) GO TO 270	INTS1000
	DO 180 I=1,NC	INTS1010
	AT(I)=AA(I)*EXP(-SLOP(I)*T)	INTS1020
	AREXP(I)=AT(I)/SLOP(I)	INTS1030
	SUM=SUM+AREXP(I)	INTS1040
180	CONTINUE	INTS1050
	IF (SUM-ARSM) 190,240,170	INTS1060
190	IF (ABS((SUM-ARSM)/ARSM).LT.0.01) GO TO 240	INTS1070
	T=T-ADD	INTS1080
	ADD=ADD/10.0	INTS1090
	IF (ADD.LE.0.0001) GO TO 200	INTS1100
	GO TO 170	INTS1110
200	T=0.0	INTS1120
	SUB=-0.1	INTS1130
210	SUM=0.0	INTS1140
	T=T+SUB	INTS1150
	DO 220 I=1,NC	INTS1160
	AT(I)=AA(I)*EXP(-SLOP(I)*T)	INTS1170
	AREXP(I)=AT(I)/SLOP(I)	INTS1180
	SUM=SUM+AREXP(I)	INTS1190
220	CONTINUE	INTS1200
	IF (SUM-ARSM) 210,240,230	INTS1210
230	IF (ABS((SUM-ARSM)/ARSM).LT.0.01) GO TO 240	INTS1220
	T=T-SUB	INTS1230
	SUB=SUB/10.0	INTS1240
	IF (SUB.LE.0.0001) GO TO 240	INTS1250
240	DO 250 I=1,NC	INTS1260
	TEN(I)=AREXP(I)/SUM	INTS1270
250	CONTINUE	INTS1280
C	OUTPUT 2	INTS1290
	WRITE (6,340)	INTS1300

	WRITE (6,390) SUM,T	INTS1310	
	WRITE (6,400) (I,TEN(I),I,AREXP(I),I,AT(I),I=1,NC)	INTS1320	
	WRITE (6,410)	INTS1330	
	DO 260 I=1,NC	INTS1340	
260	TEN(I)=TEN(I)*100.0	INTS1350	
	WRITE (6,420) (TEN(I),I=1,NC)	INTS1360	
270	RETURN	INTS1370	
C		INTS1380	
280	FORMAT (2I5)	INTS1390	
290	FORMAT (F6.0,9F7.0)	INTS1400	
300	FORMAT (F9.0,7F10.0)	INTS1410	
310	FORMAT (1H1,3X,A4,11A4)	INTS1420	
320	FORMAT (//5H NP=,I3,5H NC=,I2,6H BGD=,F10.5,6H STX=,F5.2//)	INTS1430	
330	FORMAT (5X,2HA(,I2,2H)=,F15.5,7H SLOP(,I2,2H)=,F12.8)	INTS1440	
340	FORMAT (1H0)	INTS1450	
350	FORMAT (9H YEXP(,I2,2H)=,F15.5,9H YEXP(,I2,2H)=,F15.5,9H	INTS1460	
	1YEXP(,I2,2H)=,F15.5,9H YEXP(,I2,2H)=,F15.5)	INTS1470	
360	FORMAT (//10X,5HARSM=,E15.8,7H ARTA=,E15.8,4H J=,I3,7H PTST=,F1	INTS1480	
	10.5,6H PXS=,F8.4,6H PYS=,F10.5)	INTS1490	
370	FORMAT (15X,4HSUM=,E15.8)	INTS1500	
380	FORMAT (10X,2HI(,I2,2H)=,E15.8,5X,8H AREXP(,I2,2H)=,E15.8,5X,4H	INTS1510	
	1A(,I2,2H)=,F12.5)	INTS1520	
390	FORMAT (15X,4HSUM=,E15.8,4H T=,F10.5)	INTS1530	
400	FORMAT (10X,2HI(,I2,2H)=,E15.8,5X,8H AREXP(,I2,2H)=,E15.8,5X,5H	INTS1540	
	1AT(,I2,2H)=,F12.5)	INTS1550	
410	FORMAT (////111H	1ST COMP	INTS1560
	1 2ND COMP 3RD COMP 4TH COMP 5TH COMP)	INTS1570	
420	FORMAT (//5X,15HINTENSITY (%),9X,5F16.3)	INTS1580	
	END	INTS1590	

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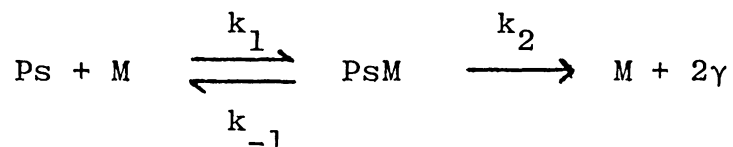
THE REACTION OF ORTHO-POSITRONIUM
WITH NITROAROMATICS

by

William Juul Madia

(ABSTRACT)

A study was made to ascertain the mechanism by which ortho-positronium reacts with nitroaromatics in solution. These reactions are two to three orders of magnitude faster than expected by simple pickoff. Conjugation of the nitro group with the aromatic ring was found to be a necessary condition for this behavior. The rate constant for each of these reactions was measured in suitable solvents over a wide temperature range, 200-500°K. Distinct departure from the Arrhenius law was found. In general, the rate constant increased linearly with increasing temperature up to a point, and then decreased linearly as the temperature was raised further. The following mechanism is postulated:



Ortho-Positronium reacts with the nitroaromatic, M, to form a complex in a reversible step. The complex may either decompose into a positronium atom and nitroaromatic molecule or it may go on to annihilate the positron. A steady-state

concentration of the complex is assumed and the observed rate constant for the process is

$$k_{\text{obs}} = \frac{k_1 k_2}{k_{-1} + k_2}$$

The two limiting cases are: (1) $k_2 \gg k_{-1}$ causing $k_{\text{obs}} \approx k_1$. Typical Arrhenius behavior is expected and observed in this region, and (2) $k_{-1} \gg k_2$ causing $k_{\text{obs}} \approx K_{\text{EQ}} k_2$. In this region $H_{\text{EQ}} + E_a(2) < 0$ and the observed rate constant decreases with increasing temperature. This corresponds to a stable but temperature sensitive PSM complex.

An approximate molecular orbital study has been made to study the possible existence of positron and positronium complexes with a select number of organic molecules. The CNDO/2 approximation is used for all electronic integrals and for electron positron Coulomb integrals. The core Hamiltonian matrix elements for the positron were estimated by combining the Wolfsberg-Helmholtz and Cusachs-Cusachs approximations. This approach shows that the positron should become bound to the molecules. Using the same criteria, the formation of stable positronium complexes is more improbable. The calculated binding energies are strongly dependent on the value of the Wolfsberg-Helmholtz proportionality constant for positrons.